

Conformal invariance of spin correlations in the planar Ising model

By DMITRY CHELKAK, CLÉMENT HONGLER, and KONSTANTIN IZYUROV

Abstract

We rigorously prove the existence and the conformal invariance of scaling limits of the magnetization and multi-point spin correlations in the critical Ising model on arbitrary simply connected planar domains. This solves a number of conjectures coming from the physical and the mathematical literature. The proof relies on convergence results for discrete holomorphic spinor observables and probabilistic techniques.

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1. Introduction

The Ising model plays a central role in equilibrium statistical mechanics, being a standard example of an order-disorder phase transition in dimensions two and above. Besides pure mathematical interest, it has found successful applications to several fields in theoretical physics and computer sciences.

The phase transition in the Ising model in two dimensions has been a subject of extensive study, both in the mathematics and in the physics literature. The value of the critical temperature on the square lattice was determined by Kramers and Wannier [KW41]. Onsager [Ons44] computed the free energy and the critical exponents of the model. Later on, many exact computations were carried out — in particular, by McCoy and Wu [MW73].

Further, a gradual understanding of the Ising model at criticality led to the conjecture by Belavin, Polyakov and Zamolodchikov that its scaling limit (as well as scaling limits of other critical models) should be conformally invariant and described by Conformal Field Theory (CFT) [BPZ84a], [BPZ84b]. Loosely speaking, this conjecture can be formulated as follows: for any conformal map $\varphi : \Omega \rightarrow \Omega'$,

(scaling limit of the model on Ω') = φ (scaling limit of the model on Ω).

In particular, if Ω_δ are discrete approximations to a continuous planar domain Ω , then various quantities (expectations, probabilities etc.) in the model, under a proper normalization, have conformally invariant or covariant limits as the mesh size δ tends to zero. Moreover, Conformal Field Theory predicts exact formulae for these limits.

In the *full-plane* case, many results were obtained following the seminal works of Onsager and Kaufman in late 1940's. The Onsager's formula for the spontaneous magnetization was proven by Yang [Yan52]. The diagonal and horizontal spin-spin correlations were explicitly computed by Wu [MW73]. A number of remarkable results were obtained for the massive limits; see [WMTB76], [SMJ80], [PT83] and references therein. Palmer [Pal07] justified the CFT predictions at criticality by taking the zero-mass limit. At criticality, the full-plane energy correlation functions (that is, the correlations of n pairs of neighboring spins) were computed on periodic isoradial graphs by Boutillier and de Tilière [BdT10], [BdT11], and the $2n$ -point full-plane spin correlation functions were treated by Dubédat, combining exact bosonization techniques [Dub11b] and results on monomer correlations in the dimer model [Dub11a].

However, the group of conformal self-maps of the full plane is only finite dimensional. Hence, in order to reveal the full strength of the conformal invariance property, it is important to consider *general planar domains with boundary*. In this setting, mathematical proofs of conformal invariance and CFT predictions at criticality have remained out of reach until recently. Smirnov [Sm06] has rigorously established conformal covariance of the *fermionic observables* in the Ising model on the square grid. Later, this result has been proven to be universal in the family of isoradial graphs [CS12], and it led to the proof of convergence of the interfaces in the Ising model to Schramm’s SLE₃ curves [CDCH⁺14]. At the same time, the scaling limit of *energy correlations* for bounded domains has been rigorously treated in [HS13], [Hon10], confirming the CFT predictions for the energy field. Nevertheless, the corresponding question about the *spin correlations* remained open.

In this paper, we rigorously prove the existence and the conformal covariance of the scaling limits of all the multi-point spin correlation functions in any simply connected planar domain with + boundary conditions at criticality.

Our main result (see [Theorem 1.2](#)) reads as follows. Let $\mathbb{E}_{\Omega_\delta}^+[\sigma_{a_1} \cdots \sigma_{a_n}]$ denote the correlation of spins at the sites a_1, \dots, a_n with + boundary conditions in a discrete domain Ω_δ of mesh size δ . Then

$$\delta^{-\frac{n}{8}} \mathbb{E}_{\Omega_\delta}^+[\sigma_{a_1} \cdots \sigma_{a_n}] \rightarrow \mathcal{C}^n \cdot \langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_\Omega^+$$

as Ω_δ approximates Ω and δ tends to zero. Here \mathcal{C} is an explicit lattice-dependent constant and $\langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_\Omega^+$ is an explicit conformally covariant tensor of degree $\frac{1}{8}$ with respect to each of the variables; see (1.1), (1.4) and (1.5).

For example, in the case of the magnetization (the expectation of a single spin), one gets

$$\delta^{-\frac{1}{8}} \mathbb{E}_{\Omega_\delta}^+[\sigma_a] \rightarrow \mathcal{C} \cdot 2^{\frac{1}{4}} \text{rad}^{-\frac{1}{8}}(a, \Omega),$$

where $\text{rad}(a, \Omega)$ denotes the conformal radius of Ω as seen from a , in other words, $\text{rad}(a, \Omega) = |\varphi'(0)|$, where φ is a conformal map from the disc $\{z \in \mathbb{C} : |z| < 1\}$ to Ω mapping the origin to a . A sketch of the proof of this result can also be found in the ICMP2012 proceedings [Hon14].

In the case of free boundary conditions, we establish a similar convergence result for the two-point function (see [Theorem 1.1](#)). In that framework, the convergence of the corresponding n -point correlations (which by symmetry are only nonzero for even n) can be obtained by our methods as well, but for conciseness it is not included in the present paper.

Our previous results [CI13] immediately allow one to treat alternating $+/-$ boundary conditions (see [Corollary 1.3](#)). The technique we use also applies to *mixed correlations*, involving spin, energy, disorder and boundary change operators, and extends to *multiply connected domains*, which will be worked out in a subsequent paper.

The explicit formulae for what we prove to be the scaling limits of $\mathbb{E}_{\Omega_\delta}^+[\sigma_{a_1} \cdots \sigma_{a_n}]$ were predicted by Conformal Field Theory methods in a number of papers originating in the seminal work [BPZ84a]. In [Car84], it was explained how to handle the half-plane case by CFT means; in particular, the two-point correlations were treated. This result was later extended to $n = 3$ [BG87] and extrapolated to larger n [BG93].

Our method is based on the extraction of information from some discrete holomorphic observables in bounded domains by means of discrete complex analysis — the approach that was firstly implemented in [Smi06], [Smi10] for the basic fermionic observables. More precisely, we use the spinor version of those that was introduced in [CI13]. Fermionic observables per se essentially go back to the Kaufman-Onsager considerations and can be written as a product $\psi_z = \sigma_z \mu_z$ of spin and disorder operators in the notation of [KC71]. The correlators $\langle \psi_z \sigma_{a_1} \mu_{a_2} \dots \mu_{a_n} \rangle$ can be found in the works of Kyoto's school [SMJ77], [SMJ79a], [SMJ79b], [SMJ80]. However, rigorous proofs of convergence results require a delicate analysis of some Riemann-type boundary value problems for discrete holomorphic functions. The analysis thereof was initiated more recently [Smi06], [Smi10], [CS12].

Simultaneously and independently of our work, Dubédat announced analogous results for $2n$ -point spin correlations in bounded domains Ω via the exact bosonization approach [Dub11b], and Camia, Garban and Newman obtained some results [CGN15] about the properly renormalized spin field seen as a random generalized function (i.e., Schwartz distribution) on Ω .

1.1. *Main results.* The Ising model on a graph \mathcal{G} is a random assignment of ± 1 spins to the vertices of \mathcal{G} . In our paper we prefer a dual setup and consider the model on the *faces* $\mathcal{F} = \mathcal{V}_{\Omega_\delta}^0$ of lattice approximations Ω_δ to a bounded simply connected planar domain $\Omega \subset \mathbb{C}$. The probability of a spin configuration $\sigma \in \{\pm 1\}^{\mathcal{F}}$ is proportional to $e^{-\beta \mathbf{H}(\sigma)}$, where $\beta > 0$ is the inverse temperature and

$$\mathbf{H}(\sigma) := - \sum_{x \sim y} \sigma_x \sigma_y$$

is the energy of the configuration. More precisely, we will work with discrete planar domains Ω_δ that are subsets of the *square grid* rotated by 45° of diagonal mesh size 2δ . (The distance between adjacent spins is thus $\sqrt{2}\delta$; see Figure 1.)

From now on we only consider the model at its *critical point*, which for the square grid corresponds to the parameter value $\beta_c = \frac{1}{2} \ln(\sqrt{2} + 1)$. We also introduce a (lattice-dependent) constant that will appear in the statements of our theorems:

$$(1.1) \quad \mathcal{C} = 2^{\frac{1}{8}} e^{-\frac{3}{2} \zeta'(-1)},$$

where ζ' denotes the derivative of Riemann's zeta function.

Below we follow the CFT notation $\langle \dots \rangle$ to represent the predicted formulae for the scaling limits of correlations on a simply connected domain Ω :

we define $\langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_{\Omega}^{\mathbf{b}}$ as some *explicit functions* of the points $a_1, \dots, a_n \in \Omega$ that may also depend on Ω and boundary conditions \mathbf{b} as parameters. In particular, we define the one-point and the two-point functions $\langle \sigma_a \rangle_{\Omega}^+$ and $\langle \sigma_a \sigma_b \rangle_{\Omega}^+$, $\langle \sigma_a \sigma_b \rangle_{\Omega}^{\text{free}}$ on the upper half-plane \mathbb{H} by the formulae

$$(1.2) \quad \langle \sigma_a \rangle_{\mathbb{H}}^+ = \frac{2^{\frac{1}{4}}}{(2\Im a)^{\frac{1}{8}}}, \quad \langle \sigma_a \sigma_b \rangle_{\mathbb{H}}^+ = \langle \sigma_a \rangle_{\mathbb{H}}^+ \langle \sigma_b \rangle_{\mathbb{H}}^+ \cdot \left[\frac{1}{2} \left(\left| \frac{b - \bar{a}}{b - a} \right|^{\frac{1}{2}} + \left| \frac{b - a}{b - \bar{a}} \right|^{\frac{1}{2}} \right) \right]^{\frac{1}{2}},$$

$$\langle \sigma_a \sigma_b \rangle_{\mathbb{H}}^{\text{free}} = \langle \sigma_a \rangle_{\mathbb{H}}^+ \langle \sigma_b \rangle_{\mathbb{H}}^+ \cdot \left[\frac{1}{2} \left(\left| \frac{b - \bar{a}}{b - a} \right|^{\frac{1}{2}} - \left| \frac{b - a}{b - \bar{a}} \right|^{\frac{1}{2}} \right) \right]^{\frac{1}{2}}$$

and on all other simply-connected domains $\Omega \neq \mathbb{C}$ by the condition that, for any conformal mapping $\varphi : \Omega \rightarrow \Omega'$, one has

$$(1.3) \quad \langle \sigma_a \rangle_{\Omega}^+ = \langle \sigma_{\varphi(a)} \rangle_{\Omega'}^+ \cdot |\varphi'(a)|^{\frac{1}{8}}, \quad \langle \sigma_a \sigma_b \rangle_{\Omega}^+ = \langle \sigma_{\varphi(a)} \sigma_{\varphi(b)} \rangle_{\Omega'} \cdot |\varphi'(a)|^{\frac{1}{8}} |\varphi'(b)|^{\frac{1}{8}},$$

$$\langle \sigma_a \sigma_b \rangle_{\Omega}^{\text{free}} = \langle \sigma_{\varphi(a)} \sigma_{\varphi(b)} \rangle_{\Omega'} \cdot |\varphi'(a)|^{\frac{1}{8}} |\varphi'(b)|^{\frac{1}{8}}.$$

Note, in particular, that this definition is consistent, as (1.3) is satisfied when $\Omega = \Omega' = \mathbb{H}$ and φ is a Möbius map since the ratio $|b - \bar{a}|/|b - a|$ is Möbius invariant. Therefore, due to the Riemann mapping theorem, the relations (1.2), (1.3) properly define $\langle \sigma_a \sigma_b \rangle_{\Omega}^+$ and $\langle \sigma_a \sigma_b \rangle_{\Omega}^{\text{free}}$ for any simply-connected domain $\Omega \neq \mathbb{C}$. Equivalent formulae for the one- and two-point functions can be given in terms of the hyperbolic distance $d_{\Omega}(a, b)$ and the conformal radius $\text{rad}(a, \Omega)$:

$$\langle \sigma_a \rangle_{\Omega}^+ = 2^{\frac{1}{4}} \text{rad}^{-\frac{1}{8}}(a, \Omega), \quad \langle \sigma_a \sigma_b \rangle_{\Omega}^+ = \langle \sigma_a \rangle_{\Omega}^+ \langle \sigma_b \rangle_{\Omega}^+ \cdot (1 - \exp(-2d_{\Omega}(a, b)))^{-\frac{1}{4}},$$

$$\langle \sigma_a \sigma_b \rangle_{\Omega}^{\text{free}} = \langle \sigma_a \sigma_b \rangle_{\Omega}^+ \cdot \exp\left(-\frac{1}{2}d_{\Omega}(a, b)\right).$$

We have the following convergence theorem for the two-point functions, both with + (the spins on the boundary of Ω_{δ} are set to +) and free (no restrictions are set for boundary spins) boundary conditions.

THEOREM 1.1. *Let Ω be a bounded simply connected domain, and for $\delta > 0$, let Ω_{δ} be a discretization of Ω by the square grid of diagonal mesh size 2δ . Then, for any $\epsilon > 0$, we have*

$$\delta^{-\frac{1}{4}} \mathbb{E}_{\Omega_{\delta}}^+ [\sigma_a \sigma_b] \rightarrow \mathcal{C}^2 \cdot \langle \sigma_a \sigma_b \rangle_{\Omega}^+ \quad \text{and} \quad \delta^{-\frac{1}{4}} \mathbb{E}_{\Omega_{\delta}}^{\text{free}} [\sigma_a \sigma_b] \rightarrow \mathcal{C}^2 \cdot \langle \sigma_a \sigma_b \rangle_{\Omega}^{\text{free}},$$

as $\delta \rightarrow 0$, uniformly over all $a, b \in \Omega$ at distance at least ϵ from $\partial\Omega$ and from each other, where the constant \mathcal{C} is given by (1.1).

Our result generalizes to multi-point correlations as follows. As for the two-point functions, define the continuous correlation functions $\langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_{\Omega}^+$

in the upper half-plane \mathbb{H} by the explicit formula

$$(1.4) \quad \langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_{\mathbb{H}}^+ = \prod_{k=1}^n \frac{1}{(2\Im a_k)^{\frac{1}{8}}} \cdot \left(2^{-\frac{n}{2}} \sum_{\mu \in \{\pm 1\}^n} \prod_{1 \leq k < m \leq n} \left| \frac{a_k - a_m}{a_k - \bar{a}_m} \right|^{\frac{\mu_k \mu_m}{2}} \right)^{\frac{1}{2}}$$

and on other domains by the condition that for any conformal map $\varphi : \Omega \rightarrow \Omega'$, we have

$$(1.5) \quad \langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_{\Omega}^+ = \langle \sigma_{\varphi(a_1)} \cdots \sigma_{\varphi(a_n)} \rangle_{\Omega'}^+ \cdot \prod_{k=1}^n |\varphi'(a_k)|^{\frac{1}{8}}.$$

Note as before that the explicit functions (1.4) are covariant under Möbius maps and thus the multi-point correlations are well defined.

THEOREM 1.2. *Let Ω be a bounded simply connected domain, and for $\delta > 0$, let Ω_δ be a discretization of Ω by the square grid of diagonal mesh size 2δ . Then, for any $\epsilon > 0$ and any $n = 1, 2, \dots$, we have*

$$\delta^{-\frac{n}{8}} \cdot \mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}] \rightarrow \mathcal{C}^n \cdot \langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_{\Omega}^+$$

as $\delta \rightarrow 0$, uniformly over all $a_1, \dots, a_n \in \Omega$ at distance at least ϵ from $\partial\Omega$ and from each other.

Using the results of [CI13], one immediately arrives at the following generalization. Let $\mathfrak{b}_\delta = \{b_1^\delta, \dots, b_{2m}^\delta\}$ be a collection of points on $\partial\Omega_\delta$ appearing in counterclockwise order and approximating a continuous collection $\mathfrak{b} = \{b_1, \dots, b_{2m}\} \subset \partial\Omega$. Denote by $\mathbb{E}_{\Omega_\delta}^{\mathfrak{b}_\delta}$ an expectation for the Ising model with $+$ boundary conditions on the counterclockwise arcs $[b_{2j-1}^\delta, b_{2j}^\delta]$ and $-$ boundary condition on the complementary arcs $[b_{2j}^\delta, b_{2j+1}^\delta]$, $j = 1, \dots, m$, where we set $b_{2m+1}^\delta := b_1^\delta$.

COROLLARY 1.3. *Suppose that the approximation Ω_δ of Ω is regular near the points of \mathfrak{b} in the sense of [CI13, Def. 3.14]. As $\delta \rightarrow 0$, one has*

$$\delta^{-\frac{n}{8}} \mathbb{E}_{\Omega_\delta}^{\mathfrak{b}_\delta} [\sigma_{a_1} \cdots \sigma_{a_n}] \rightarrow \mathcal{C}^n \cdot \langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_{\Omega}^{\mathfrak{b}},$$

where $\langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_{\Omega}^{\mathfrak{b}}$ is explicit and satisfies the conformal covariance property (1.5).

Proof. Write

$$\mathbb{E}_{\Omega_\delta}^{\mathfrak{b}_\delta} [\sigma_{a_1} \cdots \sigma_{a_n}] = (\mathbb{E}_{\Omega_\delta}^{\mathfrak{b}_\delta} [\sigma_{a_1} \cdots \sigma_{a_n}] / \mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}]) \cdot \mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}].$$

By [CI13, Cor. 5.10], the first term converges to an explicit conformally invariant limit. Thus the result follows from Theorem 1.2. \square

Remark 1.4. The condition of boundedness of Ω in Theorems 1.1 and 1.2 above is imposed for the sake of simplicity and can be removed with no essential changes in the proofs. In particular, one could take the bulk limit (i.e., $\Omega_\delta \rightarrow \mathbb{C}$

and $\delta \rightarrow 0$ while keeping the positions of a_1, \dots, a_n fixed), in which case [DFSZ87], [Dub11b]

$$\delta^{-\frac{n}{8}} \cdot \mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}] \rightarrow \mathcal{C}^n \cdot \langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_{\mathbb{C}},$$

where

$$(1.6) \quad \langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_{\mathbb{C}} = \left(2^{-\frac{n}{2}} \sum_{\substack{\mu \in \{\pm 1\}^n: \\ \mu_1 + \cdots + \mu_n = 0}} \prod_{1 \leq k < m \leq n} |a_k - a_m|^{\frac{\mu_k \mu_m}{2}} \right)^{\frac{1}{2}}.$$

(This formula can be seen as the limit of (1.4) when all a_k are replaced by $a_k + iy$ and $y \rightarrow +\infty$.) It is worth noting that, if a_1, \dots, a_n are kept at a *finite number of lattice steps* from each other, then other scaling exponents appear. For $n = 2$, the scaling limit of the energy densities $\lim_{\delta \rightarrow 0} \delta^{-1} (\mathbb{E}_{\Omega_\delta} [\sigma_a \sigma_{a'}] - \sqrt{2}/2)$, where two neighboring faces a, a' approximate the same point of Ω , has been treated in [HS13]. For $n > 2$, two terms of asymptotics have been obtained at [GHP13]. One can also wonder about the intermediate situation when a_1, \dots, a_n are at distances of order δ^β from each other for some fixed $0 < \beta < 1$. Then, the leading term in the asymptotic expansion of the discrete correlation functions has the order $\delta^{\frac{1-\beta}{8}n}$ and is provided by the bulk limit (1.6). A more sophisticated analysis is required in order to find the second term of asymptotics in this case.

1.2. *Key steps in the proof.* In this section we list the key results that allow us to prove Theorems 1.1 and 1.2. The first small step deals with the normalizing factors. It is a celebrated result of Wu [MW73] that in the unique infinite-volume limit of the critical planar Ising model (i.e., in the case $\Omega = \mathbb{C}$), one has the following asymptotics:

$$(1.7) \quad \mathbb{E}_{\mathbb{C}_\delta} [\sigma_{0_\delta} \sigma_{1_\delta}] \sim \mathcal{C}^2 \cdot \delta^{\frac{1}{4}}, \quad \delta \rightarrow 0,$$

where \mathbb{C}_δ denotes the square grid $\delta(1+i)\mathbb{Z}^2$, while 0_δ and 1_δ stand for proper approximations of the points $0, 1 \in \mathbb{C}$. (Keep in mind that our square lattice is rotated by 45° , so this is the *diagonal* spin-spin correlation.) Instead of deriving the correct normalization of spin correlations in bounded domains directly, we relate it to the behavior of the normalizing factor

$$\varrho(\delta) := \mathbb{E}_{\mathbb{C}_\delta} [\sigma_{0_\delta} \sigma_{1_\delta}].$$

Namely, we prove that, as $\delta \rightarrow 0$,

$$\begin{aligned} (\varrho(\delta))^{-\frac{n}{2}} \cdot \mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}] &\rightarrow \langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_{\Omega}^+, \\ (\varrho(\delta))^{-1} \cdot \mathbb{E}_{\Omega_\delta}^{\text{free}} [\sigma_a \sigma_b] &\rightarrow \langle \sigma_a \sigma_b \rangle_{\Omega}^{\text{free}}, \end{aligned}$$

which, combined with (1.7), readily gives Theorems 1.1 and 1.2. We point out that apart from this reduction, we never use (1.7) in this paper. On the other

hand, our methods also allow one to give a new proof of the explicit formula for the diagonal spin-spin correlations in the full-plane case as well as to derive explicit formulae for the magnetization in the half-plane; see the forthcoming work [CH].

The following theorem, concerning *discrete logarithmic derivatives* of the spin correlations with + boundary conditions, is a cornerstone for the whole paper.

THEOREM 1.5. *Let Ω be a bounded simply connected domain and Ω_δ be discretizations of Ω by the square grids $\delta(1+i)\mathbb{Z}^2$. Then, for any $\epsilon > 0$ and any $n = 1, 2, \dots$, we have*

$$(1.8) \quad \frac{1}{2\delta} \left(\frac{\mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1+2\delta}\sigma_{a_2} \cdots \sigma_{a_n}]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}]} - 1 \right) \rightarrow \Re \mathfrak{A}_\Omega(a_1, \dots, a_n),$$

$$(1.9) \quad \frac{1}{2\delta} \left(\frac{\mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1+2i\delta}\sigma_{a_2} \cdots \sigma_{a_n}]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}]} - 1 \right) \rightarrow -\Im \mathfrak{A}_\Omega(a_1, \dots, a_n)$$

as $\delta \rightarrow 0$, uniformly over all faces $a_1, \dots, a_n \in \Omega_\delta$ at distance at least ϵ from $\partial\Omega$ and from each other. The function $\mathfrak{A}_\Omega(a_1, \dots, a_n)$ is defined explicitly via the solution to some special interpolation problem (see further details in [Section 2.5](#)) and has the following covariance property under conformal mappings $\varphi : \Omega \rightarrow \Omega'$:

$$(1.10) \quad \mathfrak{A}_\Omega(a_1, \dots, a_n) = \mathfrak{A}_{\Omega'}(\varphi(a_1), \dots, \varphi(a_n)) \cdot \varphi'(a_1) + \frac{1}{8} \frac{\varphi''(a_1)}{\varphi'(a_1)}.$$

Proof. The proof is based on the convergence results for the discrete spinor observables. A rather delicate analysis is needed since we are interested in the values of observables near their singular points. See further details in [Sections 2](#) and [3](#). □

Integrating the result of [Theorem 1.5](#), we get the following weaker form of the convergence result for the spin correlations. Note that the conformal covariance degree $\frac{1}{8}$ in [\(1.5\)](#) is a direct consequence of the covariance rule [\(1.10\)](#) for $\mathfrak{A}_\Omega(a_1, \dots, a_n)$.

COROLLARY 1.6. *Under the conditions of [Theorem 1.5](#), for any $n \geq 1$, there exist some normalizing factors $\varrho_n(\delta, \Omega_\delta)$ that might depend on Ω_δ but not on the positions of the points a_1, \dots, a_n such that*

$$\mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}] \sim \varrho_n(\delta, \Omega_\delta) \cdot \langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_\Omega^+$$

as $\delta \rightarrow 0$, uniformly over all faces $a_1, \dots, a_n \in \Omega_\delta$ at distance at least ϵ from $\partial\Omega$ and from each other.

Proof. See [Section 2.7](#). □

We now focus on the special case $n = 2$. The next theorem is a crucial tool that allows us to compare the normalizing factors $\varrho_2(\delta, \Omega_\delta)$ with the full-plane case. We denote by $\mathbb{E}_{\Omega_\delta}^{\text{free}}$ the expectation for the critical Ising model defined on the *vertices* of Ω_δ (with free boundary conditions).

THEOREM 1.7. *Let Ω be a bounded simply connected domain and Ω_δ be a discretization of Ω by the square grid $\delta(1+i)\mathbb{Z}^2$. Then, for any $\epsilon > 0$, we have*

$$\frac{\mathbb{E}_{\Omega_\delta}^{\text{free}} [\sigma_{a+\delta}\sigma_{b+\delta}]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a\sigma_b]} \xrightarrow{\delta \rightarrow 0} \mathcal{B}_\Omega(a, b),$$

uniformly over all faces a, b at distance at least ϵ from $\partial\Omega$ and from each other, where $\mathcal{B}_\Omega(a, b)$ is a conformal invariant of (Ω, a, b) that can be explicitly written as $\mathcal{B}_\Omega(a, b) = \langle \sigma_a\sigma_b \rangle_\Omega^{\text{free}} / \langle \sigma_a\sigma_b \rangle_\Omega^+$.

Proof. The proof is based on the convergence results for the discrete spinor observables and the Kramers-Wannier duality; see further details in [Sections 2.4](#) and [2.6](#). □

Sketch of the proof of [Theorem 1.1](#). Having the results of [Corollary 1.6](#) and [Theorem 1.7](#), we only need to prove that $\varrho_2(\delta, \Omega_\delta) \sim \varrho(\delta)$ as $\delta \rightarrow 0$. The classical FKG inequality gives

$$\mathbb{E}_{\Omega_\delta}^{\text{free}} [\sigma_{a+\delta}\sigma_{b+\delta}] \leq \mathbb{E}_{\mathbb{C}_\delta} [\sigma_a\sigma_b] \leq \mathbb{E}_{\Omega_\delta}^+ [\sigma_a\sigma_b],$$

and it is easy to see that $\mathcal{B}_\Omega(a, b) \rightarrow 1$ as b approaches a . Since the normalizing factors $\varrho_2(\delta, \Omega_\delta)$ do not depend on the positions of $a, b \in \Omega$, we conclude that $\mathbb{E}_{\Omega_\delta}^+ [\sigma_a\sigma_b] \sim \mathbb{E}_{\mathbb{C}_\delta} [\sigma_a\sigma_b]$ in the double limit when we first let $\delta \rightarrow 0$ and then $b \rightarrow a$. This relates $\varrho_2(\delta, \Omega_\delta)$ to the full-plane normalization; see details in [Section 2.8](#). □

Sketch of the proof of [Theorem 1.2](#). Once the asymptotics $\varrho_2(\delta, \Omega_\delta) \sim \varrho(\delta)$ as $\delta \rightarrow 0$ is established, we derive the asymptotics of all the other $\varrho_n(\delta, \Omega_\delta)$, $n \neq 2$, using the following observation: as the point a_1 approaches $\partial\Omega$, the continuous correlation functions behave in the following way:

$$(1.11) \quad \langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_\Omega^+ \sim \langle \sigma_{a_1} \rangle_\Omega^+ \langle \sigma_{a_2} \cdots \sigma_{a_n} \rangle_\Omega^+,$$

and the same decorrelation result $\mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}] \sim \mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1}] \mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_2} \cdots \sigma_{a_n}]$ holds true in the double limit $\delta \rightarrow 0$ and $a_1 \rightarrow \partial\Omega$. This implies the recurrent formula $\varrho_{n+1}(\delta, \Omega_\delta) \sim \varrho_1(\delta, \Omega_\delta)\varrho_n(\delta, \Omega_\delta)$ for $n = 1, 2, \dots$ and, further, $\varrho_n(\delta, \Omega_\delta) \sim (\varrho(\delta))^{n/2}$ for all n ; see further details in [Section 2.9](#). □

1.3. Organization of the paper. [Section 2](#) contains all the main ideas. Details, especially those involving hard discrete complex analysis techniques, are mostly postponed to [Section 3](#). The readers not interested in these details may restrict themselves to [Section 2](#) only.

We fix the notation in [Section 2.1](#). The main tool of this paper, the discrete holomorphic spinor observables, is introduced and discussed in [Sections 2.2](#) and [2.3](#). In [Section 2.4](#), we prove that the ratios of spin correlations that appear in [Theorems 1.5](#) and [1.7](#) can be expressed in terms of these observables. In [Sections 2.5](#) and [2.6](#), we discuss the continuous counterparts of the discrete observables and state the convergence [Theorems 2.16](#), [2.18](#) and [2.20](#), which easily imply [Theorems 1.5](#) and [1.7](#). We do some explicit computations and prove [Corollary 1.6](#) in [Section 2.7](#). We complete the proofs of [Theorems 1.1](#) and [1.2](#) in [Sections 2.8](#) and [2.9](#), respectively.

[Section 3](#) is devoted to the proof of [Theorems 2.16](#), [2.18](#) and [2.20](#). We discuss the discrete properties of our observables and their full-plane analogue in [Sections 3.1–3.3](#), and finish the proof of the main convergence theorems in [Sections 3.4](#) and [3.5](#). In the appendix, we explicitly compute $\mathcal{A}_{\mathbb{H}}(a_1, \dots, a_n)$ for arbitrary n , which completes the proof of the formula [\(1.4\)](#).

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2. Holomorphic spinors and correlation functions

2.1. Notation. We start by fixing the notation that is used throughout the paper.

2.1.1. Graph notation. Recall that we work on the square grid rotated by 45° :

$$\mathbb{C}_\delta := \{\delta(1+i)(m+in) : m, n \in \mathbb{Z}\}.$$

The mesh size δ is hence the size of a half-diagonal of a square face. We will often identify the vertices of \mathbb{C}_δ with the corresponding complex numbers, the faces of \mathbb{C}_δ with their centers, the edges of \mathbb{C}_δ with their midpoints, etc.

By a *discrete domain of mesh size δ* we mean a union of grid faces. (See also [Figure 1](#) for notation given below.) We say that Ω_δ is *simply connected* if the corresponding polygonal domain is simply connected.

- We denote by $\text{Int}\mathcal{V}_{\Omega_\delta}^\circ$ the set of all *faces* belonging to Ω_δ (which are identified with their centers), by $\text{Int}\mathcal{V}_{\Omega_\delta}^\bullet$ the set of all *vertices* incident to these faces, and by $\text{Int}\mathcal{V}_{\Omega_\delta}^m$ the set of all *edges* incident to $\mathcal{V}_{\Omega_\delta}^\circ$, which are identified with their midpoints (or the vertices of a medial lattice).

In order to simplify the presentation, we also assume that $\mathbb{C}_\delta \setminus \Omega_\delta$ has *no fiords of a single face width*; i.e., all the edges joining vertices from $\text{Int}\mathcal{V}_{\Omega_\delta}^\bullet$ belong to $\text{Int}\mathcal{V}_{\Omega_\delta}^m$. This technical assumption can be easily relaxed, if necessary.

- We denote by $\partial\mathcal{V}_{\Omega_\delta}^\circ$, $\partial\mathcal{V}_{\Omega_\delta}^\bullet$ and $\partial\mathcal{V}_{\Omega_\delta}^m$ the sets of *boundary faces, vertices and edges*, i.e., those faces, vertices and edges that are incident but do not belong to $\text{Int}\mathcal{V}_{\Omega_\delta}^\circ$, $\text{Int}\mathcal{V}_{\Omega_\delta}^\bullet$ and $\text{Int}\mathcal{V}_{\Omega_\delta}^m$, respectively (see [Figure 1](#)).
- We set $\mathcal{V}_{\Omega_\delta}^\circ := \text{Int}\mathcal{V}_{\Omega_\delta}^\circ \cup \partial\mathcal{V}_{\Omega_\delta}^\circ$ etc.

Below we also need to work with four *corners* of a given square face separately. (In particular, see [Definition 2.1](#) below.) For a given vertex $v \in \mathbb{C}_\delta$, we identify the nearby corners with the points $v \pm \frac{1}{2}\delta$ and $v \pm \frac{1}{2}\delta i$ on the complex plane.

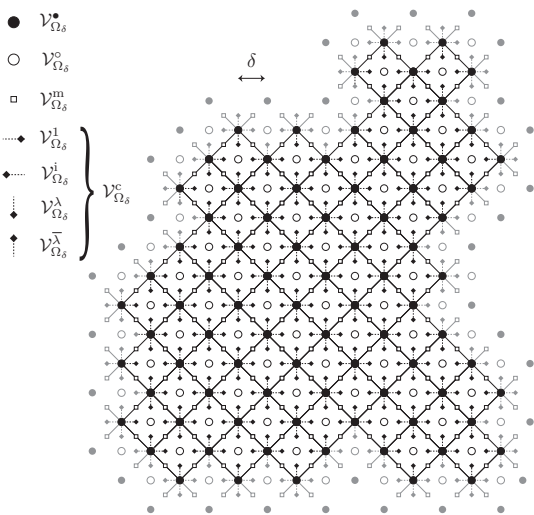


Figure 1. An example of a discrete domain Ω_δ and notation for the sets of vertices ($\mathcal{V}_{\Omega_\delta}^\bullet$), faces ($\mathcal{V}_{\Omega_\delta}^\circ$), edge midpoints ($\mathcal{V}_{\Omega_\delta}^m$) and (four types) of corners ($\mathcal{V}_{\Omega_\delta}^c = \mathcal{V}_{\Omega_\delta}^1 \cup \mathcal{V}_{\Omega_\delta}^i \cup \mathcal{V}_{\Omega_\delta}^\lambda \cup \mathcal{V}_{\Omega_\delta}^{\bar{\lambda}}$). The mesh size δ is a half-diagonal of a square face, thus the distance between adjacent spins is $\sqrt{2}\delta$. The inner vertices, faces, edges and corners are colored black, while the boundary ones are colored gray.

- We denote by $\mathcal{V}_{\Omega_\delta}^c$ the set of all corners incident to the vertices of $\text{Int}\mathcal{V}_{\Omega_\delta}^\bullet$. We also set $\mathcal{V}_{\Omega_\delta}^{\text{cm}} := \mathcal{V}_{\Omega_\delta}^c \cup \mathcal{V}_{\Omega_\delta}^{\text{m}}$.
- We use the notation $x \sim y$ if each of x, y is either a vertex, a face, an edge or a corner, and they are adjacent or incident to each other.

2.1.2. *Double covers.* In this paper we often deal with holomorphic functions (both discrete and continuous) that are defined on a double cover of a planar domain Ω and with opposite signs on opposite sheets. (That is, they have a -1 multiplicative monodromy around branching points.) We call such functions *holomorphic spinors*. The following notation will be used below:

- For a planar domain Ω and $a \in \Omega$, we denote by $[\Omega, a]$ the double cover of $\Omega \setminus \{a\}$ branching around a . All such double covers are naturally viewed as subdomains of $[\mathbb{C}, a]$. We often identify points on a double cover with their base points (so each base point is identified with the two points above it).
- For several marked points $a_1, \dots, a_n \in \Omega$, we denote by $[\Omega, a_1, \dots, a_n]$ the double cover of $\Omega \setminus \{a_1, \dots, a_n\}$ that branches around each of a_1, \dots, a_n .
- We will often compare spinors defined on $[\Omega, a_1, \dots, a_n]$ with those defined on $[\mathbb{C}, a_1]$ (e.g., with $1/\sqrt{z - a_1}$ or $\sqrt{z - a_1}$) near the common branching point a_1 . Such equations are understood to be valid in a small neighborhood of a_1 , with a proper correspondence between the sheets of both covers.

We will also consider double covers of discrete domains. In this case the following (slightly modified) notation will be convenient:

- For a discrete domain Ω_δ and a *face* $a \in \text{Int}\mathcal{V}_{\Omega_\delta}^\circ$, we set $[\Omega_\delta, a^\rightarrow] := [\Omega_\delta, a] \setminus \{a + \frac{\delta}{2}\}$, excluding both points over the corner $a + \frac{\delta}{2}$ from the natural double cover branching at a . Similarly, we set $[\Omega_\delta, a_1^\rightarrow, a_2, \dots, a_n] := [\Omega_\delta, a_1, \dots, a_n] \setminus \{a_1 + \frac{\delta}{2}\}$, if several faces $a_1, \dots, a_n \in \text{Int}\mathcal{V}_{\Omega_\delta}^\circ$ are marked.

2.1.3. *Contours.* Recall that we consider the critical Ising model on the *faces* of Ω_δ . In order to define the main tool of the paper (the discrete holomorphic spinors), we need some additional notation related to the contour representation of the model known as the *low-temperature expansion*; see [Pal07, Chap. 1].

- We denote by $\mathcal{C}_{\Omega_\delta}$ the family of all collections of closed contours on Ω_δ , i.e., the family of subsets of edges $\omega \subset \mathcal{V}_{\Omega_\delta}^{\text{m}}$ such that every vertex $v \in \mathcal{V}_{\Omega_\delta}^\bullet$ belongs to an even number of edges in ω .

The set $\mathcal{C}_{\Omega_\delta}$ is in a natural one-to-one correspondence with the spin configurations on Ω_δ with $+$ boundary conditions: trace an edge between any two adjacent faces with opposite spins. Under this mapping, the probability of a collection of interfaces $\omega \subset \Omega_\delta$ becomes proportional to $\alpha_c^{\#\text{edges}(\omega)}$, where $\alpha_c = \exp(-2\beta_c) = \sqrt{2} - 1$.

Below we also introduce families of contour collections that, besides a number of closed loops, contain a single path running from one fixed corner x to another corner or an edge midpoint y :

- For $x, y \in \mathcal{V}_{\Omega_\delta}^{\text{cm}}$, let $\pi_{x,y} = x \sim v_1 \sim \dots \sim v_n \sim y$ be some simple lattice path with $v_1, \dots, v_n \in \mathcal{V}_{\Omega_\delta}^\bullet$. We set $\mathcal{C}_{\Omega_\delta}(x, y) := \{\omega \oplus \pi_{x,y}, \omega \in \mathcal{C}_{\Omega_\delta}\}$, where \oplus denotes the XOR, or symmetric difference. It is easy to see that the set $\mathcal{C}_{\Omega_\delta}(x, y)$ does not depend on the particular choice of $\pi_{x,y}$. Note that, for any $\gamma \in \mathcal{C}_{\Omega_\delta}(x, y)$, there exists a (nonunique) decomposition of γ into a collection of disjoint, simple loops and a path $p(\gamma) \subset \gamma$ running from x to y . By a decomposition we mean that each edge in $\gamma \in \mathcal{C}_{\Omega_\delta}(x, y)$ belongs to exactly one loop (or to $p(\gamma)$) and is visited only once, and that there are no transversal intersections or self-intersections (see [Figure 2](#)).

2.2. *Construction of the discrete spinor observables.* Now we are ready to introduce the discrete spinor observables. The following definition generalizes the construction given in [\[CI13\]](#) to the case when a “source point” is inside Ω_δ .

Definition 2.1. Let Ω_δ be a discrete domain and $a_1, \dots, a_n \in \text{Int}\mathcal{V}_{\Omega_\delta}^\circ$ be inner faces. For a corner $z \in \mathcal{V}_{[\Omega_\delta, a_1^\rightarrow, \dots, a_n]}$ (below we also extend this definition to edge midpoints; see [Remark 2.2\(iii\)](#)), we define

$$(2.1) \quad F_{[\Omega_\delta, a_1, \dots, a_n]}(z) := \frac{1}{\mathcal{Z}_{\Omega_\delta}^+[\sigma_{a_1} \cdots \sigma_{a_n}]} \sum_{\gamma \in \mathcal{C}_{\Omega_\delta}(a_1 + \frac{\delta}{2}, z)} \alpha_c^{\#\text{edges}(\gamma)} \cdot \phi_{a_1, \dots, a_n}(\gamma, z),$$

where

- $\#\text{edges}(\gamma)$ is the number of full edges contained in γ and $\alpha_c = \sqrt{2} - 1$;
- the complex phase $\phi_{a_1, \dots, a_n}(\gamma, z)$ is defined by (see also [Figure 2](#))

$$\phi_{a_1, \dots, a_n}(\gamma, z) := e^{-\frac{i}{2}\text{wind}(p(\gamma))} \cdot (-1)^{\#\text{loops}_{a_1, \dots, a_n}(\gamma \setminus p(\gamma))} \cdot \text{sheet}_{a_1, \dots, a_n}(p(\gamma), z),$$

where, for a decomposition of γ mentioned in [Section 2.1.3](#),

- $\text{wind}(p(\gamma))$ is the total winding (increment of the argument of the tangent vector) of the path $p(\gamma)$ when going from $a_1 + \frac{\delta}{2}$ to z ;
- $\#\text{loops}_{a_1, \dots, a_n}(\gamma \setminus p(\gamma))$ is the number of loops in $\gamma \setminus p(\gamma)$ that contain an odd number of marked points a_1, \dots, a_n (equivalently, that do not lift to the double cover $[\Omega_\delta, a_1, \dots, a_n]$ as closed loops);
- the last factor $\text{sheet}_{a_1, \dots, a_n}(p(\gamma), z)$ is equal to $+1$ if z is on the same sheet of $[\Omega_\delta, a_1, \dots, a_n]$ as the end of the lift of $p(\gamma)$, and to -1 otherwise (more precisely, we fix one of the two points lying over the “source” $a_1 + \frac{\delta}{2}$ once forever and identify all other $z \in [\Omega_\delta, a_1, \dots, a_n]$ with paths running from *this* $a_1 + \frac{\delta}{2}$ to z modulo homotopy and an appropriate index 2 subgroup of the fundamental group);

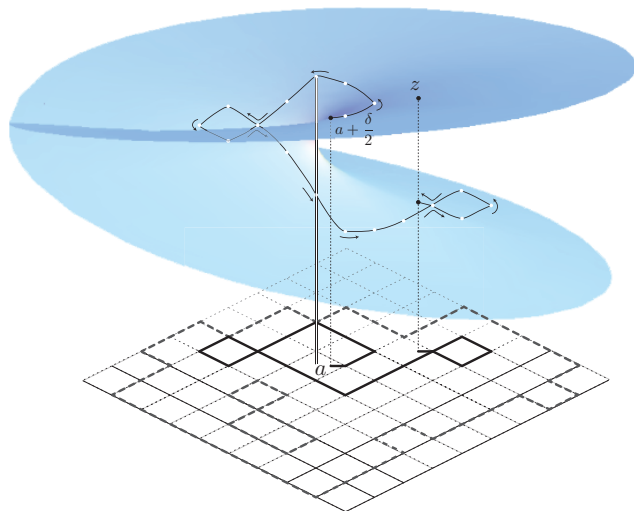


Figure 2. A contour collection $\gamma \in \mathcal{C}_{\Omega_\delta}(a + \frac{\delta}{2}, z)$ decomposed into nonintersecting loops (dashed) and a path $p(\gamma)$. Running from $a + \frac{\delta}{2}$ to the projection of z , this path makes a 3π turn counterclockwise, and thus $e^{-\frac{i}{2}\text{wind}(p(\gamma))} = i$. There is a single loop in γ surrounding a ; hence $(-1)^{\text{loops}_a(\gamma \setminus p(\gamma))} = -1$. Being lifted to the double cover $[\Omega_\delta, a]$, this path ends on the other sheet, thus $\text{sheet}_a(p(\gamma), z) = -1$, and $\phi_a(\gamma, z) = i$.

- the normalizing factor $\mathcal{Z}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}]$ is defined by

$$(2.2) \quad \mathcal{Z}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}] := \sum_{\omega \in \mathcal{C}_{\Omega_\delta}} \alpha_c^{\#\text{edges}(\omega)} (-1)^{\#\text{loops}_{a_1, \dots, a_n}(\omega)}.$$

Remark 2.2. (i) For any $\omega \in \mathcal{C}_{\Omega_\delta}$, the sign $(-1)^{\#\text{loops}_{a_1, \dots, a_n}(\omega)}$ coincides with the product of spins $\sigma_{a_1} \cdots \sigma_{a_n}$ in the corresponding Ising model configuration with $+$ boundary values. Since $\alpha_c^{\#\text{edges}(\omega)}$ is just the Ising weight of ω , one concludes that

$$\mathcal{Z}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}] = \mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}] \cdot \mathcal{Z}_{\Omega_\delta}^+ > 0,$$

where $\mathcal{Z}_{\Omega_\delta}^+ = \sum_{\omega \in \mathcal{C}_{\Omega_\delta}} \alpha_c^{\#\text{edges}}$ is the partition function of the model.

(ii) It is easy to check that the complex phase $\phi_{a_1, \dots, a_n}(\gamma, z)$ is independent of the choice of a decomposition of γ into a path $p(\gamma)$ and a collection of loops; e.g., see the discussion in [CI13]. Note that there are four types of corners: lying to the right of a nearby vertex v , below v , to the left of v , and upper v . For each of these groups, the winding of the path $p(\gamma)$ is defined uniquely modulo 2π . Therefore, the discrete spinors introduced above always are purely

real at the first group corners, are collinear to $\lambda := e^{\frac{\pi}{4}i}$ for the second group, etc. This motivates the following notation:

- we partition the set $\mathcal{V}_{\mathbb{C}_\delta}^c$ of all corners into four subsets $\mathcal{V}_{\Omega_\delta}^1, \mathcal{V}_{\Omega_\delta}^\lambda, \mathcal{V}_{\Omega_\delta}^i$ and $\mathcal{V}_{\Omega_\delta}^{\bar{\lambda}}$ depending on the position of the nearest vertex $v \in \mathcal{V}_{\mathbb{C}_\delta}^\bullet$ (to the left, upper, to the right, below) with respect to the corner.

(iii) We extend *Definition 2.1* to edge midpoints $z \in \mathcal{V}_{[\Omega_\delta, a_1^\rightarrow, \dots, a_n]}^m$ by adding the factor $(\cos \frac{\pi}{8})^{-1}$ to the formula (2.1), with $\#edges(\gamma)$ being the number of full edges contained in γ and the complex phase $\phi_{a_1, \dots, a_n}(\gamma, z)$ being defined as above. Note that each edge midpoint z can be reached by a path $p(\gamma)$ from two opposite sides, and both types of configurations are included in the sum. Thus, in this case the argument of the spinor value $F_{[\Omega_\delta, a_1, \dots, a_n]}(z)$ is no longer fixed.

(iv) The definition of $F_{[\Omega_\delta, a_1, \dots, a_n]}$ is invariant under permutations of a_2, \dots, a_n . The reader should always keep in mind, however, that the point a_1 plays a special role. The same applies to $\mathcal{A}_\Omega(a_1, \dots, a_n)$ and other related notation below.

2.3. S-holomorphicity and boundary conditions. A version of discrete holomorphicity, the notion of *s-holomorphicity* was introduced in [Smi06] together with the nonbranching version of discrete holomorphic observables as a tool to study the critical Ising model on the square lattice. The properties of such functions were further investigated in [CS12] for a more general class of graphs. On the square grid, s-holomorphic functions may be thought of as (more classical) discrete holomorphic functions whose real part is defined on $\mathcal{V}_{\Omega_\delta}^1$ and imaginary part on $\mathcal{V}_{\Omega_\delta}^i$, extended in a particular way to $\mathcal{V}_{\Omega_\delta}^\lambda, \mathcal{V}_{\Omega_\delta}^{\bar{\lambda}}$, and further to $\mathcal{V}_{\Omega_\delta}^m$ (see more details in Section 3.1). Our definitions resemble those in [Smi06].

Definition 2.3. With each corner $x \in \mathcal{V}_{\mathbb{C}_\delta}^\tau$ (with $\tau \in \{1, i, \lambda, \bar{\lambda}\}$), we associate the line $\ell(x) := \tau\mathbb{R}$ in the complex plane and denote by $P_{\ell(x)}$ the projection onto that line, defined by

$$P_{\ell(x)}[w] := \frac{1}{2} (w + \tau^2 \bar{w}), \quad w \in \mathbb{C}.$$

We say that a function $F : \mathcal{V}_{\Omega_\delta}^{cm} \rightarrow \mathbb{C}$ is *s-holomorphic* in Ω_δ if for every $x \in \mathcal{V}_{\Omega_\delta}^c$ and $z \in \mathcal{V}_{\Omega_\delta}^m$ that are adjacent, one has

$$F(x) = P_{\ell(x)}[F(z)].$$

For functions defined on double covers, we introduce the notion of s-holomorphicity exactly in the same manner.

The following proposition contains the crucial properties of $F_{[\Omega_\delta, a_1, \dots, a_n]}$ that will allow us to analyze their scaling limits. For $z \in \partial\mathcal{V}_{\Omega_\delta}^m$, let $\nu_{out}(z)$

denote the discrete “outer normal to the boundary at z ”: the edge whose midpoint is z , oriented towards the exterior of the domain and viewed as a complex number.

PROPOSITION 2.4. *The function $F_{[\Omega_\delta, a_1, \dots, a_n]}$ is s-holomorphic and has (multiplicative) monodromy -1 around each of the marked points a_1, \dots, a_n , thus being a discrete s-holomorphic spinor on $\mathcal{V}_{[\Omega_\delta, a_1^\rightarrow, \dots, a_n]}^{\text{cm}}$. Also,*

$$(2.3) \quad \Im m \left[F_{[\Omega_\delta, a_1, \dots, a_n]}(z) \sqrt{\nu_{\text{out}}(z)} \right] = 0 \text{ for all } z \in \partial \mathcal{V}_{[\Omega_\delta, a_1^\rightarrow, \dots, a_n]}^{\text{cm}}.$$

Proof. We give a proof (based on the standard XOR bijection; cf. [CS12]) in Section 3.1. \square

Remark 2.5. The boundary conditions (2.3) are *a priori* not robust enough to pass to the scaling limit: even if the limiting domain Ω has a smooth boundary, the discrete normal $\nu_{\text{out}}(z)$ can possibly admit only the values $e^{\pm \frac{\pi i}{4}}$ and $e^{\pm \frac{3\pi i}{4}}$, and so does not (pointwise) converge to its continuous counterpart. These conditions become much more tractable if one finds a way to “integrate” the square of $F_{[\Omega_\delta, a_1, \dots, a_n]}$: the real part of this primitive (antiderivative) will satisfy Dirichlet boundary conditions on $\partial\Omega$ due to (2.3). This approach is not as straightforward as in the continuum since the square of a discrete holomorphic function is, in general, not discrete holomorphic, and so does not have a well-defined discrete primitive. However, the following remarkable fact has been observed in [Smi06]: one can naturally define the real part of the integral, using the s-holomorphicity of discrete observables, which is a stronger property than the usual discrete holomorphicity. Moreover, a technique developed in [CS12] allows one to treat this real part essentially as if it were a harmonic function; see further details in Section 3.3.

2.4. From discrete spinors to ratios of correlations. The following lemma expresses ratios of spin correlations in terms of the spinor observables introduced in Section 2.2, providing a crucial ingredient for the proof of Theorems 1.5 and 1.7.

LEMMA 2.6. *For any $n = 1, 2, \dots$, we have*

$$(2.4) \quad \frac{\mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1+2\delta} \sigma_{a_2} \cdots \sigma_{a_n}]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}]} = F_{[\Omega_\delta, a_1, \dots, a_n]} \left(a_1 + \frac{3\delta}{2} \right),$$

where we take the corner $a_1 + \frac{3\delta}{2}$ on the same sheet as the “source point” $a_1 + \frac{\delta}{2}$. Moreover, in the case of just two marked points, we also have

$$(2.5) \quad \frac{\mathbb{E}_{\Omega_\delta}^{\text{free}} [\sigma_{a+\delta} \sigma_{b+\delta}]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b]} = \pm i F_{[\Omega_\delta, a, b]} \left(b + \frac{\delta}{2} \right),$$

where $\mathbb{E}_{\Omega_\delta}^{\text{free}}$ denotes the expectation in the critical Ising model defined on the vertices of Ω_δ with free boundary conditions, and the sign \pm depends on the sheet where the corner $b + \frac{\delta}{2}$ is taken.

Proof. Recall that

$$F_{[\Omega_\delta, a_1, \dots, a_n]} \left(a_1 + \frac{3\delta}{2} \right) = \frac{\sum_{\gamma \in \mathcal{C}_{\Omega_\delta}(a_1 + \frac{\delta}{2}, a_1 + \frac{3\delta}{2})} \alpha_c^{\#\text{edges}(\gamma)} \phi_{a_1, \dots, a_n} \left(\gamma, a_1 + \frac{3\delta}{2} \right)}{\mathcal{Z}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}]},$$

while, taking into account [Remark 2.2\(i\)](#),

$$\frac{\mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1+2\delta} \sigma_{a_2} \cdots \sigma_{a_n}]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}]} = \frac{\sum_{\omega \in \mathcal{C}_{\Omega_\delta}} \alpha_c^{\#\text{edges}(\omega)} (-1)^{\#\text{loops}_{a_1+2\delta, \dots, a_n}(\omega)}}{\mathcal{Z}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}]}.$$

There is a simple bijection ω between the sets $\mathcal{C}_{\Omega_\delta}(a_1 + \frac{\delta}{2}, a_1 + \frac{3\delta}{2})$ and $\mathcal{C}_{\Omega_\delta}$: removing the two corner-edges $(a_1 + \frac{\delta}{2}, a_1 + \delta)$ and $(a_1 + \delta, a_1 + \frac{3\delta}{2})$ from a given γ , we obtain a collection of closed loops $\omega(\gamma) \in \mathcal{C}_{\Omega_\delta}$ and vice versa. So, it suffices to show that

$$\phi_{a_1, \dots, a_n} \left(\gamma, a_1 + \frac{3\delta}{2} \right) = (-1)^{\#\text{loops}_{a_1+2\delta, a_2, \dots, a_n}(\omega(\gamma))}.$$

Let us pick any loop in γ and remove it. The left-hand side (respectively, the right-hand side) has changed the sign if and only if there was an odd number of points a_1, \dots, a_n (respectively, $a_1 + 2\delta, a_2, \dots, a_n$) inside the loop. However, no loop in γ separates a_1 from $a_1 + 2\delta$ (such a loop would intersect $p(\gamma)$), so the two sides can only change sign simultaneously. Thus it is sufficient to consider the case when γ is just a single non self-intersecting path $p(\gamma)$ running from $a_1 + \frac{\delta}{2}$ to $a_1 + \frac{3\delta}{2}$. This case is treated by the following observations: $\text{sheet}_{a_1, \dots, a_n}(p(\gamma), a_1 + \frac{3\delta}{2}) = -1$ if and only if there is an odd number of points a_1, \dots, a_n inside the loop $\omega(p(\gamma))$, and $\text{wind}(p(\gamma)) = 2\pi \pmod{4\pi}$ if and only if $\omega(p(\gamma))$ separates a_1 from $a_1 + 2\delta$.

For (2.5), the Kramers-Wannier duality (e.g., see [\[Pal07, Chap. 1\]](#)) implies

$$\frac{\mathbb{E}_{\Omega_\delta}^{\text{free}} [\sigma_{a+\delta} \sigma_{b+\delta}]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b]} = \frac{\sum_{\omega \in \mathcal{C}_{\Omega_\delta}(a + \frac{\delta}{2}, b + \frac{\delta}{2})} \alpha_c^{\#\text{edges}(\omega)}}{\mathcal{Z}_{\Omega_\delta}^+ [\sigma_a \sigma_b]}.$$

Hence it is sufficient to prove that the (purely imaginary) number $\phi_{a,b}(\gamma, b + \frac{\delta}{2})$ does not depend on γ . We have $\#\text{loops}_{a,b}(\gamma \setminus p(\gamma)) = 0$ since any loop in γ either surrounds both a, b or none of them, for otherwise it would intersect the path $p(\gamma)$ joining $a + \frac{\delta}{2}$ and $b + \frac{\delta}{2}$. Further, let $\pi_{ba}^\circ = b \sim v_1 \sim \dots \sim v_m \sim a$ be some simple lattice path with $v_j \in \mathcal{V}_{\Omega_\delta}^\circ$. Then, $p(\gamma) \cup (b + \frac{\delta}{2}, b) \cup \pi_{ab}^\circ \cup (a, a + \frac{\delta}{2})$ is a loop, which we denote by $l(\gamma)$. Let $n(\gamma)$ be the number of self-intersections of $l(\gamma)$ (in order words, the number of intersections of $p(\gamma)$ with π_{ba}°). Since

$\exp[-\frac{i}{2}\text{wind}(l(\gamma))] = (-1)^{n(\gamma)+1}$ (this is true for any closed loop), we see that

$$\exp[-\frac{i}{2}\text{wind}(p(\gamma))] = (-1)^{n(\gamma)+1} \cdot \exp[\frac{i}{2}\text{wind}((b + \frac{\delta}{2}, b) \cup \pi_{ba}^\circ \cup (a, a + \frac{\delta}{2}))],$$

where the second factor does not depend on γ . But we may also view π_{ba}° as a cut defining a sheet of the double cover $[\Omega, a, b]$, meaning that

$$\text{sheet}_{a,b} \left(p(\gamma), b + \frac{\delta}{2} \right) = (-1)^{n(\gamma)},$$

which proves the desired result. □

Remark 2.7. (i) Similarly to (2.4), one can check that

$$(2.6) \quad \frac{\mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1+(1\pm i)\delta}\sigma_{a_2} \cdots \sigma_{a_n}]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}]} = e^{\pm \frac{\pi i}{4}} F_{[\Omega_\delta, a_1, \dots, a_n]} \left(a_1 + \left(1 \pm \frac{i}{2}\right) \delta \right).$$

The proof boils down to the identity

$$e^{\pm \frac{\pi i}{4}} \phi_{a_1, \dots, a_n} \left(\gamma, a_1 + \left(1 \pm \frac{i}{2}\right) \delta \right) = (-1)^{\#\text{loops}_{a_1+\delta(1\pm i), a_2, \dots, a_n}(\omega(\gamma))}$$

for the natural bijection $\gamma \mapsto \omega(\gamma)$ removing the two corner-edges $(a_1 + \frac{\delta}{2}, a_1 + \delta)$ and $(a_1 + \delta, a_1 + (1 \pm \frac{i}{2})\delta)$ from a given $\gamma \in \mathcal{C}_{\Omega_\delta} \left(a_1 + \frac{\delta}{2}, a_1 + (1 \pm \frac{i}{2})\delta \right)$, which we leave to the reader.

(ii) The identity (2.5) can be extended to the case of $2n$ marked points (see [CI13, Prop. 5.6]), thus allowing one to treat $2n$ -point correlation functions with free boundary conditions.

2.5. Continuous spinors. In this section we introduce the continuous counterparts of the discrete spinor observables defined in Section 2.2: the continuous holomorphic spinors $f_{[\Omega, a_1, \dots, a_n]}$. We define them as solutions to the conformally covariant Riemann boundary value problem (2.7)–(2.9), which is a continuous analogue of the corresponding discrete boundary value problem (see Remark 2.11).

Definition 2.8. Let Ω be a bounded simply connected domain with smooth boundary, and let $a_1, \dots, a_n \in \Omega$. We define $f_{[\Omega, a_1, \dots, a_n]}$ to be the (unique) holomorphic spinor on $[\Omega, a_1, \dots, a_n]$, branching around each of a_1, \dots, a_n and satisfying the following conditions:

$$(2.7) \quad \Im \left[f_{[\Omega, a_1, \dots, a_n]}(z) \sqrt{\nu_{\text{out}}(z)} \right] = 0, \quad z \in \partial\Omega,$$

$$(2.8) \quad \lim_{z \rightarrow a_1} \sqrt{z - a_1} \cdot f_{[\Omega, a_1, \dots, a_n]}(z) = 1,$$

$$(2.9) \quad \lim_{z \rightarrow a_k} \sqrt{z - a_k} \cdot f_{[\Omega, a_1, \dots, a_n]}(z) \in i\mathbb{R}, \quad k = 2, \dots, n,$$

where $\nu_{\text{out}}(z)$ denotes the outer normal to the boundary of Ω at z .

LEMMA 2.9. *The boundary value problem (2.7)–(2.9) has a unique solution. If $\varphi : \Omega \rightarrow \Omega'$ is a conformal mapping, then one has*

$$(2.10) \quad f_{[\Omega, a_1, \dots, a_n]}(z) = f_{[\Omega', \varphi(a_1), \dots, \varphi(a_n)]}(\varphi(z)) \cdot (\varphi'(z))^{1/2}.$$

Remark 2.10. We use this covariance property as the *definition* of the continuous spinor in an arbitrary simply connected domain Ω , taking Ω' to be any smooth bounded domain.

Proof. If f_1, f_2 both satisfy (2.7)–(2.9), then the spinor $f_1 - f_2$ satisfies (2.7) and (2.9), while $\lim_{z \rightarrow a_1} \sqrt{z - a_1} \cdot (f_1(z) - f_2(z)) = 0$. Applying the Cauchy residue theorem to the *single-valued* function $(f_1(z) - f_2(z))^2$, one arrives at

$$(2.11) \quad 0 \leq i^{-1} \int_{\partial\Omega} (f_1(z) - f_2(z))^2 dz = 2\pi \sum_{k=2}^n \lim_{z \rightarrow a_k} (z - a_k)(f_1(z) - f_2(z))^2 \leq 0,$$

where the first inequality easily follows from (2.7) and the second from (2.9). Hence, $f_1 \equiv f_2$. Moreover, the conformal covariance property (2.10) now follows from an easy observation that its right-hand side satisfies all the conditions (2.7)–(2.9).

To prove the existence, consider the case $\Omega = \mathbb{H}$ and observe that if $Q(z)$ is a polynomial of degree $n - 1$ with real coefficients, then

$$(2.12) \quad f_Q(z) := e^{\frac{\pi i}{4}} \cdot \frac{Q(z)}{\sqrt{(z - a_1)(z - \bar{a}_1) \cdots (z - a_n)(z - \bar{a}_n)}}$$

is a holomorphic spinor on $[\mathbb{H}, a_1, \dots, a_n]$ satisfying (2.7) and the additional regularity condition $f_Q(z) = O(|z|^{-1})$ at infinity. Note that (2.11) shows that, if f_Q satisfies (2.9), then $\Im \lim_{z \rightarrow a_1} \sqrt{z - a_1} \cdot f_Q(z) = 0$. Therefore, conditions (2.8) and (2.9) give rise to n (real) linear equations on n unknown coefficients of Q . By the above argument, the homogeneous counterpart of this linear system has no nontrivial solutions; thus the system is nondegenerate. The solution for a given bounded smooth domain Ω is then constructed by (2.10) with $\Omega' = \mathbb{H}$; the decay of $f_Q(z)$ at infinity ensures that the right-hand side is bounded near $\varphi^{-1}(\infty)$. □

Remark 2.11. The first condition (2.7) in Definition 2.8 is a natural counterpart of (2.3). The third condition (2.9) comes from the following observation: the discrete primitive $H_{[\Omega, a_1, \dots, a_n]}$ of the “discrete differential form” $\Re [F_{[\Omega, a_1, \dots, a_n]}^2 dz]$ (which can be defined due to the s-holomorphicity property of the discrete observable; see Remark 2.5) remains bounded from below near the branching points a_2, \dots, a_n as $\delta \rightarrow 0$. Thus, we impose the same condition for the scaling limits, which means that $\Re [\int f_{[\Omega, a_1, \dots, a_n]}^2 dz]$ should behave like $c_k \log |z - a_k|$ for some *nonpositive* c_k as $z \rightarrow a_k, k = 2, \dots, n$, implying (2.9).

The second [condition \(2.8\)](#), which fixes the behavior of $f_{[\Omega, a_1, \dots, a_n]}$ near the “source point” a_1 , is the most delicate one and will be clarified later on. (See [Section 3.2](#), particularly [Lemma 3.5](#).) Note that it is sufficient to assume that $f_{[\Omega, a_1, \dots, a_n]}$ does not blow up faster than $1/\sqrt{z - a_1}$ at a_1 . Indeed, in this case the argument similar to [\(2.11\)](#) shows that $\lim_{z \rightarrow a_1} (z - a_1)(f_{[\Omega, a_1, \dots, a_n]}(z))^2 > 0$, and the rest is just a proper choice of the normalization.

Let us now introduce the quantities that play a central role in our computations of scaling limits of spin correlations, as they will turn out to be the limits of discrete logarithmic derivatives in [Theorem 1.5](#).

Definition 2.12. We define the complex number

$$\mathcal{A}_\Omega(a_1, \dots, a_n) = \mathcal{A}_{[\Omega, a_1, \dots, a_n]}$$

as the coefficient in the expansion

$$(2.13) \quad f_{[\Omega, a_1, \dots, a_n]} = \frac{1}{\sqrt{z - a_1}} + 2\mathcal{A}_{[\Omega, a_1, \dots, a_n]}\sqrt{z - a_1} + O(|z - a_1|^{3/2})$$

of $f_{[\Omega, a_1, \dots, a_n]}$ near the point a_1 . In the special case $n = 2$, we also define the quantity $\mathcal{B}_\Omega(a, b) = \mathcal{B}_{[\Omega, a, b]} > 0$ as the coefficient in the expansion of $f_{[\Omega, a, b]}$ near b :

$$(2.14) \quad f_{[\Omega, a, b]} = \pm \frac{i\mathcal{B}_{[\Omega, a, b]}}{\sqrt{z - b}} + O(|z - b|^{1/2}),$$

where the sign \pm depends on the sheet of $[\Omega, a, b]$.

Remark 2.13. Note that the covariance rule [\(1.10\)](#) for $\mathcal{A}_{[\Omega, a_1, \dots, a_n]}$ follows directly from the conformal covariance [\(2.10\)](#) of spinor observables: if $\varphi : \Omega \rightarrow \Omega'$ is a conformal mapping, then one has

$$\begin{aligned} f_{[\Omega, a_1, \dots, a_n]}(z) &= (\varphi'(z))^{1/2} f_{[\Omega', \varphi(a_1), \dots, \varphi(a_n)]}(\varphi(z)) \\ &= \left[\frac{\varphi'(z)}{\varphi(z) - \varphi(a_1)} \right]^{\frac{1}{2}} \cdot [1 + 2\mathcal{A}_\varphi \cdot (\varphi(z) - \varphi(a_1)) + \dots] \\ &= \left[\frac{1 + \frac{\varphi''(a_1)}{\varphi'(a_1)}(z - a_1) + \dots}{(z - a_1)(1 + \frac{\varphi''(a_1)}{2\varphi'(a_1)}(z - a_1) + \dots)} \right]^{\frac{1}{2}} \cdot [1 + 2\mathcal{A}_\varphi \cdot \varphi'(a_1)(z - a_1) + \dots] \\ &= \frac{1}{\sqrt{z - a_1}} \left[1 + 2 \left(\mathcal{A}_\varphi \cdot \varphi'(a_1) + \frac{1}{8} \frac{\varphi''(a_1)}{\varphi'(a_1)} \right) (z - a_1) + \dots \right], \end{aligned}$$

where $\mathcal{A}_\varphi = \mathcal{A}_{[\Omega', \varphi(a_1), \dots, \varphi(a_n)]}$. Similar arguments show that the coefficient $\mathcal{B}_{[\Omega, a, b]}$ is conformally invariant, i.e.,

$$(2.15) \quad \mathcal{B}_\Omega(a, b) = \mathcal{B}_{\Omega'}(\varphi(a), \varphi(b)).$$

We defer the further analysis of \mathcal{A}_Ω and \mathcal{B}_Ω until [Lemma 2.21](#) and the appendix, where we compute those quantities explicitly for $n \leq 2$ and $n > 2$, respectively.

2.6. *Convergence of the spinors.* The main purpose of this section is to derive [Theorems 1.5](#) and [1.7](#) from the convergence results for the discrete spinor observables $F_{[\Omega_\delta, a_1, \dots, a_n]}$ that are formulated in [Theorems 2.16](#), [2.18](#) and [2.20](#) below. The proofs of those are given in [Section 3](#), which is the most technical part of our paper. We use the following definitions concerning convergence of discrete s-holomorphic functions:

- We say that a family Ω_δ of discrete domains approximates a continuous domain $\Omega \subset \mathbb{C}$ as $\delta \rightarrow 0$ if $\partial\Omega_\delta$ converges to $\partial\Omega$ in the Hausdorff sense. (Note that our proofs can be easily generalized to the Carathéodory convergence of planar domains, which is weaker than the Hausdorff one used in this paper for simplicity.)
- We say that an s-holomorphic function (or a spinor) F_δ defined in $\Omega_\delta^{\text{cm}}$ (or its double cover) tends to a holomorphic function (or a spinor) f as $\delta \rightarrow 0$ if the “mid-edge values” $F_\delta|_{\Omega_\delta^{\text{m}}}$ approximate the values of f , while the “corner values” $F|_{\Omega_\delta^\tau}$, $\tau \in \{1, \lambda, i, \bar{\lambda}\}$, tend to the projections of f onto the corresponding lines $\tau\mathbb{R}$ (see [Definition 2.3](#)).
- We say that a convergence of discrete functions $F_\delta(z) = F_\delta(z; a_1, a_2, \dots)$ to $f(z) = f(z; a_1, a_2, \dots)$ is uniform on some compact set if and only if the differences $|F_\delta(z; a_1, a_2, \dots) - f(z; a_1, a_2, \dots)|$ are uniformly small as $\delta \rightarrow 0$, when we interpret lattice vertices (or mid-edges, corners, etc) z, a_1, a_2, \dots as the corresponding complex points when we plug them into f .

The crucial ingredient of our proofs is the interplay between (a) the values of discrete spinor observables near their branching points, which are related to the ratios of spin correlations by [Lemma 2.6](#), and (b) the mid-range behavior of these observables, which can be further related to the asymptotics expansions of their scaling limits ([2.13](#)), ([2.14](#)).

As a main tool to relate (a) and (b), we use a *full-plane version* $F_{[\mathbb{C}_\delta, a]}$ of the spinor observable. (Since we are interested in local considerations, it is sufficient to stick to the case of one marked point.) Though it could be constructed as an infinite-volume limit of the finite-domain observables, we prefer a more explicit strategy, which is outlined after the following lemma claiming the existence of $F_{[\mathbb{C}_\delta, a]}$.

LEMMA 2.14. *For $a \in \mathcal{V}_{\mathbb{C}_\delta}^{\circ}$, there exists a (unique) s-holomorphic spinor $F_{[\mathbb{C}_\delta, a]} : \mathcal{V}_{[\mathbb{C}_\delta, a]}^{\text{cm}} \rightarrow \mathbb{C}$ such that $F_{[\mathbb{C}_\delta, a]}(a + \frac{3\delta}{2}) = 1$ and $F_{[\mathbb{C}_\delta, a]}(z) = o(1)$ as $z \rightarrow \infty$. Moreover,*

$$(2.16) \quad \frac{1}{\vartheta(\delta)} F_{[\mathbb{C}_\delta, a]}(z) \xrightarrow{\delta \rightarrow 0} \frac{1}{\sqrt{z - a}} =: f_{[\mathbb{C}, a]}(z),$$

uniformly on compact subsets of $\mathbb{C} \setminus \{a\}$, where $\vartheta(\delta)$ is defined as

$$(2.17) \quad \vartheta(\delta) := F_{[\mathbb{C}_\delta, a]} \left(a + \frac{3\delta}{2} + 2\delta \lfloor \frac{1}{2\delta} \rfloor \right).$$

Proof. The detailed proof is given in [Section 3.2](#). First, we define the (real) values of $F_{[\mathbb{C}_\delta, a]}$ on $\mathcal{V}_{\mathbb{C}_\delta}^1$ as the discrete harmonic measure of the tip point $a + \frac{3\delta}{2}$ in the slit discrete plane $\mathcal{V}_{\mathbb{C}_\delta}^1 \setminus \{x + a : x \leq 0\}$. This definition is motivated by the following observation: in the continuous setup, the function $\Re[1/\sqrt{z-a}]$ can be viewed as the properly normalized harmonic measure of the tip a in the slit plane $\mathbb{C} \setminus \{x + a : x \leq 0\}$. Second, we extend $F_{[\mathbb{C}_\delta, a]}$ to $\mathcal{V}_{\mathbb{C}_\delta}^i$ by harmonic conjugation, then by symmetry to another sheet of $[\mathbb{C}_\delta, a]$, and eventually as an s -holomorphic function to $\mathcal{V}_{\mathbb{C}_\delta}^\lambda$, $\mathcal{V}_{\mathbb{C}_\delta}^{\bar{\lambda}}$ and $\mathcal{V}_{\mathbb{C}_\delta}^m$. The convergence [\(2.16\)](#) follows from known results on convergence of harmonic measures (e.g., see [\[CS11\]](#)). \square

Remark 2.15. The normalizing factor $\vartheta(\delta)$ is essentially the value of $F_{[\mathbb{C}_\delta, a]}$ at $a + 1$. In [Section 3.2](#) we show that

$$(2.18) \quad C_- \sqrt{\delta} \leq \vartheta(\delta) \leq C_+ \sqrt{\delta},$$

where $C_\pm > 0$ are some absolute constants. Note that one can compute the limit $\lim_{\delta \rightarrow 0} \vartheta(\delta)/\sqrt{\delta}$ using the recent work of Dubédat [\[Dub11a\]](#), but we do not need this sharp result.

We further use the normalizing factors $\vartheta(\delta)$ introduced above in order to formulate the following convergence theorem for discrete spinor observables away from $\partial\Omega$ and a_1, \dots, a_n .

THEOREM 2.16. *Let discrete simply connected domains Ω_δ approximate a bounded simply connected domain Ω as $\delta \rightarrow 0$. Then, for any $\epsilon > 0$ and any $n = 1, 2, \dots$, we have*

$$\frac{1}{\vartheta(\delta)} F_{[\Omega_\delta, a_1, \dots, a_n]}(z) \xrightarrow{\delta \rightarrow 0} f_{[\Omega, a_1, \dots, a_n]}(z),$$

uniformly over all $a_1, \dots, a_n \in \mathcal{V}_{\Omega_\delta}^\circ$ and $z \in \mathcal{V}_{\Omega_\delta}^m$ that are at the distance at least ϵ from $\partial\Omega$ and from each other.

Proof. See [Section 3.4](#). \square

Since we are interested in the coefficient $\mathcal{A}_{[\Omega, a_1, \dots, a_n]}$ in front of the term $\sqrt{z-a_1}$ in [\(2.13\)](#), along with the discrete analogue $F_{[\mathbb{C}_\delta, a]}$ of the function $1/\sqrt{z-a}$ given by [Lemma 2.14](#), we also need a discrete counterpart $G_{[\mathbb{C}_\delta, a]}$ of the function $\sqrt{z-a}$. We construct $G_{[\mathbb{C}_\delta, a]}$ by “discrete integration” of $F_{[\mathbb{C}_\delta, a]}$, just mimicking the continuous setup. It is sufficient to define $G_{[\mathbb{C}_\delta, a]}$ on $\mathcal{V}_{\mathbb{C}_\delta}^1$ only, as [Lemma 2.6](#) deals with the (real) values of discrete observables at the point $a_1 + \frac{3\delta}{2} \in \mathcal{V}_{\mathbb{C}_\delta}^1$.

LEMMA 2.17. *For $a \in \mathcal{V}_{\mathbb{C}_\delta}^\circ$, there exists a (unique) discrete harmonic spinor $G_{[\mathbb{C}_\delta, a]} : \mathcal{V}_{[\mathbb{C}_\delta, a]}^1 \rightarrow \mathbb{R}$ such that $G_{[\mathbb{C}_\delta, a]}$ vanishes on the half-line $\{x + a : x \leq 0\}$, $G_{[\mathbb{C}_\delta, a]}(a + \frac{3\delta}{2}) = \delta$, and $G_{[\mathbb{C}_\delta, a]}(z) = O(|z-a|^{1/2})$ as $z \rightarrow \infty$.*

Moreover,

$$(2.19) \quad \frac{1}{\vartheta(\delta)} G_{[\mathbb{C}_\delta, a]}(z) \xrightarrow{\delta \rightarrow 0} \Re \sqrt{z - a} =: g_{[\mathbb{C}, a]}(z),$$

uniformly on compact subsets of $\mathbb{C} \setminus \{a\}$.

Proof. For $z \in \mathcal{V}_{[\mathbb{C}_\delta, a]}^1$, we set $G_{[\mathbb{C}_\delta, a]}(z) := \delta \cdot \sum_{j=0}^\infty F_{[\mathbb{C}_\delta, a]}(z - 2j\delta)$. Certainly, one should check the convergence of this series and the harmonicity on the half-line $\{x + a : x \geq 0\}$; see further details in [Section 3.2](#). \square

In the continuum limit, the leading term in the expansion of $f_{[\Omega, a_1, \dots, a_n]} - f_{[\mathbb{C}, a_1]}$ near a_1 is given by $2\mathcal{A}_{[\Omega, a_1, \dots, a_n]} \sqrt{z - a_1}$. It is hence plausible to believe that the same holds true for the discrete spinors, and one has

$$(F_{[\Omega_\delta, a_1, \dots, a_n]} - F_{[\mathbb{C}_\delta, a_1]})(a_1 + \frac{3\delta}{2}) \approx 2\Re \mathcal{A}_{[\Omega_\delta, a_1, \dots, a_n]} \cdot G_{[\mathbb{C}_\delta, a_1]}(a_1 + \frac{3\delta}{2})$$

up to higher-order terms. (The real part appears due to discrete complex analysis subtleties, as real and imaginary parts of s-holomorphic functions are defined on different lattices, and $a_1 + \frac{3\delta}{2} \in \mathcal{V}_{\Omega_\delta}^1$.) We justify this heuristics in

THEOREM 2.18. *Under conditions of [Theorem 2.16](#), we have*

$$(2.20) \quad F_{[\Omega_\delta, a_1, \dots, a_n]}(a_1 + \frac{3\delta}{2}) - 1 - 2\Re \mathcal{A}_{[\Omega, a_1, \dots, a_n]} \cdot \delta = o(\delta)$$

as $\delta \rightarrow 0$, uniformly over all $a_1, \dots, a_n \in \mathcal{V}_{\Omega_\delta}^\circ$ that are at the distance at least ϵ from $\partial\Omega$ and from each other.

Proof. See [Section 3.5](#). \square

Remark 2.19. In the proof of [Theorem 2.18](#) we show that $F_{[\Omega_\delta, a_1, \dots, a_n]}$ and $F_{[\mathbb{C}_\delta, a_1]}$ are δ -close to each other at all points around a_1 — in particular, at $a_1 + (1 \pm \frac{i}{2})\delta$. Together with (2.6), this implies

$$\frac{\mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1 + (1 \pm i)\delta} \sigma_{a_2} \cdots \sigma_{a_n}]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}]} = 1 + O(\delta)$$

as $\delta \rightarrow 0$, uniformly over all $a_1, \dots, a_n \in \mathcal{V}_{\Omega_\delta}^\circ$ that are at the distance at least ϵ from $\partial\Omega$ and from each other.

The similar analysis for the quantity $\mathcal{B}_{[\Omega, a, b]}$ defined by the expansion (2.14) is simpler since we need to match the first-order coefficients instead of the second-order ones. The result is given by

THEOREM 2.20. *For $n = 2$, under conditions of [Theorem 2.16](#), we have*

$$(2.21) \quad F_{[\Omega_\delta, a, b]}(b + \frac{\delta}{2}) \pm i\mathcal{B}_{[\Omega, a, b]} = o(1)$$

as $\delta \rightarrow 0$, uniformly over all $a, b \in \mathcal{V}_{\Omega_\delta}^\circ$ that are at the distance at least ϵ from $\partial\Omega$ and from each other, where the sign \pm depends on the sheet of $[\Omega, a, b]$.

Proof. See [Section 3.5](#). \square

Proof of Theorems 1.5 and 1.7. Due to Lemma 2.6, asymptotics (1.8) is a reformulation of (2.20), while Theorem 1.7 is equivalent to Theorem 2.20. (The sign \pm is fixed due to the positivity of spin-spin correlations.) To check (1.9), we rotate our domain Ω around a by 90° clockwise. According to the conformal covariance rule (1.10), the coefficient \mathcal{A} multiplies by i , so the desired result follows from $\Re [i\mathcal{A}] = -\Im \mathcal{A}$. \square

2.7. *Integrating the logarithmic derivatives and explicit formulae.* We start this section by stating explicit formulae for $\mathcal{A}_{[\Omega, a_1, \dots, a_n]}$ and $\mathcal{B}_{[\Omega, a, b]}$.

LEMMA 2.21. *One has*

$$(2.22) \quad \mathcal{A}_{[\Omega, a_1, \dots, a_n]} = 2\partial_{a_1} \log \langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_{\Omega}^{\dagger},$$

$$(2.23) \quad \mathcal{B}_{[\Omega, a, b]} = \frac{\langle \sigma_a \sigma_b \rangle_{\Omega}^{\text{free}}}{\langle \sigma_a \sigma_b \rangle_{\Omega}^{\dagger}},$$

where $\partial_{a_1} = \frac{1}{2}(\partial_{x_1} - i\partial_{y_1})$ if $a_1 = x_1 + iy_1$, and the quantities $\langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_{\Omega}$ are defined by the conformal covariance rule (1.5) and explicit formulae (1.2) and (1.4).

Proof. Note that (1.5) readily implies that $2\partial_{a_1} \log \langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_{\Omega}^{\dagger}$ satisfies the same covariance rule (1.10) as $\mathcal{A}_{[\Omega, a_1, \dots, a_n]}$, and both sides in (2.23) are conformally invariant. Therefore, it suffices to check (2.22), (2.23) for $\Omega = \mathbb{H}$. This amounts to computing the continuous spinors $f_{[\mathbb{H}, a_1, \dots, a_n]}$, which can be done by solving a linear system for the coefficients of polynomials Q introduced in the proof of Lemma 2.9.

Let us first illustrate this procedure for $n = 1, 2$. It is straightforward to check that the spinors

$$f_{[\mathbb{H}, a]}(z) = \frac{(2i\Im a)^{\frac{1}{2}}}{\sqrt{(z-a)(z-\bar{a})}},$$

$$f_{[\mathbb{H}, a, b]}(z) = \frac{(2i\Im a)^{\frac{1}{2}}}{|b-\bar{a}| + |b-a|} \cdot \frac{[(\bar{b}-\bar{a})(\bar{b}-a)]^{\frac{1}{2}}(z-b) + [(b-a)(b-\bar{a})]^{\frac{1}{2}}(z-\bar{b})}{[(z-a)(z-\bar{a})(z-b)(z-\bar{b})]^{1/2}}$$

satisfy (2.7)–(2.9). Working out the expansions (2.13) and (2.14), one arrives at

$$\mathcal{A}_{\mathbb{H}}(a) = -\frac{1}{8i\Im a},$$

$$\mathcal{A}_{\mathbb{H}}(a, b) = -\frac{1}{8i\Im a} + \frac{|b-\bar{a}| - |b-a|}{4(|b-\bar{a}| + |b-a|)} \left(\frac{1}{b-a} - \frac{1}{\bar{b}-a} \right),$$

and

$$\mathcal{B}_{\mathbb{H}}(a, b) = \frac{(4\Im a \Im b)^{\frac{1}{2}}}{|b-\bar{a}| + |b-a|} = \frac{\langle \sigma_a \sigma_b \rangle_{\mathbb{H}}^{\text{free}}}{\langle \sigma_a \sigma_b \rangle_{\mathbb{H}}^{\dagger}},$$

proving (2.23). Taking the logarithmic derivatives with respect to a of the explicit expressions for $\langle \sigma_a \rangle_{\mathbb{H}}^+$ and $\langle \sigma_a \sigma_b \rangle_{\mathbb{H}}^+$ given by (1.2), one arrives at (2.22) for $n = 1, 2$. The proof for higher n is somewhat technical, and we defer it to the appendix. \square

The following proposition, which is a rephrasing of Corollary 1.6, now follows easily from Theorem 1.5 via integration with respect to the positions of points.

PROPOSITION 2.22. *As Ω_δ approximates Ω , one has*

$$(2.24) \quad \frac{\mathbb{E}_{\Omega_\delta}^+[\sigma_{b_1} \cdots \sigma_{b_n}]}{\mathbb{E}_{\Omega_\delta}^+[\sigma_{a_1} \cdots \sigma_{a_n}]} \xrightarrow{\delta \rightarrow 0} \frac{\langle \sigma_{b_1} \cdots \sigma_{b_n} \rangle_\Omega^+}{\langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_\Omega^+}$$

uniformly with respect to all a_k at distance at least ϵ from $\partial\Omega$ and each other, and b_k obeying the same condition.

Proof. Color the faces of Ω_δ black and white in a chessboard fashion. By Remark 2.19, the ratio of spin-spin correlations at two adjacent spins tends to 1 uniformly away from the boundary, so we can assume that all a_k, b_k are colored white. Let $a'_1 \in \Omega$ be such that $[a_1, a'_1]$ is a horizontal segment contained in Ω and disjoint with a_2, \dots, a_n . Denote by $a_1 = v_1 \sim \cdots \sim v_{m_\delta} = a'_1$ a straight horizontal lattice path approximating $[a_1, a'_1]$. Then, by Theorem 1.5, one has

$$\begin{aligned} \log \frac{\mathbb{E}_{\Omega_\delta}^+[\sigma_{v_{j+1}} \sigma_{a_2} \cdots \sigma_{a_n}]}{\mathbb{E}_{\Omega_\delta}^+[\sigma_{v_j} \sigma_{a_2} \cdots \sigma_{a_n}]} &= \left(\frac{\mathbb{E}_{\Omega_\delta}^+[\sigma_{v_{j+1}} \sigma_{a_2} \cdots \sigma_{a_n}]}{\mathbb{E}_{\Omega_\delta}^+[\sigma_{v_j} \sigma_{a_2} \cdots \sigma_{a_n}]} - 1 \right) (1 + o(1)) \\ &= 2\delta \cdot [\Re \mathcal{A}_\Omega(v_j; a_2, \dots, a_n) + o(1)] \end{aligned}$$

as $\delta \rightarrow 0$, where the $o(1)$ terms are uniform in j . Consequently,

$$(2.25) \quad \begin{aligned} \log \frac{\mathbb{E}_{\Omega_\delta}^+[\sigma_{a'_1} \sigma_{a_2} \cdots \sigma_{a_n}]}{\mathbb{E}_{\Omega_\delta}^+[\sigma_{a_1} \sigma_{a_2} \cdots \sigma_{a_n}]} &= 2\delta \cdot \sum_{j=1}^{m_\delta} \Re \mathcal{A}_\Omega(v_j, a_2, \dots, a_n) + o(1) \\ &\xrightarrow{\delta \rightarrow 0} \int_{[a_1, a'_1]} \Re \mathcal{A}_\Omega(x_1 + iy_1, a_2, \dots, a_n) dx_1. \end{aligned}$$

A similar formula with $-\Im \mathcal{A}_\Omega(x_1 + iy_1, a_2, \dots, a_n) dy_1$ in the right-hand side applies to the case when $[a_1, a'_1]$ is a vertical segment. Since by (2.22) one has

$$\begin{aligned} \Re \mathcal{A}_\Omega(x_1 + iy_1, a_2, \dots, a_n) &= \partial_{x_1} \log \langle \sigma_{x_1 + iy_1} \cdots \sigma_{a_n} \rangle_\Omega^+, \\ -\Im \mathcal{A}_\Omega(x_1 + iy_1, a_2, \dots, a_n) &= \partial_{y_1} \log \langle \sigma_{x_1 + iy_1} \cdots \sigma_{a_n} \rangle_\Omega^+, \end{aligned}$$

the proposition readily follows by applying the above computation, if necessary, to several horizontal and vertical segments and to other marked points a_2, \dots, a_n . \square

Remark 2.23. Note that whereas [Lemma 2.21](#) allows us to write the limits of correlations in [Theorem 1.2](#) explicitly, it is otherwise not essential for the proof. In fact, [Theorem 1.5](#) implies that discrete exact differential forms $d \log \mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}]$ have continuous limits $\mathcal{L}_{\Omega,n}$ given in terms of \mathcal{A}_Ω ; see [\(2.25\)](#). Therefore, one could simply *define* the functions $\langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_\Omega^+$ to be $\exp(\int \mathcal{L}_{\Omega,n})$ with properly chosen constants of integration. The only property of $\langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_\Omega^+$ that we use below (and which also defines their multiplicative normalization for $n > 2$ uniquely) is the following:

$$(2.26) \quad \frac{\langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_\Omega^+}{\langle \sigma_{a_1} \rangle_\Omega^+ \langle \sigma_{a_2} \cdots \sigma_{a_n} \rangle_\Omega^+} \rightarrow 1 \text{ as } a_1 \rightarrow \partial\Omega.$$

In principle, the existence of constants of integration such that [\(2.26\)](#) holds can be derived from the asymptotics

$$\mathcal{A}_\Omega(a_1, \dots, a_n) - \mathcal{A}_\Omega(a_1) = o(1), a_1 \rightarrow \partial\Omega,$$

$$\mathcal{A}_\Omega(a_1, \dots, a_n) - \mathcal{A}_\Omega(a_1, \dots, \hat{a}_k, \dots, a_n) = o(1), a_k \rightarrow \partial\Omega, k = 2, \dots, n,$$

which could be obtained by comparing solutions to the corresponding boundary value problems [\(2.7\)](#)–[\(2.9\)](#) as $a_1 \rightarrow \partial\Omega$.

2.8. From ratios of correlations to [Theorem 1.1](#). This section is devoted to the proof of [Theorem 1.1](#). Our goal is to relate the normalizing factors $\rho_2(\delta, \Omega_\delta)$ from [Corollary 1.6](#) (which, in principle, might depend on Ω) with the full-plane normalization $\rho(\delta)$. The proof is based on [Theorem 1.7](#), which claims the convergence of the ratios of free and + spin-spin correlations to the explicit limit $\mathcal{B}_\Omega(a, b) = \langle \sigma_a \sigma_b \rangle_\Omega^{\text{free}} / \langle \sigma_a \sigma_b \rangle_\Omega^+$. We also use the classical FKG (e.g., see [\[Gri06, Chap. 2\]](#)) and GHS (see [\[GHS70\]](#)) inequalities for the Ising model. The small additional ingredient is given by

Remark 2.24. The following property holds:

$$(2.27) \quad \mathcal{B}_\Omega(a, b) \rightarrow 1 \text{ as } \mathcal{D}_\Omega(a, b) \rightarrow 0, \quad \text{and} \quad \mathcal{B}_\Omega(a, b) \rightarrow 0 \text{ as } a \rightarrow \partial\Omega,$$

where

$$(2.28) \quad \mathcal{D}_\Omega(a, b) := \frac{|a - b|}{\text{dist}(\{a, b\}, \partial\Omega)}$$

is a quantity that measures how deeply in the bulk of Ω the points a, b are. This follows readily from conformal invariance of $\mathcal{B}_\Omega(a, b)$ and explicit formulae [\(1.4\)](#) for the half-plane. Also, if one denotes

$$\langle \sigma_a \sigma_b \rangle_{\mathbb{C}} := |a - b|^{-\frac{1}{4}},$$

then

$$(2.29) \quad \frac{\langle \sigma_a \sigma_b \rangle_\Omega^+}{\langle \sigma_a \sigma_b \rangle_{\mathbb{C}}} \rightarrow 1 \text{ as } \mathcal{D}_\Omega(a, b) \rightarrow 0.$$

Indeed, let φ be a conformal map from Ω to \mathbb{H} such that $\varphi(a) = i$. Due to standard estimates, one has $\varphi(b) \rightarrow i$ and $|\varphi'(b)| \rightarrow |\varphi'(a)|$ as $\mathcal{D}_\Omega(a, b) \rightarrow 0$. Therefore,

$$\frac{\langle \sigma_a \sigma_b \rangle_\Omega^+}{\langle \sigma_a \sigma_b \rangle_{\mathbb{C}}^+} = \frac{\langle \sigma_{\varphi(a)} \sigma_{\varphi(b)} \rangle_{\mathbb{H}}^+ |\varphi'(a)|^{\frac{1}{8}} |\varphi'(b)|^{\frac{1}{8}}}{\langle \sigma_a \sigma_b \rangle_{\mathbb{C}}^+} = \frac{|\varphi(a) - \varphi(b)|^{-\frac{1}{4}} |\varphi'(a)|^{\frac{1}{4}}}{|a - b|^{-\frac{1}{4}}} (1 + o(1)) \rightarrow 1.$$

LEMMA 2.25. *For any $\eta > 0$, there exists an $\epsilon > 0$ such that the following holds: if $\mathcal{D}_\Omega(a, b) < \epsilon$ and Ω_δ approximates Ω , then*

$$1 \leq \frac{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b]}{\mathbb{E}_{\mathbb{C}_\delta} [\sigma_a \sigma_b]} \leq 1 + \eta,$$

provided that δ is small enough.

Proof. By the FKG inequality, $\mathbb{E}_{\Lambda_\delta}^{\text{free}} [\sigma_a \sigma_b] \leq \mathbb{E}_{\mathbb{C}_\delta} [\sigma_a \sigma_b] \leq \mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b]$ for any domain Λ_δ containing a, b ; hence the left-hand side readily follows. For the right-hand side, choose $\Lambda_\delta = \Omega_\delta^\bullet - \delta$. We have, by Theorem 1.7,

$$\frac{\mathbb{E}_{\Lambda_\delta}^{\text{free}} [\sigma_a \sigma_b]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b]} = \frac{\mathbb{E}_{\Omega_\delta^\bullet}^{\text{free}} [\sigma_{a+\delta} \sigma_{b+\delta}]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b]} \xrightarrow{\delta \rightarrow 0} \mathcal{B}_\Omega(a, b).$$

Due to (2.27), one can choose ϵ so that $\mathcal{B}_\Omega(a, b) \geq 1 - \frac{\eta}{2}$, which gives the result. □

Proof of Theorem 1.1. Fix $\eta > 0$. For shortness, below we will write $a + \epsilon$ for its lattice approximations $a + 2\delta \lfloor \frac{\epsilon}{2\delta} \rfloor$ and $a + 1$ for its lattice approximation $a + 2\delta \lfloor \frac{1}{2\delta} \rfloor$. Denote

$$(2.30) \quad R_{a,b}^{\Omega_\delta} := (\varrho(\delta))^{-1} \mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b].$$

Recall that $\varrho(\delta) = \mathbb{E}_{\mathbb{C}_\delta} [\sigma_a \sigma_{a+1}]$ by definition, so one can write

$$R_{a,b}^{\Omega_\delta} = \frac{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_{a+\epsilon}]} \cdot \frac{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_{a+\epsilon}]}{\mathbb{E}_{\mathbb{C}_\delta} [\sigma_a \sigma_{a+\epsilon}]} \cdot \frac{\mathbb{E}_{\mathbb{C}_\delta} [\sigma_a \sigma_{a+\epsilon}]}{\mathbb{E}_{\mathbb{C}_\delta} [\sigma_a \sigma_{a+1}]}.$$

By Lemma 2.25, we can find a small $\epsilon > 0$ and a large domain Λ_δ containing $a, b, a + \epsilon$ and $a + 1$, such that

$$1 \leq \frac{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_{a+\epsilon}]}{\mathbb{E}_{\mathbb{C}_\delta} [\sigma_a \sigma_{a+\epsilon}]} \leq 1 + \eta \quad \text{and} \quad 1 - \eta \leq \frac{\mathbb{E}_{\mathbb{C}_\delta} [\sigma_a \sigma_{a+\epsilon}]}{\mathbb{E}_{\mathbb{C}_\delta} [\sigma_a \sigma_{a+1}]} \cdot \frac{\mathbb{E}_{\Lambda_\delta}^+ [\sigma_a \sigma_{a+1}]}{\mathbb{E}_{\Lambda_\delta}^+ [\sigma_a \sigma_{a+\epsilon}]} \leq 1 + \eta,$$

provided that δ is small enough. Consequently,

$$(1 - \eta) \frac{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_{a+\epsilon}]} \cdot \frac{\mathbb{E}_{\Lambda_\delta}^+ [\sigma_a \sigma_{a+\epsilon}]}{\mathbb{E}_{\Lambda_\delta}^+ [\sigma_a \sigma_{a+1}]} \leq R_{a,b}^{\Omega_\delta} \leq (1 + \eta)^2 \frac{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_{a+\epsilon}]} \cdot \frac{\mathbb{E}_{\Lambda_\delta}^+ [\sigma_a \sigma_{a+\epsilon}]}{\mathbb{E}_{\Lambda_\delta}^+ [\sigma_a \sigma_{a+1}]}$$

and, by convergence of the ratios of spin correlations proven in Proposition 2.22,

$$(1 - \eta)^2 \frac{\langle \sigma_a \sigma_b \rangle_\Omega^+}{\langle \sigma_a \sigma_{a+\epsilon} \rangle_\Omega^+} \cdot \frac{\langle \sigma_a \sigma_{a+\epsilon} \rangle_\Lambda^+}{\langle \sigma_a \sigma_{a+1} \rangle_\Lambda^+} \leq R_{a,b}^{\Omega_\delta} \leq (1 + \eta)^3 \frac{\langle \sigma_a \sigma_b \rangle_\Omega^+}{\langle \sigma_a \sigma_{a+\epsilon} \rangle_\Omega^+} \cdot \frac{\langle \sigma_a \sigma_{a+\epsilon} \rangle_\Lambda^+}{\langle \sigma_a \sigma_{a+1} \rangle_\Lambda^+}$$

for δ small enough. Since η can be chosen arbitrary small, and the bounds do not depend on δ , it only remains to show that we can make the factor

$$\frac{\langle \sigma_a \sigma_{a+\epsilon} \rangle_\Lambda^+}{\langle \sigma_a \sigma_{a+\epsilon} \rangle_\Omega^+ \langle \sigma_a \sigma_{a+1} \rangle_\Lambda^+}$$

as close to 1 as we wish by choosing ϵ small enough and Λ large enough. However, this follows readily from (2.29) if we multiply this factor by $\langle \sigma_a \sigma_{a+1} \rangle_{\mathbb{C}} = 1$. Thus, $\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b] \sim \varrho(\delta) \langle \sigma_a \sigma_b \rangle_\Omega^+$ as $\delta \rightarrow 0$. To derive the asymptotics of two-point correlations for free boundary conditions as $\delta \rightarrow 0$, note that Theorem 1.7 implies

$$\mathbb{E}_{\Omega_\delta^\bullet - \delta}^{\text{free}} [\sigma_a \sigma_b] \sim \mathcal{B}_\Omega(a, b) \cdot \mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b] \sim \varrho(\delta) \cdot \mathcal{B}_\Omega(a, b) \langle \sigma_a \sigma_b \rangle_\Omega^+ = \varrho(\delta) \langle \sigma_a \sigma_b \rangle_\Omega^{\text{free}}.$$

The fact that we have $\Omega_\delta^\bullet - \delta$ instead of Ω_δ plays no role, since they both approximate the same continuous domain Ω and the convergence of $(\varrho(\delta))^{-1} \mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b]$ is independent of the particular choice of lattice approximations. \square

Remark 2.26. As a simple byproduct of our analysis, we obtain the rotational invariance of the full-plane correlations recently proven by Pinson [Pin12]: by the FKG inequality, for any (large) domain Ω_δ , one has $\mathbb{E}_{\Omega_\delta}^{\text{free}} [\sigma_a \sigma_b] \leq \mathbb{E}_{\mathbb{C}_\delta} [\sigma_a \sigma_b] \leq \mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b]$ and, due to Theorem 1.1 and (2.27), both sides have the same asymptotics when Ω_δ exhausts \mathbb{C}_δ . Then, (2.29) gives the desired result:

$$(\varrho(\delta))^{-1} \mathbb{E}_{\mathbb{C}_\delta} [\sigma_a \sigma_b] \rightarrow \langle \sigma_a \sigma_b \rangle_{\mathbb{C}} = |a - b|^{-\frac{1}{4}} \text{ as } \delta \rightarrow 0.$$

2.9. *Decorrelation near the boundary and the proof of Theorem 1.2.* This section is devoted to the proof of Theorem 1.2. Note that it was already proven above in the special case $n = 2$ as a part of Theorem 1.1. Our goal is to relate the normalizing factors in Corollary 1.6 with $\varrho(\delta)$. Below we rely upon the asymptotics (2.26), which readily follow from explicit formulae (1.4) and the conformal covariance rule (1.5).

LEMMA 2.27. *Given a domain Ω with marked points a_2, \dots, a_n , $n \geq 2$, and a number $\eta > 0$, there exists $\epsilon > 0$ such that the following holds: if $a_1 \in \Omega$ is ϵ -close to the boundary, Ω_δ approximates Ω and δ is small enough, then*

$$1 - \eta \leq \frac{\mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1}] \mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_2} \cdots \sigma_{a_n}]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \cdots \sigma_{a_n}]} \leq 1.$$

Proof. The upper bound follows readily from FKG inequality. For the lower one, consider first the case $n = 2$. A celebrated application (e.g., see [DM10]) of the GHS inequality [GHS70] reads $\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b] - \mathbb{E}_{\Omega_\delta}^+ [\sigma_a] \mathbb{E}_{\Omega_\delta}^+ [\sigma_b] \leq \mathbb{E}_{\Omega_\delta}^{\text{free}} [\sigma_a \sigma_b]$ or, equivalently,

$$1 - \frac{\mathbb{E}_{\Omega_\delta}^{\text{free}} [\sigma_a \sigma_b]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b]} \leq \frac{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a] \mathbb{E}_{\Omega_\delta}^+ [\sigma_b]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b]}.$$

As $\delta \rightarrow 0$, the left-hand side converges to $1 - \mathcal{B}_\Omega(a; b)$, so (2.27) implies the claim.

To prove the result for $n \geq 3$, assume that we have already proved [Theorem 1.2](#) for all $n' < n$. (The precise description of our induction scheme is given in the proof of [Theorem 1.2](#) below; see (2.31).) Let γ be a cross-cut (simple path) in Ω separating a_1 from a_2, \dots, a_n , and let Ω' and Ω'' be the corresponding connected components. Note that the FKG inequality implies

$$\frac{\mathbb{E}_{\Omega_\delta}^+[\sigma_{a_1}]\mathbb{E}_{\Omega_\delta}^+[\sigma_{a_2} \cdots \sigma_{a_n}]}{\mathbb{E}_{\Omega_\delta}^+[\sigma_{a_1} \cdots \sigma_{a_n}]} \cdot \frac{\mathbb{E}_{\Omega'_\delta}^+[\sigma_{a_1}]\mathbb{E}_{\Omega''_\delta}^+[\sigma_{a_2} \cdots \sigma_{a_n}]}{\mathbb{E}_{\Omega_\delta}^+[\sigma_{a_1}]\mathbb{E}_{\Omega_\delta}^+[\sigma_{a_2} \cdots \sigma_{a_n}]} \geq 1$$

as $\mathbb{E}_{\Omega'_\delta}^+[\sigma_{a_1}]\mathbb{E}_{\Omega''_\delta}^+[\sigma_{a_2} \cdots \sigma_{a_n}]$ is equal to the correlation of $\sigma_{a_1}, \dots, \sigma_{a_n}$ in Ω_δ conditioned on the event that all spins neighboring γ are $+$. By the induction assumption, the second factor converges to

$$\frac{\langle \sigma_{a_1} \rangle_{\Omega'}^+}{\langle \sigma_{a_1} \rangle_{\Omega}^+} \cdot \frac{\langle \sigma_{a_2} \cdots \sigma_{a_n} \rangle_{\Omega''}^+}{\langle \sigma_{a_2} \cdots \sigma_{a_n} \rangle_{\Omega}^+}$$

as $\delta \rightarrow 0$, hence it is sufficient to show that we can make this quantity arbitrarily close to 1 by choosing a_1 and γ appropriately. We first choose the cross-cut γ in such a way that Ω'' would be Carathéodory close to Ω as seen from a_2, \dots, a_n and then put a_1 deeply inside Ω' , so that Ω' would be Carathéodory close to Ω as seen from a_1 . If two domains are Carathéodory close, then the conformal maps from these domains to \mathbb{H} mapping the marked point to i (say, with positive derivative there) are uniformly close on compact sets together with their derivatives. Thus, the lemma follows from the continuity of the half-plane functions $\langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_{\mathbb{H}}^+$ with respect to the positions of a_1, \dots, a_n . \square

Proof of Theorem 1.2. Recall that the special case $n = 2$ is already done as a part of [Theorem 1.1](#). We proceed by induction, which (together with the proof of [Lemma 2.27](#) given above) starts as follows:

$$(2.31) \quad T_2 \& L_2 \Rightarrow T_1 \Rightarrow L_3 \Rightarrow T_3 \Rightarrow L_4 \Rightarrow T_4 \Rightarrow \cdots$$

(where T_j and L_j mean the particular cases of [Theorem 1.2](#) and [Lemma 2.27](#)).

Let $n = 1$ and $\eta > 0$ be fixed. Similarly to (2.30), denote

$$R_a^{\Omega_\delta} := (\varrho(\delta))^{-\frac{1}{2}} \mathbb{E}_{\Omega_\delta}^+[\sigma_a].$$

For any $b \in \Omega$, one can write

$$(R_a^{\Omega_\delta})^2 = \frac{\mathbb{E}_{\Omega_\delta}^+[\sigma_a]}{\mathbb{E}_{\Omega_\delta}^+[\sigma_b]} \cdot \frac{\mathbb{E}_{\Omega_\delta}^+[\sigma_a \sigma_b]}{\varrho(\delta)} \cdot \frac{\mathbb{E}_{\Omega_\delta}^+[\sigma_a] \mathbb{E}_{\Omega_\delta}^+[\sigma_b]}{\mathbb{E}_{\Omega_\delta}^+[\sigma_a \sigma_b]}.$$

By [Lemma 2.27](#), if we choose b close enough to the boundary, then

$$1 - \eta \leq \frac{\mathbb{E}_{\Omega_\delta}^+[\sigma_a] \mathbb{E}_{\Omega_\delta}^+[\sigma_b]}{\mathbb{E}_{\Omega_\delta}^+[\sigma_a \sigma_b]} \leq 1,$$

provided that δ is small enough. Due to [Proposition 2.22](#) and [Theorem 1.1](#), one has

$$\frac{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a]}{\mathbb{E}_{\Omega_\delta}^+ [\sigma_b]} \xrightarrow{\delta \rightarrow 0} \frac{\langle \sigma_a \rangle_\Omega^+}{\langle \sigma_b \rangle_\Omega^+} \quad \text{and} \quad \frac{\mathbb{E}_{\Omega_\delta}^+ [\sigma_a \sigma_b]}{\varrho(\delta)} \xrightarrow{\delta \rightarrow 0} \langle \sigma_a \sigma_b \rangle_\Omega^+.$$

Consequently,

$$(1 - \eta)^3 \left(\langle \sigma_a \rangle_\Omega^+ \right)^2 \cdot \frac{\langle \sigma_a \sigma_b \rangle_\Omega^+}{\langle \sigma_a \rangle_\Omega^+ \langle \sigma_b \rangle_\Omega^+} \leq (R_a^{\Omega_\delta})^2 \leq (1 + \eta)^2 \left(\langle \sigma_a \rangle_\Omega^+ \right)^2 \cdot \frac{\langle \sigma_a \sigma_b \rangle_\Omega^+}{\langle \sigma_a \rangle_\Omega^+ \langle \sigma_b \rangle_\Omega^+}$$

provided that δ is small enough. Since η can be chosen arbitrary small, and by [\(2.26\)](#) we can make the term $\langle \sigma_a \sigma_b \rangle_\Omega^+ / \langle \sigma_a \rangle_\Omega^+ \langle \sigma_b \rangle_\Omega^+$ arbitrary close to 1 choosing b sufficiently close to $\partial\Omega$, we complete the proof for $n = 1$ by remarking that positivity of magnetization fixes the sign of $R_a^{\Omega_\delta}$.

To get the convergence of $R_{a_1, \dots, a_n}^{\Omega_\delta} := \varrho(\delta)^{-\frac{n}{2}} \mathbb{E}_{\Omega_\delta}^+ [\sigma_{a_1} \dots \sigma_{a_n}]$ for $n \geq 3$, proceed by induction and write it as

$$R_{a_1, \dots, a_n}^{\Omega_\delta} = \frac{\mathbb{E}_\Omega^+ [\sigma_{a_1} \sigma_{a_2} \dots \sigma_{a_n}]}{\mathbb{E}_\Omega^+ [\sigma_b \sigma_{a_2} \dots \sigma_{a_n}]} \cdot \frac{\mathbb{E}_\Omega^+ [\sigma_b \sigma_{a_2} \dots \sigma_{a_n}]}{\mathbb{E}_\Omega^+ [\sigma_b] \mathbb{E}_\Omega^+ [\sigma_{a_2} \dots \sigma_{a_n}]} \cdot R_b^{\Omega_\delta} R_{a_2, \dots, a_n}^{\Omega_\delta}.$$

The proof is finished similarly to the case $n = 1$: take b close to the boundary, estimate the second ratio in the right-hand side by [Lemma 2.27](#), then use convergence of the first ratio and $R_b^{\Omega_\delta} R_{a_2, \dots, a_n}^{\Omega_\delta}$ as $\delta \rightarrow 0$ (this follows from [Proposition 2.22](#) and the induction hypothesis, respectively) and asymptotics [\(2.26\)](#) for the continuous correlation functions: $\langle \sigma_b \rangle_\Omega^+ \langle \sigma_{a_2} \dots \sigma_{a_n} \rangle_\Omega^+ / \langle \sigma_b \sigma_{a_2} \dots \sigma_{a_n} \rangle_\Omega^+ \rightarrow 1$ as $b \rightarrow \partial\Omega$. □

3. Proofs of the main convergence Theorems 2.16, 2.18 and 2.20 for discrete spinors

3.1. *S*-holomorphicity of discrete observables. The notion of *s*-holomorphicity was essentially introduced by Smirnov in [\[Smi06\]](#) and used for the study of the planar critical Ising model in [\[Smi06\]](#), [\[Smi10\]](#), [\[CS12\]](#), [\[Hon10\]](#), [\[HS13\]](#), [\[CI13\]](#). Our definitions follow [\[HS13\]](#) and are equivalent to those of [\[Smi10\]](#), [\[CS12\]](#), [\[CI13\]](#) after the multiplication by \sqrt{i} ; see also [Section 3.3](#) below.

Recall that with each corner $x \in \mathcal{V}_{\mathbb{C}_\delta}^\tau$, $\tau \in \{1, i, \lambda, \bar{\lambda}\}$, we associate the line $\ell(x) = \tau\mathbb{R}$ in the complex plane, and we denote by $\mathbb{P}_{\ell(x)}$ the orthogonal projection onto that line:

$$\mathbb{P}_{\ell(x)} [w] = \frac{1}{2} (w + \tau^2 \bar{w}), \quad w \in \mathbb{C}.$$

[Definition 2.3](#) says that a function $F : \mathcal{V}_{\Omega_\delta}^{\text{cm}} \rightarrow \mathbb{C}$ is *s-holomorphic in Ω_δ* (we use the same definitions when working on double covers) if for every $x \in \mathcal{V}_{\Omega_\delta}^c$ and $z \in \mathcal{V}_{\Omega_\delta}^{\text{m}}$ that are adjacent, one has

$$F(x) = \mathbb{P}_{\ell(x)} [F(z)].$$

Remark 3.1. The set $\mathcal{V}_{\mathbb{C}_\delta}^{1,i} = \mathcal{V}_{\mathbb{C}_\delta}^1 \cup \mathcal{V}_{\mathbb{C}_\delta}^i$ may be viewed as a square lattice, divided into $\mathcal{V}_{\mathbb{C}_\delta}^1$ and $\mathcal{V}_{\mathbb{C}_\delta}^i$ in a chessboard fashion. By definition, the restriction of an s-holomorphic function F to $\mathcal{V}_{\mathbb{C}_\delta}^{1,i}$ is real on $\mathcal{V}_{\mathbb{C}_\delta}^1$ and purely imaginary on $\mathcal{V}_{\mathbb{C}_\delta}^i$. It is not difficult to check ([Smi10]) that this restriction is in fact *discrete holomorphic* in the most usual sense; that is, for any $x \in \mathcal{V}_{\mathbb{C}_\delta}^{1,i}$, one has

$$(3.1) \quad F(x + i\delta) - F(x + \delta) = i \cdot [F(x + (1+i)\delta) - F(x)].$$

The converse is also true: given a discrete holomorphic function $F : \mathcal{V}_{\mathbb{C}_\delta}^{1,i} \rightarrow \mathbb{R} \cup i\mathbb{R}$, one can first extend it to $\mathcal{V}_{\mathbb{C}_\delta}^m$ by the formula $F(z) := F(z+i\frac{\delta}{2}) + F(z-i\frac{\delta}{2})$ and then to $\mathcal{V}_{\mathbb{C}_\delta}^{\lambda,\bar{\lambda}}$ by $F(x) := \mathbf{P}_{\ell(x)} [F(x \pm \frac{\delta}{2})]$. (Due to (3.1), these projections coincide.)

We now check the s-holomorphicity of discrete spinor observables, essentially mimicking [CS12], [HS13], [CI13]. For shortness, below we use the notation

$$a_1^{\rightarrow} := a_1 + \frac{\delta}{2} \in \mathcal{V}_{\mathbb{C}_\delta}^i.$$

Proof of Proposition 2.4. Let $z \in \mathcal{V}_{[\Omega_\delta, a_1, \dots, a_n]}^m$ be a medial vertex and x be one of the four nearby corners so that $|x - z| = \frac{\delta}{2}$. We should check that

$$(3.2) \quad \mathbf{P}_{\ell(x)} [F_{[\Omega_\delta, a_1, \dots, a_n]}(z)] = F_{[\Omega_\delta, a_1, \dots, a_n]}(x),$$

where the values of $F_{[\Omega_\delta, a_1, \dots, a_n]}$ are defined as sums over the sets $\mathcal{C}_{\Omega_\delta}(a_1^{\rightarrow}, z)$ and $\mathcal{C}_{\Omega_\delta}(a_1^{\rightarrow}, x)$ respectively. There is a simple bijection $\omega_{zx} : \gamma_z \mapsto \gamma_x$ between these two sets provided by taking XOR (symmetric difference) of a configuration with the two half edges (zv) and (vx) , where v denotes the vertex that is adjacent to both x and z . Hence, it is sufficient to check that for any $\gamma_z \in \mathcal{C}_{\Omega_\delta}(a_1^{\rightarrow}, z)$, one has

$$(3.3) \quad (\cos \frac{\pi}{8})^{-1} \cdot \mathbf{P}_{\ell(x)} [\alpha_c^{\#\text{edges}(\gamma_z)} \cdot \phi_{a_1, \dots, a_n}(\gamma_z, z)] = \alpha_c^{\#\text{edges}(\gamma_x)} \cdot \phi_{a_1, \dots, a_n}(\gamma_x, x).$$

(The additional factor $(\cos \frac{\pi}{8})^{-1}$ comes from our definition of the discrete observable on medial vertices; see Remark 2.2.) There are two cases: either (zv) is contained in γ_z , which leads us to

$$\#\text{edges}(\gamma_x) = \#\text{edges}(\gamma_z), \quad \exp[-\frac{i}{2} \text{wind}(p(\gamma_x))] = e^{\pm \frac{i\pi}{8}} \exp[-\frac{i}{2} \text{wind}(p(\gamma_z))],$$

or not, which leads to

$$\begin{aligned} \#\text{edges}(\gamma_x) &= \#\text{edges}(\gamma_z) + 1, \\ \exp[-\frac{i}{2} \text{wind}(p(\gamma_x))] &= e^{\pm \frac{3i\pi}{8}} \exp[-\frac{i}{2} \text{wind}(p(\gamma_z))]. \end{aligned}$$

Let us also note that in both cases

$$\begin{aligned} (-1)^{\#\text{loops}_{a_1, \dots, a_n}(\gamma_x)} \text{sheet}_{a_1, \dots, a_n}(\gamma_x, x) \\ = (-1)^{\#\text{loops}_{a_1, \dots, a_n}(\gamma_z)} \text{sheet}_{a_1, \dots, a_n}(\gamma_z, z), \end{aligned}$$

since if ω_{zx} destroys a loop in γ_z that contributed to $\#\text{loops}_{a_1, \dots, a_n}(\gamma_z)$, leading to the change of the first factor, then this loop becomes a part of $p(\gamma_x)$, so the second factor changes simultaneously. Thus, one can factor out $\alpha_c^{\#\text{edges}(\gamma_z)}$. $\phi_{a_1, \dots, a_n}(\gamma_z, z)$ from both sides of (3.3). In the first case (3.3) readily follows, while in the second it becomes equivalent to

$$(\cos \frac{\pi}{8})^{-1} \cos \frac{3\pi}{8} = \sqrt{2} - 1 = \alpha_c.$$

Thus, $F_{[\Omega, a_1, \dots, a_n]}$ is s-holomorphic. It has multiplicative monodromy -1 around each of the marked points a_1, \dots, a_n due to the factor $\text{sheet}_{a_1, \dots, a_n}(\gamma, z)$ in (2.1). In order to prove that $F_{[\Omega, a_1, \dots, a_n]}$ obeys boundary conditions (2.3), it is sufficient to note that $\text{wind}(p(\gamma)) = \nu_{\text{out}}(z) \pmod{2\pi}$ if z is on the boundary. □

The spinor $F_{[\Omega, a_1, \dots, a_n]}$ is not defined at the corner a_1^{\rightarrow} and is not s-holomorphic there. The next lemma shows, in particular, that its values at the nearby medial vertices $a_1 + \frac{1 \pm i}{2} \delta$ have different projections onto the imaginary line $i\mathbb{R} = \ell(a_1^{\rightarrow})$, so one cannot extend $F_{[\Omega, a_1, \dots, a_n]}$ to a_1^{\rightarrow} in an s-holomorphic way.

LEMMA 3.2. *For medial vertices $a_1 + \frac{1 \pm i}{2} \delta$ taken on the same sheet of the double cover $[\Omega_\delta, a_1, \dots, a_n]$ as the “source” a_1^{\rightarrow} , one has*

$$\mathbf{P}_{i\mathbb{R}} \left[F_{[\Omega_\delta, a_1, \dots, a_n]} \left(a_1 + \frac{1 \pm i}{2} \delta \right) \right] = \mp i.$$

Proof. We consider the medial vertex $z := a_1 + \frac{1 \pm i}{2} \delta$; the opposite case is similar. Given a configuration $\gamma \in \mathcal{C}_\delta(a_1^{\rightarrow}, z)$ and applying, as above, the XOR bijection with two half-edges $(a_1^{\rightarrow}, a_1 + \delta)$ and $(a_1 + \delta, z)$, we obtain a configuration $\omega(\gamma) \in \mathcal{C}_\delta$. Since the normalizing factor $\mathcal{Z}_{\Omega_\delta}^+[a_1, \dots, a_n]$ is a sum over \mathcal{C}_δ (see (2.2)), it is sufficient to show that, for any γ ,

$$(\cos \frac{\pi}{8})^{-1} \cdot \mathbf{P}_{i\mathbb{R}} \left[\alpha_c^{\#\text{edges}(\gamma)} \phi_{a_1, \dots, a_n}(\gamma, z) \right] = -i \cdot \alpha_c^{\#\text{edges}(\omega(\gamma))} (-1)^{\#\text{loops}_{a_1, \dots, a_n}(\omega(\gamma))}.$$

Consider two cases, as in the proof of Proposition 2.4 above. If $(a_1 + \delta, z) \in \gamma$ (respectively, $(a_1 + \delta, z) \notin \gamma$), then

$$\#\text{edges}(\omega(\gamma)) = \#\text{edges}(\gamma) \quad (\text{respectively, } \#\text{edges}(\omega(\gamma)) = \#\text{edges}(\gamma) + 1).$$

We may disregard all loops in $\omega(\gamma)$ that do not contain the edge $(a_1 + \delta, a_1 + i\delta)$, as they contribute the same sign to both sides. Further, if $(a_1 + \delta, z) \in \gamma$, then $\text{wind}(p(\gamma)) = \frac{3\pi}{4} \pmod{4\pi}$, otherwise $\text{wind}(p(\gamma)) = \frac{7\pi}{4} \pmod{2\pi}$. Hence the lemma boils down to the following elementary identities:

$$(\cos \frac{\pi}{8})^{-1} \cdot \mathbf{P}_{i\mathbb{R}} \left[e^{-\frac{3\pi i}{8}} \right] = -i \quad \text{and} \quad (\cos \frac{\pi}{8})^{-1} \cdot \mathbf{P}_{i\mathbb{R}} \left[e^{-\frac{7\pi i}{8}} \right] = -i\alpha_c. \quad \square$$

3.2. *The full-plane discrete spinor $F_{[\mathbb{C}_\delta, a]}$ and its discrete primitive $G_{[\mathbb{C}_\delta, a]}$.*

This subsection is mainly devoted to the construction of the full-plane analogue of discrete spinor observables, as announced in Lemma 2.14. After this, we also prove Lemma 2.17 and the double-sided estimate (2.18) of the normalizing factor $\vartheta(\delta)$.

3.2.1. *Discrete harmonic measure in the slit plane.* We start with an important technical ingredient — the discrete Beurling estimate with optimal exponent $\frac{1}{2}$. On the square lattice, it was obtained by Kesten [Kes87] and then generalized by Lawler and Limic [LL04]. Given a face a , let $\mathbb{X}_\delta \subset \mathcal{V}_{\mathbb{C}_\delta}^1$ denote the slit discrete plane:

$$\text{Int } \mathbb{X}_\delta := \mathcal{V}_{\mathbb{C}_\delta}^1 \setminus L_a, \quad L_a := \{x + a^\rightarrow + \delta : x < 0\}.$$

(Recall that $a^\rightarrow + \delta = a + \frac{3\delta}{2} \in \mathcal{V}_{\mathbb{C}_\delta}^1$.)

LEMMA 3.3. *For all $z \in \mathbb{X}_\delta$, $A \subset \mathbb{X}_\delta$, and some absolute constant $C > 0$, the following estimates are fulfilled:*

$$(3.4) \quad \text{hm}_{\{a^\rightarrow + \delta\}}^{\mathbb{X}_\delta}(z) \leq C\delta^{\frac{1}{2}}|z - a|^{-\frac{1}{2}},$$

$$(3.5) \quad \text{hm}_A^{\mathbb{X}_\delta}(a^\rightarrow + \delta) \leq C\delta^{\frac{1}{2}}(\text{dist}(a; A))^{-\frac{1}{2}},$$

where $\text{hm}_A^{\mathbb{X}_\delta}(z)$ denotes the discrete harmonic measure of a set A in \mathbb{X}_δ viewed from z , i.e., the probability for the simple random walk on $\mathcal{V}_{\mathbb{C}_\delta}^1$ (considered as a properly scaled, rotated and shifted square grid) started at z to reach A before it hits the boundary of \mathbb{X}_δ .

Proof. This easily follows from [LL04] and simple reversibility arguments for random walks. □

Below we also need some additional estimates for the discrete harmonic measure $\text{hm}_{\{a^\rightarrow + \delta\}}^{\mathbb{X}_\delta}$ and its discrete derivatives.

LEMMA 3.4.

(i) *For $z \in \mathbb{X}_\delta$ such that $|\arg(z - a) - \pi| \leq \frac{\pi}{4}$, one has*

$$(3.6) \quad \text{hm}_{\{a^\rightarrow + \delta\}}^{\mathbb{X}_\delta}(z) \leq C\delta^{\frac{1}{2}}|\Im(z - a)||z - a|^{-\frac{3}{2}}.$$

(ii) *For all neighboring $z, z' \in \mathbb{X}_\delta$, one has*

$$(3.7) \quad \delta^{-1} \left| \text{hm}_{\{a^\rightarrow + \delta\}}^{\mathbb{X}_\delta}(z') - \text{hm}_{\{a^\rightarrow + \delta\}}^{\mathbb{X}_\delta}(z) \right| \leq C\delta^{\frac{1}{2}}|z - a|^{-\frac{3}{2}}.$$

(iii) *Being normalized by the value $\vartheta(\delta)$ at the proper lattice approximation of the point $a+1$, the functions $(\vartheta(\delta))^{-1}\text{hm}_{\{a^\rightarrow + \delta\}}^{\mathbb{X}_\delta}$ converge to $\Re[1/\sqrt{z - a}]$ as $\delta \rightarrow 0$, uniformly on compact subsets of $\mathbb{C} \setminus L_a$. Moreover, this convergence holds true in the C^1 sense meaning that the discrete derivatives (3.7) converge to the corresponding partial derivatives as well.*

Proof. (i) In order to prove (3.6), note that the probability for the random walk started at z to leave the ball of radius $\frac{1}{2}|z - a|$ around z before hitting $\partial\mathbb{X}_\delta$ is $O(|\Im(z - a)| \cdot |z - a|^{-1})$, and once this has happened the probability to hit $a^\rightarrow + \delta$ is uniformly bounded by $O(\delta^{\frac{1}{2}}|z - a|^{-\frac{1}{2}})$ due to the Beurling estimate (3.4).

(ii) The estimate (3.7) for the discrete derivatives follows from (3.4), (3.6) and the (discrete) Harnack estimate (e.g., see [CS11, Prop. 2.7]), applied to the ball of radius $\frac{1}{2}|z - a|$ (or $\frac{1}{2}|\Im(z - a)|$, if z is close to L_a) around z .

(iii) This is essentially a special case of [CS11, Th. 3.13], which claims the C^1 -convergence of discrete Poisson kernels normalized at some inner point to their continuous counterparts. The fact that our domain is unbounded plays no role here as, for any $r > 0$, the positive discrete harmonic functions $f_a^\delta := (\vartheta(\delta))^{-1} \text{hm}_{\{a^\rightarrow + \delta\}}^{\mathbb{X}_\delta}$ are uniformly bounded in the annulus $\{z : |z - a| \geq r\}$. Indeed, if f_a^δ is big at some point v in this annulus, then, by the maximum principle, f_a^δ is also big along some nearest-neighbor path running from v to $a^\rightarrow + \delta$. Since the discrete harmonic measure of such a path as seen from $a + 1$ is bounded from below (by a constant depending on r but not on δ), this leads to a contradiction with $f_a^\delta(a + 1) = 1$. \square

3.2.2. *Construction of the full-plane spinor $F_{[\mathbb{C}_\delta, a]}$.* Now we are ready to construct $F_{[\mathbb{C}_\delta, a]}$, as announced in Lemma 2.14.

Proof of Lemma 2.14. Let \mathbb{X}_δ^\pm denote two copies of the slit plane $\mathbb{X}_\delta \subset \mathcal{V}_{\mathbb{C}_\delta}^1$. We first define the real component $F_{[\mathbb{C}_\delta, a]}^1$ on \mathbb{X}_δ^\pm as

$$(3.8) \quad F_{[\mathbb{C}_\delta, a]}^1(z) = \pm \text{hm}_{\{a^\rightarrow + \delta\}}^{\mathbb{X}_\delta^\pm}(z), \quad z \in \mathbb{X}_\delta^\pm.$$

Since $F_{[\mathbb{C}_\delta, a]}^1$ vanishes identically on L_a , by identifying the upper side of this cut in \mathbb{X}_δ^+ with the lower side in \mathbb{X}_δ^- and vice versa, we obtain a function $F_{[\mathbb{C}_\delta, a]}^1$ that is defined and harmonic everywhere on $\mathcal{V}_{[\mathbb{C}_\delta, a]}^1$ except at the (two) points over $a^\rightarrow + \delta$. Recall that Lemma 3.4(iii) ensures the convergence

$$(3.9) \quad (\vartheta(\delta))^{-1} F_{[\mathbb{C}_\delta, a]}^1(z) \xrightarrow{\delta \rightarrow 0} \Re[1/\sqrt{z - a}]$$

as well as the convergence of discrete derivatives of $F_{[\mathbb{C}_\delta, a]}^1(z)$ to partial derivatives of $\Re[1/\sqrt{z - a}]$, uniformly on compact subsets of $\mathbb{C} \setminus L_a$.

We then define the imaginary component $F_{[\mathbb{C}_\delta, a]}^i : \mathcal{V}_{[\mathbb{C}_\delta, a^\rightarrow]}^i = \mathcal{V}_{[\mathbb{C}_\delta, a]}^i \setminus \{a^\rightarrow\} \rightarrow i\mathbb{R}$ as discrete harmonic conjugate of $F_{[\mathbb{C}_\delta, a]}^1$ — that is, by integrating the identity (3.1) along paths on $\mathcal{V}_{[\mathbb{C}_\delta, a^\rightarrow]}^i$ starting from, say, one of the two fibers of the point $a^\rightarrow + 2\delta$. Since $F_{[\mathbb{C}_\delta, a]}^1$ is harmonic, its discrete harmonic conjugate is well defined at least on the universal cover of $[\mathbb{C}_\delta, a^\rightarrow]$. Further, let ϖ denote some simple loop in $\mathcal{V}_{\mathbb{C}_\delta}^i$, starting and ending at $a^\rightarrow + 2\delta$, symmetric

with respect to the horizontal line $\{x : \Im \mathfrak{m}(x - a) = 0\}$, and surrounding the singularity (so, it lifts to $[\mathbb{C}_\delta, a^\rightarrow]$ as a path connecting two different fibers of $a^\rightarrow + 2\delta$). It follows from the antisymmetry of $F_{[\mathbb{C}_\delta, a]}^1$ with respect to L_a (also, note that ϖ changes the sheet once passing across the cut), that the total increment of $F_{[\mathbb{C}_\delta, a]}^i$ along ϖ is zero. Thus, $F_{[\mathbb{C}_\delta, a]}^i$ vanishes at both fibers of the point $a^\rightarrow + 2\delta$, and hence inherits the spinor property of its discrete harmonic conjugate $F_{[\mathbb{C}_\delta, a]}^1$. Moreover, for $b := a^\rightarrow + 2j\delta$, $j \geq 1$, the discrete holomorphicity [equation \(3.3\)](#) and the symmetry of $F_{[\mathbb{C}_\delta, a]}^1$ with respect to the half-line $R_a := \{x + a^\rightarrow : x > 0\}$ imply

$$F_{[\mathbb{C}_\delta, a]}^i(b + 2\delta) - F_{[\mathbb{C}_\delta, a]}^i(b) = \mp i[F_{[\mathbb{C}_\delta, a]}^1(b \pm i\delta) - 2F_{[\mathbb{C}_\delta, a]}^1(b + \delta) + F_{[\mathbb{C}_\delta, a]}^1(b + (2 \pm i)\delta)] = 0.$$

Therefore, $F_{[\mathbb{C}_\delta, a]}^i$ vanishes everywhere on R_a . Further, the estimate [\(3.7\)](#) guarantees that $F_{[\mathbb{C}_\delta, a]}^i$ is uniformly bounded: just take a path of discrete integration in the definition of $F_{[\mathbb{C}_\delta, a]}^i$ running from R_a to z along the circular arc centered at 0.

It is worth noting that $F_{[\mathbb{C}_\delta, a]}^i$ also admits a discrete harmonic measure representation similar to [\(3.8\)](#). Namely, let \mathbb{Y}_δ^\pm denote the two sheets of $\mathbb{V}_{[\mathbb{C}_\delta, a^\rightarrow]} \setminus R_a$, where signs in the notation are chosen so that $F_{[\mathbb{C}_\delta, a]}^1 > 0$ in the upper half-plane on \mathbb{Y}_δ^+ and in the lower half-plane on \mathbb{Y}_δ^- . Then,

$$(3.10) \quad F_{[\mathbb{C}_\delta, a]}^i(z) = \mp i \cdot \text{hm}_{\{a^\rightarrow\}}^{\mathbb{Y}_\delta^\pm}(z), \quad z \in \mathbb{Y}_\delta^\pm.$$

Indeed, on each of the sheets \mathbb{Y}_δ^\pm one can further extend $F_{[\mathbb{C}_\delta, a]}^i$ (as a harmonic conjugate of $F_{[\mathbb{C}_\delta, a]}^1$) to the point a^\rightarrow : the only obstruction to do so on $[\mathbb{C}_\delta, a]$ was that the increment of $F_{[\mathbb{C}_\delta, a]}^i$ along the smallest loop surrounding $a^\rightarrow + \delta$ would be nonzero as $F_{[\mathbb{C}_\delta, a]}^1$ is not harmonic at $a^\rightarrow + \delta$, but now this loop intersects the cut R_a . Since the bounded function $F_{[\mathbb{C}_\delta, a]}^i$ is harmonic on \mathbb{Y}_δ^\pm and vanishes on R_a , it is proportional to $\text{hm}_{\{a^\rightarrow\}}^{\mathbb{Y}_\delta^\pm}$. Finally, it follows from symmetry arguments that

$$\begin{aligned} \pm \text{hm}_{\{a^\rightarrow\}}^{\mathbb{Y}_\delta^\pm}(a^\rightarrow + (1+i)\delta) - (\pm 1) &= \pm \text{hm}_{\{a^\rightarrow + \delta\}}^{\mathbb{X}_\delta^\pm}(a^\rightarrow + i\delta) - (\pm 1) \\ &= F_{[\mathbb{C}_\delta, a]}^1(a^\rightarrow + i\delta) - F_{[\mathbb{C}_\delta, a]}^1(a^\rightarrow + \delta) \\ &= i \cdot [F_{[\mathbb{C}_\delta, a]}^i(a^\rightarrow + (1+i)\delta) - F_{[\mathbb{C}_\delta, a]}^i(a^\rightarrow)], \end{aligned}$$

which fixes the multiplicative constant $\mp i$ in [\(3.10\)](#).

Similarly to [\(3.9\)](#), we have

$$(3.11) \quad (\vartheta(\delta))^{-1} F_{[\mathbb{C}_\delta, a]}^i(z) \xrightarrow[\delta \rightarrow 0]{} \Im \mathfrak{m}[1/\sqrt{z - a}]$$

together with the convergence of discrete derivatives, uniformly on compact subsets of $\mathbb{C} \setminus R_a$. Now, following Remark 3.1, we extend $F_{[\mathbb{C}_\delta, a]}$ to $\mathcal{V}_{[\mathbb{C}, a]}^m$ and $\mathcal{V}_{[\mathbb{C}, a]}^{\lambda, \bar{\lambda}}$ as an s-holomorphic function. The convergence (2.16) readily follows from (3.9) and (3.11). (For points lying near L_a , one can approximate $F_{[\mathbb{C}_\delta, a]}^1$ by summing discrete derivatives of $F_{[\mathbb{C}_\delta, a]}^i$ and vice versa near R_a .) \square

The following lemma explains the future role of $F_{[\mathbb{C}_\delta, a]}$: it has the same “discrete singularity” (whatever it means) at a^\rightarrow as $F_{[\Omega_\delta, a_1, \dots, a_n]}$ has at a_1^\rightarrow . This allows one to handle this singularity by taking the difference $F_{[\Omega_\delta, a_1, \dots, a_n]} - F_{[\mathbb{C}_\delta, a_1]}$

LEMMA 3.5. *The spinor $F_{[\Omega_\delta, a_1, \dots, a_n]}^\dagger := F_{[\Omega_\delta, a_1, \dots, a_n]} - F_{[\mathbb{C}_\delta, a_1]}$, extended to be zero at a_1^\rightarrow , is s-holomorphic in any simply connected neighborhood of a_1 that does not contain other branching points a_2, \dots, a_n .*

Proof. By Lemma 3.2, it suffices to check that $P_{i\mathbb{R}} [F_{[\mathbb{C}_\delta, a_1]}(a_1 + \frac{1\pm i}{2}\delta)] = \mp i$ on the sheet \mathbb{X}_δ^+ . The considerations given in the previous proof show that, being considered on \mathbb{Y}_δ^\pm instead of \mathbb{X}_δ^+ , the spinor $F_{[\mathbb{C}_\delta, a_1]}$ can be extended at a_1^\rightarrow in an s-holomorphic way and

$$P_{i\mathbb{R}} [F_{[\mathbb{C}_\delta, a_1]}(a_1 + \frac{1\pm i}{2}\delta)] = F_{[\mathbb{C}_\delta, a_1]}^i(a_1^\rightarrow) = \mp i \text{ on } \mathbb{Y}_\delta^\pm.$$

However, by definition, \mathbb{X}_δ^+ coincides with \mathbb{Y}_δ^+ in the upper half-plane and with \mathbb{Y}_δ^- in the lower half-plane. \square

3.2.3. *Construction of the harmonic spinor $G_{[\mathbb{C}_\delta, a]}$.* Now we “integrate” the real component of $F_{[\mathbb{C}_\delta, a]}$ in order to construct a discrete counterpart of the harmonic spinor $\Re \sqrt{z - a}$, as announced in Lemma 2.17. Along the way, we also prove the double-sided estimate (2.18) of the normalizing factor $\vartheta(\delta)$. Note that the upper bound $\vartheta(\delta) \leq C_+ \sqrt{\delta}$ directly follows from the discrete Beurling estimate (3.4), so we need to prove the lower bound only.

Proof of Lemma 2.17 and the estimate (2.16). We use the notation for the sheets and the cuts of $[\mathbb{C}_\delta, a]$ introduced above. For $z \in \mathbb{X}_\delta^+ \subset \mathcal{V}_{[\mathbb{C}_\delta, a]}^1$, define

$$(3.12) \quad G_{[\mathbb{C}_\delta, a]}(z) := \delta \cdot \sum_{j=0}^\infty F_{[\mathbb{C}_\delta, a]}^1(z - 2j\delta).$$

Due to the estimate (3.6) for $F_{[\mathbb{C}_\delta, a]}^1(\cdot) = \text{hm}_{\{a^\rightarrow + \delta\}}^{\mathbb{X}_\delta^1}(\cdot)$, this sum always converges. Note that $F_{[\mathbb{C}_\delta, a]}^1 = 0$ on the cut L_a , so $G_{[\mathbb{C}_\delta, a]}$ vanishes on L_a too.

We are going to prove that $G_{[\mathbb{C}_\delta, a]}$ is harmonic *everywhere* inside \mathbb{X}_δ^+ , including the point $a^\rightarrow + \delta$ right near the cut L_a . For z outside R_a , this harmonicity readily follows from the harmonicity of $F_{[\mathbb{C}_\delta, a]}^1$, so let $z \in R_a$. By

$\mathcal{S}_N(z)$ denote

$$\text{Int } \mathcal{S}_N(z) := \mathcal{V}_{\mathbb{C}_\delta}^1 \cap \{z : |\Re(w - z)|, |\Im(w - z)| \leq 2N\delta\}$$

a sufficiently large square centered at z , and write the discrete Green formula:

$$\begin{aligned} \sum_{j=\lfloor z/2\delta \rfloor}^N \Delta F_{[\mathbb{C}_\delta, a]}^1(z - 2j\delta) &= \sum_{w \in \mathcal{S}_N(z)} \Delta F_{[\mathbb{C}_\delta, a]}^1(w) \\ &= \sum_{(ww') \in \partial \mathcal{S}_N(z)} (F_{[\mathbb{C}_\delta, a]}^1(w') - F_{[\mathbb{C}_\delta, a]}^1(w)). \end{aligned}$$

Let N be large enough so that $|w - a| \geq N\delta$ for all boundary edges of the square $\mathcal{S}_N(z)$. Then, it immediately follows from the estimate (3.7) that the last sum is $O(N \cdot \delta^{\frac{3}{2}}(N\delta)^{-\frac{3}{2}}) = O(N^{-\frac{1}{2}})$. Passing to the limit as $N \rightarrow \infty$, we conclude that

$$\Delta G_{[\mathbb{C}_\delta, a]}(z) = \delta \cdot \sum_{j=\lfloor z/2\delta \rfloor}^\infty \Delta F_{[\mathbb{C}_\delta, a]}^1(z - 2j\delta) = 0.$$

Thus, $G_{[\mathbb{C}_\delta, a]}$ is indeed discrete harmonic in \mathbb{X}_δ^+ . Moreover, since it vanishes on L_a , one can harmonically extend $G_{[\mathbb{C}_\delta, a]}$ to the double cover $\mathcal{V}_{[\mathbb{C}_\delta, a]}^1$ by symmetry.

Let $\nu(\delta)$ denote the value of $G_{[\mathbb{C}_\delta, a]}$ in a lattice approximation of the point $a + 1$. Arguing as in the proof of Lemma 3.4(iii), we see that, uniformly on compact subsets of $\mathbb{C} \setminus L_a$,

$$(3.13) \quad (\nu(\delta))^{-1} G_{[\mathbb{C}_\delta, a]}(z) \xrightarrow{\delta \rightarrow 0} \Re \sqrt{z - a}.$$

Indeed, since $G_{[\mathbb{C}_\delta, a]}$ vanishes everywhere on L_a and remains bounded near 0 by the maximum principle, in this case the limiting positive harmonic function should be proportional to $\Re \sqrt{z - a}$, and the multiplicative normalization is fixed at $a + 1$. Moreover, a similar convergence holds true for discrete derivatives, yielding

$$(\nu(\delta))^{-1} \cdot \frac{1}{2} F_{[\mathbb{C}_\delta, a]}^1(z) \xrightarrow{\delta \rightarrow 0} \partial_x \Re \sqrt{z - a} = \frac{1}{2} \Re \left[\frac{1}{\sqrt{z - a}} \right].$$

In particular, $\nu(\delta) \sim \vartheta(\delta)$ as $\delta \rightarrow 0$, which allows us to give a simple proof of the lower bound in (2.18): as discrete harmonic functions $(\nu(\delta))^{-1} G_{[\mathbb{C}_\delta, a]}$ are uniformly bounded near the unit circle around a and vanish identically on L_a , the discrete Beurling estimate (3.5) implies

$$(\nu(\delta))^{-1} \cdot \delta = (\nu(\delta))^{-1} G_{[\mathbb{C}_\delta, a]}(a^\rightarrow + \delta) \leq C\sqrt{\delta}.$$

Finally, the convergence (3.13) near L_a follows from the convergence of $F_{[\mathbb{C}_\delta, a]}^1(z)$, since the tails in (3.12) are uniformly small due to estimate (3.6). \square

3.3. *The boundary value problem for spinors.* In this section we reformulate the Riemann-type boundary value problem for holomorphic spinors (both discrete and continuous) using primitives (antiderivatives) of their squares. Note that this approach is not completely straightforward, since the square of a discrete holomorphic function, in general, does not have a discrete primitive. However, it was noted in [Smi06] that one can naturally define the real part of this primitive using the *s-holomorphicity* of observables. Moreover, a technique developed in [CS12] (see, in particular, Sections 3.4 and 3.5 therein) allows one to treat this real part essentially as if it were a harmonic function. Below, we summarize the tools we will use. We warn the reader that all our definitions are equivalent to those of [CS12] *after the multiplication of the spinor by \sqrt{i}* , which means that imaginary parts, sub-/super-harmonicity and positivity of functions and their (inner) normal derivatives used in [CS12] should be replaced by real parts, super-/sub-harmonicity and negativity, respectively.

3.3.1. *Discrete integration of the squared spinor observables.* Let Δ_δ° be the standard (unnormalized) discrete Laplacian acting on functions $H_\delta^\circ : \mathcal{V}_{\Omega_\delta}^\circ \rightarrow \mathbb{R}$ (which are defined on faces of Ω_δ):

$$\Delta_\delta^\circ H_\delta^\circ(z) := \sum_{w \sim z} (H_\delta^\circ(w) - H_\delta^\circ(z)), \quad z \in \text{Int}\mathcal{V}_{\Omega_\delta}^\circ,$$

where the sum is over the four neighbors $w \in \mathcal{V}_{\Omega_\delta}^\circ$ of z . Similarly, for functions $H_\delta^\bullet : \mathcal{V}_{\Omega_\delta}^\bullet \rightarrow \mathbb{R}$ defined on vertices of Ω_δ , let Δ_δ^\bullet denote the slightly modified discrete Laplacian:

$$\Delta_\delta^\bullet H_\delta^\bullet(z) := \sum_{w \sim z} c_{zw} \cdot (H_\delta^\bullet(w) - H_\delta^\bullet(z)), \quad z \in \text{Int}\mathcal{V}_{\Omega_\delta}^\bullet,$$

where the conductance c_{zw} is equal to 1 for inner edges (i.e., for $w \in \text{Int}\mathcal{V}_{\Omega_\delta}^\bullet$) and $c_{zw} = 2(\sqrt{2} - 1)$ for the boundary edges. (See Section 3.6 in [CS12] or [DCHN11] for the reason of this “boundary modification” of Δ_δ^\bullet .)

PROPOSITION 3.6. *For an s-holomorphic spinor observable*

$$F_\delta = F_{[\Omega_\delta, a_1, \dots, a_n]} : [\Omega_\delta, a_1^\rightarrow, \dots, a_n] \rightarrow \mathbb{C}$$

satisfying boundary conditions (2.3), one can define a real-valued function $H_\delta = H_{[\Omega_\delta, a_1, \dots, a_n]} : \mathcal{V}_{\Omega_\delta}^\circ \rightarrow \mathbb{R}$ that is a discrete analogue of the primitive $\Re \int (F_\delta(z))^2 dz$, so that the following properties are fulfilled:

- For any adjacent $w \in \mathcal{V}_{\Omega_\delta}^\circ$ and $v \in \mathcal{V}_{\Omega_\delta}^\bullet$, one has

$$(3.14) \quad H_\delta^\circ(w) - H_\delta^\bullet(v) = 2\delta \left| F_\delta \left(\frac{1}{2}(w+v) \right) \right|^2,$$

where $\frac{1}{2}(w+v) \in \mathcal{V}_{\Omega_\delta}^\circ$ is the corner between v and w . (In accordance with Lemma 3.2, we set $|F_\delta(a_1^\rightarrow)| := 1$ in the case $w = a_1$ and $v = a_1 + \delta$.)

- H_δ satisfies Dirichlet boundary conditions: $H_\delta^\circ(w) = 0$ for any $w \in \partial\mathcal{V}_{\Omega_\delta}^\circ$, and $H_\delta^\bullet(v) = 0$ for any $v \in \partial\mathcal{V}_{\Omega_\delta}^\bullet$.
- H_δ^\bullet has a “nonpositive inner normal derivative,” i.e., $H_\delta(v) \leq 0$ for any vertex $v \in \mathcal{V}_{\Omega_\delta}^\bullet$ adjacent to a boundary vertex.
- H_δ° is Δ_δ° -subharmonic on $\mathcal{V}_{\Omega_\delta}^\circ \setminus \{a_1, \dots, a_n\}$, while H_δ^\bullet is Δ_δ^\bullet -superharmonic on $\mathcal{V}_{\Omega_\delta}^\bullet \setminus \{a_1 + \delta\}$.

Proof. All the claims follow directly from the results of Section 3.3 in [CS12]. Since all listed properties are local, the spinor nature of F_δ plays no role here. (Note that the right-hand side of (3.14) does not depend on the sheet.) □

Remark 3.7. By construction, $H_\delta^\circ(w) \geq H_\delta^\bullet(v)$ for adjacent w and v . Combined with the sub-/super-harmonicity, this implies a maximum principle for H_δ : if $\Omega'_\delta \subset \Omega_\delta$ does not contain $a_1 + \delta$ (respectively, any of a_1, \dots, a_n), then

$$\min_{\Omega'_\delta} H_\delta = \min_{\partial\Omega'_\delta} H_\delta^\bullet \quad (\text{respectively, } \max_{\Omega'_\delta} H_\delta = \max_{\partial\Omega'_\delta} H_\delta^\circ).$$

Moreover, if $\text{hm}_A(z)$ denotes the discrete harmonic measure of a set A in Ω'_δ viewed from z , then

$$\begin{aligned} H_\delta(z) &\geq (1 - \text{hm}_A^\bullet(z)) \min_{\partial\Omega'_\delta} H_\delta^\bullet + \text{hm}_A^\bullet(z) \min_A H_\delta^\bullet, \\ H_\delta(z) &\leq (1 - \text{hm}_A^\circ(z)) \max_{\partial\Omega'_\delta} H_\delta^\circ + \text{hm}_A^\circ(z) \max_A H_\delta^\circ. \end{aligned}$$

Remark 3.8. The subharmonicity of H_δ° fails at the faces a_1, \dots, a_n because $F_{[\Omega_\delta, a_1 \dots a_n]}$ branches at those points, while the superharmonicity of H_δ^\bullet fails at the vertex $a_1 + \delta$ because of the discrete singularity of $F_{[\Omega_\delta, a_1 \dots a_n]}$, which is not defined at a_1^\rightarrow . Due to Lemma 3.5, one can remove the latter singularity by subtracting the full-plane observable $F_{[\mathbb{C}_\delta, a_1]}$. Then, the function

$$H_\delta^\dagger = H_{[\Omega_\delta, a_1 \dots, a_n]}^\dagger := \Re \int (F_\delta^\dagger(z))^2 dz, \quad F_\delta^\dagger := F_{[\Omega_\delta, a_1 \dots, a_n]} - F_{[\mathbb{C}_\delta, a_1]}$$

(defined in the same way as H_δ accordingly to Proposition 3.6) is subharmonic on $\mathcal{V}_{\Omega_\delta}^\circ \setminus \{a_1\}$ and superharmonic on $\mathcal{V}_{\Omega_\delta}^\bullet$ everywhere in a neighborhood of a_1 . Moreover, since $F_\delta^\dagger(a_1^\rightarrow) = 0$ on both sheets, the values of F_δ^\dagger at the nearby corners $a_1 \pm \frac{\delta}{2}, a_1 \pm \frac{i\delta}{2}$ and midedges $a_1 \pm \frac{1 \pm i}{2}\delta$ satisfy the s-holomorphicity conditions as if F_δ^\dagger were nonbranching at a_1 . Thus, H_δ^\dagger is subharmonic at the point $a_1 \in \mathcal{V}_{\Omega_\delta}^\circ$ too.

3.3.2. Integration of squared spinors in the continuous setup. We now give a characterization of the continuous spinors solving the boundary value problem (2.7)–(2.9) in terms of the primitives of their squares. In the next section, this characterization will be used in the proofs of the main convergence results.

PROPOSITION 3.9. *Let Ω be a simply connected domain, and suppose a holomorphic spinor $f = f_{[\Omega, a_1, \dots, a_n]}$ solves the boundary value problem (2.7)–(2.9) (or is defined according to Remark 2.10 if Ω is not smooth). Define two harmonic functions*

$$h := \Re \int (f(z))^2 dz \quad \text{and} \quad h^\dagger := \Re \int (f(z) - f_{[\mathbb{C}, a_1]}(z))^2 dz,$$

where $f_{[\mathbb{C}, a_1]}(z) := 1/\sqrt{z - a_1}$. Then, the following holds true:

- (1) h is a single-valued function in $\Omega \setminus \{a_1, \dots, a_n\}$, continuous up to $\partial\Omega$, which satisfies Dirichlet boundary conditions $h \equiv \text{const}$ on $\partial\Omega$. (Since h is defined up to an additive constant, below we assume that $h \equiv 0$ on $\partial\Omega$.)
- (2) There is no point z_0 on $\partial\Omega$ such that $h(z) \geq 0$ in a neighborhood of z_0 .
- (3) h is bounded from below in a neighborhood of each a_2, \dots, a_n .
- (4) h^\dagger is single-valued and bounded in a neighborhood of a_1 .

Moreover, if h and h^\dagger satisfy (1)–(4), then f solves the problem (2.7)–(2.9).

Proof. Note that, being integrated, the covariance property (2.10) claims the conformal invariance of both h and h^\dagger . As the properties (1)–(4) are preserved under conformal mappings too, we will further assume that $\partial\Omega$ is smooth. Note that the property (2.7) is equivalent to (1) and (2): it states that $f^2(z)\nu_{\text{out}}(z)$ is nonnegative on the boundary, which yields $\partial_{\tau(z)}h \equiv 0$ (where $\tau(z)$ denotes a tangent vector) and $\partial_{\nu_{\text{out}}(z)}h \geq 0$ everywhere on $\partial\Omega$. A straightforward integration of the asymptotics of f near a_k given by the conditions (2.8) and (2.9) yields

$$(3.15) \quad \begin{aligned} h^\dagger(z) &= O(1), & z &\rightarrow a_1, \\ h(z) &= -c_k \log|z - a_k| + O(1), & z &\rightarrow a_k, \quad 2 \leq k \leq n \end{aligned}$$

for some $c_k \geq 0$, giving (3) and (4). Vice versa, (3), (4) imply (2.9), (2.8) by differentiating and taking the square root. □

3.4. Convergence of discrete observables away from singularities. We now prove the convergence of the (properly normalized) discrete spinor observables to their continuous counterparts on compact subsets of $\Omega \setminus \{a_1, \dots, a_n\}$.

Proof of Theorem 2.16. Let the discrete integrals

$$H_\delta := \Re \int \left((\vartheta(\delta))^{-1} F_\delta(z) \right)^2 dz, \quad F_\delta = F_{[\Omega_\delta, a_1, \dots, a_n]},$$

be defined on $\mathcal{V}_{\Omega_\delta}^\bullet$ accordingly to Proposition 3.6. Given $\epsilon > 0$, denote

$$\Omega_\delta(\epsilon) := \Omega_\delta \cap \{z : \text{dist}(z; \{a_1, \dots, a_n\}) > \epsilon\}.$$

Assume for a moment that

$$(3.16) \quad \begin{aligned} &\text{for any } \epsilon > 0, \text{ the functions } H_\delta \text{ remain uniformly} \\ &\text{bounded on } \Omega_\delta(\epsilon) \text{ by some constant } C(\epsilon) \text{ as } \delta \rightarrow 0. \end{aligned}$$

Then by [CS12, Th. 3.12], the functions $(\vartheta(\delta))^{-1}F_\delta$ are equicontinuous on $\Omega_\delta(\epsilon)$ for any $\epsilon > 0$. Therefore, by passing to a subsequence and applying the diagonal process, we can assume that $(\vartheta(\delta))^{-1}F_\delta$ tends to a limit \tilde{f} and $H_\delta \rightarrow \tilde{h} := \Re \int \tilde{f}^2$ uniformly on compact subsets of $\Omega \setminus \{a_1, \dots, a_n\}$. Our goal is to check that \tilde{f} satisfies the properties (1)–(4) given in Proposition 3.9. Then, the uniqueness of a solution to the boundary value problem (2.7)–(2.9) proven in Lemma 2.9 implies $\tilde{f} = f_{[\Omega, a_1, \dots, a_n]}$.

Clearly, \tilde{f} is a holomorphic spinor on $[\Omega, a_1, \dots, a_n]$. By superharmonicity of H_δ^\bullet , for $2 \leq k \leq n$, we have

$$\min_{|z-a_k| \leq \epsilon} H_\delta(z) \geq \min_{\epsilon < |z-a_k| \leq 2\epsilon} H_\delta^\bullet(z);$$

thus \tilde{h} is bounded from below in the neighborhoods of a_2, \dots, a_n , so property (3) holds true. By the maximum principle for H_δ (see Remark 3.7), taking into account that $H_\delta \equiv 0$ on $\partial\Omega_\delta$, we have that

$$(3.17) \quad \begin{aligned} H_\delta^\bullet(z) &\geq -C(\epsilon)(1 - \text{hm}_{\partial\Omega_\delta}^{\Omega_\delta(\epsilon)}(z)), \\ H_\delta^\circ(z) &\leq C(\epsilon)(1 - \text{hm}_{\partial\Omega_\delta}^{\Omega_\delta(\epsilon)}(z)) \end{aligned}$$

Since $\text{hm}_{\partial\Omega_\delta}^{\Omega_\delta(\epsilon)}(z) \rightarrow 1$ uniformly in δ as z approaches the boundary of Ω_δ , this implies $\tilde{h} \equiv 0$ on $\partial\Omega$, giving (1). Moreover, thanks to Remark 6.3 in [CS12], we also have (2): there is no point on $\partial\Omega$ such that $\tilde{h} \geq 0$ in its neighborhood.

Now consider the function discussed (up to normalization) in Remark 3.8:

$$H_\delta^\dagger := \Re \int ((\vartheta(\delta))^{-1}(F_\delta(z) - F_{[\mathbb{C}, a_1]}(z)))^2 dz,$$

which is well defined in the disc $\{z : |z - a_1| < r\}$ provided that r is small enough. Since $(\vartheta(\delta))^{-1}F_\delta$ and $(\vartheta(\delta))^{-1}F_{[\mathbb{C}, a_1]}$ converge to \tilde{f} and $f_{[\mathbb{C}, a_1]}$, respectively (see Lemma 2.14), uniformly on compact subsets of $A_r := \{z : 0 < |z - a_1| < r\}$, we conclude that H_δ^\dagger converges to $\tilde{h}^\dagger := \Re \int (\tilde{f}(z) - f_{[\mathbb{C}, a_1]}(z))^2 dz$ everywhere in A_r . By Remark 3.8 (sub-/super-harmonicity of H_δ^\dagger on $\mathcal{V}^\circ/\mathcal{V}^\bullet$ near the point a_1), the functions H_δ^\dagger are uniformly bounded in A_r , so \tilde{h}^\dagger is bounded in A_r too, which concludes the proof of the property (4). Therefore, $\tilde{f} = f_{[\Omega, a_1, \dots, a_n]}$.

It remains to justify (3.16). On the contrary, suppose that

$$M_\delta(\epsilon) := \max_{\Omega_\delta(\epsilon)} |H_\delta| \xrightarrow{\delta \rightarrow 0} \infty$$

for some $\epsilon > 0$ and along some subsequence of δ 's. Then the re-normalized functions $(M_\delta(\epsilon))^{-1}H_\delta$ are uniformly bounded in $\Omega_\delta(\epsilon)$, and thus $(M_\delta(\epsilon))^{-1/2} \cdot (\vartheta(\delta))^{-1}F_\delta$ and $(M_\delta(\epsilon))^{-2}H_\delta$ have (taking subsequences) the limits \tilde{f} and $\tilde{h} = \Re \int \tilde{f}^2$, which are holomorphic and harmonic in $\Omega_\delta(\epsilon)$, respectively. An

important observation, proven in [Lemma 3.10](#) below, is that \tilde{h} cannot be identically zero. In particular, for any $0 < \epsilon' < \epsilon$, we have $M_\delta(\epsilon') \leq CM_\delta(\epsilon)$ with some $C = C(\epsilon', \epsilon)$ independent of δ .

Applying the diagonal procedure, we may assume $(M_\delta(\epsilon))^{-1/2} \cdot (\vartheta(\delta))^{-1} F_\delta$ tends to a limit \tilde{f} (and $(M_\delta(\epsilon))^{-1} H_\delta$ tends to $\tilde{h} = \Re \int \tilde{f}^2$) uniformly on *each* of $\Omega_\delta(\epsilon')$. Arguing as above, we see that \tilde{h} is harmonic in $\Omega \setminus \{a_1, \dots, a_n\}$, satisfies the Dirichlet boundary conditions, has nonnegative outer normal derivative, and is bounded from below near a_2, \dots, a_n . Moreover, repeating the last step of the proof given above we see that the function

$$\begin{aligned} \tilde{h}^\dagger &= \lim_{\delta \rightarrow 0} (M_\delta(\epsilon))^{-1} \Re \int ((\vartheta(\delta))^{-1} (F_\delta(z) - F_{[\mathbb{C}_\delta, a_1]}(z)))^2 dz \\ &= \lim_{\delta \rightarrow 0} (M_\delta(\epsilon))^{-1} \Re \int ((\vartheta(\delta))^{-1} F_\delta(z))^2 dz = \tilde{h} \end{aligned}$$

is also bounded in a neighborhood of a_1 . (One can neglect $F_{[\mathbb{C}_\delta, a_1]}(z)$ in the last expression since $(\vartheta(\delta))^{-1} F_{[\mathbb{C}_\delta, a_1]}(z)$ tends to $f_{[\mathbb{C}, a_1]}$ and $M_\delta(\epsilon) \rightarrow \infty$.) Thus, \tilde{h} is bounded from below near *all* a_1, \dots, a_n and has a nonnegative outer normal derivative. This contradicts the maximum principle since \tilde{h} does not vanish identically. \square

LEMMA 3.10. *In the notation of the proof above, none of the subsequential limits of $(M_\delta(\epsilon))^{-1} H_\delta$ is identically zero in $\Omega(\epsilon)$.*

Proof. Suppose by contradiction that $(M_\delta(\epsilon))^{-1} H_\delta \rightarrow 0$ uniformly on compact subsets of $\Omega_\delta(\epsilon)$. Let z_δ^{\max} be the point of $\Omega_\delta(\epsilon)$ where the maximum of $|H_\delta|$ is attained. Since H_δ vanishes on $\partial\Omega_\delta$, the sub-/super-harmonicity of H_δ implies that z_δ^{\max} belongs to one of the “discrete circles”

$$\varpi_k(\epsilon) := \{z : \epsilon \leq |z - a_k| \leq \epsilon + 5\delta\}$$

of radius ϵ around a_k . Passing to a subsequence, one can assume that $z_\delta^{\max} \rightarrow z^{\max}$. Recall that $H^\circ \geq H^\bullet$ at adjacent points (see [\(3.14\)](#)). Hence, either $z_\delta^{\max} \in \mathcal{V}_{\Omega_\delta}^\circ$ and $M_\delta(\epsilon) = H_\delta^\circ(z_\delta^{\max})$, or $z_\delta^{\max} \in \mathcal{V}_{\Omega_\delta}^\bullet$ and $M_\delta(\epsilon) = -H_\delta^\bullet(z_\delta^{\max})$.

Suppose that $z_\delta^{\max} \in \varpi_k(\epsilon)$ for some $2 \leq k \leq n$. Denote

$$m_\delta(2\epsilon) := \min_{z: |z - a_k| \leq 2\epsilon} H_\delta^\bullet(z).$$

As H_δ^\bullet is superharmonic inside of the disc bounded by $\varpi_k(2\epsilon)$ and $(M_\delta(\epsilon))^{-1} H_\delta^\bullet$ tends to zero uniformly on $\varpi_k(2\epsilon)$ by our assumption, we have

$$(3.18) \quad (M_\delta(\epsilon))^{-1} \cdot m_\delta(2\epsilon) \xrightarrow{\delta \rightarrow 0} 0.$$

Therefore, $z_\delta^{\max} \in \mathcal{V}_{\Omega_\delta}^\circ$. By the subharmonicity of H_δ° , we can find a discrete nearest-neighbor path $\gamma^\circ := \{z_\delta^{\max} = z_1 \sim z_2 \sim \dots\} \subset \mathcal{V}_{\Omega_\delta}^\circ$ with $M_\delta(\epsilon) \leq \dots \leq H_\delta^\circ(z_j) \leq H_\delta^\circ(z_{j+1}) \leq \dots$, which may only end up at a_k , where the subharmonicity fails. By [\[CS12, Rem. 3.10\]](#), the functions $H_\delta^\bullet - m_\delta(2\epsilon)$

and $H_\delta^\circ - m_\delta(2\epsilon)$ are uniformly comparable at adjacent points inside $\varpi_k(2\epsilon)$. Taking into account (3.18), this implies $H_\delta^\bullet(z) \geq cM_\delta(\epsilon)$ for some absolute constant $c > 0$ and all $z \in \gamma^\bullet$, where γ^\bullet is the set of vertices adjacent to γ° . Further, the maximum principle gives

$$H_\delta^\bullet(z) \geq cM_\delta(\epsilon)\text{hm}_\gamma^\bullet(z) + (1 - \text{hm}_\gamma^\bullet(z))m_\delta(2\epsilon) \text{ for } z : |z - a_k| \leq 2\epsilon,$$

where hm_γ^\bullet denotes the discrete harmonic measure of γ in $\{z : |z - a_k| \leq 2\epsilon\}$. It follows from the discrete Beurling estimate that $\text{hm}_\gamma^\bullet(z) \geq \frac{1}{2}$ if z is chosen close enough (but at fixed distance that is independent of δ) to z^{\max} where the path γ starts. For such z , one has $(M_\delta(\epsilon))^{-1}H_\delta^\bullet(z) \geq \frac{1}{2}c + o(1)$ as $\delta \rightarrow 0$, and hence the limit of $(M_\delta(\epsilon))^{-1}H_\delta(z)$ cannot be identically zero in $\Omega(\epsilon)$.

It remains to treat the case $z_\delta^{\max} \in \varpi_1(\epsilon)$. Consider the function

$$(M_\delta(\epsilon))^{-1}H_\delta^\dagger = (M_\delta(\epsilon))^{-1}\Re \int ((\vartheta(\delta))^{-1}(F_\delta(z) - F_{[\mathbb{C}_\delta, a_1]}(z)))^2 dz.$$

Note that it tends to zero on $\varpi_1(2\epsilon)$ since one can neglect the term $F_{[\mathbb{C}_\delta, a_1]}(z)$. (Recall that $(\vartheta(\delta))^{-1}F_{[\mathbb{C}_\delta, a_1]}(z) \rightarrow f_{[\mathbb{C}, a_1]}$ and $M_\delta(\epsilon) \rightarrow \infty$.) Therefore, by Remark 3.8 and the maximum principle, it also tends to zero in a neighborhood of $\varpi_1(\epsilon)$, so we consequently derive that each of the functions

$$\frac{F_\delta - F_{[\mathbb{C}_\delta, a_1]}}{(M_\delta(\epsilon))^{1/2}\vartheta(\delta)}, \quad \frac{F_\delta}{(M_\delta(\epsilon))^{1/2}\vartheta(\delta)} \quad \text{and} \quad \frac{H_\delta}{M_\delta(\epsilon)}$$

tends to zero uniformly on $\varpi_1(\epsilon)$. In particular, $1 = (M_\delta(\epsilon))^{-1}|H_\delta(z_\delta^{\max})| \rightarrow 0$, which is a contradiction. □

3.5. Analysis near the singularities. We now pass to the most delicate part of our analysis: matching the second-order terms in the values $F_\delta(a_1 + \frac{3\delta}{2})$ with the second coefficient in the expansion of the continuous spinor near a_1 . For shortness, below we use the notation $a := a_1$, $F_\delta = F_{[\Omega_\delta, a, a_2, \dots, a_n]}$ and $\mathcal{A} = \mathcal{A}_{[\Omega, a, a_2, \dots, a_n]}$.

Proof of Theorem 2.18. Let \mathcal{R} denote the reflection with respect to the horizontal line $\{x : \Im(x - a) = 0\}$, and let Λ_δ be a small neighborhood of a in $\Omega_\delta \cap \mathcal{R}(\Omega_\delta)$. Recall the notation $L_a = \{x + a + \frac{3\delta}{2} : x < 0\}$, and denote by $\Lambda_\delta^+ \subset \mathcal{V}_{\Lambda_\delta}^1$ one of two sheets of $[\Lambda_\delta, a] \setminus L_a$ such that $F_\delta(a + \frac{3\delta}{2}) > 0$ on Λ_δ^+ . We define a real-valued function $S_\delta : \Lambda_\delta^+ \rightarrow \mathbb{R}$ by

$$S_\delta := (\vartheta(\delta))^{-1}(\frac{1}{2}(F_\delta + F_\delta^{(\mathcal{R})}) - F_{[\mathbb{C}_\delta, a]} - 2\Re \mathcal{A} \cdot G_{[\mathbb{C}_\delta, a]}),$$

where $F_\delta^{(\mathcal{R})} = F_{[\mathcal{R}(\Omega_\delta), a, \mathcal{R}(a_2), \dots, \mathcal{R}(a_n)]}$ and the functions $F_{[\mathbb{C}_\delta, a]}$, $G_{[\mathbb{C}_\delta, a]}$ were constructed in Section 3.2. By symmetry, one has $F_\delta^{(\mathcal{R})}(a + \frac{3\delta}{2}) = F_\delta(a + \frac{3\delta}{2})$. Thus,

$$(3.19) \quad S_\delta(a + \frac{3\delta}{2}) = (\vartheta(\delta))^{-1} \left(F_\delta(a + \frac{3\delta}{2}) - 1 - 2\Re \mathcal{A} \cdot \delta \right)$$

and our goal is to estimate this value. Note that S_δ vanishes on the cut L_a : both $F_{[\mathbb{C}_\delta, a]}$ and $G_{[\mathbb{C}_\delta, a]}$ vanish by construction, and $F_\delta^{(\mathcal{R})} = -F_\delta$ on L_a due to the spinor property. (F_δ changes the sign between opposite sides of L_a since they belong to different sheets.) It is clear that S_δ is discrete harmonic everywhere in Λ_δ^+ except at the point $a + \frac{3\delta}{2}$ since all terms are discrete harmonic there. Moreover, due to Lemma 3.5, it is discrete harmonic at $a + \frac{3\delta}{2}$ also. Therefore, for any (small, but fixed) $r > 0$, the discrete Beurling estimate (3.4) implies

$$(3.20) \quad |S_\delta(a + \frac{3\delta}{2})| \leq C\delta^{\frac{1}{2}} \cdot r^{-\frac{1}{2}} \max_{\varpi(r)} |S_\delta|,$$

where $\varpi(r) := \{z : r \leq |z - a| \leq r + 5\delta\}$ denotes the “discrete circle” of radius r around a . Further, it follows from Theorem 2.16 and Lemmas 2.14 and 2.17 that

$$S_\delta \xrightarrow{\delta \rightarrow 0} s := \Re \left[\frac{1}{2}(f + f^{(\mathcal{R})}) - f_{[\mathbb{C}, a]} \right] - 2\Re \mathcal{A} \cdot g_{[\mathbb{C}, a]},$$

uniformly on $\varpi(r)$, where $f = f_{[\Omega, a_1, \dots, a_n]}$, $f^{(\mathcal{R})} = f_{[\mathcal{R}(\Omega), a, \mathcal{R}(a_2), \dots, \mathcal{R}(a_n)]}$, and $f_{[\mathbb{C}, a]}(z) = \Re [1/\sqrt{z - a}]$ and $g_{[\mathbb{C}, a]}(z) = \Re \sqrt{z - a}$. Recall that, by definition of the coefficient \mathcal{A} , one has

$$f - f_{[\mathbb{C}, a]} = 2\mathcal{A}\sqrt{z - a} + O(|z - a|^{3/2}), \quad z \rightarrow a.$$

It is easy to check that $f^{(\mathcal{R})}(z) \equiv \overline{f(\mathcal{R}(z))}$ (since this spinor solves the corresponding boundary value problem), and hence

$$f^{(\mathcal{R})} - f_{[\mathbb{C}, a]} = 2\overline{\mathcal{A}}\sqrt{z - a} + O(|z - a|^{3/2}), \quad z \rightarrow a.$$

Thus, we arrive at $s(z) = O(|z - a|^{3/2})$ as $z \rightarrow a$, which means

$$(3.21) \quad r^{-\frac{1}{2}} \max_{\varpi(r)} |S_\delta| \xrightarrow{\delta \rightarrow 0} r^{-\frac{1}{2}} \cdot O(r^{3/2}) = O(r).$$

Combining (3.19)–(3.21), one concludes that, for any given $r > 0$,

$$|F_\delta(a + \frac{3\delta}{2}) - 1 - 2\Re \mathcal{A} \cdot \delta| \leq C\vartheta(\delta)\delta^{\frac{1}{2}}r$$

if δ is small enough. Since $\vartheta(\delta) = O(\delta^{\frac{1}{2}})$ by (2.18), and r can be chosen arbitrary small, this yields (2.20). All estimates are uniform with respect to a_1, \dots, a_n at definite distance from the boundary and from each other. \square

To prove Theorem 2.20, we need to analyze the discrete spinor $F_{[\Omega_\delta, a, b]}$ near the point b . In contrast to Theorem 2.18, where the second-order information near a was extracted, here we only need to match the first-order coefficients. Note that the situation is slightly different from the first-order analysis near a , as $F_{[\Omega_\delta, a, b]}$ does not have an explicit discrete singularity, remaining s-holomorphic near the branching point b , while its limit blows up at b . Still, the strategy for the identification resembles the one used above and appeals to the symmetrization with respect to the horizontal line passing through b . Let $L'_b := \{x + b + \frac{\delta}{2} : x < 0\}$.

Proof of Theorem 2.20. Let \mathcal{R} denote the reflection with respect to the line $\{x : \Im m(x - b) = 0\}$, and let $\Lambda_\delta \subset \Omega_\delta \cap \mathcal{R}(\Omega_\delta)$ be a small neighborhood of b . Fix the sheet $\Omega_\delta^+ \subset \mathcal{V}_{\Lambda_\delta}^i$ of $[\Omega_\delta, a, b] \setminus L'_b$ so that $F_\delta(b + \frac{\delta}{2}) = i\mathcal{B}_\delta$ with $\mathcal{B}_\delta > 0$ (recall that the values of discrete spinors on $\mathcal{V}_{\mathbb{C}_\delta}^i$ are purely imaginary), and do the same for the continuous spinor: fix a sheet so that $\Im m f_{[\Omega, a, b]}(b + x) > 0$ as $x \downarrow 0$ (see (2.14)).

Let $W_\delta := \text{hm}_{\{b + \frac{\delta}{2}\}}(\cdot)$ denote the harmonic measure of the point $b + \frac{\delta}{2}$ in the slit discrete plane $\mathbb{C}_\delta^i \setminus L'_b$ and

$$T_\delta := (\vartheta(\delta))^{-1}(\frac{1}{2}(F_\delta + F_\delta^{(\mathcal{R})}) - i\mathcal{B}_\delta \cdot W_\delta) : \Lambda_\delta^+ \rightarrow i\mathbb{R},$$

where $F_\delta = F_{[\Omega_\delta, a, b]}$ and $F_\delta^{(\mathcal{R})} = F_{[\mathcal{R}(\Omega), \mathcal{R}(a), b]}$. By symmetry, $F_\delta^{(\mathcal{R})}(b + \frac{\delta}{2}) = i\mathcal{B}_\delta$, where the sheet for $F_\delta^{(\mathcal{R})}$ is fixed by the same condition $\Im m F_\delta^{(\mathcal{R})}(b + \frac{\delta}{2}) > 0$. We have to prove that $\mathcal{B}_\delta \rightarrow \mathcal{B} = \mathcal{B}_{[\Omega, a, b]}$ as $\delta \rightarrow 0$. Note that, by passing to a subsequence, one may assume that $\mathcal{B}_\delta \rightarrow \tilde{\mathcal{B}}$ for some $\tilde{\mathcal{B}} \in [0, +\infty]$.

Suppose that $\tilde{\mathcal{B}}$ is finite. Then, for any fixed (small) $r > 0$, Theorem 2.16 and Lemma 3.4(iii) imply the uniform convergence

$$(3.22) \quad T_\delta(z) \xrightarrow{\delta \rightarrow 0} t(z) := i \cdot \Im m [\frac{1}{2}(f(z) + f^{(\mathcal{R})}(z))] - i\tilde{\mathcal{B}} \cdot \Re e [1/\sqrt{z - b}]$$

on compact subsets of $\{z : 0 < |z - b| \leq r\}$, where $f = f_{[\Omega, a, b]}$ and $f^{(\mathcal{R})}(z) = f_{[\mathcal{R}(\Omega), \mathcal{R}(a), b]}(z) = -\overline{f(\mathcal{R}(z))}$ due to our conventions about sheets. Note that

$$(3.23) \quad t(z) = i(\mathcal{B} - \tilde{\mathcal{B}}) \cdot \Re e [1/\sqrt{z - b}] + O(|z - b|^{1/2}), \quad z \rightarrow b.$$

Clearly, the function T_δ is harmonic everywhere in Λ_δ^+ except at the point $b + \frac{\delta}{2}$, and $T_\delta = 0$ on L'_b : indeed, $F_\delta^{(\mathcal{R})} = -F_\delta$ on L'_b due to the spinor property of F_δ and symmetry reasons. Moreover, $T_\delta(b + \frac{\delta}{2}) = 0$ as $W_\delta(b + \frac{\delta}{2}) = 1$. Therefore, $\tilde{\mathcal{B}} \neq \mathcal{B}$ would contradict the maximum principle for the discrete harmonic function T_δ : in this case (3.22) and (3.23) imply that, for sufficiently small δ , the values of T_δ near b are bigger than those near the circle $\{z : |z - b| = r\}$.

The similar argument works if $\tilde{\mathcal{B}} = +\infty$: in this case one would have

$$\mathcal{B}_\delta^{-1}T_\delta(z) \xrightarrow{\delta \rightarrow 0} -i \cdot \Re e [1/\sqrt{z - b}]$$

which contradicts the maximum principle for T_δ , just as before. □

A. Appendix

The goal of this appendix is to derive the explicit formulae (1.4) for the scaling limits of n -point spin correlations in the upper half-plane \mathbb{H} . Recall that we have reduced this problem to the computation of the coefficients $2\mathcal{A}_{[\mathbb{H}, a_1, \dots, a_n]}$ in front of $\sqrt{z - a_1}$ in the asymptotic expansion (2.13) of the spinor $f_{[\mathbb{H}, a_1, \dots, a_n]}$. The latter is defined to be the unique solution to the boundary value problem (2.7)–(2.9) with additional regularity condition $f = O(|z|^{-1})$

at infinity. What remains to be proven is the equality (2.22), which claims that those coefficients $\mathcal{A}_{[\mathbb{H}, a_1, \dots, a_n]}$ are logarithmic derivatives of the explicit quantities $\langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_{\Omega}^+$ given by (1.4).

Proof of (2.22) for $n \geq 3$. Recall from the proof of Lemma 2.9 that the function $f_{[\mathbb{H}, a_1, \dots, a_n]}$ can be written in the form

$$(A.1) \quad f_Q(z) = e^{\frac{\pi i}{4}} \cdot \frac{Q(z)}{\sqrt{(z - a_1)(z - \bar{a}_1) \cdots (z - a_n)(z - \bar{a}_n)}},$$

where Q is a polynomial of degree $n - 1$ with real coefficients. This polynomial is to be determined from the conditions

$$\begin{aligned} \Re \left(\lim_{z \rightarrow a_1} \sqrt{z - a_1} \cdot f_{[\Omega, a_1, \dots, a_n]}(z) \right) &= 1, \\ \Re \left(\lim_{z \rightarrow a_k} \sqrt{z - a_k} \cdot f_{[\Omega, a_1, \dots, a_n]}(z) \right) &= 0, \quad k = 2, \dots, n, \end{aligned}$$

which are equivalent to

$$(A.2) \quad \frac{Q(a_k)}{\prod_{j \neq k} \sqrt{(a_k - a_j)(a_k - \bar{a}_j)}} + \frac{Q(\bar{a}_k)}{\prod_{j \neq k} \sqrt{(\bar{a}_k - a_j)(\bar{a}_k - \bar{a}_j)}} = 2\Im a_k \cdot \delta_{1k},$$

where $k = 1, \dots, n$ and δ_{1k} is the Kronecker delta. We can write $Q(z)$ in the form

$$(A.3) \quad Q(z) = \sum_{j=1}^n \frac{q_j}{z - a_j} \cdot \prod_{j=1}^n (z - a_j),$$

with some unknown coefficients q_1, \dots, q_n . Then, (A.2) is equivalent to

$$q_k \cdot \frac{\prod_{j \neq k} \sqrt{a_k - a_j}}{\prod_{j \neq k} \sqrt{a_k - \bar{a}_j}} + \sum_{s=1}^n q_s \cdot \frac{(\bar{a}_k - a_k) \prod_{j \neq k} \sqrt{\bar{a}_k - a_j}}{(\bar{a}_k - a_s) \prod_{j \neq k} \sqrt{\bar{a}_k - \bar{a}_j}} = 2\Im a_k \cdot \delta_{1k},$$

which can be further rewritten as

$$(A.4) \quad \sum_{s=1}^n (D_{ks} + C_{ks}) q_s = c(a_1, \dots, a_k) \cdot \delta_{1k}, \quad k = 1, \dots, n,$$

for a constant $c(a_1, \dots, a_k) \neq 0$ that only affects the overall normalization of the spinor $f_{[\mathbb{H}, a_1, \dots, a_n]}$, where D is the diagonal matrix with entries

$$D_{kk} = \frac{1}{\bar{a}_k - a_k} \prod_{j \neq k} \chi_{kj}^{\frac{1}{2}}, \quad \chi_{kj} = \chi_{jk} = \frac{(a_k - a_j)(\bar{a}_k - \bar{a}_j)}{(a_k - \bar{a}_j)(\bar{a}_k - a_j)} = \left| \frac{a_k - a_j}{a_k - \bar{a}_j} \right|^2,$$

and C is the Cauchy matrix

$$C_{ks} = \frac{1}{\bar{a}_k - a_s}, \quad 1 \leq k, s \leq n.$$

Solving (A.4) by Cramer's rule, one gets $q_j/q_1 = (-1)^{j+1} \det A_{[1j]} / \det A_{[11]}$, where $\det A_{[1j]}$ stands for the minors of the matrix $A = D + C$. Plugging (A.1)

and (A.3) into the expansion (2.13) of $f_{[\mathbb{H}, a_1, \dots, a_n]}(z) = f_Q(z)$ as $z \rightarrow a_1$, we arrive at the formula

$$\begin{aligned} 2\mathcal{A}_{[\mathbb{H}, a_1, \dots, a_n]} &= \partial_z(\log[f_Q(z)\sqrt{z - a_1}])\Big|_{z=a_1} \\ &= \sum_{j=2}^n \frac{(-1)^{j+1} \det A_{[1j]}}{(a_1 - a_j) \det A_{[11]}} - \frac{1}{2(a_1 - \bar{a}_1)} \\ &\quad + \frac{1}{2} \sum_{j=2}^n \left(\frac{1}{a_1 - a_j} - \frac{1}{a_1 - \bar{a}_j} \right). \end{aligned}$$

Recall that principal minors of the Cauchy matrix C are given by

$$\det(C_S) = \prod_{k \in S} \frac{1}{\bar{a}_k - a_k} \prod_{\substack{k < m \\ k, m \in S}} \chi_{km} = \prod_{k \in S} \frac{1}{\bar{a}_k - a_k} \prod_{\substack{k \neq m \\ k, s \in S}} \chi_{km}^{\frac{1}{2}}.$$

Hence, expanding $\det A_{[11]}$ by linearity and encoding a subset $S \subset \{2, \dots, n\}$ by the collection of signs $\mu \in \{\pm 1\}^n$, where $\mu_1 = -1$ and $\mu_s = 2\mathbb{1}_{s \in S} - 1$ for $s = 2, \dots, n$, we find

$$\begin{aligned} \det A_{[11]} &= \sum_{S \subset \{2, \dots, n\}} (\det C_S \cdot \det D_{\bar{S}}) \\ &= \prod_{k=2}^n \frac{1}{\bar{a}_k - a_k} \cdot \sum_{S \subset \{2, \dots, n\}} \left(\prod_{\substack{k, m \in S \\ k \neq m}} \chi_{km}^{\frac{1}{2}} \cdot \prod_{\substack{k \in \bar{S} \\ j=1 \\ j \neq k}} \prod_{j=1}^n (\chi_{kj} \chi_{jk})^{\frac{1}{4}} \right) \\ &= \prod_{k=2}^n \frac{1}{\bar{a}_k - a_k} \cdot \sum_{\substack{\mu \in \{\pm 1\}^n \\ \mu_1 = -1}} \left(\prod_{\substack{k, m=2 \\ k \neq m}}^n \chi_{km}^{\frac{3+\mu_k \mu_m}{8}} \prod_{k=2}^n (\chi_{k1} \chi_{1k})^{\frac{1-\mu_k}{8}} \right) \\ &= \prod_{k=2}^n \frac{1}{\bar{a}_k - a_k} \prod_{\substack{k, m=2 \\ k \neq m}}^n \chi_{km}^{\frac{3}{8}} \prod_{k=2}^n \chi_{1k}^{\frac{1}{4}} \cdot \sum_{\substack{\mu \in \{\pm 1\}^n \\ \mu_1 = -1}} \prod_{\substack{k, m=1 \\ k \neq m}}^n \chi_{km}^{\frac{\mu_k \mu_m}{8}}. \end{aligned}$$

Note that we can also write

$$\sum_{j=2}^n \frac{(-1)^{j+1} \det A_{[1j]}}{a_1 - a_j} = \det[\tilde{D} + \tilde{C}],$$

where \tilde{D} and \tilde{C} are n -by- n matrices, \tilde{D} is diagonal with entries $\tilde{D}_{11} = 0$ and $\tilde{D}_{kk} = D_{kk}$ for $2 \leq k \leq n$, and the entries of \tilde{C} are given by $\tilde{C}_{11} = 0$, $\tilde{C}_{1s} = 1/(a_1 - a_s)$ for $2 \leq s \leq n$ and $\tilde{C}_{ks} = C_{ks} = 1/(\bar{a}_k - a_s)$ if $2 \leq k \leq n$.

We now compute the principal minors of \tilde{C} . If $1 \notin S$, then $\det \tilde{C}_S = \det C_S$. Otherwise, denote $S := \{1\} \cup S_1$, where $S_1 \subset \{2, \dots, n\}$. Treating \bar{a}_1

and a_1 as independent variables (note that \tilde{C} does not contain \bar{a}_1), we obtain

$$\begin{aligned} \det \tilde{C}_S &= \lim_{\bar{a}_1 \rightarrow a_1} \left(\det C_S - \frac{1}{\bar{a}_1 - a_1} \det C_{S_1} \right) \\ &= \prod_{k \in S_1} \frac{1}{\bar{a}_k - a_k} \prod_{\substack{k < m \\ k, m \in S_1}} \chi_{km} \cdot \lim_{\bar{a}_1 \rightarrow a_1} \frac{\prod_{s \in S_1} \chi_{1s} - 1}{\bar{a}_1 - a_1} \\ &= \prod_{k \in S_1} \frac{1}{\bar{a}_k - a_k} \prod_{\substack{k \neq m \\ k, m \in S_1}} \chi_{km}^{\frac{1}{2}} \cdot \sum_{s \in S_1} \left(\frac{1}{a_1 - \bar{a}_s} - \frac{1}{a_1 - a_s} \right). \end{aligned}$$

Note that $1/(a_1 - \bar{a}_s) - 1/(a_1 - a_s) = -\partial_{a_1} \log \chi_{1s}$. Similarly to the computation of $\det A_{[11]}$ given above, we have

$$\begin{aligned} \det[\tilde{D} + \tilde{C}] &= \sum_{S \subset \{1, \dots, n\}} (\det \tilde{C}_S \cdot \det \tilde{D}_{\bar{S}}) = \sum_{\substack{S = \{1\} \cup S_1, \\ S_1 \subset \{2, \dots, n\}}} (\det \tilde{C}_S \cdot \det \tilde{D}_{\bar{S}}) \\ &= \prod_{k=2}^n \frac{1}{\bar{a}_k - a_k} \prod_{\substack{k, m=2 \\ k \neq m}}^n \chi_{km}^{\frac{3}{8}} \prod_{k=2}^n \chi_{1k}^{\frac{1}{4}} \\ &\quad \cdot \sum_{\substack{\mu \in \{\pm 1\}^n \\ \mu_1 = -1}} \left(\prod_{\substack{k, m=1 \\ k \neq m}}^n \chi_{km}^{\frac{\mu_k \mu_m}{8}} \cdot \sum_{s=2}^n \frac{1 + \mu_s}{2} (-\partial_{a_1} \log \chi_{1s}) \right), \end{aligned}$$

where we set $\mu_s = 2\mathbb{1}_{s \in S_1} - 1$ for $s = 2, \dots, n$. Taking into account the symmetry of $\mu_k \mu_m$ and $\mu_1 \mu_s$ under the simultaneous sign flip of *all* μ_j , $j = 1, \dots, n$, we obtain

$$\frac{\det[\tilde{D} + \tilde{C}]}{\det A_{[11]}} = \frac{\sum_{\mu \in \{\pm 1\}^n} \left(\prod_{1 \leq k < m \leq n} \chi_{km}^{\frac{\mu_k \mu_m}{4}} \cdot \sum_{s=2}^n \frac{1 - \mu_1 \mu_s}{2} (-\partial_{a_1} \log \chi_{1s}) \right)}{\sum_{\mu \in \{\pm 1\}^n} \prod_{1 \leq k < m \leq n} \chi_{km}^{\frac{\mu_k \mu_m}{4}}}.$$

Therefore,

$$\frac{\det[\tilde{D} + \tilde{C}]}{\det A_{[11]}} = -\frac{1}{2} \sum_{m=2}^n \partial_{a_1} \log \chi_{1m} + \frac{2\partial_{a_1} \left(\sum_{\mu \in \{\pm 1\}^n} \prod_{1 \leq k < m \leq n} \chi_{km}^{\frac{\mu_k \mu_m}{4}} \right)}{\sum_{\mu \in \{\pm 1\}^n} \prod_{1 \leq k < m \leq n} \chi_{km}^{\frac{\mu_k \mu_m}{4}}},$$

and we arrive at

$$\mathcal{A}_{[\mathbb{H}, a_1, \dots, a_n]} = -\frac{1}{4(a_1 - \bar{a}_1)} + \partial_{a_1} \log \left(\sum_{\mu \in \{\pm 1\}^n} \prod_{1 \leq k < m \leq n} \chi_{km}^{\frac{\mu_k \mu_m}{4}} \right),$$

thus concluding the proof of the identity $\mathcal{A}_{[\mathbb{H}, a_1, \dots, a_n]} = 2\partial_{a_1} \log \langle \sigma_{a_1} \cdots \sigma_{a_n} \rangle_{\mathbb{H}}^{\dagger}$. \square

Remark A.1. When preparing the first version of our paper, we were only able to check the explicit formula (1.4) in some special cases; thus, the exposition went along the lines of Remark 2.23. It later turned out that “interpolation problems” similar to (A.2) also appear in the analysis of the critical Ising model with mixed free/fixed boundary conditions, in which case the collection $a_1, \bar{a}_1, \dots, a_n, \bar{a}_n$ should be replaced by $2n$ boundary points where these boundary conditions change. In particular, a computation similar to the one given above, combined with the Edwards-Sokal coupling, yields general $2n$ -point crossing formulae in the random cluster (Fortuin-Kasteleyn) representation of the critical Ising model; see [Izy14].

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ST. PETERSBURG DEPARTMENT OF STEKLOV MATHEMATICAL INSTITUTE,
ST. PETERSBURG, RUSSIA

CHEBYSHEV LABORATORY, DEPARTMENT OF MATHEMATICS AND MECHANICS,
ST. PETERSBURG STATE UNIVERSITY, ST. PETERSBURG, RUSSIA

E-mail: dchelkak@pdmi.ras.ru

Current address: ETH ZÜRICH, INSTITUTE FOR THEORETICAL STUDIES,
ZÜRICH, SWITZERLAND

COLUMBIA UNIVERSITY, NEW YORK, NY

E-mail: hongler@math.columbia.edu

Current address: ÉCOLE POLYTECHNIQUE FÉDÉRALE DE LAUSANNE,
LAUSANNE, SWITZERLAND

CHEBYSHEV LABORATORY, DEPARTMENT OF MATHEMATICS AND MECHANICS,
ST. PETERSBURG STATE UNIVERSITY, ST. PETERSBURG, RUSSIA

Current address: UNIVERSITY OF HELSINKI, HELSINKI, FINLAND

E-mail: konstantin.izyurov@helsinki.fi