

On geometric transfer in real twisted endoscopy

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Abstract

We prove the existence of a transfer of orbital integrals in endoscopy for real reductive groups when there is twisting by an automorphism defined over the reals and by a character on the real points of the group. Our proof contains a relatively short self-contained argument for the already known case of standard endoscopy.

1. Introduction

Endoscopy concerns conjugacy classes and irreducible representations for reductive groups: conjugacy classes within a stable class and irreducible representations within a packet. We consider just real groups. In the case of standard endoscopy, where there is no twisting, geometric and spectral *transfer identities* have been proved. These identities display structure on both packets and stable classes. In particular, the structure on a packet of discrete series representations reflects that on the set of conjugacy classes in a regular elliptic stable class. As is well known, this structure plays a role in various comparisons of trace formulas and in multiplicity formulas for automorphic representations. In the present paper we consider the broader setting of twisted endoscopy, again for real groups. Our purpose is to present a complete argument for the main geometric transfer identity. This identity shows that sums of integrals over suitably regular twisted conjugacy classes, when weighted by the transfer factors introduced in [KS99] (see also [KS12]), may be interpreted as integrals over stable conjugacy classes in an endoscopic group. The precise result has two immediate applications. First, locally (i.e., for real groups), it establishes the underlying structure for a functorial dual transfer of stable traces on a twisted endoscopic group to virtual twisted traces on the ambient group. In a separate paper [She11] we begin the description of an explicit form for the dual transfer via compatible spectral transfer factors. This extends the standard, or untwisted, case [She10], [She08b] and appears useful in the global theory; see, for example, [Art11, Th. 2.2.4]. Second, in the global picture,

our geometric transfer identity is of course one ingredient for stabilization of the geometric side of the general twisted version of the Arthur-Selberg trace formula.

Suppose G is a connected reductive algebraic group defined over \mathbb{R} . There are two familiar types of twisting we will consider for an admissible representation π of the reductive Lie group $G(\mathbb{R})$: composing π with an \mathbb{R} -automorphism θ of G and multiplying π by a character ϖ of $G(\mathbb{R})$. An isomorphism A_π between $\pi \circ \theta$ and $\varpi \otimes \pi$, if it exists, provides us with a distribution $f \rightarrow \text{Trace}(\pi(f)A_\pi)$, a (θ, ϖ) -twisted character for π , on a suitable space of test functions f . Comparing these twisted traces with ordinary stable traces for a lower dimensional group, an endoscopic group $H_1(\mathbb{R})$ for (G, θ, ϖ) , requires a correspondence on test functions. That is provided by the main geometric transfer identity which displays weighted sums of (θ, ϖ) -twisted orbital integrals of test functions on $G(\mathbb{R})$ as stable orbital integrals of the corresponding test functions on $H_1(\mathbb{R})$. For the remainder of Section 1 we will discuss in some detail our setting for this and related results. The results themselves will then be described in more detail in Section 2.

Our setting is based on the constructions and results of [KS99] for the case of real groups. For the norm correspondence of [KS99, Chap. 3, §5.4] between points of $G(\mathbb{R})$ and points of an endoscopic group $H_1(\mathbb{R})$ for (θ, ϖ) , it is an associated outer automorphism θ^* of a quasi-split inner form G^* that is significant. If θ is inner, then θ^* is the identity, and we have a slight variant of the setting for standard endoscopy [LS87, §1.3]. To simplify the presentation we will carry a minor assumption on the norm correspondence for most of the paper. Fix an inner twist $\psi : G \rightarrow G^*$, where G^* is quasi-split over \mathbb{R} . There is an \mathbb{R} -automorphism θ^* of G^* which preserves a given \mathbb{R} -splitting of G^* and for which θ^* and $\psi \circ \theta \circ \psi^{-1}$ differ by an inner automorphism of G^* . We then say (G, θ, ψ) is an inner twist of (G^*, θ^*) , as in [KS99, App. B]. Start now with the pair (G^*, θ^*) . We will consider those (isomorphism classes of) inner twists (G, θ, ψ) for which there is a norm correspondence from twisted conjugacy classes in $G(\mathbb{R})$ to the ordinary, i.e., untwisted, conjugacy classes in an endoscopic group $H_1(\mathbb{R})$. See Section 6 for a precise version of the assumption. If θ^* is the identity, this excludes certain inner automorphisms θ . In these cases the twist θ persists to conjugacy classes in the endoscopic group according to the formalism of [KS99, §5.4]. The general excluded case is a variant of this, and we use a *slightly twisted* norm correspondence. It can be handled by a straightforward extension of our arguments, as we will describe in Section 12.

An endoscopic group H_1 comes with more data. First we assume that we are given, rather than the twisting character ϖ itself, a 1-cocycle a_ϖ (of the Weil group $W_{\mathbb{C}/\mathbb{R}}$ in the center of the connected complex dual group G^\vee of G or

G^*) to which ϖ is attached by Langlands' construction [Bor79, 10.1]. A set ϵ of endoscopic data for (G, θ, a_ϖ) or $(G^*, \theta^*, a_\varpi)$ is a tuple $(H, \mathcal{H}, \mathfrak{s}, \xi)$ as in [KS99, §2.1]. There is no harm in assuming that ξ , an embedding of the group \mathcal{H} in the L -group ${}^L G = G^\vee \rtimes W_{\mathbb{C}/\mathbb{R}}$, is the inclusion map incl , so that \mathcal{H} is given as a subgroup of ${}^L G$. We do so, and drop ξ entirely from notation. This subgroup \mathcal{H} is, by definition, a split extension of $W_{\mathbb{C}/\mathbb{R}}$ by H^\vee . In some cases, there is an L -isomorphism $\xi_1 : \mathcal{H} \rightarrow {}^L H$. This provides us then with an L -embedding $\text{incl} \circ (\xi_1)^{-1}$ of ${}^L H$ in ${}^L G$, and H itself may serve as an endoscopic group. The L -embedding $\text{incl} \circ (\xi_1)^{-1}$ determines both a term for geometric transfer factors and a shift in infinitesimal character for the dual spectral transfer from $H(\mathbb{R})$ to $G(\mathbb{R})$. The shift is necessary for the existence of a transfer identity satisfying the functoriality principle; for some examples, see [She, Part B, §2]. Existence of ξ_1 as isomorphism, however, excludes many cases; quick examples can be found for an outer automorphism of $\text{SU}(2, 1)$ or for base change in Sp_4 . (In standard endoscopy, examples are harder to find.) To avoid these exclusions, we add to the endoscopic data $\epsilon = (H, \mathcal{H}, \mathfrak{s})$ a z -pair (H_1, ξ_1) as in [KS99, Chap. 2], and then H_1 , rather than H , serves as endoscopic group. This group H_1 is quasi-split over \mathbb{R} with simply-connected derived group, and there is an exact sequence $1 \rightarrow Z_1 \rightarrow H_1 \rightarrow H \rightarrow 1$ defined over \mathbb{R} , where Z_1 is an induced central torus in H_1 . Then $H_1(\mathbb{R}) \rightarrow H(\mathbb{R})$ is surjective and ${}^L H$ is naturally embedded in ${}^L H_1$. Now ξ_1 is an injective L -homomorphism of \mathcal{H} in ${}^L H_1$ (see [KS99, §2.2] for proof of existence), and ξ_1 determines, in particular, a character ϖ_1 on $Z_1(\mathbb{R})$. For example, in the $\text{SU}(2, 1)$ case we may pass from the problematic $H = \text{PGL}(2)$ to the group $H_1 = \text{GL}(2)$ with sign character ϖ_1 on $Z_1(\mathbb{R}) = \mathbb{R}^\times$. Spectral transfer from $H_1(\mathbb{R})$ to $G(\mathbb{R})$ involves just those representations π_1 of $H_1(\mathbb{R})$ for which $Z_1(\mathbb{R})$ acts by ϖ_1 . We will assume this property for a representation π_1 without further mention. The L -embedding $\text{incl} \circ (\xi_1)^{-1}$, now defined on a subgroup of ${}^L H_1$, plays essentially the same role in spectral transfer as before, but of course with H_1 in place of H .

We will prove transfer for test functions on $G(\mathbb{R})$ that are smooth and either compactly supported or rapidly decreasing on $G(\mathbb{R})$. (Passage to functions with prescribed behavior under the action of twisted conjugation by the center is then routine.) In the case of smooth functions of compact support this provides a direct analogue of Waldspurger's results in the nonarchimedean case [Wal08]. In particular, we use the same normalization of twisted orbital integrals [Wal08, §§1.5, 3.10]. We no longer need the technical assumption on the central behavior of θ from an earlier draft; see Lemma 8.1 and its preparation from Sections 6 and 7. Following the formalism of z -pairs [KS99, §2.2], we do prescribe behavior of test functions on $H_1(\mathbb{R})$ under translation by the central subgroup $Z_1(\mathbb{R}) = \text{Ker}(H_1(\mathbb{R}) \rightarrow H(\mathbb{R}))$. Namely, we require that a

test function f_1 on $H_1(\mathbb{R})$ satisfy

$$f_1(z_1 h_1) = \varpi_1(z_1)^{-1} f_1(h_1)$$

for all $z_1 \in Z_1(\mathbb{R})$, $h_1 \in H_1(\mathbb{R})$.

For our test functions we could go directly to C_c^∞ -spaces and then obtain, as a corollary of the geometric transfer, the dual transfer of stable admissible traces to twisted-invariant distributions. Instead we prefer to start with a more general space of functions of Harish-Chandra Schwartz type and then later pass to C_c^∞ -functions using a well-known result of Bouaziz [Bou94, Th. 6.2.1]. Thus from our main theorem we obtain first a dual spectral transfer of stable *tempered* traces to *tempered* twisted-invariant distributions. There has been recent progress by Mezo [Mez12] on identifying these distributions as weighted sums of tempered irreducible twisted traces. For standard endoscopy, this program has been completed [She10], with the weights identified as the predefined canonical spectral transfer factors of [She10]. Then, for standard endoscopy, we conclude from the existence of geometric transfer that a spectrally defined transfer identity for a pair (f, f_1) of test functions of any type also yields a geometric transfer identity for the pair if and only if it is correct on the tempered spectrum, i.e., it has the spectral transfer factors as weights. For progress with twisted spectral factors and their relation to Mezo's constants, see [She11].

To define a θ -Schwartz function f on $G(\mathbb{R})$ we consider, as usual, the manifold $G(\mathbb{R})\theta$ within $G(\mathbb{R}) \rtimes \text{Aut}_{\mathbb{R}}(G)$. On $G(\mathbb{R})\theta$ there is an action of $G(\mathbb{R})$ by conjugation: $x\theta \cdot g = g^{-1}(x\theta)g = g^{-1}x \theta(g) \theta$. To a smooth complex-valued function f on $G(\mathbb{R})$ we attach the smooth function f_θ on $G(\mathbb{R})\theta$ given by $f_\theta(x\theta) = f(x)$. We call f a θ -Schwartz function on $G(\mathbb{R})$ if f_θ is Schwartz on $G(\mathbb{R})\theta$. This requires a straightforward generalization of Harish-Chandra's definition; see the appendix for details and references. Write $\mathcal{C}(G(\mathbb{R}), \theta)$ for the space of all such functions. On $H_1(\mathbb{R})$ we consider the space $\mathcal{C}(H_1(\mathbb{R}), \varpi_1)$ of functions that are ϖ_1 -Schwartz in the usual sense. As mentioned already, for the fully general case there is a twist also on $H_1(\mathbb{R})$ by an inner automorphism θ_1 . In that setting, $H_1(\mathbb{R})\theta_1$ may be replaced by an appropriate coset of $H_1(\mathbb{R})$ in $H_1(\mathbb{C})$ (see Section 12), and we again require that test functions transform by ϖ_1^{-1} under the translation action of $Z_1(\mathbb{R})$.

To specify a correspondence (f, f_1) it will be sufficient to consider those twisted conjugacy classes of elements δ in $G(\mathbb{R})$ that are strongly θ -regular and have a (strongly G -regular) norm γ_1 in $H_1(\mathbb{R})$ in the sense of [KS99, §§3.3, 5.4]. Then the θ -twisted centralizer $\text{Cent}_\theta(\delta, G)$ of δ is reductive and abelian, but is not necessarily connected (as complex group). Because δ has a norm in $H_1(\mathbb{R})$, ϖ is trivial on $\text{Cent}_\theta(\delta, G)(\mathbb{R})$; see [KS99], where we use Theorem 5.1.D to strengthen Lemma 4.4.C. The ordinary centralizer $\text{Cent}(\gamma_1, H_1)$ is a torus which we write as H_{γ_1} . Recall we assume no twisting in $H_1(\mathbb{R})$ until Section 12.

There is a simple notion of compatibility for normalization of Haar measures on $\text{Cent}_\theta(\delta, G)(\mathbb{R})$ and $H_{\gamma_1}(\mathbb{R})$; see Section 11. We fix Haar measures dg on $G(\mathbb{R})$ and dh_1 on $H_1(\mathbb{R})$. This choice can be avoided if we work instead with Schwartz measures fdg and f_1dh_1 . In any case, it plays no significant role provided we insist on compatible measures dt_δ and dt_{γ_1} for $\text{Cent}_\theta(\delta, G)(\mathbb{R})$ and $H_{\gamma_1}(\mathbb{R})$ when γ_1 is a norm of strongly θ -regular δ . For $f \in \mathcal{C}(G(\mathbb{R}), \theta)$ and quotient measure $\frac{dg}{dt_\delta}$, we have the well-defined (θ, ϖ) -twisted orbital integral

$$O^{\theta, \varpi}(\delta, f) = \int_{\text{Cent}_\theta(\delta, G)(\mathbb{R}) \backslash G(\mathbb{R})} f(g^{-1}\delta\theta(g))\varpi(g) \frac{dg}{dt_\delta}$$

(see the appendix). Finally, the familiar stable orbital integral $\text{SO}(\gamma_1, f_1)$ is defined for $f_1 \in \mathcal{C}(H_1(\mathbb{R}), \varpi_1)$ and the quotient measure $\frac{dh_1}{dt_{\gamma_1}}$. If strongly θ -regular δ does not have a norm in $H_1(\mathbb{R})$, we may still define a (θ, ϖ) -twisted orbital integral $O^{\theta, \varpi}(\delta, f)$, but it plays no role in the transfer to $H_1(\mathbb{R})$. There will be other endoscopic groups that do account for it [KS99, Chap. 6].

The last ingredient for our transfer identity is the transfer factor $\Delta(\gamma_1, \delta)$ from [KS99] (see also [KS12]). While its definition is complicated in general, it has the property that the relative factor

$$\Delta(\gamma_1, \delta) / \Delta(\overline{\gamma_1}, \overline{\delta}) = \Delta(\gamma_1, \delta; \overline{\gamma_1}, \overline{\delta})$$

is canonical [KS99, Th. 4.6.A]. This means that the relative factor depends only on the data we have prescribed: the inner twist (G, θ, ψ) , cocycle a_ϖ defining the twisting character ϖ , endoscopic data ϵ with z -pair (H_1, ξ_1) for ϵ , and of course the pairs $(\gamma_1, \delta), (\overline{\gamma_1}, \overline{\delta})$. When ϖ is trivial, it is only the appropriate conjugacy classes of these pairs that matter: the stable (slightly twisted) conjugacy classes of $\gamma_1, \overline{\gamma_1}$ in $H_1(\mathbb{R})$ and the ordinary $(G(\mathbb{R})$ -) twisted conjugacy classes of $\delta, \overline{\delta}$ in $G(\mathbb{R})$. In general, there is a further twist by ϖ over twisted conjugacy classes in $G(\mathbb{R})$ in the sense of [KS99, Th. 5.1.D (2)].

The canonicity property motivates our approach to proving transfer and is critical to our arguments, reducing the difficulties in establishing the main transfer identity to simply stated problems at various walls in the endoscopic group. We are free to make convenient choices for the data determining the individual terms in transfer factors at each wall, and we thereby avoid the long consistency arguments for various local choices over on the ambient group in our original approach to the case of standard endoscopy for real groups [She82]. In particular, given the definitions of the transfer factors in [LS87] and the alternate characterization of stable orbital integrals we use here (see Section 4 and Theorem 12.1, where we may set g_0 to be the identity), the present paper offers a relatively short proof of the transfer for standard endoscopy. Indeed, we may go directly to Section 9 since the results of Sections 6–8 for ordinary conjugacy are known [She79a] and Section 5 is essentially just a statement of

the main jump formula from which the transfer follows quickly. The argument for this jump formula is a special case of the arguments in Sections 9–11. There we reduce easily to questions about the terms in transfer factors. Then, in loosely technical terms, our choice of a -data from Section 3, which is different from but in the same spirit as that in [Kal11], makes the previously intractable term Δ_I easy to handle (Lemma 9.5). The term Δ_{II} is trivial to handle and so the burden is on Δ_{III} . Our choice of χ -data from Section 3 allows us to deal with this term in our main lemma (Lemma 9.3) by a sequence of cohomological calculations based on results in [LS87], [LS90] and [KS99], and we are done. In particular, we avoid the convoluted arguments needed in Section 13 of [She08a] for the proof of standard transfer sketched there.

In some cases there are particular normalizations for the absolute factor $\Delta(\gamma_1, \delta)$ which simplify its form, but these do not play a direct role in the arguments of the present paper. In fact, since the choice of normalization does not matter for existence of the transfer identity, in Section 5 we simply fix a pair $(\overline{\gamma}_1, \overline{\delta})$ and specify $\Delta(\overline{\gamma}_1, \overline{\delta})$ in a way that allows us to avoid carrying various constants in our calculations.

Finally, we note that Waldspurger has pointed out two corrections ([Wal09], personal communication) needed for the definition of twisted transfer factors in [KS99]. These have been addressed in [KS12]. The first does not affect our archimedean setting; see Remark 1 of Section 9. The second involves the choice of a sign in the Galois hypercohomology pairing of Appendix A of [KS99] used to define the term Δ_{III} in transfer factors. In the archimedean case we may simply invert the pairing without further change, as explained in Remark 2 of Section 9.

2. Statement of the main theorem

We fix a set \mathfrak{e} of endoscopic data, along with a z -pair (H_1, ξ_1) for \mathfrak{e} , and study geometric transfer for $G(\mathbb{R})$ and $H_1(\mathbb{R})$ under the transfer factor Δ . Until Section 12 we assume that the norm correspondence involves no twisting of the conjugacy classes in $H_1(\mathbb{R})$.

Suppose f is a θ -Schwartz function on $G(\mathbb{R})$, i.e., $f \in \mathcal{C}(G(\mathbb{R}), \theta)$. We have attached to \mathfrak{e} and (H_1, ξ_1) the shift character ϖ_1 on the central subgroup $Z_1(\mathbb{R})$ of $H_1(\mathbb{R})$. Define the subset

$$\text{Trans}(f)$$

of $\mathcal{C}(H_1(\mathbb{R}), \varpi_1)$ to consist of those ϖ_1 -Schwartz functions f_1 on $H_1(\mathbb{R})$ whose strongly G -regular stable orbital integrals match, through the norm correspondence for $G(\mathbb{R})$ and $H_1(\mathbb{R})$ attached to θ , Δ -weighted combinations of (θ, ϖ) -twisted orbital integrals of f :

$$\text{SO}(\gamma_1, f_1) = \sum_{\delta, \theta\text{-conj}} \Delta(\gamma_1, \delta) O^{\theta, \varpi}(\delta, f)$$

for all strongly G -regular γ_1 in $H_1(\mathbb{R})$. The summation is over θ -conjugacy classes of strongly θ -regular elements in $G(\mathbb{R})$; for fixed γ_1 , the product $\Delta(\gamma_1, \delta) O^{\theta, \varpi}(\delta, f)$ depends only on the θ -conjugacy class of strongly θ -regular δ and is nonvanishing on finitely many such classes (see Section 5).

This *transfer identity* for the pair (f, f_1) says, in particular, that if strongly G -regular γ_1 is not a norm, then

$$\text{SO}(\gamma_1, f_1) = 0$$

since, by definition, we then have $\Delta(\gamma_1, \delta) = 0$ for all strongly θ -regular δ in $G(\mathbb{R})$. Moreover, the stable orbital integrals of f_1 have relatively simple behavior around semiregular semisimple elements. One requirement of the identity is thus that the weights Δ provide a great deal of cancellation in the singularities of the individual (θ, ϖ) -twisted orbital integrals of f .

Notice that $f_1 \in \text{Trans}(f)$ is determined uniquely modulo the annihilator in $\mathcal{C}(H_1(\mathbb{R}), \varpi_1)$ of the space of stable tempered characters on $H_1(\mathbb{R})$. The strongly G -regular elements are dense in the set of all regular semisimple elements in $H_1(\mathbb{R})$, and so functions f_1 and f_2 in $\text{Trans}(f)$ have the same stable orbital integrals on all regular semisimple elements. Then, by Harish-Chandra’s regularity theorem for characters (see [HC75, §11, Th. 1]) and a simple application of a stable Weyl integration formula, those integrals generate all stable tempered characters on $H_1(\mathbb{R})$. Hence, f_1 and f_2 agree on such characters, as asserted.

We may consider instead $f \in C_c^\infty(G(\mathbb{R}), \theta)$, by which we mean that f_θ lies in $C_c^\infty(G(\mathbb{R})\theta)$, and define the set $\text{Trans}_c(f)$ of functions $f_1 \in C_c^\infty(H_1(\mathbb{R}), \varpi_1)$ such that f and f_1 have Δ -matching orbital integrals in the same manner. Embedding $C_c^\infty(G(\mathbb{R}), \theta)$ in $\mathcal{C}(G(\mathbb{R}), \theta)$, we may adapt the argument above to see that $f_1 \in \text{Trans}_c(f)$ is determined uniquely modulo the annihilator in $C_c^\infty(H_1(\mathbb{R}), \varpi_1)$ of the space of all stable tempered characters on $H_1(\mathbb{R})$.

THEOREM 2.1 (Main Theorem). *For all $f \in \mathcal{C}(G(\mathbb{R}), \theta)$, the subset $\text{Trans}(f)$ of $\mathcal{C}(H_1(\mathbb{R}), \varpi_1)$ is nonempty.*

We conclude from this theorem that the correspondence (f, f_1) , where $f \in \mathcal{C}(G(\mathbb{R}), \theta)$ and $f_1 \in \text{Trans}(f)$, is well defined. This correspondence determines a map from $\mathcal{C}(G(\mathbb{R}), \theta)$ to the quotient of $\mathcal{C}(H_1(\mathbb{R}), \varpi_1)$ by the annihilator of stable tempered characters on $H_1(\mathbb{R})$. If we switch from Schwartz functions to Schwartz measures, then the map is determined uniquely up to normalization of transfer factors. In standard endoscopy, where the dual tempered spectral transfer is available (see [She10] and [She08b] for the form needed), we may normalize the tempered spectral factors $\Delta(\pi_1, \pi)$ first if we wish. For example, for certain inner forms there is a common Whittaker normalization that has desirable properties [She08b, §§11, 13]. Then for simultaneous geometric and

spectral transfer identities the geometric factors must be normalized so that $\Delta(\pi_1, \pi)/\Delta(\gamma_1, \delta)$ coincides with a predefined, and canonical, compatibility factor $\Delta(\pi_1, \pi; \gamma_1, \delta)$ [She10, §12]. In the Whittaker case, this brings us back to the geometric version of the Whittaker normalization in [KS99, §5.3] for $\Delta(\gamma_1, \delta)$ [She10, §12]. Similar results are expected for the twisted case; see [She11].

There is an analogue for C_c^∞ -functions.

COROLLARY 2.2. *For all $f \in C_c^\infty(G(\mathbb{R}), \theta)$, the subset $\text{Trans}_c(f)$ is non-empty.*

Proof. Let $f \in C_c^\infty(G(\mathbb{R}), \theta)$. Using the main theorem we first find f'_1 in the subset $\text{Trans}(f)$ of $\mathcal{C}(H_1(\mathbb{R}), \varpi_1)$. Then because the stable orbital integrals of f'_1 vanish off the conjugacy classes meeting a set in $H_1(\mathbb{R})$ that is bounded modulo $Z_1(\mathbb{R})$, Bouaziz's characterization of stable orbital integrals of C_c^∞ -functions shows that there exists $f_1 \in C_c^\infty(H_1(\mathbb{R}), \varpi_1)$ such that

$$\text{SO}(\gamma_1, f_1) = \text{SO}(\gamma_1, f'_1)$$

for all strongly G -regular γ_1 in $H_1(\mathbb{R})$. Here, a slight extension of [Bou94, Th. 6.2.1] is needed; see [Ren03, §5.3]. Then $f_1 \in \text{Trans}_c(f)$. \square

Let K, K_1 be maximal compact subgroups of $G(\mathbb{R}), H_1(\mathbb{R})$ respectively. If $f \in C_c^\infty(G(\mathbb{R}), \theta)$ is K -finite, then spectral methods are expected to show that there is K_1 -finite f_1 in $\text{Trans}_c(f)$, as for standard endoscopy. In the standard setting, if $\Delta(\pi_1, \pi)$ is the spectral transfer factor compatible with given geometric factor $\Delta(\gamma_1, \delta)$, then the Paley-Wiener argument of Clozel in an appendix to [CD84] shows that there is K_1 -finite f_1 satisfying tempered spectral transfer for f with weights $\Delta(\pi_1, \pi)$. Thus $f_1 \in \text{Trans}_c(f)$.

Sections 3–11 are dedicated to a proof of the main theorem which, after some preparation, hinges almost entirely on Theorem 5.1. In Sections 3 and 4, we introduce a variant of Harish-Chandra's $'F_f$ transform that fits better with transfer factors. In particular, we obtain the limit formulas of Theorem 4.2 for ordinary stable orbital integrals. These are simpler; for example, the troublesome fourth root of unity that appears in the jump formulas for stable $'F_f$ (see [She08a, §3]) is gone. In Sections 5–10, our main goal is to prove Theorem 5.1, which amounts to analogous limit formulas for the right side of the transfer identity, i.e., for sums of twisted orbital integrals weighted by the transfer factors. At this stage we ignore the limit formulas for derivatives that will be required later in the paper and focus instead on the needed analysis of terms in the transfer factors.

The main lemma (Lemma 9.3) in the proof of Theorem 5.1 is a simple wall-crossing property of the term Δ_{III} in the transfer factor $\Delta = \Delta_I \Delta_{II} \Delta_{III} \Delta_{IV}$ that we deduce from a detailed examination of constructions from [LS87],

[LS90], and [KS99]. Two features are crucial to the cancellations that yield this result: use of the *s-compatible data sets* introduced in Section 3 and precise control of data attached to the abstract norm map; see *toral descent data at γ_0* in Section 7. The term Δ_{II} then contributes trivially at the wall, apart from the piece needed for descent to a neighborhood of the identity in a twisted centralizer of Dynkin type A_1 , while analysis of Δ_I may be avoided if we use known results for standard endoscopy. Since we plan to deduce that case as well, we also give an independent analysis of Δ_I as an exercise with descent formulas from [LS90]. The term Δ_{IV} is, as usual, absorbed into the definition of normalized integrals.

Once we have finished the proof of Theorem 5.1, we extend the limit formulas to derivatives. Again, use of the alternative transform simplifies both statements and arguments. We then complete our proof of the main theorem in Section 11. In Section 12, the theorem is extended to the general case, i.e., to the case of *slightly twisted* norms.

Our notation will follow this pattern: O for unnormalized integrals, Φ for normalized integrals, and $\Psi_{a,\chi}$ for our variant of the stabilized $'F_f$ transform.

We should mention the work of Renard [Ren97], [Ren03], which offers insight into the difficulties of local analysis for twisted transfer. In [Ren03], however, the focus is different from ours; certain choices are made there that we expressly exclude here by the symmetry (*s-compatibility*) requirements of the next section. Those choices are reminiscent of our initial approach to standard endoscopy [She82], and unfortunately the reference [Sh6] in [Ren03] consists only of some personal notes which make no attempt to address the remaining problems for making the method work. In the example of base change, we note that the consistency problems in [She81a] were resolved only by the new approach of [She84]. With the dual spectral transfer in mind (see [She11, §11]), we also need the slightly more general setting of [KS99], and we start with Schwartz functions to capture the dual tempered transfer first. Some of our early results from Section 6 have analogues in [Ren03], but our paths soon separate since we bundle transfer factors with the twisted integrals from the start and then focus on the space of (abstract) norms and the endoscopic group. This leads us to local problems for transfer factors directly related to descent arguments from [LS87] and [LS90]. Those are the problems we propose to describe and solve here since, as we have already mentioned for the special case of standard endoscopy, the desired transfer then follows quite quickly.

3. Generalized Weyl denominators

A stabilized version of Harish-Chandra's $'F_f$ transform was introduced in [She79a] to characterize stable orbital integrals. We prepare in the present section to introduce a variant of this transform based on the generalized Weyl

denominators from [She08a, §9] (see also [She10, §7c]) that depend on the a -data and χ -data of [LS87, §2] rather than on a choice of positive roots.

Let G be a connected reductive algebraic group defined over \mathbb{R} , and let T be a maximal torus in G defined over \mathbb{R} . The familiar skew-symmetric Weyl denominator on the Lie algebra $\mathfrak{t}_{\mathbb{R}}$ of $T(\mathbb{R})$ does not in general lift to $T(\mathbb{R})$. Harish-Chandra introduced the closely related function Δ' on $T(\mathbb{R})$ defined by

$$\Delta'(\gamma) = \left| \det(\text{Ad}(\gamma) - I)_{\mathfrak{g}/\mathfrak{m}} \right|^{1/2} \prod_{\alpha > 0, \text{imag}} (\alpha(\gamma) - 1),$$

where \mathfrak{m} is Lie algebra of the centralizer M in G of the split component of T . The product is over those imaginary roots, i.e., roots of T in M , which are positive for some specified ordering. See Section 17 of [HC75]; this paper has the final version of $'F_f$. An earlier definition, which differs by a sign that depends on the ordering, is recognized by the presence of a term $\epsilon_{\mathbb{R}}$. Note also that we have modified the definition to accommodate the use of the right action of conjugation in prescribing orbital integrals. Following Harish-Chandra [HC75], we partition roots of T in G as real ($\sigma\alpha = \alpha$), imaginary ($\sigma\alpha = -\alpha$), or complex ($\sigma\alpha \neq \pm\alpha$). Here, and throughout, σ denotes the action of the nontrivial element of $\Gamma = \text{Gal}(\mathbb{C}/\mathbb{R})$ on T , on the rational characters $X^*(T)$, etc. Then

$$\Delta'(\gamma) = \prod_{\alpha > 0, \text{imag}} (\alpha(\gamma) - 1) \prod_{\alpha \text{ real, cplx}} \left| \alpha(\gamma)^{1/2} - \alpha(\gamma)^{-1/2} \right|^{1/2},$$

where $|\alpha(\gamma)^{1/2} - \alpha(\gamma)^{-1/2}|$ is convenient notation for $|(\alpha(\gamma) - 1)(\alpha(\gamma)^{-1} - 1)|^{1/2}$. If γ is regular as an element of M , we may further write

$$\Delta'(\gamma) = \prod_{\alpha > 0, \text{imag}} \frac{(\alpha(\gamma) - 1)}{|\alpha(\gamma) - 1|} \prod_{\alpha} \left| \alpha(\gamma)^{1/2} - \alpha(\gamma)^{-1/2} \right|^{1/2}.$$

Let \mathcal{O}_{α} denote the Galois orbit of the root α of T in G . If α is imaginary, then \mathcal{O}_{α} is symmetric: $\mathcal{O}_{\alpha} = -\mathcal{O}_{\alpha} = \{\pm\alpha\}$. Otherwise \mathcal{O}_{α} is asymmetric. Then \mathcal{O}_{α} and $-\mathcal{O}_{\alpha}$ are disjoint and \mathcal{O}_{α} consists of one or two roots according as α is real or complex. Recall that we define a -data $\{a_{\alpha}\}$ and χ -data $\{\chi_{\alpha}\}$ as follows [LS87, 2.2 and 2.5]. For each root α , a_{α} is a nonzero complex number and

$$a_{\sigma\alpha} = \bar{a}_{\alpha}, a_{-\alpha} = -a_{\alpha}.$$

In particular, if α is real, then a_{α} is a real number, while if α imaginary, then a_{α} is purely imaginary. Turning to χ -data, if α is imaginary or complex, then χ_{α} is a character on \mathbb{C}^{\times} . Further, if α is imaginary, then χ_{α} must be an extension to \mathbb{C}^{\times} of the sign character on \mathbb{R}^{\times} . Finally,

$$\chi_{\sigma\alpha} = \chi_{\alpha} \circ \sigma, \chi_{-\alpha} = \chi_{\alpha}^{-1}.$$

If α is real, then χ_{α} is a character on \mathbb{R}^{\times} and $\chi_{-\alpha} = \chi_{\alpha}^{-1}$.

If \mathcal{O}_{α} is asymmetric, then χ_{α} may be the trivial character, in which case the choice of a_{α} will not matter for the objects we construct (for the sake of

completeness, we will often pick $a_\alpha = \pm 1 = -a_{-\alpha}$, and we say that such data are *trivial*.

The associated (right) generalized Weyl denominator is

$$\begin{aligned} \Delta_{a,\chi,\text{right}}(\gamma) &= \prod_{\mathcal{O}} \chi_\alpha \left(\frac{(\alpha(\gamma) - 1)}{a_\alpha} \right) \prod_{\alpha} \left| \alpha(\gamma)^{1/2} - \alpha(\gamma)^{-1/2} \right|^{1/2} \\ &= \left| \det(\text{Ad}(\gamma) - I)_{\mathfrak{g}/\mathfrak{t}} \right|^{1/2} \prod_{\mathcal{O}} \chi_\alpha \left(\frac{(\alpha(\gamma) - 1)}{a_\alpha} \right), \end{aligned}$$

where the product is over all Galois orbits \mathcal{O} , symmetric or not. Notice that the choice of representative α for \mathcal{O} does not matter.

We may also define $\Delta_{a,\chi,\text{left}}(\gamma)$ by replacing each term $\chi_\alpha \left(\frac{(\alpha(\gamma) - 1)}{a_\alpha} \right)$ with the term

$$\chi_\alpha(-a_\alpha(1 - \alpha(\gamma)^{-1})).$$

A useful property for computing the dual transfer of characters is that the product

$$\Delta_{a,\chi,\text{left}}(\gamma) \Delta_{a,\chi,\text{right}}(\gamma)$$

coincides with the term $\left| \det(\text{Ad}(\gamma) - I)_{\mathfrak{g}/\mathfrak{t}} \right|$ appearing in the Weyl integration formula [She10, Lemma 7.3]. In the present paper we are interested only in $\Delta_{a,\chi,\text{right}}(\gamma)$ and will write it simply as $\Delta_{a,\chi}(\gamma)$.

To return to the Harish-Chandra factor $\Delta'(\gamma)$, we choose a positive system for the imaginary roots and then set

$$\chi_\alpha(z) = (z/\bar{z})^{\frac{1}{2}} = \frac{z}{|z|}$$

for α positive imaginary. We also set χ_α trivial for all real roots and all complex roots. Then for any choice $\{a_\alpha\}$ of a -data, we have

$$\Delta'(\gamma) = \Delta_{a,\chi}(\gamma) \prod_{\alpha > 0, \text{imag}} \frac{a_\alpha}{|a_\alpha|}.$$

Notice that the product on the right is a fourth root of unity.

Suppose (arbitrarily chosen) χ -data $\{\chi_\alpha\}$ are replaced by another such set $\{\chi'_\alpha = \eta_\alpha \chi_\alpha\}$. Then

$$\Delta_{a,\chi'}(\gamma) = \Delta_{a,\chi}(\gamma) \prod_{\mathcal{O}, \text{symm}} \eta_\alpha \left(\frac{(\alpha(\gamma) - 1)}{a_\alpha} \right) \prod_{\pm \mathcal{O}, \text{asymm}} \eta_\alpha(\alpha(\gamma)).$$

Suppose α is imaginary, and choose a square root $\alpha(\gamma)^{1/2}$ for $\alpha(\gamma)$. Then $\eta_\alpha(\alpha(\gamma)^{1/2})$ is independent of this choice, and the last formula may be rewritten as

$$\Delta_{a,\chi'}(\gamma) = \Delta_{a,\chi}(\gamma) \prod_{\mathcal{O}, \text{symm}} \eta_\alpha(\alpha(\gamma)^{1/2}) \prod_{\pm \mathcal{O}, \text{asymm}} \eta_\alpha(\alpha(\gamma)),$$

showing that the change is independent of the choice of a -data. Replacing $\{a_\alpha\}$ by another set $\{a'_\alpha = a_\alpha b_\alpha\}$ yields

$$\Delta_{a',\chi}(\gamma) = \Delta_{a,\chi}(\gamma) \prod_{\mathcal{O}, \text{symm}} \text{sign}(b_\alpha),$$

and then that change is independent of the choice of χ -data.

Let α be an imaginary root of T . By a Cayley transform with respect to α we mean the restriction to T of an inner automorphism of G , written $\gamma \rightarrow \gamma^s = s^{-1}\gamma s$ or $T \rightarrow T^s$, for which $s\sigma(s)^{-1}$ acts on T as the Weyl reflection ω_α with respect to α . Then T^s is defined over \mathbb{R} . This is a generalization of the usual Cayley transform (see [She79a], [She79b, §3], also a review in [She83, §2]) that works well for stable conjugacy. Such a transform exists if and only if the orbit of α under the imaginary Weyl group, i.e., the Weyl group of T in M , contains a noncompact root (see [She79a, Prop. 4.11]). In the terminology of [She83, §2] this says that α is not *totally compact*. For each root β of T , we denote by β^s its transport by s to a root of T^s .

Suppose that $\{a_\beta\}, \{\chi_\beta\}$ are a -data and χ -data for T , and fix an imaginary root α . Assume that α is not totally compact so that we may choose a Cayley transform s with respect to α . Then we call $\{a_\beta\}, \{\chi_\beta\}$ together with a -data and χ -data $\{a_{\beta^s}\}, \{\chi_{\beta^s}\}$ for T^s an *s-compatible data set* if

$$a_{\omega_\alpha(\beta)} = a_\beta, \chi_{\omega_\alpha(\beta)} = \chi_\beta$$

for all $\beta \neq \pm\alpha$, and

$$a_{\beta^s} = a_\beta, \chi_{\beta^s} = \chi_\beta$$

for all roots $\beta \neq \pm\alpha$ of T except those complex β for which β^s is real, while for such β , we require

$$a_\beta = a_{\beta^s}, \chi_\beta = \chi_{\beta^s} \circ Nm_{\mathbb{R}}^{\mathbb{C}}.$$

This definition places no additional restrictions on the data $a_\alpha, a_{\alpha^s}, \chi_\alpha$ or χ_{α^s} corresponding to the Cayley roots α, α^s . On the other hand, we are *not* free to make the usual assumption that the data are trivial on all asymmetric orbits for T^s : the data must be nontrivial on those asymmetric (complex) orbits for T^s which bifurcate into symmetric orbits on passage back to T ; see the last step in the proof of Lemma 3.1. In the case of bifurcation of an asymmetric (complex) orbit for T into asymmetric (real) orbits for T^s , mentioned in the definition, we may choose trivial data, but if we do not, then only real (Galois-invariant) a_β, χ_β are allowed. The requirements in this last case are made with the proofs of Lemmas 4.1 and 9.1 in mind.

LEMMA 3.1. *Suppose that s is a Cayley transform. Then s-compatible data sets exist.*

Proof. Write σ, σ^s for the Galois actions on T, T^s respectively. By construction,

$$\sigma^s(\beta^s) = (\omega_\alpha \sigma \beta)^s$$

for all roots β of T . Thus, as in the case of the standard Cayley transform, the roots $\pm\alpha^s$ are real. If β is real, then so is β^s . If β is complex, then either

$\omega_\alpha\beta \neq \pm\sigma\beta$ and β^s is complex or $\omega_\alpha\beta = \sigma\beta$ and β^s is real. Here the case $\omega_\alpha\beta = -\sigma\beta$ (equivalently, β^s imaginary) has been excluded since that implies β^s is orthogonal to α^s , so that β must be imaginary and orthogonal to α .

First we pick a -data and χ -data for T . Clearly we may adjust the data to satisfy the conditions that $a_{\omega_\alpha(\beta)} = a_\beta$ and $\chi_{\omega_\alpha(\beta)} = \chi_\beta$ for all imaginary $\beta \neq \pm\alpha$. Suppose that β is real. Then we may take χ_β trivial and arrange that $a_\beta = \pm 1 = -a_{-\beta}$. Suppose that β is complex. Then we again take χ_β to be trivial and arrange that $a_\beta = \pm 1 = -a_{-\beta}$. We may also require that $a_{\omega_\alpha(\beta)} = a_\beta = a_{\sigma\beta}$. For this we observe that the orbit of β under the group generated by σ and ω_α is asymmetric and, moreover, disjoint from its negative; if β^s is real, then $\omega_\alpha\beta = \sigma\beta$ and the orbit is $\{\beta, \sigma\beta\}$, whereas if β^s is complex, then $\omega_\alpha\beta \neq \pm\sigma\beta$ and the orbit is $\{\beta, \sigma\beta, \omega_\alpha\beta, \omega_\alpha\sigma\beta\}$. The disjointness property is then clear.

To complete the proof of the lemma we show that we may define a -data and χ -data for T^s as follows. First use the formulas

$$a_{\beta^s} = a_\beta, \chi_{\beta^s} = \chi_\beta$$

for all roots β of T except $\pm\alpha$ and those complex β for which β^s is real. Suppose β is complex and β^s is real. We pick $a_{\beta^s} = a_\beta$ and take χ_{β^s} trivial on \mathbb{R}^\times . We choose $\chi_{\pm\alpha^s}$ trivial on \mathbb{R}^\times and $a_{\alpha^s} = 1 = -a_{-\alpha^s}$.

There is nothing left to show for $a_{\beta^s}, \chi_{\beta^s}$ unless β is imaginary and $\beta \neq \pm\alpha$. Then β^s is imaginary or complex according as β is orthogonal to α or not. If β is orthogonal to α , then $\sigma^s(\beta^s) = (\sigma\beta)^s$ and so it is clear that our chosen $a_{\pm\beta^s} = a_\beta, \chi_{\pm\beta^s} = \chi_\beta$ are appropriate. If β is not orthogonal to α , then

$$\sigma^s(\beta^s) = (-\omega_\alpha\beta)^s.$$

Using the additional requirement

$$a_{\omega_\alpha(\beta)} = a_\beta, \chi_{\omega_\alpha(\beta)} = \chi_\beta,$$

we see that

$$a_{\sigma^s(\beta^s)} = a_{-\omega_\alpha\beta} = \overline{a_{\omega_\alpha\beta}} = \overline{a_\beta} = \overline{a_{\beta^s}}$$

and

$$\chi_{\sigma^s(\beta^s)} = \chi_{-\omega_\alpha\beta} = \chi_{\omega_\alpha\beta} \circ \sigma = \chi_\beta \circ \sigma = \chi_{\beta^s} \circ \sigma.$$

Since clearly $a_{-\beta^s} = -a_{\beta^s}$ and $\chi_{-\beta^s} = \chi_{\beta^s}^{-1}$, this finishes the proof. □

4. A limit formula for stable orbital integrals

We continue with the setting of the last section. Suppose that SO is an unnormalized stable orbital integral on the regular semisimple set of $G(\mathbb{R})$, i.e., that there is a Schwartz function f on $G(\mathbb{R})$ such that, for each regular semisimple γ in $G(\mathbb{R})$, $\text{SO}(\gamma)$ is the stable orbital integral $\text{SO}(\gamma, f)$. Suppose

also that γ lies in $T(\mathbb{R})$. Then we use the factors Δ' and $\Delta_{a,\chi}$ from the last section to define the transforms

$$\Psi(\gamma) = \Delta'(\gamma)\text{SO}(\gamma)$$

for a given choice of positive imaginary roots for T and

$$\Psi_{a,\chi}(\gamma) = \Delta_{a,\chi}(\gamma)\text{SO}(\gamma)$$

for a given choice of a -data and χ -data for T . The choice of measures has been suppressed in notation; we follow [She79a] (see also Section 11). Our purpose in the present section is to deduce simple limit formulas for $\Psi_{a,\chi}$ from the limit formulas for Ψ ; see [She79a] for a detailed proof of the latter.

We confine our attention to the behavior of orbital integrals near *semiregular* semisimple elements of $G(\mathbb{R})$, those elements with centralizer of type A_1 . Suppose then that γ_0 is a semiregular element of $T(\mathbb{R})$, that $\alpha(\gamma_0) = 1$, where α is an imaginary root which is not totally compact, and that s is a Cayley transform with respect to α . We may regard the coroot α^\vee as an element of the Lie algebra of T , and then $a_\alpha\alpha^\vee$ lies in the real Lie algebra: $\sigma(a_\alpha\alpha^\vee) = a_{-\alpha}(-\alpha^\vee) = a_\alpha\alpha^\vee$. For a sufficiently small nonzero real number ν , the element $\gamma_\nu = \gamma_0 \exp(\nu a_\alpha\alpha^\vee)$ is a regular element in $T(\mathbb{R})$. Moreover, it is unchanged if α is replaced by $-\alpha$. At the same time, the element γ_0^s lies in $T^s(\mathbb{R})$ and is annihilated only by the real roots $\pm\alpha^s$. Then $\Psi_{a^s,\chi^s}(\gamma_0^s)$ is prescribed by smooth extension [HC75, §17, Th. 1]. In particular, if we set $\gamma_{s,\nu} = \gamma_0^s \exp(\nu a_{\alpha^s}(\alpha^s)^\vee)$, then

$$\Psi_{a^s,\chi^s}(\gamma_0^s) = \lim_{\nu \rightarrow 0} \Psi_{a^s,\chi^s}(\gamma_{s,\nu}).$$

We note first a lemma that simplifies our argument for the next theorem (and motivates the definition of s -compatibility).

LEMMA 4.1. *For any s -compatible data set $\{a_\beta\}, \{\chi_\beta\}, \{a_{\beta^s}\}, \{\chi_{\beta^s}\}$, we have*

$$\prod_{\mathcal{O} \neq \mathcal{O}_\alpha} \chi_\beta \left(\frac{(\beta(\gamma_0) - 1)}{a_\beta} \right) = \prod_{\mathcal{O}^s \neq \pm \mathcal{O}_{\alpha^s}} \chi_{\beta^s} \left(\frac{(\beta^s(\gamma_0^s) - 1)}{a_{\beta^s}} \right).$$

On the left, the product is over all Galois orbits \mathcal{O} for T except $\mathcal{O}_\alpha = \{\pm\alpha\}$. Each term is independent of the choice of representative β for \mathcal{O} . The right side is defined by using all Galois orbits for T^s except $\{\alpha^s\}$ and $\{-\alpha^s\}$, and again the choice of representative has no effect on the terms.

Proof. If \mathcal{O} is orthogonal to \mathcal{O}_α , then we find immediately a matching term for \mathcal{O} on the right side of the equation. For the remaining cases, if β is imaginary and $\beta' = \omega_\alpha\beta$ is distinct from β , then the contributions to the left from $\{\pm\beta\}$ and $\{\pm\beta'\}$ are equal and moreover they each equal the contribution to the right from each of the two orbits $\{\beta^s, -(\beta')^s\}$ and $\{-\beta^s, (\beta')^s\}$. If β

is complex and β^s is complex, then we clearly have matching terms. If β is complex and β^s is real, then $(\sigma\beta)^s = \omega_{\alpha^s}\beta^s$. The product of the terms for $\{\beta, \sigma\beta\}$ and $\{-\beta, -\sigma\beta\}$ is $\chi_\beta(\beta(\gamma_0))$. The product of the terms for $\{\beta^s\}$, $\{-\beta^s\}$ is $\chi_{\beta^s}(\beta^s(\gamma_0^s))$, which equals the product for $\{\omega_{\alpha^s}\beta^s\}$, $\{-\omega_{\alpha^s}\beta^s\}$. Since $\beta(\gamma_0) = \beta^s(\gamma_0^s)$ is real, s -compatibility ensures that

$$\chi_\beta(\beta(\gamma_0)) = \chi_{\beta^s}(\beta^s(\gamma_0^s)^2) = \chi_{\beta^s}(\beta^s(\gamma_0^s)) \cdot \chi_{\omega_{\alpha^s}\beta^s}(\omega_{\alpha^s}\beta^s(\gamma_0^s)),$$

and the lemma is proved. □

THEOREM 4.2. *For any s -compatible data set, we have*

$$\lim_{\nu \rightarrow 0^-} \Psi_{a,\chi}(\gamma_\nu) = - \lim_{\nu \rightarrow 0^+} \Psi_{a,\chi}(\gamma_\nu)$$

and

$$\lim_{\nu \rightarrow 0^+} \Psi_{a,\chi}(\gamma_\nu) = \Psi_{a^s,\chi^s}(\gamma_0^s).$$

Proof. As a first step, we check that it is sufficient to verify these limits for one s -compatible data set. Suppose then that the result is true for the choice $\{a_\beta\}$, $\{\chi_\beta\}$ and $\{a_{\beta^s}\}$, $\{\chi_{\beta^s}\}$. We now use another set, which we write as $\{a_\beta b_\beta\}$, $\{\chi_\beta \eta_\beta\}$ and $\{a_{\beta^s} b_{\beta^s}\}$, $\{\chi_{\beta^s} \eta_{\beta^s}\}$, and we consider the effect on $\Psi_{a,\chi}(\gamma_\nu)$ and $\Psi_{a^s,\chi^s}(\gamma_0^s)$. We may argue orbit by orbit.

Notice that only the data for $\mathcal{O}_\alpha = \{\pm\alpha\}$ affect γ_ν . The characters $\eta_{\pm\alpha} = \eta_\alpha^{\pm 1}$ are trivial on \mathbb{R}^\times , while $b_\alpha = b_{-\alpha}$ may be any nonzero real number. Then γ_ν is replaced by $\gamma_{b_\alpha\nu}$ and $\Delta_{a,\chi}(\gamma_\nu)$ is multiplied by

$$\chi_\alpha(b_\alpha)^{-1} \eta_\alpha \left(\frac{\alpha(\gamma_\nu) - 1}{a_\alpha} \right) = \text{sign}(b_\alpha) \eta_\alpha(e^{\nu a_\alpha})$$

since $\alpha(\gamma_\nu) = e^{2\nu a_\alpha}$ and $(e^{\nu a_\alpha} - e^{-\nu a_\alpha})/a_\alpha$ is real. Thus the first limit statement remains true (each side is replaced with the negative of the other if b_α is negative), and then the second limit statement follows also. Next we observe that $\eta_{\pm\alpha^s} = \eta_{\alpha^s}^{\pm 1}$ and $b_{\alpha^s} = b_{-\alpha^s}$ contribute no change to $\Psi_{a^s,\chi^s}(\gamma_0^s)$ since

$$\eta_{\alpha^s} \left(\frac{\alpha^s(\gamma) - 1}{a_{\alpha^s}} \right) \eta_{-\alpha^s} \left(\frac{\alpha^s(\gamma)^{-1} - 1}{a_{-\alpha^s}} \right) \chi_{\alpha^s}(b_{\alpha^s})^{-1} \chi_{-\alpha^s}(b_{-\alpha^s})^{-1} = \eta_{\alpha^s}(\alpha^s(\gamma))$$

for any regular γ in $T^s(\mathbb{R})$, and so has limit 1 as γ approaches γ_0^s . Thus we are done with the orbits $\mathcal{O}_\alpha, \pm\mathcal{O}_{\alpha^s}$.

For the remaining orbits, we could do a calculation for each symmetric orbit \mathcal{O} and each asymmetric pair $\pm\mathcal{O}$ individually. Instead we appeal to Lemma 4.1 to see that the (nonzero) total contribution can be cancelled from the limit formulas. This finishes the first step.

The second step in our proof is to compare the proposed limit formulas with the limit formulas for the stable version $\Psi = \Delta'.\text{SO}$ of Harish-Chandra's $'F_f$ transform ([She79a], recalled in Section 3 of [She08a]). It is convenient to assume first that α itself is noncompact and then drop this assumption later.

We pick a system of positive imaginary roots for T that is adapted to α . This means that α is positive and that if β is positive imaginary and not orthogonal to α , then $\beta_1 = -\omega_\alpha(\beta)$ is also positive. For convenience, we will choose χ_β to be the standard character $z \rightarrow \frac{z}{|z|}$ if β is positive imaginary and orthogonal to α . This is also assumed for $\beta = \alpha$. In each of these cases we set $a_\beta = i$. For each pair of positive roots $\beta, \beta_1 = -\omega_\alpha(\beta)$ not orthogonal to α and distinct from α , we pick one, labeling it β , and make χ_β the standard character. Then χ_{β_1} must be its inverse. Also we set $a_\beta = i$, so that a_{β_1} must be $-i$. We assume that χ_β is the identity character if β is real or complex.

Now we compare $\Delta_{a,\chi}(\gamma)$ with $\Delta'(\gamma)$ at $\gamma_\nu = \gamma_0 \exp(i\nu\alpha^\vee)$, as well as $\Delta_{a^s,\chi^s}(\gamma_0^s)$ with $\Delta'(\gamma_0^s)$. We proceed orbit by orbit, considering the contribution of \mathcal{O} to the change for $\Delta_{a,\chi}$ and of \mathcal{O}^s to the change for Δ_{a^s,χ^s} . Real or complex orbits for T contribute no change to either $\Delta'(\gamma_\nu)$ or $\Delta'(\gamma_0^s)$. Consider the imaginary orbits orthogonal to α . Suppose there are N such orbits. Then passage to $\Delta_{a,\chi}(\gamma_\nu)$ multiplies $\Delta'(\gamma_\nu)$ by $(i)^{-N}$. Since N is the number of imaginary orbits for T^s and we use s -compatible data for T^s , the term $\Delta'(\gamma_0^s)$ is also multiplied by $(i)^{-N}$. Consider next the orbits of a pair of positive imaginary roots β, β_1 not orthogonal to α and distinct from α . Then we replace

$$A(\gamma_\nu) = \frac{\beta(\gamma_\nu) - 1}{|\beta(\gamma_\nu) - 1|} \cdot \frac{\beta_1(\gamma_\nu) - 1}{|\beta_1(\gamma_\nu) - 1|}$$

by

$$B(\gamma_\nu) = \frac{\beta(\gamma_\nu) - 1}{|\beta(\gamma_\nu) - 1|} \cdot \frac{|\beta_1(\gamma_\nu) - 1|}{-(\beta_1(\gamma_\nu) - 1)/i} = \frac{\beta(\gamma_\nu) - 1}{1 - \beta_1(\gamma_\nu)} \cdot \frac{|\beta_1(\gamma_\nu) - 1|}{|\beta(\gamma_\nu) - 1|}.$$

Because $\beta_1(\gamma_0) = \beta(\gamma_0)^{-1} = \overline{\beta(\gamma_0)}$, we have

$$\lim_{\nu \rightarrow 0^+} A(\gamma_\nu) = \lim_{\nu \rightarrow 0^-} A(\gamma_\nu) = 1,$$

whereas

$$\lim_{\nu \rightarrow 0^+} B(\gamma_\nu) = \lim_{\nu \rightarrow 0^-} B(\gamma_\nu) = \beta(\gamma_0).$$

Thus we have to multiply all limits by $\beta(\gamma_0)$. Consider now the change to $\Delta'(\gamma_0^s)$. This term is multiplied by

$$\begin{aligned} \chi_{\beta^s} \left(\frac{\beta^s(\gamma_0^s) - 1}{a_{\beta^s}} \right) \cdot \chi_{-\beta^s} \left(\frac{\beta^s(\gamma_0^s)^{-1} - 1}{-a_{\beta^s}} \right) &= \chi_{\beta^s} \left(\frac{\beta^s(\gamma_0^s) - 1}{1 - \beta^s(\gamma_0^s)^{-1}} \right) \\ &= \chi_\beta \left(\frac{\beta(\gamma_0) - 1}{1 - \beta_1(\gamma_0)} \right) = B(\gamma_0) = \beta(\gamma_0), \end{aligned}$$

and so we are done with this case.

There is one remaining orbit, that of α . Its contribution multiplies $\Delta'(\gamma_\nu)$ by i^{-1} , but there is no change to $\Delta'(\gamma_0^s)$. This is exactly what we need to deduce the claimed limits from the analogous limits for the Harish-Chandra

type function Ψ (see [She08a, §3, Property (vi)]). Step 2 is thus complete and the assertions of the theorem proved for the case that α is noncompact.

Suppose that α is compact and that ω is an element of the imaginary Weyl group for which $\alpha^\dagger = \omega^{-1}\alpha$ is noncompact. Assume that ω acts on T as $\text{Int}(w)$. Then if s is a Cayley transform relative to α , $s^\dagger = w^{-1}s$ is a Cayley transform relative to α^\dagger . Also if γ_0 is a semiregular element of $T(\mathbb{R})$ such that $\alpha(\gamma_0) = 1$, then γ_0^w is a semiregular element of $T(\mathbb{R})$, $\alpha^\dagger(\gamma_0^w) = 1$, and $(\gamma_0^w)^{s^\dagger} = \gamma_0^s$. Finally, to obtain an s^\dagger -compatible data set from an s -compatible data set $\{a_\beta\}, \{\chi_\beta\}$ and $\{a_{\beta^s}\}, \{\chi_{\beta^s}\}$, we may replace $\{a_\beta\}, \{\chi_\beta\}$ by $\{a'_\beta\}, \{\chi'_\beta\}$, where $a'_\beta = a_{\omega\beta}$ and $\chi'_\beta = \chi_{\omega\beta}$, and leave $\{a_{\beta^s}\}, \{\chi_{\beta^s}\}$ unchanged. Then

$$\gamma_\nu^w = \gamma_0^w \exp \nu a_{\alpha^\dagger} (\alpha^\dagger)^\vee,$$

and because SO is stable, we have

$$\Psi_{a,\chi}(\gamma_\nu) = \Psi_{a',\chi'}(\gamma_\nu^w).$$

The limit formulas at γ_0 now follow immediately from those at γ_0^w , and this completes the proof of Theorem 4.2. \square

Notice that Lemma 4.1 allows us to use Δ_α in place of $\Delta_{a,\chi}$ in the statement of Theorem 4.2, where

$$\Delta_\alpha(\gamma) = \chi_\alpha\left(\frac{(\alpha(\gamma) - 1)}{a_\alpha}\right) \left| \det(\text{Ad}(\gamma) - I)_{\mathfrak{g}/\mathfrak{t}} \right|^{1/2}.$$

Here $\Delta_{-\alpha} = \Delta_{\sigma\alpha} = \Delta_\alpha$, and so only the (symmetric) orbit \mathcal{O} of α matters. We then write $\Delta_{\mathcal{O}}$ in place of Δ_α .

We end this section with a remark on the *normalized orbital integral*

$$\Phi(\gamma) = \left| \det(\text{Ad}(\gamma) - I)_{\mathfrak{g}/\mathfrak{t}} \right|^{1/2} \text{SO}(\gamma).$$

Set

$$\Psi_{\mathcal{O}}(\gamma) = \Delta_{\mathcal{O}}(\gamma) \text{SO}(\gamma) = \chi_\alpha\left(\frac{(\alpha(\gamma)-1)}{a_\alpha}\right) \Phi(\gamma).$$

Assume, as in the theorem, that α is not totally compact. Notice that if we write a_α as ib_α , where b_α is real, then for $|\nu|$ small and nonzero, we have

$$\begin{aligned} \chi_\alpha\left(\frac{(\alpha(\gamma_\nu) - 1)}{a_\alpha}\right) &= \chi_\alpha(e^{i\nu b_\alpha}) \chi_\alpha\left(\frac{e^{i\nu b_\alpha} - e^{-i\nu b_\alpha}}{ib_\alpha}\right) \\ &= \chi_\alpha(e^{i\nu b_\alpha}) \chi_\alpha\left(\frac{2 \sin(\nu b_\alpha)}{\nu b_\alpha} \nu\right) = \chi_\alpha(e^{\nu a_\alpha}) \text{sign}(\nu). \end{aligned}$$

Because s defines an inner twist between the identity components of their respective centralizers, the elements γ_0 and γ_0^s are stably conjugate in $G(\mathbb{R})$ in the sense introduced by Kottwitz in Section 3 of [Kot82] for the untwisted setting. Comparing limits for $\Psi_{\mathcal{O}}$ with limits for Φ , we see, by an argument along the lines of Section 2 that the assertions of Theorem 4.2 may be rephrased as the existence and equality of the limits of $\Phi(\gamma)$ as (i) γ approaches γ_0

through the regular elements of $T(\mathbb{R})$ and (ii) γ approaches the stable conjugate γ_0^s of γ_0 through the regular elements of $T^s(\mathbb{R})$ (see Section 2 of [She83]). This suggests another approach to the proof of transfer; we simply found our present approach quicker. Our preference for working with $\Psi_{\mathcal{O}}$ rather than Φ is explained by the formulas of Section 10 for derivatives.

It is now a short exercise to modify the characterization theorem for stable orbital integrals in [She79a] using the statement of Theorem 4.2 or, more precisely, its generalization to derivatives. As mentioned in Section 1, we will need eventually to introduce a slight twist in the stable integrals. Thus we will wait until Section 12, and then write a slightly more general characterization theorem (Theorem 12.1).

5. A limit formula for twisted orbital integrals

We return to the statement of the main theorem in Section 2 and follow the notation established in that setting. In particular, we will consider (θ, ϖ) -twisted integrals for G , while the endoscopic group H_1 will now assume the role of the group of the last two sections. Recall that, because of our assumption on the inner twist (G, θ, ψ) , we consider completely untwisted integrals on $H_1(\mathbb{R})$. To commence the proof of the main theorem, we assume that $f \in \mathcal{C}(G(\mathbb{R}), \theta)$ and define a function Φ_1 on the strongly G -regular elements γ_1 of $H_1(\mathbb{R})$ by

$$\Phi_1(\gamma_1) = \left| \det(\text{Ad}(\gamma_1) - I)_{\mathfrak{h}_1/\mathfrak{t}_1} \right|^{1/2} \sum_{\delta, \theta\text{-conj}} \Delta(\gamma_1, \delta) O^{\theta, \varpi}(\delta, f).$$

We must show Φ_1 is a normalized stable orbital integral on $H_1(\mathbb{R})$. Our primary concern will be an analogue of the limit formulas of the last section.

Consider Φ_1 near a semiregular element γ_0 in $H_1(\mathbb{R})$ annihilated by an imaginary root α_1 of a maximal torus T_1 in H_1 . Because H_1 is quasi-split over \mathbb{R} , the root α_1 is not totally compact [She79b, Lemma 9.2]. We then have a Cayley transform s_1 in the sense of Section 3 for α_1 , along with the semiregular element $\gamma_0^{s_1}$ in the adjacent Cartan subgroup $T_1^{s_1}(\mathbb{R})$ annihilated by the real root $\alpha_1^{s_1}$. We will choose an s_1 -compatible data set in Section 9 based on compatible twisted data. We make the additional requirement that γ_0 be G -semiregular; see Section 6 for definition. For all nonzero real ν with $|\nu|$ sufficiently small, we will see that both $\gamma_\nu = \gamma_0 \exp(\nu a_{\alpha_1} \alpha_1^\vee)$ in $T_1(\mathbb{R})$ and $\gamma_{s_1, \nu} = \gamma_0^{s_1} \exp(\nu a_{\alpha_1^{s_1}} (\alpha_1^{s_1})^\vee)$ in $T_1^{s_1}(\mathbb{R})$ are G -regular and then that $\Phi_1(\gamma_\nu)$, $\Phi_1(\gamma_{s_1, \nu})$ are defined.

THEOREM 5.1. *All relevant limits exist, and the assertions of Theorem 4.2 are true for the group H_1 when Φ (normalized stable orbital integral on $H_1(\mathbb{R})$) is replaced by Φ_1 (normalized transport of a weighted sum of twisted integrals*

on $G(\mathbb{R})$:

$$\lim_{\nu \rightarrow 0^-} \Psi_{a,\chi}(\gamma_\nu) = - \lim_{\nu \rightarrow 0^+} \Psi_{a,\chi}(\gamma_\nu)$$

and

$$\lim_{\nu \rightarrow 0^+} \Psi_{a,\chi}(\gamma_\nu) = \lim_{\nu \rightarrow 0} \Psi_{a^{s_1}, \chi^{s_1}}(\gamma_{s_1, \nu}).$$

We will gather ingredients for a proof of the theorem over the next four sections, completing the argument in Section 10. Later the theorem will be strengthened to include derivatives (see Lemmas 10.1, 10.2) and all semiregular γ_0 (see Section 11). Often we will write γ'_ν for $\gamma_{s_1, \nu}$ and a', χ' for a^{s_1}, χ^{s_1} .

To begin, we replace $O^{\theta, \varpi}(\delta, f)$ by the normalized integral

$$\Phi^{\theta, \varpi}(\delta, f) = \left| \det(\text{Ad}(\delta) \circ \theta - I)_{\mathfrak{g}/\text{Cent}(\mathfrak{g}_\delta^{\theta, \varpi})} \right|^{1/2} O^{\theta, \varpi}(\delta, f).$$

Assume strongly G -regular γ_1 is a norm of δ . Then the term $\Delta_{IV}(\gamma_1, \delta)$ in the transfer factor is the quotient of the normalizing term above by that for ordinary orbital integrals on $H_1(\mathbb{R})$. Thus our proposed normalized stable orbital integral is given on γ_1 by

$$\Phi_1(\gamma_1) = \sum_{\delta, \theta\text{-conj}} \frac{\Delta(\gamma_1, \delta)}{\Delta_{IV}(\gamma_1, \delta)} \Phi^{\theta, \varpi}(\delta, f).$$

We may as well assume for the rest of the paper that there exists a strongly G -regular element in $H_1(\mathbb{R})$ that is a norm, for otherwise the zero function lies in $\text{Trans}(f)$ and the main theorem is proved. We then fix a pair $(\bar{\gamma}, \bar{\delta})$, with strongly G -regular $\bar{\gamma} \in H_1(\mathbb{R})$ a norm of strongly θ -regular $\bar{\delta} \in G(\mathbb{R})$, in order to normalize transfer factors as mentioned in Section 1. We gather all terms involving only $(\bar{\gamma}, \bar{\delta})$ as

$$\Delta^*(\bar{\gamma}, \bar{\delta}) = \Delta(\bar{\gamma}, \bar{\delta}) [\Delta_I(\bar{\gamma}) \Delta_{II}(\bar{\gamma}) \Delta_{IV}(\bar{\gamma})]^{-1}.$$

Here we have dropped the second argument in our notation for $\Delta_I, \Delta_{II}, \Delta_{IV}$ since it plays no role. There is no harm for the proof of Theorem 5.1 in assuming that transfer factors are normalized so that

$$\Delta(\bar{\gamma}, \bar{\delta}) = \Delta_I(\bar{\gamma}) \Delta_{II}(\bar{\gamma}) \Delta_{IV}(\bar{\gamma}),$$

and then

$$\Delta^*(\bar{\gamma}, \bar{\delta}) = 1.$$

This allows us to rewrite $\Phi_1(\gamma_1)$, for any strongly G -regular $\bar{\gamma} \in H_1(\mathbb{R})$, as

$$\Delta_I(\gamma_1) \Delta_{II}(\gamma_1) \sum_{\delta, \theta\text{-conj}} \Delta_{III}(\gamma_1, \delta; \bar{\gamma}, \bar{\delta}) \Phi^{\theta, \varpi}(\delta, f),$$

where the summation is over θ -conjugacy classes of strongly θ -regular elements δ in $G(\mathbb{R})$. Here we declare the contribution of the class of δ to be zero if γ_1 is not a norm of δ .

If γ_1 is a norm of δ , then the torus $\text{Cent}(\gamma_1, H_1)$ is a norm group (in the sense of the next section) which, as noted in Section 1, implies that the character ϖ is trivial on $\text{Cent}_\theta(\delta, G)(\mathbb{R})$. The transformation rule (2) of Theorem 5.1.D of [KS99] further allows us to write $\Delta_{III}(\gamma_1, \delta; \bar{\gamma}, \bar{\delta})O^{\theta, \varpi}(\delta, f)$ in the form

$$\int_{\text{Cent}_\theta(\delta, G)(\mathbb{R}) \backslash G(\mathbb{R})} \Delta_{III}(\gamma_1, g^{-1}\delta\theta(g); \bar{\gamma}, \bar{\delta})f(g^{-1}\delta\theta(g))dg/dt.$$

As a function of δ , this is constant on θ -conjugacy classes, as is the normalizing factor $|\det(\text{Ad}(\delta) \circ \theta - I)_{\mathfrak{g}/\text{Cent}(\mathfrak{g}_\delta^\theta, \mathfrak{g})}|^{1/2}$ for $\Phi^{\theta, \varpi}(\delta, f)$. The set of elements with γ_1 as norm forms a single stable θ -conjugacy class of elements in $G(\mathbb{R})$, as will be reviewed in Sections 6 and 7. Thus the summation in $\Phi_1(\gamma_1)$ may be taken over the (finite) set of θ -conjugacy classes in this stable class.

In Section 7 we will define $\Phi_1(\gamma_1)$ for G -regular elements γ_1 that are not strongly G -regular in the same way as for the untwisted case, i.e., by smooth extension. First, we need to describe our choice of stable θ -conjugacy class with norm γ_1 in that setting. At the same time we prepare for the more delicate analysis of $\Phi_1(\gamma_1)$ when γ_1 is near semiregular γ_0 .

6. Norm groups and semiregular elements

To view semisimple elements of the endoscopic group $H_1(\mathbb{R})$ as norms, we adapt the definition of *image* in standard endoscopy (see (1.2) of [LS90]) to our twisted setting. Recall that we have made an assumption to avoid any twisting in $H_1(\mathbb{R})$. Namely, we have fixed quasi-split data (G^*, θ^*) and inner twist (G, θ, ψ) :

$$\psi\sigma(\psi)^{-1} = \text{Int}(u(\sigma))$$

and

$$\psi \circ \theta \circ \psi^{-1} = \text{Int}(g_\theta)^{-1} \circ \theta^*,$$

where $u(\sigma), g_\theta$ lie in G_{sc}^* . We write $u(\sigma), g_\theta$ also for the images of these two elements in G^* under the natural map $G_{\text{sc}}^* \rightarrow G^*$. Define $m : G \rightarrow G^*$ by $m(\delta) = \psi(\delta)g_\theta^{-1}$. Then our assumption is that we may choose $u(\sigma), g_\theta$ so that

$$\sigma(m)(\delta) = u(\sigma)^{-1}m(\delta)\theta^*(u(\sigma)).$$

See Lemma 3.1.A and Appendix B of [KS99] for its (hyper)cohomological significance. It is not difficult to drop the assumption, as we will check in Section 12.

We start our discussion of norms with the correspondence of [KS99] between the set of stable conjugacy classes of strongly G -regular elements in $H_1(\mathbb{R})$ and the set of stable θ -conjugacy classes of strongly θ -regular elements in $G(\mathbb{R})$. Recall from the last section that we may as well assume this correspondence is nonempty. It is uniquely determined by the choice of g_θ ; see [She11] for a related discussion. If the class of strongly θ -regular δ in $G(\mathbb{R})$ corresponds to the class of strongly G -regular γ_1 in $H_1(\mathbb{R})$, then γ_1 is a norm of δ . We will call a maximal torus T_1 over \mathbb{R} in H_1 a *norm group* for (G, θ) if

$T_1(\mathbb{R})$ contains strongly G -regular elements that are norms of strongly θ -regular elements in $G(\mathbb{R})$; this generalizes a definition in [KS99, §3.3].

Let T_1 be a maximal torus over \mathbb{R} in H_1 . Then by Lemma 3.3.B of [KS99] there exist a θ^* -stable maximal torus T in G^* defined over \mathbb{R} and an admissible homomorphism $T_1 \rightarrow T_{\theta^*}$ from T_1 to the coinvariants of θ^* in T . In more detail: there exist a θ^* -stable maximal torus T in G^* defined over \mathbb{R} and a θ^* -stable Borel subgroup B containing T , along with Borel subgroup B_1 containing T_1 such that the homomorphism

$$T_1 \rightarrow T_1/Z_1 \rightarrow T_{\theta^*}$$

attached to the pairs (B_1, T_1) and (B, T) is defined over \mathbb{R} . Here the map $T_1 \rightarrow T_1/Z_1$ is the natural projection, and the construction of $T_1/Z_1 \rightarrow T_{\theta^*}$ comes from the definition of endoscopic data. The strongly θ^* -regular elements of $T(\mathbb{R})$, which include a dense subset of $T(\mathbb{R})^0$, have strongly G^* -regular norms in $T_1(\mathbb{R})$, and so the cited lemma shows that any maximal torus over \mathbb{R} in H_1 is a norm group for the pair (G^*, θ^*) .

Assume now that T_1 is a norm group for (G, θ) . Suppose that γ_1 is a strongly G -regular element of $T_1(\mathbb{R})$ and that γ_1 is a norm of strongly θ -regular δ in $G(\mathbb{R})$. First we take an admissible homomorphism $T_1 \rightarrow T_{\theta^*}$ mapping γ_1 to an element, say γ^* , of $T_{\theta^*}(\mathbb{R})$. Because γ_1 is a norm of δ , there is also an associated isomorphism

$$\text{Int}(g) \circ \psi : G_{\delta}^{\theta} \rightarrow (T^{\theta^*})^0$$

defined over \mathbb{R} , where g is chosen in G_{sc}^* so that

$$\delta^* = gm(\delta)\theta^*(g)^{-1}$$

lies in T and $N(\delta^*) = \gamma^*$; see [KS99, §§3.3, 4.4]. Here, as in [KS99, §3.2], N denotes the *abstract norm* map, i.e., the projection $T \rightarrow T_{\theta^*}$ to coinvariants, while G_{δ}^{θ} denotes $\text{Cent}_{\theta}(\delta, G)^0$, a torus defined over \mathbb{R} . In the equation $\delta^* = gm(\delta)\theta^*(g)^{-1}$, the element g has been identified with its image in G^* (we will do this repeatedly, often without mention) and m is the modification of the inner twist $\psi : G \rightarrow G^*$ defined in the first paragraph. Because of the strong regularity condition, g is unique up to an element of T_{sc} once $T_1 \rightarrow T_{\theta^*}$ has been fixed. Also, changing $T_1 \rightarrow T_{\theta^*}$ changes g in a simple manner [KS99, §4.4].

In summary, if strongly G -regular γ_1 in $H_1(\mathbb{R})$ is a norm of strongly θ -regular δ in $G(\mathbb{R})$, we identify the quotient of $\text{Cent}(\gamma_1, H_1) = T_1$ by Z_1 with the group of θ^* -coinvariants in T . Here T is provided by the data for an admissible homomorphism $T_1 \rightarrow T_{\theta^*}$. We also identify $G_{\delta}^{\theta} = \text{Cent}_{\theta}(\delta, G)^0$ with the identity component of the group of θ^* -invariants in T .

Recall that the strong θ -regularity of δ ensures only that $\text{Cent}_{\theta}(\delta, G)$ is abelian and diagonalizable. The isomorphism $\text{Int}(g) \circ \psi$ above maps $\text{Cent}_{\theta}(\delta, G)$ onto the full group of θ^* -invariants in T .

Now we drop the assumption of *strong G-regularity* on a semisimple element in $H_1(\mathbb{R})$. Then the ambient norm group is not unique unless the element is *G-regular*, and so we proceed torus by torus.

Suppose that γ_0 is an element in the norm group $T_1(\mathbb{R})$, and assume that δ_0 is a θ -semisimple element of $G(\mathbb{R})$. Then, by definition [KS99, §3.2], $\text{Int}(\delta_0) \circ \theta$ preserves some pair (B^\dagger, T^\dagger) . Write T^{δ_0} for the identity component of the fixed points of $\text{Int}(\delta_0) \circ \theta$ in T^\dagger . Then T^{δ_0} is a maximal torus in the reductive group $G_{\delta_0}^\theta$ defined over \mathbb{R} , and we may assume T^{δ_0} is defined over \mathbb{R} . (Otherwise replace (B^\dagger, T^\dagger) by a suitable $G_{\delta_0}^\theta$ -conjugate pair.) Fix an admissible homomorphism $T_1 \rightarrow T_{\theta^*}$, and write γ_0^* for the image of γ_0 . Then there is an isomorphism $\text{Int}(g) \circ \psi$ carrying (B^\dagger, T^\dagger) to (B, T) , where $g \in G_{\text{sc}}^*$. This implies that $\delta_0^* = gm(\delta_0)\theta^*(g)^{-1}$ lies in T .

Definition. We call γ_0 a T_1 -norm of δ_0 if we may choose $g \in G_{\text{sc}}^*$ so that (i) $N(\delta_0^*) = \gamma_0^*$ and (ii) the isomorphism $\text{Int}(g) \circ \psi : T^{\delta_0} \rightarrow (T^{\theta^*})^0$ is defined over \mathbb{R} .

In the case that γ_0 is strongly *G-regular* (ii) follows from (i) [KS99, (3.3.6)]. In general, for given T_1 , the choice of admissible homomorphism $T_1 \rightarrow T_{\theta^*}$ does not affect the existence of g .

Next, we consider together all elements in the γ_0 -component $\gamma_0 T_1(\mathbb{R})^0$ of $T_1(\mathbb{R})$.

LEMMA 6.1. *The following are equivalent for $\gamma_0 \in T_1(\mathbb{R})$:*

- (i) γ_0 is a T_1 -norm,
- (ii) some strongly *G-regular* element in the γ_0 -component is a norm,
- (iii) every element of the γ_0 -component is a T_1 -norm.

Proof. Fix an admissible homomorphism $T_1 \rightarrow T_{\theta^*}$, and assume that $\gamma_0 \in T_1(\mathbb{R})$ is a T_1 -norm of a θ -semisimple $\delta_0 \in G(\mathbb{R})$. Choose elements g, δ_0^* as in the definition. Take ε in the identity component of the Cartan subgroup $T^{\delta_0}(\mathbb{R})$ of $G_{\delta_0}^\theta(\mathbb{R})$, and consider $\delta = \varepsilon\delta_0$. Then δ is θ -semisimple since $\text{Int}(\delta) \circ \theta$ preserves the same pair (B^\dagger, T^\dagger) as $\text{Int}(\delta_0) \circ \theta$. Also, by results of Steinberg (see Theorem 1.1.A in [KS99]), $G_\delta^\theta(\mathbb{R})$ contains $T^{\delta_0}(\mathbb{R})$ as Cartan subgroup. Further we may choose ε so that δ is strongly θ -regular; the elements ε with this property are dense in $T^{\delta_0}(\mathbb{R})^0$. Set

$$\delta^* = gm(\delta)\theta^*(g)^{-1} = gm(\varepsilon\delta_0)\theta^*(g)^{-1} = g\psi(\varepsilon)g^{-1}.gm(\delta_0)\theta^*(g)^{-1} = \varepsilon^*\delta_0^* = \delta_0^*\varepsilon^*,$$

where $\varepsilon^* = g\psi(\varepsilon)g^{-1}$ lies in $(T^{\theta^*}(\mathbb{R}))^0$. The image of the γ_0 -component in $T_1(\mathbb{R})$ under $T_1 \rightarrow T_{\theta^*}$ then contains $N(\delta^*) = \gamma_0^*N(\varepsilon^*)$, where γ_0^* is, as before, the image of γ_0 under $T_1 \rightarrow T_{\theta^*}$. Since δ^* is strongly θ^* -regular, each element in the γ_0 -component which maps to $N(\delta^*)$ under $T_1 \rightarrow T_{\theta^*}$ is strongly *G-regular*, and (ii) now follows.

Assume (ii), and suppose strongly G -regular γ_1 in the γ_0 -component of $T_1(\mathbb{R})$ is a norm of δ . Choose δ^*, g as in the definition of norm for strongly G -regular elements. By our assumption that the restriction of θ to the center of G is (strongly) semisimple, the homomorphism $N : T^{\theta^*}(\mathbb{R})^0 \rightarrow T_{\theta^*}(\mathbb{R})^0$ is surjective. Thus the image of $\gamma_1 T_1(\mathbb{R})^0$ under $T_1 \rightarrow T_{\theta^*}$ coincides with the image under N of $\delta^* T^{\theta^*}(\mathbb{R})^0$. We write an element γ_2 of $\gamma_1 T_1(\mathbb{R})^0 = \gamma_0 T_1(\mathbb{R})^0$ as the image under N of some element δ_2^* in $\delta^* T^{\theta^*}(\mathbb{R})^0$. Then, as in Lemma 4.4.A of [KS99],

$$\sigma(\delta_2^*)\delta_2^{*-1} = \sigma(\delta^*)\delta^{*-1} = (\theta^* - 1)v(\sigma),$$

where the cochain $v(\sigma)$ is (the image in T of) the cochain $gu(\sigma)\sigma(g)^{-1}$ in T_{sc} . Thus

$$\delta_2 = m^{-1}(g^{-1}\delta_2^*\theta^*(g))$$

is θ -semisimple, lies in $G(\mathbb{R})$, and has norm γ_2 , so that (iii) follows. The rest is immediate. \square

We now expand on the argument for (i) \Rightarrow (ii) in the last lemma. Write the element ε defined there as $\exp Y$, where Y belongs to the Cartan subalgebra $\mathfrak{t}^{\delta_0}(\mathbb{R})$ of the Lie algebra $\mathfrak{g}_{\delta_0}^\theta(\mathbb{R})$ of $G_{\delta_0}^\theta(\mathbb{R})$. Let Y map to Y^* , where $Y^* \in \mathfrak{t}^{\theta^*}(\mathbb{R})$, under the bijection provided by $\text{Int}(g) \circ \psi$. Recall from the definition of z -pair that we have the exact sequence $1 \rightarrow Z_1 \rightarrow H_1 \rightarrow H \rightarrow 1$, with Z_1 central in H_1 . We split the corresponding sequence for Lie algebras in the usual manner and identify, over \mathbb{R} , the Lie algebra \mathfrak{h} as a subalgebra of \mathfrak{h}_1 complementary to \mathfrak{z}_1 . Then the Lie algebra \mathfrak{t}_H of T_1/Z_1 is a subspace of \mathfrak{t}_1 complementary to \mathfrak{z}_1 . There is a linear isomorphism

$$\mathfrak{t}^{\theta^*}(\mathbb{R}) \rightarrow \mathfrak{t}_{\theta^*}(\mathbb{R}) \rightarrow \mathfrak{t}_H(\mathbb{R})$$

determined by the restriction of $N : T \rightarrow T_{\theta^*}$ to θ^* -invariants and the chosen admissible isomorphism $T_{\theta^*} \rightarrow T_1/Z_1$. Write Y_H for the image of Y^* , so that we have

$$\mathfrak{t}^{\delta_0}(\mathbb{R}) \ni Y \leftrightarrow Y^* \leftrightarrow N(Y^*) \leftrightarrow Y_H \in \mathfrak{t}_H(\mathbb{R}).$$

Write $Y_1 \in \mathfrak{t}_1(\mathbb{R})$ as $Y_1 = Y_H + Y_{\mathfrak{z}_1}$. Then the following is immediate.

LEMMA 6.2. *Assume that γ_0 is a T_1 -norm of δ_0 and $Y_1 \in \mathfrak{t}_1(\mathbb{R})$. Then the element*

$$\gamma_0(Y_1) = \gamma_0 \cdot \exp Y_1 = \exp Y_1 \cdot \gamma_0$$

in the γ_0 -component of $T_1(\mathbb{R})$ is a T_1 -norm of the element

$$\delta_0(Y) = \exp Y \cdot \delta_0 = \delta_0 \cdot \exp \theta Y.$$

The cochain $v(\sigma)$ attached to γ_0 also serves for $\gamma_0(Y_1)$, while the attached element of T is

$$\delta^*(Y) = \delta_0^* \cdot \exp Y^* = \exp Y^* \cdot \delta_0^*.$$

Now we consider all tori T_1 containing a given semisimple element γ_0 in $H_1(\mathbb{R})$. Let δ_0 be a θ -semisimple element of $G(\mathbb{R})$. We call γ_0 a *norm* of δ_0 (or, for emphasis on the ambient group, a *G-norm* of δ_0) if there exists a norm group T_1 such that γ_0 is a T_1 -norm of δ_0 . Otherwise we say that γ_0 is not a (G -) norm. The following will be proved after Lemma 6.6.

LEMMA 6.3. *Let γ_0 be a semisimple element in $H_1(\mathbb{R})$ and let δ_0, δ'_0 be θ -semisimple elements of $G(\mathbb{R})$. Then (i) if γ_0 is a G -norm of δ_0 , then so are all stable conjugates of γ_0 in $H_1(\mathbb{R})$, and (ii) if γ_0 is a G -norm of both δ_0 and δ'_0 , then δ_0 and δ'_0 are stably θ -conjugate.*

Remark. By δ'_0 is stably θ -conjugate to δ_0 we mean that we may write $\delta'_0 \in G(\mathbb{R})$ as $x\delta_0\theta(x)^{-1}$, where $x \in G$ and $\text{Int}(x): G_{\delta_0}^\theta \rightarrow G_{\delta'_0}^\theta$ is an inner twist.

Remark. As pointed out by a referee, the converse statement for (ii) in Lemma 6.3 is false in general.

Assume now that semisimple $\gamma_0 \in H_1(\mathbb{R})$ is a T_1 -norm of $\delta_0 \in G(\mathbb{R})$. Fix admissible $T_1 \rightarrow T_{\theta^*}$ and choose g, δ_0^* as in the definition of T_1 -norm. Then $\text{Int}(g) \circ \psi$ is an isomorphism of $G_{\delta_0}^\theta$ with $(G^*)_{\delta_0^*}^{\theta^*}$. We will abbreviate (slightly) the notation for the latter group as $G_{\delta_0^*}^{\theta^*}$. We have required that $\text{Int}(g) \circ \psi$ maps, over \mathbb{R} , the maximal torus T^{δ_0} over \mathbb{R} in $G_{\delta_0}^\theta$ to the maximal torus $(T^{\theta^*})^0$ in $G_{\delta_0^*}^{\theta^*}$. In the case that $G_{\delta_0}^\theta$ is of Dynkin type A_1 we claim that this requirement ensures first that $G_{\delta_0^*}^{\theta^*}$ is defined over \mathbb{R} and then that $\text{Int}(g) \circ \psi : G_{\delta_0}^\theta \rightarrow G_{\delta_0^*}^{\theta^*}$ is an inner twist. Indeed, $\text{Int}(g) \circ \psi$ transports the two roots of T^{δ_0} in $G_{\delta_0}^\theta$, either both imaginary or both real, to the roots of $(T^{\theta^*})^0$ in $G_{\delta_0^*}^{\theta^*}$ which must be of the same type. An argument with root vectors then finishes the proof.

With no restriction on the Dynkin type of $G_{\delta_0}^\theta$, we will prove the next lemma at the end of this section.

LEMMA 6.4. *Suppose that semisimple $\gamma_0 \in H_1(\mathbb{R})$ is a T_1 -norm of $\delta_0 \in G(\mathbb{R})$ and that $T_1 \rightarrow T_{\theta^*}$ is an admissible homomorphism. Then we may choose the elements g, δ_0^* so that (i) $\sigma(\delta_0^*)\delta_0^{*-1}$ is central in G^* , and (ii) $v(\sigma) = gu(\sigma)\sigma(g)^{-1}$ lies in the product of the torus $(T_{\text{sc}})^{\theta^*}$ with the center of G_{sc}^* .*

In particular, if G is of adjoint type, then we may arrange that δ_0^* lies in $T(\mathbb{R})$. In general, for any g, δ_0^* as in this lemma, the group $G_{\delta_0^*}^{\theta^*}$ is defined over \mathbb{R} and $\text{Int}(g) \circ \psi : G_{\delta_0}^\theta \rightarrow G_{\delta_0^*}^{\theta^*}$ is an inner twist.

Before continuing with the case that $G_{\delta_0}^\theta$ is of Dynkin type A_1 , we record an explicit analysis of the roots of $G_{\delta_0}^\theta$ and $G_{\delta_0^*}^{\theta^*}$ following Steinberg (see [KS99, Chap. 1]). By a *restricted root* we will mean the restriction α_{res} of a root α of T in G^* to the torus $(T^{\theta^*})^0$. This torus is maximal in each of the reductive

groups $(G^{*\theta^*})^0$ and $G_{\delta_0^*}^{\theta^*}$. The set of all restricted roots forms a nonreduced root system in general. As in Section 1.3 of [KS99], we call α of type R_1 if neither $2\alpha_{\text{res}}$ nor $\frac{1}{2}\alpha_{\text{res}}$ is a restricted root, of type R_2 if $2\alpha_{\text{res}}$ is a restricted root, or of type R_3 if $\frac{1}{2}\alpha_{\text{res}}$ is a restricted root. Also following [KS99], we may identify a root $\alpha_1 = ((\alpha^\vee)_{\text{res}})^\vee$ of T_1 in H_1 , or of $T_1/Z_1 \simeq T_{\theta^*}$ in H_1/Z_1 , as $N\alpha$ or $2N\alpha$. If α is of type R_1, R_3 , then $\alpha_1 = N\alpha$, and if α is of type R_2 , then $\alpha_1 = 2N\alpha$. Recall that $N\alpha$ denotes the sum of all distinct roots in the θ^* -orbit of α . Assume α_1 is a root of T_1 in the identity component $(H_1)_{\gamma_0}$ of the centralizer of γ_0 in H_1 . The identification of roots then implies that

$$\alpha_1(\gamma_0) = N\alpha(\delta_0^*) = 1$$

if α is of type R_1 or R_3 and that

$$\alpha_1(\gamma_0) = N\alpha(\delta_0^*)^2 = 1$$

if α is of type R_2 . Write this second case as $R_{2,\pm}$ according as $N\alpha(\delta_0^*) = \pm 1$.

We use $\text{Int}(g) \circ \psi$ to identify roots of T^{δ_0} in $G_{\delta_0}^\theta$ with roots of $(T^{\theta^*})^0$ in $G_{\delta_0^*}^{\theta^*}$. Let α be a root of T in G^* . Then α_{res} is a root of T^{δ_0} in $G_{\delta_0}^\theta$ if and only if $N\alpha(\delta_0^*) = 1$ in the cases α is of type R_1, R_2 , or if and only if $N\alpha(\delta_0^*) = -1$ in the case α is of type R_3 . We conclude the following.

LEMMA 6.5. *Assume that $\alpha_1 = ((\alpha^\vee)_{\text{res}})^\vee$ is a root of T_1 in $(H_1)_{\gamma_0}$, i.e., that $\alpha_1(\gamma_0) = 1$. Then (i) $\alpha_0 = r_\alpha \alpha_{\text{res}}$ is a root of T^{δ_0} in $G_{\delta_0}^\theta$, where $r_\alpha = 1$ if α is of type R_1 or $R_{2,+}$, $r_\alpha = 2$ if α is of type $R_{2,-}$, and $r_\alpha = \frac{1}{2}$ if α is of type R_3 . Also, (ii) if α is of any type except $R_{2,-}$, then $N\alpha(\delta_0^*) = 1$ and α_0 is a root of $(T^{\theta^*})^0$ in $(G^{*\theta^*})^0$. Finally, (iii) if α is of type $R_{2,-}$, then $N\alpha(\delta_0^*) = -1$ and $\alpha_{\text{res}} = \frac{1}{2}\alpha_0$ is a root of $(T^{\theta^*})^0$ in $(G^{*\theta^*})^0$.*

Remark. We will often write $N\alpha(\delta_0)$ for $N\alpha(\delta_0^*)$. Notice we may make a definition of $N\alpha$ that is intrinsic to G by using the automorphism $\text{Int}(\delta_0) \circ \theta$ and the maximal torus $T^\dagger = \text{Cent}(T^{\delta_0}, G)$.

Next, we assume also that γ_0 is semiregular; i.e., $\pm\alpha_1$ are the only roots of T_1 in $(H_1)_{\gamma_0}$. We will say that γ_0 is G -semiregular if $\pm\alpha_0$ are the only roots of T^{δ_0} in $G_{\delta_0}^\theta$; i.e., both $(H_1)_{\gamma_0}$ and $G_{\delta_0}^\theta$ are of Dynkin type A_1 . Explicitly, the extra condition is that if root β of T is not in the \mathbb{Q} -span of the θ^* -orbit of α , then $N\beta(\delta_0) \neq 1$ if β is of type R_1 or R_2 , and $N\beta(\delta_0^*) \neq -1$ if β is of type R_3 . Notice that if β is of type R_2 , then $N\beta(\delta_0^*) = -1$ implies that $2\beta_{\text{res}}$ is a root of $G_{\delta_0}^\theta$, and so we conclude that for β of type R_2 , the extra condition can be rewritten as $\beta_1(\gamma_0) = N\beta(\delta_0^*)^2 \neq 1$, and then that the condition for β of type R_3 is redundant. We may now write the G -semiregularity condition directly in terms of γ_0 as

$\alpha_1(\gamma_0) = 1$ and $\beta_1(\gamma_0) \neq 1$ for all roots β of type R_1 or R_2 not in the \mathbb{Q} -span of the θ^* -orbit of α .

If semiregular $\gamma_0 \in T_1(\mathbb{R})$ is not a norm, we will use this condition as our definition of G -semiregularity (which coincides with the more natural definition using the map $\mathcal{A}_{G/H}$ of [KS99, Th. 3.3.A]).

We return to the setting of Theorem 5.1, where α_1 is imaginary and s_1 is a Cayley transform with respect to α_1 . Because of the stability of the transfer factor $\Delta(\gamma_1, \delta)$ in its first argument γ_1 [KS99, Lemma 5.1.B], the argument of the last paragraph of the proof of Theorem 4.2 shows that there is no harm (for the proof of Theorem 5.1) in assuming α_1 itself is noncompact and that s_1 is a Cayley transform within $(H_1)_{\gamma_0}$. Then also $\gamma_0^{s_1} = \gamma_0$; i.e., γ_0 lies in $T_1 \cap T_1^{s_1}$.

An element γ_1 in $T_1(\mathbb{R})$ is G -regular in the sense of [KS99] if and only if $\beta_1(\gamma_1) \neq 1$ for all roots β of type R_1 or R_2 . Because γ_0 is assumed G -semiregular, the elements

$$\gamma_\nu = \gamma_0 \exp(\nu a_{\alpha_1}(\alpha_1)^\vee)$$

in the γ_0 -component of $T_1(\mathbb{R})$ and the elements

$$\gamma_{s_1, \nu} = \gamma_0 \exp(\nu a_{\alpha'_1}(\alpha'_1)^\vee),$$

where $\alpha'_1 = \alpha_1^{s_1}$, in the γ_0 -component of $T_1^{s_1}(\mathbb{R})$ are easily checked to be G -regular for all real nonzero ν with $|\nu|$ sufficiently small. We gather the following observations with some special cases of Theorem 5.1 in mind (see Lemma 7.2).

LEMMA 6.6. *Suppose γ_0 is a G -semiregular element in a Cartan subgroup $T_1(\mathbb{R})$ of $H_1(\mathbb{R})$ annihilated by a noncompact imaginary root α_1 . Suppose that s_1 is a Cayley transform for α_1 in $(H_1)_{\gamma_0}$. Then (i) if γ_0 is not a G -norm, then the G -regular elements γ_ν and $\gamma_{s_1, \nu}$ are not norms, (ii) if $T_1^{s_1}$ is a norm group for (G, θ) , then T_1 is also a norm group for (G, θ) , (iii) if γ_0 is a G -norm of δ_0 in $G(\mathbb{R})$, then γ_0 is a T_1 -norm of δ_0 , (iv) if γ_0 is a G -norm of δ_0 in $G(\mathbb{R})$, then γ_0 is a $T_1^{s_1}$ -norm of δ_0 if and only if $G_{\delta_0}^\theta$ is split modulo center, and (v) if $T_1^{s_1}$ is not a norm group for (G, θ) , then the group $G_{\delta_0}^\theta$ is compact modulo center for each δ_0 in $G(\mathbb{R})$ with T_1 -norm γ_0 .*

Proof. For (i), assume γ_0 is not a G -norm. We then apply Lemma 6.1 to γ_0 as element of T_1 to conclude that γ_ν is not a norm, and to γ_0 as element of $T_1^{s_1}$ to conclude that $\gamma_{s_1, \nu}$ is not a norm.

For (ii), assume that $T_1^{s_1}$ is a norm group for (G, θ) . By Lemma 6.1, there is a component of $T_1^{s_1}(\mathbb{R})$ consisting of $T_1^{s_1}$ -norms. Choose a G -semiregular element γ_2 of this component annihilated by the real root $\alpha_1^{s_1}$, and suppose it is a $T_1^{s_1}$ -norm of δ_2 . There are θ^* -stable maximal tori T, T' in G^* defined over \mathbb{R} and admissible homomorphisms $T_1 \rightarrow T_{\theta^*}, T_1^{s_1} \rightarrow T'_{\theta^*}$. Since $T_1^{s_1}$ is a

norm group for (G, θ) , there is also an isomorphism $\text{Int}(g_2) \circ \psi : T'^{\delta_2} \rightarrow (T'^{\theta^*})^0$ defined over \mathbb{R} , where $\delta_2^* = g_2 m(\delta_2) \theta^*(g_2)^{-1}$ lies in T' and $N(\delta_2^*)$ is the image of γ_2 under $T_1^{s_1} \rightarrow T'_{\theta^*}$. Recall that $G_{\delta_2^*}^{\theta^*}$ is defined over \mathbb{R} and $\text{Int}(g_2) \circ \psi : G_{\delta_2^*}^{\theta} \rightarrow G_{\delta_2^*}^{\theta^*}$ is an inner twist. The root α'_0 of T'^{δ_2} in $G_{\delta_2^*}^{\theta}$ corresponding to $\alpha'_1 = \alpha_1^{s_1}$ is also real: $\sigma \alpha'_0$ corresponds to $\sigma \alpha'_1 = \alpha'_1$ and so equals α'_0 . Let $t \in G_{\text{sc}}$ define an inverse Cayley transform in $(G_{\delta_2^*}^{\theta})_{\text{sc}}$ for α'_0 . On the other hand, the θ^* -stable pairs (B, T) and (B', T') defining $T_1 \rightarrow T_{\theta^*}$, $T_1^{s_1} \rightarrow T'_{\theta^*}$ are conjugate under $(G_{\text{sc}}^*)^{\theta^*_{\text{sc}}}$ (by Steinberg's structure results; see [KS99, Th. 1.1.A]) and so they determine an element t^* of $(G_{\text{sc}}^*)^{\theta^*_{\text{sc}}}$ such that $\text{Int}(t^*)$ maps T' to T , T'_{θ^*} to T_{θ^*} and completes a commutative diagram with $\text{Int}(s_1)^{-1} : T_1^{s_1} \rightarrow T_1$ and the admissible homomorphisms $T_1 \rightarrow T_{\theta^*}$, $T_1^{s_1} \rightarrow T'_{\theta^*}$. Then t^* is an inverse Cayley transform for the real root $r\alpha'_0$ of $(T'^{\theta^*})^0$ in $(G^{*\theta^*})^0$, where $r = 1$ or $\frac{1}{2}$ since the action of $\sigma(t^*)^{-1}t^*$ on $(T'^{\theta^*})^0$ coincides with the dual transport of $\sigma(s_1)s_1^{-1}$ which acts on $T_1^{s_1}$ as the Weyl reflection for α'_1 ; this dual transport coincides with the Weyl reflection for $r\alpha'_0$. Here we define dual transport using the bijection (1.3.8) of [KS99]. We may arrange the choices so that t^* is standard; i.e., t^* lies in the image of SL_2 in $(G_{\text{sc}}^*)^{\theta^*_{\text{sc}}}$ corresponding to the root $r\alpha'_0$. The action of $\sigma(t^*)^{-1}t^*$ on $(T'^{\theta^*})^0$ coincides with the transport by $\text{Int}(g_2) \circ \psi$ of the action of $\sigma(t)^{-1}t$ on T'^{δ_2} (t is the inverse Cayley transform defined earlier in the present paragraph) since again each acts as the same Weyl reflection. Let T^{δ_2} be the image of T'^{δ_2} under t . This property of t, t^* (via our definition of Cayley transform in Section 3) implies that if $g_3 = t^* \cdot g_2 \cdot \psi(t^{-1})$, then the composition

$$\text{Int}(g_3) \circ \psi : T^{\delta_2} \rightarrow (T^{\theta^*})^0$$

is defined over \mathbb{R} and that

$$g_3 m(\delta_2) \theta^*(g_3)^{-1} = \text{Int}(t^*)(\delta_2^*) = \delta_3^*$$

lies in T . Finally, $N(\delta_3^*)$ is the image of $(\gamma_2)^{s_1^{-1}} \in T_1(\mathbb{R})$, so that $(\gamma_2)^{s_1^{-1}}$ is a T_1 -norm of δ_3 . In particular, T_1 is a norm group for (G, θ) , and (ii) is proved.

For (iii), assume γ_0 is a G -norm. Then because $(H_1)_{\gamma_0}$ is of type A_1 , we see that γ_0 must be either a T_1 -norm or $T_1^{s_1}$ -norm. The argument for (ii) with $\gamma_2 = \gamma_0$ shows that if γ_0 is a $T_1^{s_1}$ -norm of an element δ_0 , then it is also a T_1 -norm of δ_0 .

For (iv), we return to the argument for (ii), except that now T_1 in place of $T_1^{s_1}$ is assumed a norm group for (G, θ) . We replace the element δ_2 by δ_0 and, as usual, write T^{δ_0} for the image of $(T^{\theta^*})^0$ under the embedding into $G_{\delta_0}^{\theta}$. If $G_{\delta_0}^{\theta}$ is split modulo center, which implies that the root α_0 of T^{δ_0} is noncompact imaginary, then we may construct a Cayley transform s^1 in $(G_{\delta_0}^{\theta})_{\text{sc}}$ and argue along the same lines as (ii) to write γ_0 as a $T_1^{s_1}$ -norm of δ_0 . For the converse,

assume γ_0 is also a $T_1^{s_1}$ -norm of δ_0 . Then the argument for (ii) shows that $G_{\delta_0}^\theta$ contains a torus T^{δ_0} which has a real root and so is split modulo center.

Lemma 6.1 shows that (v) is a consequence of (iv) and the lemma follows. \square

Proof of Lemma 6.3. Suppose semisimple γ_0 is a T_1' -norm of δ_0 , where T_1' is arbitrary. By definition, T_1' lies in $(H_1)_{\gamma_0}$. If T_1' is not fundamental in $(H_1)_{\gamma_0}$, then it has a real root. Now we argue similarly as for (ii) in Lemma 6.6, with γ_0 in place of γ_2 and δ_0 in place of δ_2 , to display γ_0 as a T_1 -norm of δ_0 , with T_1 of split rank one less than that of T_1' . Repeating this argument until real roots are exhausted, we conclude that if γ_0 is a G -norm of δ_0 , then γ_0 is a T_1 -norm of δ_0 , where $T_1 = T_{\text{fund}}$ is fundamental in $(H_1)_{\gamma_0}$. Recall that a stable conjugate of γ_0 in $H_1(\mathbb{R})$ may be written as $w\gamma_0w^{-1}$, where the restriction of $\text{Int}(w)$ to T_{fund} is defined over \mathbb{R} [She83, Lemma 2.5.1]. Then (i) follows.

To prove (ii), let semisimple γ_0 be a G -norm of δ_0, δ'_0 . Then by the last paragraph we may use an admissible homomorphism $T_1 \rightarrow T_{\theta^*}$, with T_1 fundamental in $(H_1)_{\gamma_0}$, to attach $g, g' \in G_{\text{sc}}^*$ and $\delta_3^*, (\delta'_3)^* \in T$ to δ_0, δ'_0 respectively. Following the proof for (i) \Rightarrow (ii) in Lemma 6.1, we use the elements g, g' to define strongly θ -regular δ_3, δ'_3 and corresponding elements $\delta_3^*, (\delta'_3)^*$ in $\delta_0^* \cdot (T^{\theta^*})^0(\mathbb{R})$ and $(\delta'_0)^* \cdot (T^{\theta^*})^0(\mathbb{R})$ respectively, such that $N(\delta_3^*) = N((\delta'_3)^*)$. That construction allows us to assume $(\delta'_3)^* = \delta_3^* t \theta^*(t)^{-1}$, where $t \in T$ satisfies $(\delta'_0)^* = \delta_0^* t \theta^*(t)^{-1}$. Set $x = \psi^{-1}(g'^{-1}tg)$. Then $x\delta_3\theta(x)^{-1} = \delta'_3$ and $x\delta_0\theta(x)^{-1} = \delta'_0$. From the first of these two equations (the strongly regular case) we conclude that $\sigma(x)^{-1}x$ lies in the product of $G_{\delta_3}^\theta = T^{\delta_0}$ with θ -invariants in the center of G , and so δ'_0 is stably θ -conjugate to δ_0 . \square

We turn now to the proof of Lemma 6.4. Our first remark is that the elements $v(\sigma) = gu(\sigma)\sigma(g)^{-1}$ and $\delta_0^* = gm(\delta_0)\theta^*(g)^{-1}$ from the statement of the lemma are unchanged when the inner twist $\psi : G \rightarrow G^*$ is replaced by $\text{Int}(x) \circ \psi$, where $x \in G_{\text{sc}}^*$, provided we replace $u(\sigma)$ by $xu(\sigma)\sigma(x)^{-1}$ and g_θ by $\theta^*(x)g_\theta x^{-1}$. Recall that $u(\sigma), g_\theta$ were discussed in the first paragraph of the present section. Notice also that the change in ψ does not affect our assumption there about $u(\sigma), g_\theta$. We are thus free to choose ψ as we wish within its inner class. Our choice will use fundamental splittings, as in [She11] but without the cuspidality assumption. The definitions are as follows.

Let T_G be a fundamental maximal torus over \mathbb{R} in G and B_G be a Borel subgroup of G containing T_G . Then we call the pair (B_G, T_G) fundamental if the set of B_G -simple roots of T_G in G is preserved by the action of $-\sigma_T$ on $X^*(T)$. Such pairs exist; see [Kot82, §10.4]. We will review this below as we use it. Consider a splitting $\text{spl}_G = (B_G, T_G, \{X_\alpha\})$ for G . Here X_α is a root vector for the B_G -simple root α . Denote by $X_{-\alpha}$ the root vector for $-\alpha$ completing X_α and the coroot H_α to a simple triple. There are two possibilities: α is

complex and $|\{\pm\alpha, \pm\sigma_T\alpha\}| = 4$ or α is imaginary and $|\{\pm\alpha, \pm\sigma_T\alpha\}| = 2$. We call spl_G fundamental if the pair (B_G, T_G) is fundamental and $\sigma X_\alpha = X_{\sigma_T\alpha}$ for all B_G -simple roots that are complex or noncompact imaginary, $\sigma X_\alpha = -X_{\sigma_T\alpha}$ for all B_G -simple roots that are compact imaginary. A fundamental pair may be extended to a fundamental splitting; see [She11, §3] regarding imaginary roots. Suppose that η is an automorphism of G that preserves the fundamental splitting spl_G . If the restriction of η to T_G is defined over \mathbb{R} , then an argument with root vectors shows that η is defined over \mathbb{R} as automorphism of G .

The automorphism θ^* of G^* preserves a (fixed) \mathbb{R} -splitting $(B^*, T^*, \{X_{\alpha^*}\})$. Here T^* is a maximally split maximal torus defined over \mathbb{R} and B^* is also defined over \mathbb{R} . We may construct a θ^* -stable fundamental pair (B, T) for G^* as follows. Consider the identity component G^1 of the group of fixed points of θ^* in G^* . Then G^1 has an \mathbb{R} -splitting that extends the pair $(G^1 \cap B^*, G^1 \cap T^*)$. Following Sections 10.3, 10.4 of [Kot86], we apply a rationality theorem of Steinberg to find a fundamental pair (B^1, T^1) for G^1 : choose $h \in (G_{\text{sc}}^*)^{\theta_{\text{sc}}^*}$ such that $h\sigma(h)^{-1}$ preserves $G^1 \cap T^*$ and acts on $G^1 \cap T^*$ as the longest element of the Weyl group of $G^1 \cap T^*$ in G^1 , and then set $B^1 = h^{-1}B^*h, T^1 = h^{-1}T^*h$. Let (B, T) be the corresponding θ^* -stable pair for G^* . Then T is fundamental since a real root would provide a real root for the fundamental torus T^1 , and further the pair (B, T) is fundamental, again by Steinberg’s structure theorem. We extend (B, T) to a fundamental splitting spl . Then θ^* preserves spl up to an inner automorphism by an element of T_{sc} ; this inner automorphism is defined over \mathbb{R} .

Returning to the inner twist $\psi : G \rightarrow G^*$, we adjust ψ within its inner class so that the restriction of ψ^{-1} to T is defined over \mathbb{R} . Set $B_G = \psi^{-1}(B), T_G = \psi^{-1}(T)$. Then (B_G, T_G) is a fundamental pair. We may further adjust ψ by an inner automorphism by an element of T_{sc} so that ψ^{-1} transports the fundamental splitting spl of G^* to a fundamental splitting spl_G of G extending (B_G, T_G) . With these adjustments to ψ we now conclude that $\theta_G = \psi^{-1} \circ \theta^* \circ \psi$ is defined over \mathbb{R} . Then $\theta = \text{Int}(h_\theta) \circ \theta_G$, where $h_\theta = \psi_{\text{sc}}^{-1}(g_\theta^{-1})$. Both $\text{Int}(h_\theta)$ and $\text{Int}(g_\theta)$ are defined over \mathbb{R} , and we may take $u(\sigma)$ to be fixed by θ_{sc}^* since $(T_{\text{sc}})^{\theta_{\text{sc}}^*} \rightarrow (T_{\text{ad}})^{\theta_{\text{ad}}^*}$ is surjective. Then the cocycle z_σ of Lemma 3.1.A of [KS99] is simply $\psi_{\text{sc}}(h_\theta^{-1}\sigma(h_\theta))$. Returning to the assumption of the first paragraph of this section, we adjust the choice of $g_\theta, u(\sigma)$ by central elements in G_{sc}^* to arrange that $z_\sigma = 1$ [KS99, p. 26].

Remark. Since θ_G has finite order, it follows that θ may be written as the product of an inner automorphism and an automorphism of finite order, where each automorphism is defined over \mathbb{R} . This result was pointed out by a referee who also supplied another proof.

We will also make use of connectivity properties of real points of fundamental tori. We continue with the same setting. From Sections 10.3, 10.4

of [Kot86] we see that $T_{\text{sc}}(\mathbb{R})$ is connected: because (B, T) is a fundamental pair, $X_*(T_{\text{sc}})$ has a base preserved by $-\sigma_T$, namely the coroots of the B -simple roots of T , and so each σ_T -invariant element of $X_*(T_{\text{sc}})$ lies in $(1 + \sigma_T)X_*(T_{\text{sc}})$, which implies that $T_{\text{sc}}(\mathbb{R})$ has one component. The same argument for $X_*(T_{\text{ad}})$, using fundamental coweights in place of coroots, shows that $T_{\text{ad}}(\mathbb{R})$ is connected. Finally, recall that (B, T) is θ^* -stable. The image in $X_*(T_{\text{ad}})/(1 - \theta_{\text{ad}}^*)X_*(T_{\text{ad}}) = X_*((T_{\text{ad}})_{\theta_{\text{ad}}^*})$ of the chosen base for $X_*(T_{\text{ad}})$ is a base for $X_*((T_{\text{ad}})_{\theta_{\text{ad}}^*})$ since it has the correct cardinality, by Steinberg's structure theorem. Thus $(T_{\text{ad}})_{\theta_{\text{ad}}^*}(\mathbb{R})$ is connected.

Proof of Lemma 6.4. First observe that (ii) follows once we have proved (i). The equation $\sigma(\delta^*)\delta^{*-1} = (\theta^* - 1)v(\sigma)$ from [KS99, Lemma 4.4.A] (see the proof of Lemma 6.1) implies that the image $v(\sigma)_{\text{ad}}$ of $v(\sigma)$ in T_{ad} is an element, in fact a cocycle, in $(T_{\text{ad}})_{\theta_{\text{ad}}^*}$. Since $(T_{\text{sc}})_{\theta_{\text{sc}}^*}$ and $(T_{\text{ad}})_{\theta_{\text{ad}}^*}$ are both connected, the natural projection $T_{\text{sc}} \rightarrow T_{\text{ad}}$ projects $(T_{\text{sc}})_{\theta_{\text{sc}}^*}$ onto $(T_{\text{ad}})_{\theta_{\text{ad}}^*}$, and (ii) follows.

For the proof of (i), it is sufficient to consider the case that the endoscopic group is basic, i.e., attached to the trivial endoscopic data $(G_1, G_1^\vee \rtimes W_{\mathbb{R}}, 1)$ for the pair (G, θ) , where G_1^\vee denotes the identity component of the fixed points of θ^\vee in G^\vee : if H_1 is any endoscopic group and $H = H_1/Z_1$, then an admissible embedding $T_H \rightarrow T_{\theta^*}$ determines an admissible embedding $T_{G_1} \rightarrow T_{\theta^*}$ (see [KS99, §3.3]), with same data g, δ_0^* attached to the same (strongly G -regular) element in $T_{\theta^*}(\mathbb{R})$.

Assume then that H_1 is basic. There exists an admissible embedding $T_H \rightarrow T_{\theta^*}$, where (B, T) is a θ^* -stable fundamental pair, and thus there exist strongly G -regular T_1 -norms; here $T_H = T_1/Z_1$. Suppose strongly G -regular $\gamma_1 \in T_1(\mathbb{R})$ is a norm of $\delta \in G(\mathbb{R})$. Attach $g \in G_{\text{sc}}^*$ and $\delta^* \in T$ as usual. Then $N\delta^* \in T_{\theta^*}(\mathbb{R})$. Passing to the adjoint form G_{ad}^* of G^* , we have that δ_{ad}^* has image in $(T_{\text{ad}})_{\theta_{\text{ad}}^*}(\mathbb{R})$ under N_{ad} . Since $N_{\text{ad}} : T_{\text{ad}}(\mathbb{R}) \rightarrow (T_{\text{ad}})_{\theta_{\text{ad}}^*}(\mathbb{R})$ is surjective (domain and target are connected), we may then find $\delta^{**} \in T$ such that $\sigma(\delta^{**})\delta^{**,-1}$ is central in G^* and $\delta^{**} \equiv \delta^*(1 - \theta^*)T$. Multiplying δ^{**} by a suitable central element allows us to replace $(1 - \theta^*)T$ by the image of $(1 - \theta_{\text{sc}}^*)T_{\text{sc}}$. Then multiplying g by a suitable element of T_{sc} , we obtain a replacement for the pair g, δ^* with the desired property (i). Lemma 6.1 shows that the assumption of strongly G -regularity is unnecessary.

We remove the assumption that T is fundamental using induction on the split rank of T^{θ^*} . By Lemma 6.1 we may assume that γ_0, δ_0 are the elements γ_2, δ_2 of the proof of (ii) in Lemma 6.6, with attached g_2, δ_2^* . We construct g_3, δ_3^* and adjust them using the induction hypothesis, then replace g_2, δ_2^* accordingly. Then $\delta_2^* = \text{Int}(t^*)^{-1}(\delta_3^*)$. Recall that $\sigma(\delta_2^*)\delta_2^{*-1}$ is the image in G^* of the element $(\theta_{\text{sc}}^* - 1)v_2(\sigma)$ in G_{sc}^* and $\sigma(\delta_3^*)\delta_3^{*-1}$ is the image of $(\theta_{\text{sc}}^* - 1)v_3(\sigma)$. We claim that we can adjust g_2 again to arrange that $(\theta_{\text{sc}}^* - 1)v_2(\sigma)$ and $(\theta_{\text{sc}}^* - 1)v_3(\sigma)$ are the same central element of G_{sc}^* . This will both complete our

inductive proof of (i) and provide a modification of the hypercocycle property that is useful for the proof of the main lemma of Section 9.

To justify the claim, we return to Lemma 4.4.A of [KS99] and argue with G_{sc} instead of G . We replace δ_0 by $\delta_{sc} \in G_{sc}$ with image δ_0 up to a central element. Then $\sigma(\delta_{sc}) = z_0\delta_{sc}$, where z_0 is central in G_{sc} . Passing to a suitable strongly G_{sc} -regular element in each case, we find that

$$(\theta_{sc}^* - 1)v_2(\sigma) = \psi_{sc}(z_0)\sigma(\delta_{2,sc}^*)(\delta_{2,sc}^*)^{-1}$$

and

$$(\theta_{sc}^* - 1)v_3(\sigma) = \psi_{sc}(z_0)\sigma(\delta_{3,sc}^*)(\delta_{3,sc}^*)^{-1},$$

where $\delta_{2,sc}^* = g_2m_{sc}(\delta_{sc})\theta_{sc}^*(g_2)^{-1}$ has image δ_2^* up to a central element in G^* and $\delta_{2,sc}^* = g_3m_{sc}(\delta_{sc})\theta_{sc}^*(g_3)^{-1}$ has image δ_3^* up to the same central element in G^* . Moreover, $\delta_{2,sc}^* = \text{Int}(t^*)^{-1}(\delta_{3,sc}^*)$ and $\sigma(\delta_{3,sc}^*)(\delta_{3,sc}^*)^{-1}$ is central. To prove the claim, we observe that $\text{Int}(t^*)^{-1}(\sigma(\delta_{3,sc}^*)(\delta_{3,sc}^*)^{-1})$ coincides with $\sigma(\delta_{2,sc}^*)(\delta_{2,sc}^*)^{-1}$ up to an element of $(1 - \theta_{sc}^*)T'_{sc}$, so that we may adjust g_2 as desired. \square

Definition. Choose g, δ_0^* satisfying (i), and thus (ii), of Lemma 6.4. Then we will call $(T_1 \rightarrow T_{\theta^*}, g)$ *toral data at γ_0* .

7. Application to Theorem 5.1

We return to the normalized sum of twisted integrals

$$\Phi_1(\gamma_1) = \left| \det(\text{Ad}(\gamma_1) - I)_{\mathfrak{h}_1/\mathfrak{t}_1} \right|^{1/2} \sum_{\delta, \theta\text{-conj}} \Delta(\gamma_1, \delta) O^{\theta, \varpi}(\delta, f)$$

for γ_1 strongly G -regular. This was rewritten in Section 5 as

$$\Delta_I(\gamma_1)\Delta_{II}(\gamma_1) \sum_{\delta, \theta\text{-conj}} \Delta_{III}(\gamma_1, \delta; \bar{\gamma}, \bar{\delta}) \Phi^{\theta, \varpi}(\delta, f),$$

where the twisted integrals themselves are now normalized, and the terms Δ_I, Δ_{II} , and Δ_{III} come from the twisted transfer factor Δ .

Fix a maximal torus T_1 over \mathbb{R} in H_1 , a G -semiregular element γ_0 in the Cartan subgroup $T_1(\mathbb{R})$ annihilated by an imaginary root α_1 , and a Cayley transform s_1 for α_1 .

Our next step is to write $\Phi_1(\gamma_1)$ for strongly G -regular γ_1 in the γ_0 -component of $T_1(\mathbb{R})$ in a way that will be useful both for extending Φ_1 to all G -regular elements and for jump analysis around γ_0 .

If γ_0 is not a T_1 -norm, then $\Phi_1(\gamma_1) = 0$ for all strongly G -regular γ_1 in the γ_0 -component of $T_1(\mathbb{R})$, and so we define $\Phi_1(\gamma_1) = 0$ also for the remaining G -regular elements in the component. Assume then that γ_0 is a T_1 -norm of the θ -semisimple element δ_0 of $G(\mathbb{R})$. Let $(T_1 \rightarrow T_{\theta^*}, g)$ be toral data at γ_0 . As in Lemma 6.2, we have the element $\delta = \delta_0(Y) = (\exp Y)\delta_0$ with given norm

$\gamma_1 = \gamma_1(Y_1) = \gamma_0 \exp(Y_H + Y_{\delta_1})$ in the γ_0 -component of $T_1(\mathbb{R})$. Suppose γ_1 is strongly G -regular, so that δ is strongly θ -regular. We fix representatives δ' for the θ -conjugacy classes in the stable θ -conjugacy class of δ and then define $\text{inv}(\delta, \delta')$ and κ_δ as in the preamble to Theorem 5.1.D of [KS99], which also describes how these two objects are paired. (For more on the definitions, see the proof of Lemma 9.6.) Then by (1) of that theorem, $\Phi_1(\gamma_1)$ may be rewritten as

$$\Delta_I(\gamma_1)\Delta_{II}(\gamma_1)\Delta_{III}(\gamma_1, \delta; \bar{\gamma}, \bar{\delta}) \sum_{\delta'} \langle \text{inv}(\delta, \delta'), \kappa_\delta \rangle \Phi^{\theta, \varpi}(\delta', f).$$

Suppose, slightly more generally, that δ' is stably θ -conjugate to strongly θ -regular $\delta = \varepsilon\delta_0$, where $\varepsilon \in T^{\delta_0}(\mathbb{R})$. We may write

$$\delta' = \delta(w) = w^{-1}\delta\theta(w) = w^{-1}\varepsilon w.w^{-1}\delta_0\theta(w),$$

where $w \in G(\mathbb{C})$ (we stress \mathbb{C} in notation just for this paragraph) and $\sigma(w)w^{-1}$ lies in $\text{Cent}_\theta(\delta, G(\mathbb{C}))$. As earlier, let $T^\dagger = \text{Cent}(T^{\delta_0}, G)$. Then strong θ -regularity implies that $\text{Cent}_\theta(\delta, G)$ coincides with the group T_{δ_0} of fixed points of $\text{Int}(\delta_0) \circ \theta$ in T^\dagger . Set

$$\mathfrak{A}_\theta(T^{\delta_0}) = \{w \in G(\mathbb{C}) : \sigma(w)w^{-1} \in T_{\delta_0}(\mathbb{C})\}.$$

Then, via the map $w \rightarrow \delta(w)$,

$$\mathfrak{D}_\theta(T^{\delta_0}) = T_{\delta_0}(\mathbb{C}) \backslash \mathfrak{A}_\theta(T^{\delta_0}) / G(\mathbb{R})$$

parametrizes the θ -conjugacy classes in the stable θ -conjugacy class of δ .

If now we assume only that $\delta = \delta_0(Y)$ is θ -regular, then by definition (see the remark after Lemma 6.3),

$$\{\delta(w) : w \in \mathfrak{A}_\theta(T^{\delta_0})\}$$

is the stable θ -conjugacy class of δ . We will define $\Phi_1(\gamma_1)$ to be

$$\left| \det(\text{Ad}(\gamma_1) - I)_{\mathfrak{h}_1/\mathfrak{t}_1} \right|^{1/2} \sum_w \Delta(\gamma_1, \delta(w)) O^{\theta, \varpi}(\delta(w), f),$$

where \sum_w indicates summation over a set of representatives w for $\mathfrak{D}_\theta(T^{\delta_0})$ and

$$\Delta(\gamma_1, \delta(w)) = \lim_{\gamma_1^\dagger \rightarrow \gamma_1} \Delta(\gamma_1^\dagger, \delta^\dagger(w)).$$

In this limit, the variable $\gamma_1^\dagger = \gamma_1 \exp Y^\dagger$ is a strongly G -regular element in the γ_0 -component of $T_1(\mathbb{R})$. This element γ_1^\dagger is a norm of each (strongly θ -regular) element $\delta^\dagger(w)$, where $\delta^\dagger = (\exp Y^{\dagger\dagger}) \delta$. Here $Y^{\dagger\dagger} \leftrightarrow Y_H$, where $Y^\dagger = Y_H + Y_{\delta_1}$ as in Lemma 6.2. To see that the limit exists, we just have to recall how the term $\Delta(\gamma_1^\dagger, \delta^\dagger(w))$ depends on Y^\dagger . First,

$$\Delta(\gamma_1^\dagger, \delta^\dagger(w)) = \langle \text{inv}(\delta^\dagger, \delta^\dagger(w)), \kappa_{\delta^\dagger} \rangle \Delta(\gamma_1^\dagger, \delta^\dagger).$$

The first term is a constant sign and so can be ignored. The term $\Delta(\gamma_1^\dagger, \delta^\dagger)$ is a product

$$\Delta_I(\gamma_1 \exp Y^\dagger) \Delta_{II}(\gamma_1 \exp Y^\dagger) \Delta_{III}(\gamma_1 \exp Y^\dagger, (\exp Y^{\dagger\dagger})\delta; \bar{\gamma}, \bar{\delta}) \Delta_{IV}(\gamma_1 \exp Y^\dagger).$$

The new first term is a constant sign. The term $\Delta_{II}\Delta_{IV}$ is a quotient of generalized Weyl denominators for G and H_1 (see [KS99, §4.3]). It is well defined, smooth, and nonzero on the subset of all G -regular elements in $T_1(\mathbb{R})$. It remains then to examine $\Delta_{III}(\gamma_1 \exp Y^\dagger, (\exp Y^{\dagger\dagger})\delta; \bar{\gamma}, \bar{\delta})$. A check of definitions shows that it is the product of a constant and a character on $T^{\delta_0}(\mathbb{R})$ evaluated at $\exp Y^{\dagger\dagger}$; see the beginning of the proof of Lemma 9.3 where we introduce more detailed notation for an analysis of Δ_{III} . We conclude then that $\lim_{\gamma_1^\dagger \rightarrow \gamma_1} \Delta(\gamma_1^\dagger, \delta^\dagger(w))$ is well defined, which completes our (smooth) extension of Φ_1 to the full G -regular set in $T_1(\mathbb{R})$.

Let $w \in \mathfrak{A}_\theta(T^{\delta_0})$, and write $w^{-1}\delta_0\theta(w)$ as $\delta_0(w)$. Then $\text{Int}(w^{-1}) : G_{\delta_0}^\theta \rightarrow G_{\delta_0(w)}^\theta$ is an inner twist and $\delta_0(w)$ is stably conjugate to δ_0 . The inner type of the group $G_{\delta_0(w)}^\theta$ of Dynkin type A_1 , either split modulo center or compact modulo center, depends only on the double coset of w in $\mathfrak{D}_\theta(T^{\delta_0})$. We may ignore those w for which $G_{\delta_0(w)}^\theta$ is compact modulo center, as they contribute nothing to the final limit formula (see Section 8). We have the following generalization of Lemma 4.2 of [She79a]. Again α_0 denotes the multiple of α_{res} that is a root of T^{δ_0} in $G_{\delta_0}^\theta$.

LEMMA 7.1. *If both $G_{\delta_0}^\theta$ and $G_{\delta_0(w)}^\theta$ are split modulo center (i.e., both α_0 and $w\alpha_0$ are noncompact imaginary roots), then there exists $g \in G(\mathbb{R})$ such that $\text{Int}(g)$ maps $G_{\delta_0}^\theta$ to $G_{\delta_0(w)}^\theta$ and T^{δ_0} to $T^{\delta_0(w)}$, and $w^{-1}\alpha_0 = \pm g\alpha_0$.*

Proof. We follow the proof of Lemma 4.2 in [She79a]. First, a simple argument with root vectors shows that we can arrange that $\text{Int}(w^{-1}) : G_{\delta_0}^\theta \rightarrow G_{\delta_0(w)}^\theta$ is defined over \mathbb{R} ; see the first paragraph of the cited proof. Let s be the standard Cayley transform in $(G_{\delta_0}^\theta)_{\text{sc}} = \text{SL}(2)$ relative to the root α_0 of T^{δ_0} in $G_{\delta_0}^\theta$, and set $T'^{\delta_0} = (T^{\delta_0})^s$. We may argue in the untwisted setting with $w \in \mathfrak{A}(T'_G)$, where T'_G is the maximal torus $\text{Cent}(T'^{\delta_0}, G)$ in G , to choose g_1 in $G(\mathbb{R})$ so that $\text{Int}(g_1)$ maps T'^{δ_0} to $w^{-1}T'^{\delta_0}w$ and acts on the maximal split torus in T'^{δ_0} as $\text{Int}(w^{-1})$. Then $\text{Int}(g_1^{-1}w^{-1})$ normalizes the derived group of $G_{\delta_0}^\theta$ (by another argument with root vectors) as well as T'^{δ_0} . Then $\text{Int}(g_1^{-1}w^{-1})$ normalizes $G_{\delta_0}^\theta$ itself. Multiplying g_1 by a suitable element of $G_{\delta_0}^\theta(\mathbb{R})$, we obtain g in $G(\mathbb{R})$ such that $\text{Int}(g^{-1}w^{-1})$ normalizes both $G_{\delta_0}^\theta$ and T^{δ_0} . Then $w^{-1}\alpha_0$ coincides with $\pm g\alpha_0$. \square

We will need a twisted version of Proposition 4.6 of [She79a] in order to match the elements of $\mathfrak{D}_\theta(T^{\delta_0})$ contributing to jumps with the elements of $\mathfrak{D}_\theta(T'^{\delta_0})$, where $T'^{\delta_0} = (T^{\delta_0})^s$. Assume α_0 is noncompact and that s is

standard in $(G_{\delta_0}^\theta)_{\text{sc}} = \text{SL}(2)$. Let \mathbf{w} be an element of $\mathfrak{D}_\theta(T^{\delta_0})$ such that $G_{\delta_0(w)}^\theta$ is split modulo center for some, and hence every, w representing \mathbf{w} . Then, following the last lemma, we may choose w so that w normalizes $G_{\delta_0}^\theta$ and T^{δ_0} , and $w^{-1}\alpha_0 = \pm\alpha_0$. Now consider those \mathbf{w} with representative w such that $w^{-1}\alpha_0 = \alpha_0$. Suppose w_0 is an element of $\text{Cent}_\theta(\delta_0, G)$ normalizing T^{δ_0} for which the action of $\text{Int}(w_0)$ on T^{δ_0} realizes the Weyl reflection relative to α_0 . Then w and ww_0 represent the same element of $\mathfrak{D}_\theta(T^{\delta_0})$ if and only if we may choose w_0 in $G(\mathbb{R})$, i.e., in $\text{Cent}_\theta(\delta_0, G)(\mathbb{R})$. If that is so, then we say that the Weyl reflection relative to α_0 is realized in $G(\mathbb{R})$, keeping in mind that this notion depends on the choice of δ_0 . The elements \mathbf{w} of $\mathfrak{D}_\theta(T^{\delta_0})$ with a representative w such that $w^{-1}\alpha_0 = \alpha_0$ are then exactly those \mathbf{w} such that $G_{\delta_0(w)}^\theta$ is split modulo center for each representative w . We denote this subset of $\mathfrak{D}_\theta(T^{\delta_0})$ by $\mathfrak{D}_\theta(\alpha_0)$. On the other hand, if the Weyl reflection relative to α_0 is not realized in $G(\mathbb{R})$, then for each element \mathbf{w} of $\mathfrak{D}_\theta(T^{\delta_0})$ with representative w such that $w^{-1}\alpha_0 = \alpha_0$, there is an element \mathbf{w}_- , distinct from \mathbf{w} , with representative $w_- = ww_0$ such that $w_-^{-1}\alpha_0 = -\alpha_0$. In this case, $\mathfrak{D}_\theta(\alpha_0)$ will denote the set of pairs $\{\mathbf{w}, \mathbf{w}_-\}$.

Consider \mathbf{w}' in $\mathfrak{D}_\theta(T'^{\delta_0})$. Again following on from the proof of Lemm 7.1, since α_0^s is a real root, we may find a representative w' for \mathbf{w}' such that w' normalizes both T'^{δ_0} and $G_{\delta_0}^\theta$, and $w'^{-1}\alpha_0^s = \alpha_0^s$. We can then further arrange that w' centralizes $(G_{\delta_0}^\theta)_{\text{der}}$. Thus $w = s^{-1}w's = w'$ lies in $\mathfrak{A}_\theta(T^{\delta_0})$ and $w^{-1}\alpha_0 = \alpha_0$. Let \mathbf{w} be the class of w in $\mathfrak{D}_\theta(T^{\delta_0})$. Then another argument with root vectors shows that $\mathbf{w}' \rightarrow \mathbf{w}$ is a well-defined bijective map of $\mathfrak{D}_\theta(T'^{\delta_0})$ to those $\mathbf{w} \in \mathfrak{D}_\theta(T^{\delta_0})$ with representative w such that $w^{-1}\alpha_0 = \alpha_0$. This provides us with a bijection of $\mathfrak{D}_\theta(T'^{\delta_0})$ with $\mathfrak{D}_\theta(\alpha_0)$.

Before continuing with the analysis, we finish the proof of Theorem 5.1 for some special cases.

LEMMA 7.2. *All limits in Theorem 5.1 are zero if*

- (i) γ_0 is not a norm, or if
- (ii) γ_0 is a norm but $T_1^{s_1}$ is not a norm group for G , or if
- (iii) γ_0 is a norm, $T_1^{s_1}$ is a norm group for G , but γ_0 is not a $T_1^{s_1}$ -norm.

Proof. For (i), we have only to apply (i) of Lemma 6.6: $\Phi_1(\gamma_\nu) = 0$ and $\Phi_1(\gamma_{s_1, \nu}) = 0$ for $|\nu|$ sufficiently small and nonzero. On the other hand, for (ii) and (iii) we have, in general, only that $\Phi_1(\gamma_{s_1, \nu}) = 0$ for $|\nu|$ sufficiently small. Thus it remains to show $\lim_{\nu \rightarrow 0} \Phi_1(\gamma_\nu) = 0$. By Lemma 6.6, each group $G_{\delta_0(w)}^\theta(\mathbb{R})$ is compact modulo center and so each unnormalized integral $O^{\theta, \varpi}(\delta(w), f)$ appearing in $\Phi_1(\gamma_\nu)$ is bounded as $\nu \rightarrow 0$ (see Section 8). Thus the limit of $\Phi_1(\gamma_\nu)$ exists and is zero. □

For our analysis of the limits in Theorem 5.1, we may now assume that both T_1 and $T_1^{s_1}$ are norm groups and that γ_0 is both a T_1 -norm and a $T_1^{s_1}$ -norm. Recall that we assume also that the root α_1 of T_1 annihilating γ_0 is noncompact and that s_1 is a Cayley transform in $((H_1)_{\gamma_0})_{sc} = \text{SL}(2)$.

We return to the setting established at the end of the proof of Lemma 6.4. We may suppose γ_0 is both a $T_1^{s_1}$ -norm and a T_1 -norm of an element δ_0 of $G(\mathbb{R})$ for which $G_{\delta_0}^\theta$ is split modulo center, as there. We have admissible homomorphisms $T_1 \rightarrow T_{\theta^*}, T_1^{s_1} \rightarrow T'_{\theta^*}$ and inverse Cayley transform t^* in $(G_{sc}^*)^{\theta^*}$ which maps $(T'^{\theta^*})^0$ to $(T^{\theta^*})^0, T'$ to T, T'_{θ^*} to T_{θ^*} and completes a commutative diagram with $\text{Int}(s_1)^{-1} : T_1^{s_1} \rightarrow T_1$ and $T_1 \rightarrow T_{\theta^*}, T_1^{s_1} \rightarrow T'_{\theta^*}$. Also $t \in G_{sc}$ defines an inverse Cayley transform in $(G_{\delta_0}^\theta)_{sc}$ for α'_0 , where α'_0 is the root of T'^{δ_0} in $G_{\delta_0}^\theta$ corresponding to $\alpha'_1 = \alpha_1^{s_1}$. Then, with g_2, g_3 as at the end of the proof of Lemma 6.6, we choose $g = g_3$ and $g' = g_2$. There is another requirement that will be useful since it makes the limits we consider for Δ_I, Δ_{III} in Lemmas 9.3 and 9.5 both equal to one. Namely we insist that if a complex root of $(T'^{\theta^*})^0$ is positive in the ordering determined by the toral data and our choice of \mathbb{R} -splitting for $(G^{\theta^*})^0$, then its complex conjugate is also positive. That this is possible follows from a familiar argument using a suitable lexicographic ordering of roots for the \mathbb{R} -splitting. (Start with toral data for a maximally split torus in H_1 , identify inverse Cayley transforms needed to reach $(T'^{\theta^*})^0$ through H_1 , adjust the \mathbb{R} -splitting accordingly via Cayley transforms from the torus attached to the maximally split torus in H_1 , and prescribe toral data for $T_1^{s_1}$ using the inverse transforms.)

We call the data of the last paragraph *toral descent data at γ_0* .

8. Jump analysis for twisted orbital integrals

The limit formulas for the individual twisted orbital integrals guide our analysis of the transfer factors, and so we write them next. Formulas of this type are well known. We need only to extend the setting and to write the results in a way that fits well with our transforms.

We continue with the toral descent data at γ_0 from the end of the last section: γ_0 is both a $T_1^{s_1}$ -norm and a T_1 -norm of an element δ_0 of $G(\mathbb{R})$ for which $G_{\delta_0}^\theta$ is split modulo center and of Dynkin type A_1 . Now s will be the Cayley transform t^{-1} in $(G_{\delta_0}^\theta)_{sc}$. Fix an element of $\mathfrak{D}_\theta(T'^{\delta_0})$. Our choice in the last section of representative w' , along with w and w_0 , ensures that $G_{\delta_0(w)}^\theta = G_{\delta_0(ww_0)}^\theta = G_{\delta_0}^\theta$ and that the points $\delta_0(w), \delta_0(ww_0), \delta_0(w')$ all coincide. We will make a descent from $G(\mathbb{R})$ to $\text{Cent}_\theta(\delta_0, G)(\mathbb{R})$, then into $G_{\delta_0}^\theta(\mathbb{R})$, around $\delta_0(w)$. This generalizes the descent used in Section 4 of [She79a] for the untwisted case. Notice that because the twisting character ϖ is trivial on both Cartan subgroups $T^{\delta_0}(\mathbb{R}), T'^{\delta_0}(\mathbb{R})$ in $G_{\delta_0}^\theta(\mathbb{R})$ [KS99, Lemma 4.4.C] (more generally,

ϖ is trivial on both $T_{\delta_0}(\mathbb{R})$, $T'_{\delta_0}(\mathbb{R})$ by [KS99, Th. 5.1.D]), we have that ϖ is trivial on $G_{\delta_0}^\theta(\mathbb{R})$.

We may write α_0 as $r_\alpha \alpha_{\text{res}}$, where $\alpha_1 = ((\alpha^\vee)_{\text{res}})^\vee$ and the coefficient r_α is described in Lemma 6.5. As in [KS99] (see Section 9 also), we use the same a -data and χ -data for all multiples of α_{res} and write χ, a, χ' , and a' for data $\chi_{\alpha_{\text{res}}}, a_{\alpha_{\text{res}}}, \chi_{\alpha_{\text{res}}^\vee}$, and $a_{\alpha_{\text{res}}^\vee}$.

Assume δ is a θ -regular element in $T^{\delta_0}(\mathbb{R})^0 \delta_0$. For α_0 of type R_1 , set

$$\Delta_{\alpha_0}(\delta) = \chi \left(\frac{N\alpha(\delta) - 1}{a} \right) |N\alpha(\delta) - 1|^{1/2} |N\alpha(\delta)^{-1} - 1|^{1/2},$$

which we abbreviate as

$$\chi \left(\frac{N\alpha(\delta) - 1}{a} \right) |N\alpha(\delta)^{1/2} - N\alpha(\delta)^{-1/2}|.$$

For α_0 of type R_2 or R_3 , we include the contribution from the orbits of all multiples of α_{res} to the numerators of Δ_{II}, Δ_{IV} :

$$\Delta_{\alpha_0}(\delta) = \chi \left(\frac{N\alpha(\delta)^2 - 1}{a} \right) |N\alpha(\delta) - N\alpha(\delta)^{-1}|.$$

On the other hand, the roots $\pm\alpha'_0$ of T'^{δ_0} form two Galois orbits and we include them both. Thus if δ' is a θ -regular element in $T'^{\delta_0}(\mathbb{R})^0 \delta_0$, then we define $\Delta_{\alpha'_0}(\delta')$ as we have $\Delta_{\alpha_0}(\delta)$, but using only the contribution from $+\alpha'_0$ for the absolute value term. Set

$$\begin{aligned} \Delta_{\pm\alpha'_0}(\delta') &= \Delta_{\alpha'_0}(\delta') \cdot \Delta_{-\alpha'_0}(\delta') \\ &= \chi' \left(\frac{N\alpha'(\delta')^r - 1}{a'} \right) \cdot |N\alpha'(\delta')^r - 1|^{1/2} \\ &\quad \cdot (\chi')^{-1} \left(\frac{N\alpha'(\delta')^{-r} - 1}{-a'} \right) \cdot |N\alpha'(\delta')^{-r} - 1|^{1/2} \\ &= \chi'(N\alpha'(\delta')^r) \cdot |(N\alpha'(\delta')^{r/2} - N\alpha'(\delta')^{-r/2})|, \end{aligned}$$

where $r = 1$ if α_0 is of type R_1 and $r = 2$ if α_0 is of type R_2 or R_3 .

For $\nu \in \mathbb{R}$, set $\delta_\nu = \exp(\nu Y(a\alpha_1^\vee)) \cdot \delta_0$, where $Y(a\alpha_1^\vee) \in \mathfrak{t}^{\delta_0}(\mathbb{R})$ corresponds under the bijection of Lemma 6.2 to the multiple $a\alpha_1^\vee$ of the coroot α_1^\vee regarded as an element of $\mathfrak{t}_H(\mathbb{R})$. Then δ_ν has as T_1 -norm the element γ_ν from the statement of Theorem 5.1. Also,

$$\delta_\nu(w) = w^{-1} \delta_\nu \theta(w) = \exp(\nu Y(a\alpha_1^\vee)) \cdot w^{-1} \delta_0 \theta(w)$$

since $w\alpha_0 = \alpha_0$ implies that $w^{-1} \cdot Y(a\alpha_1^\vee) = Y(a\alpha_1^\vee)$. Again starting with δ_0 , define $\delta_{s,\nu}$ with $\gamma_{s,\nu}$ as T'_1 -norm, and $\delta_{s,\nu}(w')$ similarly. For $|\nu|$ sufficiently small but nonzero, the elements $\delta_\nu(w), \delta_{s,\nu}(w')$ are θ -regular.

Since s is a Cayley transform mapping T^{δ_0} to T'^{δ_0} within the group $G_{\delta_0}^\theta$, we require that the Haar measures on $T^{\delta_0}(\mathbb{R})$ and $T'^{\delta_0}(\mathbb{R})$ are compatible in the

sense of [She79a]. (Also see Section 1.4 of [LS87]; we may start with differential forms, attach measures and define compatibility using $a/|a|$ in place of i .)

LEMMA 8.1. *Let $f \in \mathcal{C}(G(\mathbb{R}), \theta)$. Then for any choice of χ, a, χ_s and a_s we have*

$$\begin{aligned} & \lim_{\nu \rightarrow 0^+} \Delta_{\alpha_0}(\delta_\nu) O^{\theta, \varpi}(\delta_\nu(w), f) - \lim_{\nu \rightarrow 0^-} \Delta_{\alpha_0}(\delta_\nu) O^{\theta, \varpi}(\delta_\nu(w), f) \\ &= d(\alpha_0) \lim_{\nu \rightarrow 0} \Delta_{\pm \alpha'_0}(\delta_{s, \nu}) O^{\theta, \varpi}(\delta_{s, \nu}(w'), f), \end{aligned}$$

where $d(\alpha_0) = 2$ if w_0 is realized in $G(\mathbb{R})$ in the sense of Section 7, and $d(\alpha_0) = 1$ otherwise.

For the proof, we first replace the version of Harish-Chandra’s compactness principle in Section 4 of [She79a] by the following.

LEMMA 8.2. *If C is a compact subset of $G(\mathbb{R})$, then there exist a neighborhood \mathcal{Y} of 0 in $\mathfrak{g}_{\delta_0}^\theta(\mathbb{R})$ and a compact subset \overline{C} of $\text{Cent}_\theta(\delta_0, G)(\mathbb{R}) \setminus G(\mathbb{R})$ such that if $g \in G(\mathbb{R})$, $Y \in \mathcal{Y}$, and $g^{-1}(\exp Y)\delta_0 \theta(g) \in C$, then $\text{Cent}_\theta(\delta_0, G)(\mathbb{R})g \in \overline{C}$.*

Proof. We follow the argument for Theorem 8.1.4.1 of [War72] in our setting, noting the arguments for Proposition 3.1 of [Ren97]. □

Proof of Lemma 8.1. Notice that the choice of χ, a, χ' , and a' does not matter, by an argument as in the first step of the proof of Theorem 4.2. In particular, there is no harm in taking χ' trivial and $a' = 1$.

By a continuity argument (see the appendix), it is enough to consider the case that $f \in C_c^\infty(G(\mathbb{R}), \theta)$. Using Lemma 8.2 with δ_0 replaced by $\delta_0(w)$, we may then apply a variant of Harish-Chandra’s descent argument (specifically, we generalize step by step the arguments of [She79a, §4]) to write the normalized twisted integrals $\Phi^{\theta, \varpi}(\delta_\nu(w), f)$ and $\Phi^{\theta, \varpi}(\delta_{s, \nu}(w'), f)$ as the normalized ordinary orbital integrals of a function ϕ in $C_c^\infty(G(\delta_0)^+)$, evaluated at $\exp(\nu Y(a\alpha_1^\vee))$ and $\exp(\nu Y(a'(\alpha_1')^\vee))$, respectively. Here $G(\delta_0)^+$ denotes the identity component of the derived group of $G_{\delta_0}^\theta(\mathbb{R})$. In the descent we may replace $\delta_\nu(w)$ by an element $\delta = \varepsilon\delta_0$ with ε sufficiently close to the identity in $\exp \mathfrak{t}^{\delta_0}(\mathbb{R})$ so that δ is θ -regular. There will be no harm in assuming further that δ is strongly θ -regular, so that $\text{Cent}_\theta(\delta, G) = T_{\delta_0}$. (Otherwise we use $T_{\delta_0}(\mathbb{R})$ in place of $\text{Cent}_\theta(\delta, G)(\mathbb{R})$ in the definition of twisted orbital integral.) We may do the same in T'_{δ_0} , replacing $\delta_{s, \nu}(w')$ by an appropriate element $\delta' = \varepsilon'\delta_0$.

The constant $d(\alpha_0)$ appears when we generalize Proposition 4.4 of [She79a]. We have $T_{\delta_0} = Z^\theta T^{\delta_0}$, $T'_{\delta_0} = Z^\theta T'^{\delta_0}$, and an argument with root vectors shows that we also have $\text{Cent}_\theta(\delta_0, G) = Z^\theta G_{\delta_0}^\theta$. (Here Z^θ denotes the θ -invariants in the center of G .) Denote the center of $G_{\delta_0}^\theta$ by Z_{δ_0} , and write \mathcal{G} for the product $Z_{\delta_0}(\mathbb{R}).G(\delta_0)^+$. Then the three indices, all finite, that concern us are

$[\text{Cent}_\theta(\delta_0, G)(\mathbb{R}) : \mathcal{G}]$, $[T_{\delta_0}(\mathbb{R}) : T_{\delta_0} \cap \mathcal{G}]$, and $[T'_{\delta_0}(\mathbb{R}) : T'_{\delta_0} \cap \mathcal{G}]$. We use them to replace the three indices in the statement of Proposition 4.4. Arguing as in [She79a], we see that a coset of \mathcal{G} in $\text{Cent}_\theta(\delta_0, G)(\mathbb{R})$ has a representative g which normalizes T^{δ_0} and $T^{\delta_0} \cap G(\delta_0)^+$, so that $g\alpha_0 = \pm\alpha_0$. Suppose $w_0 \in \text{Cent}_\theta(\delta_0, G)$ realizes the Weyl reflection for α_0 . Either g or w_0g lies in T_{δ_0} , and $T_{\delta_0} \cap \mathcal{G} = T^{\delta_0}(\mathbb{R})$. Suppose we cannot choose w_0 in $\text{Cent}_\theta(\delta_0, G)(\mathbb{R})$; i.e., w_0 is not realized in $G(\mathbb{R})$ in the sense of Section 7. Then we conclude that all three indices are the same. Suppose we may choose w_0 in $\text{Cent}_\theta(\delta_0, G)(\mathbb{R})$. Then the first index is twice the second, and further the first equals the third. Now we can proceed with the descent along the same lines as in Section 4 of [She79a], and the constant $d(\alpha_0)$ will persist to the final jump formula in the statement of Lemma 8.1.

Let $Y_0(a) = a\alpha_0^\vee \in \mathfrak{t}^{\delta_0}(\mathbb{R})$ and $Y'_0(a') = a'(\alpha'_0)^\vee \in \mathfrak{t}^{\delta_0}(\mathbb{R})$. (We could drop a' from notation since we have assumed $a' = 1$.) Then the familiar jump formula at the identity element for the ordinary orbital integrals of ϕ may be rewritten as

$$\begin{aligned} \lim_{\nu \rightarrow 0^+} \Delta(\exp \nu Y_0(a)) O(\exp \nu Y_0(a), \phi) &= \lim_{\nu \rightarrow 0^-} \Delta(\exp \nu Y_0(a)) O(\exp \nu Y_0(a), \phi) \\ &= \lim_{\nu \rightarrow 0} \Delta'(\exp \nu Y'_0(a')) O(\exp \nu Y'_0(a'), \phi), \end{aligned}$$

where $\Delta(\exp \nu Y_0(a))$ is given by

$$\begin{aligned} \chi \left(\frac{\alpha_0(\exp \nu Y_0(a)) - 1}{a} \right) &\left| (\alpha_0(\exp \nu Y_0(a))^{1/2} - \alpha_0(\exp \nu Y_0(a))^{-1/2}) \right| \\ &= \chi \left(\frac{e^{2\nu a} - 1}{a} \right) \left| e^{\nu a} - e^{-\nu a} \right|, \end{aligned}$$

and

$$\Delta'(\exp \nu Y'_0(a')) = \left| e^\nu - e^{-\nu} \right|.$$

The vectors $Y(a\alpha_1^\vee)$, $Y(a'(\alpha'_1)^\vee)$ are positive multiples of $Y_0(a)$, $Y'_0(a')$, and so it remains to check that the (germs at the identity of the) normalizing factors $\Delta, \Delta', \Delta_{\alpha_0}, \Delta_{\pm\alpha'_0}$ behave correctly under a rescaling of the variable ν . Rather than write down the evident general principle, we record explicit calculations for each of the three types for α .

Assume first that $\alpha_0 = \alpha_{\text{res}}$, where α is a root of T in G^* of type R_1 . Here, as in Section 6, we have transported the root α_0 of T^{δ_0} in $G_{\delta_0}^\theta$ to $(T^{\theta^*})^0$ by the twist $\text{Int}(g) \circ \psi : G_{\delta_0}^\theta \rightarrow G_{\delta_0}^{\theta^*}$, without change in notation. We similarly identify the elements Y and Y^* of Lemma 6.2. The coroot of α_0 is $N(\alpha^\vee)$, the sum of the coroots in the θ^\vee -orbit of α^\vee , so that $Y_0(a) = aN(\alpha^\vee)$. The root α_1 of T_1 in H_1 has coroot $(\alpha^\vee)_{\text{res}}$. In the Lie algebra $\mathfrak{t}_1 \simeq \mathfrak{t}_{\theta^*} = \mathfrak{t}/(\theta^* - 1)\mathfrak{t}$, we identify $(\alpha^\vee)_{\text{res}}$ with the coset of $\alpha^\vee \in \mathfrak{t}$. Then $Y(a\alpha_1^\vee)$ must be the real θ^* -invariant $\frac{a}{l_\alpha} N(\alpha^\vee) = \frac{1}{l_\alpha} Y_0(a)$, where l_α is the cardinality of the θ^\vee -orbit of

α^\vee (or θ^* -orbit of α). Since $N\alpha(\delta_0) = N\alpha(\delta_0^*) = 1$ and $\langle N\alpha, N(\alpha^\vee) \rangle = 2l_\alpha$, we have that

$$\begin{aligned} \Delta_{\alpha_0}(\delta_\nu) &= \Delta_{\alpha_0}(\exp(\nu Y(a\alpha_1^\vee))) \\ &= \chi\left(\frac{N\alpha(\exp(\frac{\nu a}{l_\alpha} N(\alpha^\vee))) - 1}{a}\right) \\ &\quad \cdot \left| N\alpha\left(\exp\left(\frac{\nu a}{l_\alpha} N(\alpha^\vee)\right)\right)^{1/2} - N\alpha\left(\exp\left(\frac{\nu a}{l_\alpha} N(\alpha^\vee)\right)\right)^{-1/2} \right| \\ &= \chi\left(\frac{e^{2\nu a} - 1}{a}\right) |e^{\nu a} - e^{-\nu a}|. \end{aligned}$$

By the same argument, $Y(a'(\alpha_1')^\vee) = \frac{1}{l_\alpha} Y_0'(a')$ and

$$\Delta_{\pm\alpha_0'}(\delta_{s,\nu}) = \Delta_{\pm\alpha_0'}(\exp(Y(a'(\alpha_1')^\vee))) = |e^\nu - e^{-\nu}|.$$

We can now finish the proof for the case α_0 is of type R_1 . In the limit formula for the orbital integrals of ϕ , replace the variable ν throughout by $\frac{1}{l_\alpha}\nu$. Rewrite the quotient of

$$\chi\left(\frac{e^{2\nu a} - 1}{a}\right) |e^{\nu a} - e^{-\nu a}|$$

by

$$\chi\left(\frac{e^{2\nu a/l_\alpha} - 1}{a}\right) |e^{\nu a/l_\alpha} - e^{-\nu a/l_\alpha}|,$$

as

$$\chi(e^{\nu a(1-1/l_\alpha)}) \chi\left(\frac{\sin(\nu b)}{\sin(\nu b/l_\alpha)}\right) \frac{\sin(\nu b)}{\sin(\nu b/l_\alpha)},$$

where $a = ib$. Since χ is trivial on positive real numbers, the second term in this product is trivial, and so the quotient extends continuously at $\nu = 0$ with nonzero value l_α . The same is true, with same value l_α , for the analogue

$$|e^\nu - e^{-\nu}| |e^{\nu/l_\alpha} - e^{-\nu/l_\alpha}|^{-1} = \frac{\sinh(\nu)}{\sinh(\nu/l_\alpha)}$$

on the other Cartan subgroup. This allows us to replace $\Delta(\exp \frac{1}{l_\alpha} \nu Y_0(a))$ by $\Delta_{\alpha_0}(\delta_\nu)$ and $\Delta'(\exp \frac{1}{l_\alpha} \nu Y_0'(a'))$ by $\Delta_{\pm\alpha_0'}(\delta_{s,\nu})$ when computing limits, and so we get the desired formula.

Suppose that α_0 is of type R_3 , so that α_0 again has coroot $N(\alpha^\vee)$, and $Y_0(a) = aN(\alpha^\vee)$. Here the root α_1 of T_1 in H_1 has coroot $(\beta^\vee)_{\text{res}}$ in the notation of Section 1.3 of [KS99], where $(\alpha^\vee)_{\text{res}} = 2(\beta^\vee)_{\text{res}}$ (see Lemma 6.2). Thus

$$Y^*(a\alpha_1^\vee) = \frac{a}{l_\beta} N(\beta^\vee) = \frac{a}{2l_\alpha} N(\alpha^\vee) = \frac{1}{2l_\alpha} Y_0(a).$$

Again

$$\Delta_{\alpha_0}(\delta_\nu) = \Delta_{\alpha_0}(\exp(\nu Y(a\alpha_1^\vee))),$$

since $N\alpha(\delta_0^*)^2 = (-1)^2 = 1$. Also, $\langle N\alpha, N(\alpha^\vee) \rangle = 2l_\alpha$ and so we again get the formula

$$\Delta_{\alpha_0}(\delta_\nu) = \chi \left(\frac{e^{2\nu a} - 1}{a} \right) |e^{\nu a} - e^{-\nu a}|.$$

After the substitution of $\frac{1}{2l_\alpha}\nu$ for ν , we have to examine the quotient of

$$\chi \left(\frac{e^{2\nu a} - 1}{a} \right) |e^{\nu a} - e^{-\nu a}|$$

by

$$\chi \left(\frac{e^{\nu a/l_\alpha} - 1}{a} \right) |e^{\nu a/2l_\alpha} - e^{-\nu a/2l_\alpha}|,$$

and we may proceed as for R_1 .

Suppose that α_0 is of type R_2 . In keeping with the notation of the last paragraph, we write the coroot of α_0 as $2N(\beta^\vee)$ and $Y_0(a) = 2aN(\beta^\vee)$. Now the coroot α_1^\vee may be either $(\beta^\vee)_{\text{res}}$ or $(\alpha^\vee)_{\text{res}} = 2(\beta^\vee)_{\text{res}}$. Suppose $\alpha_1^\vee = (\beta^\vee)_{\text{res}}$. Then

$$Y^*(a\alpha_1^\vee) = \frac{a}{l_\beta} N(\beta^\vee) = \frac{1}{2l_\beta} Y_0(a).$$

Also,

$$\langle N\beta, N(\beta^\vee) \rangle = 2l_\alpha = l_\beta,$$

so that

$$\Delta_{\alpha_0}(\delta_\nu) = \chi \left(\frac{e^{2\nu a} - 1}{a} \right) |e^{\nu a} - e^{-\nu a}|.$$

Suppose $\alpha_1^\vee = (\alpha^\vee)_{\text{res}}$. Then

$$Y^*(a\alpha_1^\vee) = \frac{a}{l_\alpha} N(\alpha^\vee) = \frac{1}{2l_\alpha} Y_0(a) = \frac{1}{l_\beta} Y_0(a),$$

$$\Delta_{\alpha_0}(\delta_\nu) = \chi \left(\frac{e^{4\nu a} - 1}{a} \right) |e^{2\nu a} - e^{-2\nu a}|,$$

and once again we finish the argument the same way. □

For any $w \in \mathfrak{A}_\theta(T^{\delta_0})$, we may also do a similar descent (i.e., find ϕ as in the proof above) around $\delta_0(w)$ in $\text{Cent}(\delta_0(w), G)(\mathbb{R})$. If $\text{Cent}(\delta_0(w), G)(\mathbb{R})$ is compact modulo center, then we conclude that $O^{\theta, \varpi}(\delta_\nu(w), f)$, like the ordinary orbital integral for ϕ , is bounded as $\nu \rightarrow 0$ and so contributes nothing to the jump formula for Φ_1 . This remark also applies to the proof of Lemma 7.2 for the setting where every $O^{\theta, \varpi}(\delta_\nu(w), f)$ is of this type.

9. Twisted transfer factors

We now examine the various terms $\Delta_I, \dots, \Delta_{IV}$ of the twisted transfer factor $\Delta(\gamma_1, \delta)$ in the setting of toral descent data at γ_0 (last paragraph of Section 7). For the relative analysis, we have three associated Cayley transforms. First, there is $s_1 : T_1 \rightarrow T_1^{s_1}$ associated with the root α_1 in H_1 . Second, there is $t^{*-1} = s^* : T \rightarrow T'$ in $(G^{\theta^*})^0$ associated with the least positive multiple of α_{res} that is a root, and, finally, there is $t^{-1} = s : T^{\delta_0} \rightarrow T'^{\delta_0}$ for the root α_0 in $G_{\delta_0}^\theta$. Details of the construction of the terms Δ_I and Δ_{III} will be included where they are used in proofs. There is a last ingredient for our setting, a twisted analogue of the s -compatible data sets of Sections 5–7. The results for Δ_{II} and Δ_{IV} then follow quickly (Corollary 9.2), while the analysis for Δ_I and Δ_{III} takes longer. The proof of the main lemma, Lemma 9.3, will consist of several steps to remove parts (which we show to be trivial) of a particular Δ_{III} term until we arrive finally at a term we can compute explicitly and also show to be trivial.

We choose a -data and χ -data following Section 1.3 of [KS99]. These are data for the system of restricted roots β_{res} of T in G^* . We use the same pair $a_{\beta_{\text{res}}}, \chi_{\beta_{\text{res}}}$ for any positive multiple of β_{res} that is also a restricted root and the same data for coroots of the restrictions and for the restrictions of coroots

$$a_{\beta_{\text{res}}} = a_{(\beta^\vee)_{\text{res}}} = a_{(\beta_{\text{res}})^\vee}$$

and

$$\chi_{\beta_{\text{res}}} = \chi_{(\beta^\vee)_{\text{res}}} = \chi_{(\beta_{\text{res}})^\vee}.$$

This provides us then with data for the roots and coroots of T_1 in H_1 . We make the same choices for the torus T' and define s^* -compatibility for the twisted data set $\{a_{\beta_{\text{res}}}\}, \{\chi_{\beta_{\text{res}}}\}, \{a_{\beta'_{\text{res}}}\}, \{\chi_{\beta'_{\text{res}}}\}$ as in Section 3. Our constructions ensure that s^* -compatible data (which we also call s -compatible) provide data for $T_1, T_1^{s_1}$ that are s_1 -compatible.

Following p. 36 of [KS99], we write Δ_{II} in quotient form

$$\Delta_{II} = \Delta_{II}^{\text{num}} / \Delta_{II}^{\text{denom}},$$

where Δ_{II}^{num} is a term attached to (G, θ) and $\Delta_{II}^{\text{denom}}$ is from standard endoscopy for the group H_1 . We now prefer to index the contributions to Δ_{II}^{num} by the orbits \mathcal{O} of *reduced* restricted roots α_{res} . Thus the formulas of p. 36 of [KS99] yield

$$\Delta_{II}^{\text{num}}(\gamma_1, \delta) = \prod_{\mathcal{O}} \chi_{\alpha_{\text{res}}} \left(\frac{N\alpha(\delta^*)^r - 1}{a_{\alpha_{\text{res}}}} \right),$$

where α_{res} represents \mathcal{O} , and $r = 1$ or $r = 2$ according as α_{res} is of type R_1 or of type R_2 .

Remark 1. Waldspurger [Wal09] has pointed out that a correction is needed in the definition of twisted transfer factors in the nonarchimedean case and that it can be made by the insertion of 2 in certain contributions to Δ_{II} when the system of restricted roots α_{res} is not reduced. This has no effect in our present archimedean case; see [KS12, §1] for details. An alternate way of making the correction, which involves Δ_I instead and makes sense in all characteristics, is presented in [KS12]. It also has no effect on the definitions in the archimedean case [KS12, Prop. 3.5.2].

Remark 2. First we observe an error on p. 137 of [KS99] pointed out to us by Waldspurger. The exponent -1 in the formula (A.3.13) does not belong there. We emphasize that by the term *Langlands' pairing* in the statement of (A.3.13) we mean the pairing from [Lan97]. The source of this error is on p. 131, where what is described as the *Langlands map* is the reciprocal of that defined in [Lan97]. To be explicit in the case at hand, if T is a torus defined over \mathbb{R} , then the isomorphism $H_1(\mathbb{C}^\times, X_*(T)) \rightarrow X_*(T) \otimes \mathbb{C}^\times = T(\mathbb{C})$ defined in the middle of p. 131 of [KS99] has an exponent -1 not present in the isomorphism defined in [Lan97]. (An explicit formula is found on p. 243 of [Lan97] after the first commutative diagram.) We resolve this by inverting the formula for the pairing in (A.3.9). Then the formula (A.3.13) is true as stated in [KS99]. Now, in principle, we should insert an exponent -1 in the formula (A.3.14) involving Tate-Nakayama duality, but here in the archimedean case the term is simply a sign, and so we may use the formula as stated in [KS99]. Our resolution agrees with that suggested to us by Waldspurger for the general case; i.e., our Δ coincides with the term Δ' of (5.4.1) in [KS12]. It also gives the correct shift in infinitesimal character for Langlands functoriality of the dual spectral transfer [She11].

Returning now to our analysis of the various terms $\Delta_I, \dots, \Delta_{IV}$, we observe the following generalization of Lemma 4.1.

LEMMA 9.1. *For any s -compatible twisted data set $\{a_{\beta_{\text{res}}}\}, \{\chi_{\beta_{\text{res}}}\}, \{a_{\beta'_{\text{res}}}\}, \{\chi_{\beta'_{\text{res}}}\}$ we have*

$$\prod_{\mathcal{O}} \chi_{\beta_{\text{res}}} \left(\frac{(N\beta(\delta_0)^r - 1)}{a_{\beta_{\text{res}}}} \right) = \prod_{\mathcal{O}'} \chi_{\beta'_{\text{res}}} \left(\frac{(N\beta'(\delta_0)^r - 1)}{a_{\beta'_{\text{res}}}} \right).$$

On the left, the product is over all Galois orbits \mathcal{O} of reduced restricted roots for T^{δ_0} (i.e., of types R_1 or R_2) *except* those containing a multiple of α_0 . Each term is independent of the choice of representative β_{res} for \mathcal{O} ; $r = 1$ if β_{res} is of type R_1 and $r = 2$ if β_{res} is of type R_2 . The right side is defined analogously, using all Galois orbits \mathcal{O}' of reduced restricted roots for T'^{δ_0} *except* those containing a multiple of α'_0 . For the precise meaning of $N\beta(\delta_0)$, see the remark after Lemma 6.5.

Proof. We match contributions to each side of the formula orbit by orbit as in the proof of Lemma 4.1. \square

Because we will eventually consider derivatives of the transforms $\Psi_{a,\chi}$ and $\Psi_{a',\chi'}$, we use the variables $\delta_0(Y), \gamma_0(Y_1)$, etc. from Lemma 6.2 in our limit formulas for terms of the transfer factors. Each of Δ_{II} and Δ_{IV} is defined as a quotient of a term associated with G and a term associated with H_1 (Sections 4.3 and 4.5 of [KS99]). Each denominator cancels with an identical term in one of the transforms $\Psi_{a,\chi}$ and $\Psi_{a',\chi'}$ of Theorem 5.1. Denote the numerators as $\Delta_{II,\text{num}}$ and $\Delta_{IV,\text{num}}$. These numerators contribute the factors $\Delta_{\alpha_0}, \Delta_{\pm\alpha'_0}$ from the orbits in $\mathbb{Q}\alpha_0, \mathbb{Q}\alpha'_0$ for the twisted transforms in the jump formulas of the last section, and so these terms will also be removed. In the case of Δ_{II} , what remains is each side of the equation in Lemma 9.1. There is a similar assertion for Δ_{IV} . Thus

COROLLARY 9.2. *For an s -compatible data set and toral descent data at γ_0 we have*

$$\begin{aligned} \lim_{Y_1 \rightarrow 0} \Delta_{II,\text{num}}(\gamma_0(Y_1)) \Delta_{IV,\text{num}}(\delta_0(Y)) \Delta_{\alpha_0}(\delta_0(Y))^{-1} \\ = \lim_{Y'_1 \rightarrow 0} \Delta_{II,\text{num}}(\gamma_0(Y'_1)) \Delta_{IV,\text{num}}(\delta_0(Y')) \Delta_{\pm\alpha'_0}(\delta_0(Y'))^{-1}. \end{aligned}$$

LEMMA 9.3 (Main lemma). *For an s -compatible data set and toral descent data at γ_0 , we have*

$$\lim_{Y_1, Y'_1 \rightarrow 0} \Delta_{III}(\gamma_0(Y_1), \delta_0(Y); \gamma_0(Y'_1), \delta_0(Y')) = 1.$$

Transitivity of the relative transfer factor (Lemma 5.1.A of [KS99]) then implies immediately the following about the terms of type Δ_{III} which appear in the limit formulas of Theorem 5.1 and Lemmas 10.1 and 10.2.

COROLLARY 9.4. *In the same setting, we have*

$$\lim_{Y_1 \rightarrow 0} \Delta_{III}(\gamma_0(Y_1), \delta_0(Y); \bar{\gamma}, \bar{\delta}) = \lim_{Y'_1 \rightarrow 0} \Delta_{III}(\gamma_0(Y'_1), \delta_0(Y'); \bar{\gamma}, \bar{\delta}).$$

Proof of Lemma 9.3. We start by showing that

$$\Delta_{III}(\gamma_0(Y_1), \delta_0(Y); \gamma_0(Y'_1), \delta_0(Y')),$$

defined as the term $\langle \mathbf{V}_1, \mathbf{A}_1 \rangle$ on p. 43 of [KS99], is the product of a term independent of Y_1, Y'_1 , which we will denote

$$\Delta_{III}(\gamma_0, \delta_0; T_1, T'_1),$$

and a term which has limit 1 as Y_1, Y'_1 approach 0. A longer argument will then show that

$$\Delta_{III}(\gamma_0, \delta_0; T_1, T'_1) = 1.$$

Recall Remark 2 earlier in this section: the pairing $\langle -, - \rangle$ is now defined by the reciprocal of the formula displayed on p. 135 of [KS99]. First we factor \mathbf{V}_1 as $\mathbf{V}_0 \cdot \mathbf{V}(Y)$. The tori U, S , and S_1 are attached to T, T' in Section 4.4 of [KS99]. Notice that our T_1, T'_1 are labeled T_{H_1}, T'_{H_1} there. The element \mathbf{V}_1 belongs to the hypercohomology group denoted $H^1(\Gamma, U \xrightarrow{1-\theta} S_1)$. It is the class of the pair (V, D_1) , where $V = V(\sigma)$ is a Galois 1-cocycle in U and D_1 is an element in S_1 , and the hypercocycle identity $(1-\theta)V = \sigma(D_1)D_1^{-1}$ is satisfied. We have defined 1-cochains $v(\sigma), v'(\sigma)$ in Section 6. The pair $(v(\sigma)^{-1}, v'(\sigma))$ lies in $T_{\text{sc}} \times T'_{\text{sc}}$. Its image under the projection to

$$U = T_{\text{sc}} \times T'_{\text{sc}} / \{(z^{-1}, z) : z \in Z_{\text{sc}}\}$$

is, by definition, V . (Our modification at the end of Section 6 does not affect V .) To describe D_1 we start with the elements $\delta_0^*(Y), (\delta_0^*)'(Y')$ of T, T' (Lemma 6.2). To resolve a notational conflict with [KS99], we write the pull-back torus T_1 of p. 42 of [KS99] as T_2 . Then $(\delta_0^*(Y), \gamma_0(Y_1))$ lies in T_2 and $((\delta_0^*)'(Y'), \gamma_0(Y'_1))$ lies in T'_2 . The element

$$((\delta_0^*(Y), \gamma_0(Y_1))^{-1}, ((\delta_0^*)'(Y'), \gamma_0(Y'_1)))$$

of $T_2 \times T'_2$ factors as

$$((\delta_0^*, \gamma_0)^{-1}, ((\delta_0^*)', \gamma_0)) \cdot ((\exp Y^*, \exp Y_1)^{-1}, (\exp Y^{*'}, \exp Y'_1)).$$

This factoring persists for images in the quotient S_1 (defined on p. 42 of [KS99]), and we write the factoring in S_1 as $D_1 = D_0 \cdot D(Y_1, Y'_1)$. Because Y^* , etc., lie in the real Lie algebras of the relevant tori, we also have a factoring of hypercocycles:

$$(V, D_1) = (V, D_0) \cdot (1, D(Y_1, Y'_1)).$$

Then \mathbf{V}_0 will denote the class of (V, D_0) and $\mathbf{V}(Y_1, Y'_1)$ will denote the class of

$$(1, D(Y_1, Y'_1)).$$

Define

$$\Delta_{III}(\gamma_0, \delta_0; T_1, T'_1) = \langle \mathbf{V}_0, \mathbf{A}_1 \rangle,$$

so that

$$\langle \mathbf{V}_1, \mathbf{A}_1 \rangle = \Delta_{III}(\gamma_0, \delta_0; T_1, T'_1) \cdot \langle \mathbf{V}(Y_1, Y'_1), \mathbf{A}_1 \rangle.$$

To see that the complementary term $\langle \mathbf{V}(Y_1, Y'_1), \mathbf{A}_1 \rangle$ has limit 1 as Y_1, Y'_1 approach 0, we recall that the element \mathbf{A}_1 in the hypercohomology group $H^1(W_{\mathbb{R}}, S_1^{\vee} \xrightarrow{1-\theta^{\vee}} U^{\vee})$ is represented by the pair (A^{-1}, \mathfrak{s}_U) specified on p. 45 of [KS99]. In particular, A is a 1-cocycle of $W_{\mathbb{R}}$ in S_1^{\vee} . The pairing for hypercohomology is compatible with the Langlands parametrization of characters on $S_1(\mathbb{R})$ ([KS99, A.3.13] , as corrected in Remark 2). This allows us to compute

$\langle \mathbf{V}(Y_1, Y'_1), \mathbf{A}_1 \rangle$ as the value of the character attached to the class of A^{-1} in $H^1(W_{\mathbb{R}}, S_1^{\vee})$ on the image $D(Y_1, Y'_1)$ of

$$((\exp Y^*, \exp Y_1), (\exp Y^{*'}, \exp Y'_1))$$

in $S_1(\mathbb{R})$. The limit assertion is now immediate.

Thus it remains to show that $\langle \mathbf{V}_0, \mathbf{A}_1 \rangle = 1$. We factor each of (V, D_0) and (A^{-1}, \mathfrak{s}_U) further, and so reduce to calculations with familiar pairings in cohomology.

Recall from the end of Section 7 that we have arranged that the 1-cochains $v(\sigma), v'(\sigma)$ are such that $(\theta^* - 1)v(\sigma)$ and $(\theta^* - 1)v'(\sigma)$ are the same central element, so that $V(\sigma)$ is θ^* -invariant. Thus V is a 1-cocycle in U^{θ^*} . We then have (or may check directly) that $D_0 \in S_1(\mathbb{R})$, so that (V, D_0) factors as $(V, 1) \cdot (1, D_0)$. Turning to the dual side, we have from the hypercocycle equation that the element \mathfrak{s}_U determines a Γ -invariant element in $U_{\theta}^{\vee} = U^{\vee} / (1 - \theta^{\vee})U^{\vee}$, and hence an element $\mathfrak{s}_{U, \theta}$ in $\pi_0((U_{\theta}^{\vee})^{\Gamma})$. The group U^{θ^*} is a torus since the usual isomorphism of U with $T_{\text{sc}} \times T'_{\text{ad}}$ (see p. 38 of [KS99]) is θ^* -equivariant and the invariants for each factor in the product torus are connected. The dual of U^{θ^*} is U_{θ}^{\vee} . Write $\langle V, \mathfrak{s}_{U, \theta} \rangle$ for the Tate-Nakayama pairing of the class of V in $H^1(\Gamma, U^{\theta^*})$ with $\mathfrak{s}_{U, \theta} \in \pi_0((U_{\theta}^{\vee})^{\Gamma})$ and Λ for the character on $S_1(\mathbb{R})$ attached to A^{-1} by the Langlands correspondence. Then we may compute $\langle \mathbf{V}_0, \mathbf{A}_1 \rangle$ as the product $\langle V, \mathfrak{s}_{U, \theta} \rangle \cdot \Lambda(D_0)$ (see p. 135 of [KS99]). We check now that each term in this product equals 1.

The image $v'_{\text{ad}}(\sigma)$ of the cochain $v'(\sigma)$ in T'_{ad} is a cocycle in the torus $(T'_{\text{ad}})^{\theta^*}$. As usual, we identify the cocharacters of this torus with the θ^* -invariant coweights of T'_{ad} . Under the Tate-Nakayama isomorphism

$$H^{-1}(\Gamma, X_*((T'_{\text{ad}})^{\theta^*})) \rightarrow H^1(\Gamma, (T'_{\text{ad}})^{\theta^*}),$$

$v'_{\text{ad}}(\sigma)$ is cohomologous to the cup product of the fundamental 2-cocycle for \mathbb{C}/\mathbb{R} with a θ^* -invariant coweight x'_{cw} for T'_{ad} such that $\sigma x'_{cw} = -x'_{cw}$, i.e., to $(-1)^{x'_{cw}}$. Write

$$v'_{\text{ad}}(\sigma) = (-1)^{x'_{cw}}(\sigma(t')t'^{-1})_{\text{ad}},$$

where t' lies in the torus $(T'_{\text{sc}})^{\theta^*}$. Extend the root α'_0 trivially to Z_{sc} . Then our assumptions on g' ensure that

$$\alpha'_0(v'(\sigma)) = \alpha'_0(v'_{\text{ad}}(\sigma)) = 1.$$

Thus $\alpha'_0(\sigma(t')t'^{-1}) = 1$. Apply the inverse Cayley transform t^* to x'_{cw} to obtain a θ^* -invariant coweight x_{cw} for T_{ad} . Then $\sigma x_{cw} = -x_{cw}$ and a calculation shows that

$$v_{\text{ad}}(\sigma) = (-1)^{x_{cw}}((-1)^{\epsilon \alpha_0^{\vee}} \sigma(t'')t''^{-1})_{\text{ad}},$$

where $\epsilon \in \{0, 1\}$ and $t'' \in (T_{\text{sc}})^{\theta^*}$. To recall the characters and cocharacters of U , we use t^* to identify T' with T over \mathbb{C} . The characters may be identified as

pairs (λ, μ) , where each of λ and μ is a weight of T_{ad} and $\lambda - \mu$ is an integral combination of roots, while the cocharacters may be identified as pairs (λ^\vee, μ^\vee) of coweights such that $\lambda^\vee + \mu^\vee$ is an integral combination of coroots. The canonical pairing is

$$\langle (\lambda, \mu), (\lambda^\vee, \mu^\vee) \rangle = \langle \lambda - \mu, \lambda^\vee \rangle + \langle \mu, \lambda^\vee + \mu^\vee \rangle.$$

Set

$$x = (-x_{cw} - \epsilon\alpha_0^\vee, x_{cw}).$$

(Recall x'_{cw} has now been identified with x_{cw} .) Then x lies in $X_*(U^{\theta^*})$, $\sigma x = -x$, and, by evaluating characters on both sides of the following, we see that $(-1)^x = V(\sigma)\sigma(u)u^{-1}$, where u is the image in U^{θ^*} of $(t'', t')^{-1}$. Thus $\sigma \rightarrow (-1)^x$ is cohomologous to V .

We may now compute $\langle V, \mathfrak{s}_{U,\theta} \rangle$ by evaluating x , as character on $(U^{\theta^*})^\vee = U_\theta^\vee$, at the element $\mathfrak{s}_{U,\theta}$. In the notation of p. 39 of [KS99] where \mathfrak{s}_U is defined, we have arranged that $\tilde{\mathfrak{s}}_T = \tilde{\mathfrak{s}}_{T'}$, so that to show $\langle V, \mathfrak{s}_{U,\theta} \rangle = 1$, it is enough to show that $\alpha_0^\vee(\mathfrak{s}_T) = 1$; i.e., $N(\alpha^\vee)(\mathfrak{s}_T) = 1$ if α_0 is of type R_1 or R_3 , or $N(\alpha^\vee)(\mathfrak{s}_T)^2 = 1$ if α_0 is of type R_2 . But if α_0 is of type R_1 or R_3 , then the corresponding root α_1 of H_1 is of type R_1 or R_2 only, so that $N(\alpha^\vee)(\mathfrak{s}_T) = 1$, as desired. If α_0 is of type R_2 , then the corresponding root α_1 is of type R_2 or R_3 , and $N(\alpha^\vee)(\mathfrak{s}_T) = \pm 1$ accordingly. Since we need only $N(\alpha^\vee)(\mathfrak{s}_T)^2 = 1$, we are done. This remark, namely that $\alpha_0^\vee(\mathfrak{s}_T) = 1$, will be useful again. Also, a partial converse result (see the proof of Lemma 11.1) provides a crucial cancellation in the final steps of our proof of the main theorem.

It remains then to show that $\Lambda(D_0) = 1$. Here s -compatibility of the χ -data plays a key role, along with an extension of the comparison arguments of Section 4 of [LS90] already used in the definition of A in Section 4 of [KS99]. Our (second) argument for Lemma 9.5 below will have a similar structure, using the first lemma of comparison from [LS90] in place of the second.

The element D_0 of $S_1(\mathbb{R})$ is the image of $((\delta_0^*, \gamma_0)^{-1}, ((\delta_0^*)', \gamma_0))$ under $T_2 \times T'_2 \rightarrow S_1$. As before, we will use t^* to identify T' with T , and then T'_2 with T_2 , over \mathbb{C} . The element $((\delta_0^*)', \gamma_0)$ is thus identified with (δ_0^*, γ_0) . As on p. 42 of [KS99], we identify S_1 as $T'_2 \times T_{\text{ad}}$ (T'_2 is labelled T_1 in the reference) and then as $T_2 \times T_{\text{ad}}$. The Galois action on the first component is the transport σ' of that on T'_2 , while on the second component we use the twisted action

$$(1, t_{\text{ad}}) \rightarrow (\psi_{w_0}(\sigma'(t_{\text{ad}})), \sigma(t_{\text{ad}})).$$

Here $\psi_{w_0} : T_{\text{ad}} \rightarrow T_2$ is defined as follows. Pick $t_2 \in T_2$ in the inverse image of t_{ad} under the surjection $T_2 \rightarrow T \rightarrow T_{\text{ad}}$. Then $\psi_{w_0}(t_{\text{ad}}) = w_0(t_2)t_2^{-1}$ is independent of the choice for t_2 . The chosen Galois action makes

$$T_2 \times T'_2 \rightarrow S_1 \rightarrow T_2 \times T_{\text{ad}} : (t_2, t'_2) \rightarrow (t_2 t'_2, (t_2)_{\text{ad}})$$

defined over \mathbb{R} . Write δ_{ad} for the image of $(\delta_0^*, \gamma_0)^{-1}$ in T_{ad} . Then δ_{ad} is fixed by σ and σ' (recall our assumptions on $\delta_0^*, (\delta_0^*)'$, and $\psi_{w_0}(\delta_{\text{ad}}) = 1$, also because of our assumptions on $\delta_0^*, (\delta_0^*)'$). Notice that $D_0 \in S_1(\mathbb{R})$ is identified with

$$(1, \delta_{\text{ad}}) \in (T_2 \times T_{\text{ad}})(\mathbb{R}).$$

As in [KS99], we identify S_1^\vee as $T_2^\vee \times T_{\text{sc}}^\vee$, with Galois action

$$(t_2, t_{\text{sc}}) \rightarrow (\sigma'(t_2), \varphi_{w_0}(\sigma'(t_2))\sigma(t_{\text{sc}})).$$

(Recall we have chosen T_2' rather than T_2 to be the torus T_1 in [KS99].) Here $\varphi_{w_0} : T_2^\vee \rightarrow T_{\text{sc}}^\vee$ is the dual of ψ_{w_0} . (This is the variant of the definition of $\alpha(w_0)$ in [KS99] needed when U is replaced by S_1 , and we will recall how to compute it when needed below.) The 1-cocycle $A(w)$ of $W_{\mathbb{R}}$ in S_1^\vee is constructed as the element $(a_{T_2'}(w), x_{\text{sc}}(w))$ of $T_2^\vee \times T_{\text{sc}}^\vee$, where $x_{\text{sc}}(w)$ is a product

$$\hat{\tau}(w_0, \sigma') \cdot \hat{b}(w_0)^{-1} \cdot w_0(c'(w)) \cdot c(w)^{-1} \cdot \varphi_{w_0}(a_{T_2'}(w)).$$

To begin examining these terms, we observe that we may replace the cocycles $a_{T_2}(w), a_{T_2'}(w)$ by cocycles $a_-(w), a'_-(w)$ for which $\varphi_{w_0}(a'_-(w)) = 1$. Then $A(w)$ will be replaced by

$$A_-(w) = (a'_-(w), \hat{\tau}(w_0, \sigma') \cdot \hat{b}(w_0)^{-1} \cdot w_0(c'_-(w)) \cdot c_-(w)^{-1}),$$

and there is now no twist in the Galois action on the first component. This ensures that the second component is a cocycle for the action by σ . Our strategy then will be to examine that cocycle and see that the attached character on $T_{\text{ad}}(\mathbb{R})$ takes the value 1 at δ_{ad} , which is sufficient to complete the proof of the lemma.

To define $a_-(w), a'_-(w)$, it is more convenient to view S_1^\vee as a subtorus of $T_2^\vee \times T_2^\vee$, with Galois actions σ', σ on the first and second components respectively. The cocycle $A(w) = (a_{T_2'}(w), a_{T_2}(w))$ lies in S_1^\vee . By construction, $T^\vee \times T^\vee$ embeds in $T_2^\vee \times T_2^\vee$, and S_1^\vee contains the image of the standard homomorphism $T_{\text{sc}}^\vee \times T_{\text{sc}}^\vee \rightarrow T^\vee \times T^\vee$. Consider a cocycle in S_1^\vee which is the image of a cocycle $(a'_+(w), a_+(w))$ in $T_{\text{sc}}^\vee \times T_{\text{sc}}^\vee$. We will write this image also as $(a'_+(w), a_+(w))$. To evaluate the corresponding character on $S_1(\mathbb{R})$ under the Langlands correspondence on the element D_0 we may, by functoriality of the correspondence, evaluate at $(\delta_{\text{ad}}, \delta_{\text{ad}})$ the character on $T_{\text{ad}}(\mathbb{R}) \times T_{\text{ad}}(\mathbb{R})$ (σ' is the action for the first component, σ for the second) attached to $(a'_+(w), a_+(w))$ as cocycle in $T_{\text{sc}}^\vee \times T_{\text{sc}}^\vee$. We will choose $(a'_+(w), a_+(w))$ so that the resulting value is 1, and thus $\Lambda(D_0)$ is unchanged when we divide $A(w)$ by $(a'_+(w), a_+(w))$. The cocycles $a'_+(w), a_+(w)$ will come from $c'(w), c(w)$.

The cochain $c'(w)$ is defined as a quotient $r'_1(w)/r'_s(w)$ of terms from constructions in Section 2.5 of [LS87]. First, $r'_1(w)$ is the term $r_p(w)$ for the group $G_*^\vee = ((G^\vee)^{\theta^\vee})^0$, Galois action σ' and gauge p associated to our choice

of positive roots (that determined by our fixed Γ -splitting of G^\vee preserved by θ^\vee and our choice of toral data). Then

$$r'_1(w) = s_{p/p_0}(w) \prod r_{\pm\mathcal{O}'}(w),$$

where the product is over pairs $\pm\mathcal{O}'$ of orbits for σ' in the roots of $T_*^\vee = ((T^\vee)^{\theta^\vee})^0$ in G_*^\vee . The term $r'_s(w)$ is defined similarly, using the roots of T_*^\vee in H^\vee . In the next paragraph, we will keep track of contributions after cancellation, using the pairs $\pm\mathcal{O}'$ of orbits of roots in G_*^\vee (the reduced restricted roots) to index them.

We claim that there are nontrivial contributions from $\pm\mathcal{O}'$, $\pm 2\mathcal{O}'$ to $c'(w)$ only in the following two cases: (i) neither $\pm\mathcal{O}'$, $\pm 2\mathcal{O}'$ belongs to H^\vee , and (ii) $\pm 2\mathcal{O}'$ belongs to H^\vee . Recall we have fixed the root $\alpha_0 = \alpha_{\text{res}}$ of $G_{\delta_0}^\theta$ and (reduced) α_* is the multiple of α_0 that is a root of $(G^{\theta^*})^0$. Now on the dual side, we set α_{**} to be the multiple of $(\alpha^\vee)_{\text{res}}$ that is a root of T_*^\vee in G_*^\vee and denote by $\beta_{**} = (\beta^\vee)_{\text{res}}$ a root of T_*^\vee in G_*^\vee distinct from $\pm\alpha_{**}$. The coroot of β_{**} is $rN\beta$, where $r = 1$ unless β (and hence also β^\vee) is of type R_2 , in which case $r = 2$. The term $r_{\pm\mathcal{O}'}(w)$ is constructed in Section 2.5 of [LS87]. We will need its explicit form only for symmetric orbits. Then

$$r_{\pm\mathcal{O}'}(w) = \chi_{\beta_{**}}(u_0(w))^{rN\beta},$$

where β_{**} represents \mathcal{O}' and $u_0(w)$ is defined in Section 2.5 of [LS87]. This applies also if β_{**} is not reduced (as in case (ii)). Now to check the claim we examine the various possibilities as in the argument on p. 49 of [KS99]. We see that the contribution in case (i) is $r_{\pm\mathcal{O}'}(w)$, while in case (ii) it is $r_{\pm 2\mathcal{O}'}(w)^{-1}$. In the remaining cases, it is 1, as asserted. We of course define $c(w)$ in the same way as $c'(w)$, using instead the action σ .

The terms $s_{p/p_0}(w)$ are, in general, different for G_*^\vee and H^\vee . We have assumed that our toral data have the property that complex conjugates (relative to σ' only) of positive complex roots are positive. Then both terms $s_{p/p_0}(w)$ contributing to $c'(w)$, but not necessarily those contributing to $c(w)$, are trivial (see Section 2.4 of [LS87] for their definition) and will be deleted in notation. We will deal with $s_{p/p_0}(w)$ for the action defined by σ in the last paragraph of our proof.

Suppose $\mathcal{O}' \neq \{\pm\alpha_{**}\}$ is asymmetric and not orthogonal to α_{**} . Our plan is to remove a cocycle for each $\pm\mathcal{O}'$ contributing to $c'(w)$ and then to remove a matching cocycle from $c(w)$. Because there exist trivial χ -data for $\pm\mathcal{O}'$, the contribution $r_{\pm\mathcal{O}'}(w)$ must be a cocycle (see also Corollary 2.5.B of [LS87]), and we may compute the corresponding character $\Lambda_{\pm\mathcal{O}'}$ on $T_{\text{ad}}(\mathbb{R})$ as in Section 3.3 of [LS87]. Assume β_{**} belongs to $\pm\mathcal{O}'$. Suppose first that $\sigma'\beta_{**} = -w_0\beta_{**} \neq \pm\beta_{**}$ (i.e., β_{**} is complex for σ' and imaginary for σ). Then according to Lemma 3.3.D of [LS87], $\Lambda_{\pm\mathcal{O}'}(\delta_{\text{ad}}) = \chi_{\beta_{**}}(N\beta(\delta_{\text{ad}})^r)$. To extract

a matching cocycle from $c(w)$ we may simply write down any cocycle that gives the correct character value. We will, however, take time to motivate our construction, as we will use the result later. Namely, we consider the (distinct, symmetric) σ -orbits \mathcal{O} and $w_0\mathcal{O}$ of β_{**} and $w_0\beta_{**}$. The contributions $r_{\pm\mathcal{O}}(w)$ and $r_{\pm w_0\mathcal{O}}(w)$ to $c(w)$ are not cocycles. However, because we use compatible χ -data, $r_{\pm\mathcal{O}}(w)r_{\pm w_0\mathcal{O}}(w)$ is of the form

$$\chi_{\beta_{**}}(u_0(w))^{rN\beta} \chi_{w_0\beta_{**}}(u_0(w))^{rw_0N\beta} = \chi_{\beta_{**}}(u_0(w))^{r(N\beta + w_0N\beta)}.$$

But

$$N\beta + w_0N\beta \equiv 2N\beta \pmod{N\alpha}.$$

We extract $\chi_{\beta_{**}}(u_0(w))^{2rN\beta}$ from $c(w)$. This is a cocycle since $\chi_{\beta_{**}}^2$ is trivial on \mathbb{R}^\times (Lemma 2.5.B of [LS87]). The value of the corresponding character at δ_{ad} is $\chi_{\beta_{**}}(x^2)$, where $x/\bar{x} = N\beta(\delta_{\text{ad}})^r$; see the calculations of Section 3.3 of [LS87]. Since

$$\chi_{\beta_{**}}(x^2) = \chi_{\beta_{**}}(x/\bar{x}.x\bar{x}) = \chi_{\beta_{**}}(x/\bar{x}) = \Lambda_{\pm\mathcal{O}'}(\delta_{\text{ad}}),$$

we have removed an appropriate pair of cocycles from $c'(w)$, $c(w)$.

In the next step of the definition of $a'_+(w)$, $a_+(w)$ we consider the asymmetric orbits \mathcal{O}' not orthogonal to α_{**} for which the σ -orbit \mathcal{O} of $\beta_{**} \in \mathcal{O}'$ is also asymmetric. Then both $r_{\pm\mathcal{O}'}(w)$ and $r_{\pm\mathcal{O}}(w)$ are cocycles. If $\mathcal{O}', \mathcal{O}$ are of the same cardinality (i.e., both consist of a complex root and its conjugate), then the corresponding characters have the same value at δ_{ad} , and so we remove $r_{\pm\mathcal{O}'}(w), r_{\pm\mathcal{O}}(w)$ from $c'(w)$, $c(w)$ respectively. It remains to consider the case that $\sigma'\beta_{**} = \beta_{**}$ and β_{**} is not orthogonal to α_{**} . Then $w_0\beta_{**}$ also has this property, is distinct from β_{**} , and has same σ -orbit as β_{**} . Here we remove both $r_{\pm\mathcal{O}'}(w)$ and $r_{\pm w_0\mathcal{O}'}(w)$ from $c'(w)$, and $r_{\pm\mathcal{O}}(w)$ from $c(w)$. The requirement of s -compatibility that $\chi_{\beta_{**}} = \chi'_{\beta_{**}} \circ Nm$ ensures that the corresponding characters match at δ_{ad} (see Section 3.3 of [LS87]), and so we are done.

The final step in the definition of $a'_+(w)$, $a_+(w)$ is needed only for the case where α_{**} is of type R_2 and $2\alpha_{**}$ is a root of H^\vee , so that $\{\pm\alpha_{**}\}$ satisfies the requirements of case (ii) above. Then $r_{\pm 2\mathcal{O}'}(w)^{-1}$ is a cocycle which we include in $a'_+(w)$, i.e., discard from $c'(w)$, since the method of Section 3.3 of [LS87] shows that the value of the corresponding character at δ_{ad} is 1.

We observe next that $c'_-(w) = c'(w)/a'_+(w)$ has contributions only from orbits which are orthogonal to α_{**} , so that $c'_-(w)$ is fixed by w_0 . Moreover, by construction [LS87], each contribution is the image of an w_0 -invariant in T_{sc}^\vee .

Turning now to $A_-(w)$, we verify that $\varphi_{w_0}(a'_-(w)) = 1$. The cocycle $a_{T_2'}(w)$ takes values in the torus T_2^\vee , which is the quotient of $T_1^\vee \times T^\vee$ by the diagonally embedded torus T_H^\vee . It may be written as the image of the pair $(t_1(w), t(w)^{-1})$ on p. 45 of [KS99]. To compute φ_{w_0} on this image, we may choose an element

$t_H(w)$ of T_H^\vee so that $t_1(w)t_H(w)$ lies in the center of H_1^\vee and then compute $\varphi_{w_0}(t(w)^{-1}t_H(w))$. In this last formula, φ_{w_0} denotes the standard homomorphism $T^\vee \rightarrow T_{sc}^\vee : t \rightarrow w_0(t_{sc})t_{sc}^{-1}$, where t_{sc} has same image as t in T_{ad}^\vee . Then

$$\varphi_{w_0}(a'_-(w)) = \varphi_{w_0}(a_{T'_2}(w).a'_+(w)^{-1}) = \varphi_{w_0}(t(w)^{-1}t_H(w).a'_+(w)^{-1}).$$

Notice that $\varphi_{w_0}(c'_-(w)) = 1$. Our claim now is that

$$\varphi_{w_0}(t(w)^{-1}t_H(w).a'_+(w)^{-1}) = 1.$$

To provide a more explicit description of $t(w)$, and of our choice for $t_H(w)$, we review the construction of $a_{T'_2}(w)$. We fix a Γ -splitting of G^\vee that is preserved by θ^\vee and assume that the endoscopic datum \mathfrak{s} lies in the maximal torus of this splitting (which we identify with T^\vee using our chosen toral data). Then we use the attached Γ -splittings for G_*^\vee and H^\vee . Let w_H denote the action of $1 \times w \in {}^L H$ on H^\vee . Then w_H acts on $T_H^\vee = T_*^\vee$ and thus on $T^\vee = \text{Cent}(T_*^\vee, G^\vee)$ as $\omega(w_H)w_G$, where w_G is the action of $1 \times w \in {}^L G$ (or ${}^L G_*$) and $\omega(w_H)$ lies in the Weyl group of G_*^\vee . Let M_*^\vee be the Levi group in G_*^\vee containing T_*^\vee and with root system consisting of those β_{**} for which $\sigma\beta_{**} = -\beta_{**}$. (Recall we use σ as an abbreviation for σ_T .) Then $\omega(w_H)$ lies in the Weyl group of M_*^\vee , and so we construct $n(\omega(w_H))$ in M_*^\vee acting as $\omega(w_H)$ as in [LS87]. Further, we may find $t_H^1(w)$ in $T_*^\vee \cap (M_*^\vee)_{\text{der}}$ so that $h(w) = t_H^1(w)n(\omega(w_H)) \times w \in {}^L G$ lies in \mathcal{H} (part of the endoscopic data \mathfrak{e}) and acts on H^\vee as w_H . Then for the embedding $\xi_1 : \mathcal{H} \rightarrow {}^L H_1$ (part of the z -pair), we have $\xi_1(h(w)) = z_1(w) \times w$, where $z_1(w)$ is central in H_1^\vee . The embedding $\xi_{T'_1} : {}^L T_1 \rightarrow {}^L H_1$ has the property

$$\begin{aligned} \xi_{T'_1}(1 \times w) &= r'_\mathfrak{s}(w)n_\mathfrak{s}(\omega'_H(w)) \times w \\ &= z_1(w)^{-1}r'_\mathfrak{s}(w).n_\mathfrak{s}(\omega'_H(w)).\xi_1(h(w)) \\ &= z_1(w)^{-1}r'_\mathfrak{s}(w).\xi_1(n_\mathfrak{s}(\omega'_H(w)).h(w)). \end{aligned}$$

Here σ' acts as $\omega'_H(\sigma).\sigma_H$ on T_H^\vee , and $\omega'_H(w) = \omega'_H(\sigma)$ if $w \rightarrow \sigma$ under $W_{\mathbb{R}} \rightarrow \Gamma$, while $\omega'_H(w) = 1$ if $w \rightarrow 1$. Notice that $\omega'(\sigma)$ also lies in the Weyl group of M_*^\vee (although we construct $n_\mathfrak{s}(\omega'_H(w))$ in H^\vee). We have to compare $\xi_{T'_1}$ with the embedding $\xi_{T'_*} : {}^L T'_* \rightarrow {}^L G_*$ which extends naturally to $\xi_{T'} : {}^L T' \rightarrow {}^L G$. Write the action of σ' on $T_H^\vee = T_*^\vee$ as $\omega'_G(\sigma).\sigma_G$. Construct $n(\omega'_G(w))$ in M_*^\vee , and notice that $\omega'_G(w) = \omega'_H(w).\omega(w_H)$. Then

$$\begin{aligned} \xi_{T'_*}(1 \times w) &= r'_1(w)n(\omega'_G(w)) \times w \\ &= r'_1(w)\omega'_H(w)(t_H^1)^{-1}n(\omega'_G(\sigma))n(\omega(w_H))^{-1}h(w). \end{aligned}$$

We claim that

$$n(\omega'_G(\sigma))n(\omega(w_H))^{-1} = t_H^2(w)n_\mathfrak{s}(\omega'_H(w)),$$

where $t_H^2(w) \in T_*^\vee \cap (M_*^\vee)_{\text{der}}$. To prove this, we compare the left side to $n(\omega'_H(w))$ using Lemma 2.1.A of [LS87]. For the right side, there is a routine

generalization of Lemma 4.3.A of [LS90] to the twisted case that allows us to compare $n_{\mathfrak{s}}(\omega'_H(w))$, an element in the Levi group of H^\vee corresponding to the appropriate multiples of roots in M_*^\vee , to $n(\omega'_H(w))$, an element of M_*^\vee . The claim then follows. Thus

$$\xi_{T_*'}(1 \times w) = r'_1(w)\omega'_H(w)(t_H^1)^{-1}t_H^2(w).n_{\mathfrak{s}}(\omega'_H(w))h(w).$$

Turning now to $a_{T_2'}(w)$, we set

$$t_1(w) = z_1(w)^{-1}r'_s(w), t_H(w) = r'_s(w)^{-1}$$

and

$$t(w) = r'_1(w)\omega'_H(w)(t_H^1)^{-1}t_H^2(w).$$

Then

$$\varphi_{w_0}(t(w)^{-1}t_H(w).a'_+(w)^{-1}) = \varphi_{w_0}(\omega'_H(w)(t_H^1)^{-1}t_H^2(w)) = 1$$

since

$$\omega'_H(w)(t_H^1)^{-1}t_H^2(w) \in T_*^\vee \cap (M_*^\vee)_{\text{der}}$$

and φ_{w_0} is trivial on $T_*^\vee \cap (M_*^\vee)_{\text{der}}$.

Our last step is to examine the second component

$$\hat{\tau}(w_0, \sigma').\hat{b}(w_0)^{-1}.w_0(c'_-(w)).c_-(w)^{-1}$$

of $A_-(w)$. Consider

$$w_0(c'_-(w)).c_-(w)^{-1} = c'_-(w)c_-(w)^{-1}.$$

If \mathcal{O}' is orthogonal to α_{**} , then \mathcal{O}' is also a σ -orbit \mathcal{O} , and $r_{\pm\mathcal{O}'}(w) = r_{\pm\mathcal{O}}(w)$. Thus all that remains in $w_0(c'_-(w)).c_-(w)^{-1}$ is a term in $(\mathbb{C}^\times)^{rN\alpha}$ and the term $s_{p/p_0}(w)$ for the action σ . We compare $s_{p/p_0}(w)$ with $\hat{\tau}(w_0, \sigma').\hat{b}(w_0)^{-1}$. Recall our assumption that if $\beta_{**} > 0$, then $\sigma'\beta_{**} > 0$ unless $\sigma'\beta_{**} = -\beta_{**}$. Then $\beta_{**} > 0$ and $\sigma\beta_{**} > 0$ requires $w_0\beta_{**} = \sigma'\sigma\beta_{**} > 0$. Thus the sum defining $\hat{\tau}(w_0, \sigma')$ is empty, so that $\hat{\tau}(w_0, \sigma') = 1$. Next, we use a routine generalization of Lemma 4.3.B in [LS90]. This shows that the term $\hat{b}(w_0)$ is a product of an element of order two and an element in $(\mathbb{C}^\times)^{rN\alpha}$. The element of order two is of the form $\prod_{\beta_{**}} (-1)^{\beta_{**}^\vee}$, where the product is over representatives for the pairs $\{\beta_{**}, -w_0\beta_{**}\}$ with the property that $\beta_{**} > 0$ and $-w_0\beta_{**} > 0$. If we consider just those pairs where $\beta_{**}, -w_0\beta_{**}$ are also complex roots (if one is, the other is), then we obtain $s_{p/p_0}(w)$. See Section 2.5 of [LS87], and cancel terms for G^\vee, H^\vee appropriately. Assume now that β_{**} is imaginary. Then $(-1)^{\beta_{**}^\vee}$ is a Galois cocycle which inflates to a cocycle of $W_{\mathbb{C}/\mathbb{R}}$ of order at most two. To evaluate the corresponding character at δ_{ad} , we use the method of Section 3.2 of [LS87] to reduce the calculation to evaluation at the element $N\beta(\delta_{\text{ad}})^r$ of a character of order at most two on the real points of a 1-dimensional torus $T^{\beta_{**}}$. Since $T^{\beta_{**}}(\mathbb{R})$ is compact, the (cocycle and) character must be trivial. Notice that here the canonical constructions of [LS87] have allowed us to avoid

the more complicated setting in Theorem 6.1.1 of [She81b], where case-by-case computations were needed. Thus we may discard the pairs $\{\beta_{**}, -w_0\beta_{**}\}$ for which β_{**} , and hence also $-w_0\beta_{**}$, is imaginary. Since no real roots contribute, we conclude that, after the discard, $\widehat{\tau}(w_0, \sigma') \cdot \widehat{b}(w_0)^{-1} \cdot w_0(c'_-(w)) \cdot c_-(w)^{-1}$ is a cocycle with values in $(\mathbb{C}^\times)^{rN\alpha}$. It remains to evaluate the corresponding character at δ_{ad} . We again use the method of Section 3.2 of [LS87] to reduce this to the value of a character on a 1-dimensional torus $T^{\alpha_{**}}$ at the element $N\alpha(\delta_{\text{ad}})^r$. If α_{**} is of type R_1 , then $N\alpha(\delta_{\text{ad}})^r = N\alpha(\delta_{\text{ad}}) = 1$, and if α_{**} is of type R_2 , then $N\alpha(\delta_{\text{ad}})^r = N\alpha(\delta_{\text{ad}})^2 = (\pm 1)^2 = 1$ also. Thus the value is 1, and we have finished the proof of Lemma 9.3. Notice that we could have based our last calculation on the coroot of the root α_1 of H_1 in place of the reduced α_{**} . Then we arrive at the evaluation of a character at $\alpha_1(\gamma_0) = 1$. \square

LEMMA 9.5. *For an s -compatible data set and toral descent data at γ_0 , we have*

$$\Delta_I(\gamma_0(Y_1), \delta_0(Y)) = \Delta_I(\gamma_0(Y'_1), \delta_0(Y'))$$

for all $Y_1 \in \mathfrak{t}_1(\mathbb{R})$ and $Y'_1 \in \mathfrak{t}'_1(\mathbb{R})$.

Proof. Given our choices, the sign Δ_I depends only on the torus T_1 or T'_1 to which the first argument, $\gamma_0(Y_1)$ or $\gamma_0(Y'_1)$, belongs. The lemma asserts that not even that matters. There are two ways we can argue this. The first is to observe that the proof in [She82], [LS90] of geometric transfer (with the transfer factors of [LS87]) for untwisted endoscopy avoids Lemma 9.5, using instead regular unipotent analysis and the local hypothesis. We deduce Lemma 9.5 in the untwisted case from the cited proof together with Corollaries 9.2 and 9.4 above: if transfer exists and all terms but Δ_I are known to match correctly then Δ_I must match correctly. We then prove Lemma 9.5 in the general case with the observation from Section 4.2 of [KS99] that Δ_I for the twisted case may be interpreted as Δ_I for a case of standard endoscopy.

Our second proof for Lemma 9.5 is a direct argument, allowing us to complete a proof for geometric transfer that works as well for, rather than assumes, standard endoscopy. The starting point is the observation cited above for twisted Δ_I . We consider standard endoscopy for the quasi-split group $G^{\theta_{\text{sc}}^*} = (G_{\text{sc}}^*)^{\theta_{\text{sc}}^*}$ (denoted G^x in [KS99]) and the datum $\mathfrak{s}_{T, \theta}$ defined on p. 32 of [KS99]. The two maximal tori $T^{\theta_{\text{sc}}^*}, T'^{\theta_{\text{sc}}^*}$ in $G_{\text{sc}}^{\theta_{\text{sc}}^*}$ are norm (image) groups for the endoscopic group J . Our toral data and a -data provide data for this setting also. Write α_* for the multiple of α_0 that is a root of $T^{\theta_{\text{sc}}^*}$ and define α'_* similarly. Recall that the inverse Cayley transform t^* carries α'_* to α_* . Pick a $G^{\theta_{\text{sc}}^*}$ -semiregular element ε of $T'^{\theta_{\text{sc}}^*}(\mathbb{R})$ with image ε_J in $J(\mathbb{R})$. Then we make an endoscopic descent around the pair $(\varepsilon, \varepsilon_J)$ as in [LS90]. By construction, the connected centralizers of $\varepsilon, \varepsilon_J$ are isomorphic over \mathbb{R} , so that the base endoscopy is trivial up to passage to z -extensions. In particular, each

Δ_I term is trivial. Our setting satisfies the requirements for the comparison formulas of Section 3.3 of [LS90], including the condition (3.3.2). The formula of Lemma 9.5 is the same as the corresponding formula for $G_{sc}^{\theta^*}$ relative to the tori $T^{\theta^*}_{sc}, T'^{\theta^*}_{sc}$. Thus it is enough to show that the quotient of the two terms in the formula divided by the (trivial) quotient for the centralizers, or the quotient of the terms Θ_I of [LS90] for $T'^{\theta^*}_{sc}$ and $T^{\theta^*}_{sc}$, is trivial. Lemma 3.3.D of [LS90] describes a class v in $H^1(\Gamma, T^{\theta^*}_{sc})$ with which we may pair $\mathfrak{s}_{T,\theta}$, by the Tate-Nakayama pairing, to obtain this quotient of the Θ_I . It remains thus to examine v (which we will write as v_*) and conclude that, because of our particular choice of a -data, this class is represented by a cocycle $(-1)^{\epsilon\alpha^{\vee}}$, where $\epsilon \in \{0, 1\}$. Since α^{\vee}_* is a root of J^{\vee} , the pairing yields 1, and the lemma is then proved.

We use, just for this paragraph, α to denote a reduced root of $T^{\theta^*}_{sc}$ different from $\pm\alpha_*$. (We argue in $G_{sc}^{\theta^*}$ with no reference to H^{\vee} or the endoscopic data.) Identify $T'^{\theta^*}_{sc}$ with $T^{\theta^*}_{sc}$ via t^* , and write σ for the Galois action on $T^{\theta^*}_{sc}$, σ' for the transport of the Galois action on $T'^{\theta^*}_{sc}$, and a'_α for the a -data for $T'^{\theta^*}_{sc}$. Then $\sigma = w_0\sigma'$, where w_0 is the Weyl reflection for α_* , and

$$v_*(\sigma) = \tau(w_0, \sigma').b(w_0)^{-1}.w_0(y'(\sigma)).y(\sigma)^{-1}.$$

Here

$$\tau(w_0, \sigma') = \prod_{\alpha>0, w_0\alpha<0, \sigma\alpha>0} (-1)^{\alpha^{\vee}}.$$

Up to multiplication by an element of $(\mathbb{C}^\times)^{\alpha^{\vee}_*}$, the term $b(w_0)$ is $\prod (-1)^{\alpha^{\vee}}$, where the product is over representatives for pairs $\{\alpha, -w_0\alpha\}$ such that $\alpha > 0$, $w_0\alpha < 0$ (see Lemma 4.3.A of [LS90]). Here the order on the roots is obtained by transport of that determined by our choice of an \mathbb{R} -splitting. The choice of splitting does not affect the quotient of Δ_I terms, and there is no harm in our assumption that if $\alpha > 0$ and $\sigma'\alpha \neq -\alpha$, then $\sigma'\alpha > 0$. (Or see Lemma 2.3.A of [LS87], and note that the assumption (3.3.2) of [LS90] is retained.) Finally,

$$y'(\sigma) = \prod_{\alpha>0, \sigma'\alpha<0} (a'_\alpha)^{\alpha^{\vee}}$$

and

$$y(\sigma) = \prod_{\alpha>0, \sigma\alpha<0} (a_\alpha)^{\alpha^{\vee}}.$$

Suppose $\alpha > 0$, $\sigma'\alpha < 0$, so that α contributes to $w_0(y'(\sigma))$. Then $\alpha = -\sigma'\alpha$ and $w_0\alpha = \alpha = -\sigma'\alpha = -\sigma\alpha$, so that α is imaginary for both $T'^{\theta^*}_{sc}$ and $T^{\theta^*}_{sc}$. By t^* -compatibility of our a -data, we have $a'_\alpha = a_\alpha$, and the contribution from α to $w_0(y'(\sigma))$ cancels that to $y(\sigma)$. There are two remaining types of contribution to $y(\sigma)$. The first is for $\alpha > 0$, such that $\alpha = -\sigma\alpha$ and $w_0\alpha \neq \alpha$. Then $w_0\alpha = -w_0\sigma\alpha = -\sigma'\alpha < 0$ since $-\sigma'\alpha \neq \alpha$. Thus we also have $-w_0\alpha > 0$,

and $-w_0\alpha$ is of same type as α . The contribution to $y(\sigma)$ from $\{\alpha, -w_0\alpha\}$ is

$$(a_\alpha)^{\alpha^\vee} (a_{-w_0\alpha})^{-w_0\alpha^\vee} = (a_\alpha)^{\alpha^\vee - w_0\alpha^\vee} (-1)^{-w_0\alpha^\vee}$$

since $a_{-w_0\alpha} = -a_{w_0\alpha} = -a_\alpha$. The first term in the product lies in $(\mathbb{C}^\times)^{\alpha^\vee}$ and the second cancels with a term in $b(w_0)$ up to multiplication by an element of $(\mathbb{C}^\times)^{\alpha^\vee}$. The second type of contribution to $y(\sigma)$ is from $\alpha > 0$ such that $\sigma\alpha < 0$ and $\sigma\alpha \neq -\alpha$. Then each of α and $-\sigma\alpha$ contributes and their joint contribution is

$$(a_\alpha)^{\alpha^\vee} (a_{-\sigma\alpha})^{-\sigma\alpha^\vee} = (a_\alpha)^{\alpha^\vee} (\overline{a_\alpha})^{-\sigma\alpha^\vee} (-1)^{-\sigma\alpha^\vee}.$$

Since $(a_\alpha)^{\alpha^\vee} (\overline{a_\alpha})^{-\sigma\alpha^\vee}$ is a coboundary, we may ignore it. Let $\beta = -\sigma\alpha$. Then $\beta > 0, \sigma\beta < 0$. Also $w_0\beta = -\sigma'\alpha < 0$ since $\alpha > 0$ and $\sigma'\alpha \neq -\alpha$. Thus $(-1)^{-\sigma\alpha^\vee} = (-1)^{\beta^\vee}$ cancels with the corresponding term in $b(w_0)$, and so we conclude that, up to coboundaries, the cocycle $v_*(\sigma)$ lies in $(\mathbb{C}^\times)^{\alpha^\vee}$. The lemma now follows. \square

Finally, the following equalities will be used in assembling the jump formulas in the next section. The terms were introduced in Section 7.

LEMMA 9.6. *Under the assumptions of the present section, we have*

$$\begin{aligned} &\langle \text{inv}(\delta_0(Y), \delta_0(Y)(w)), \kappa_{\delta_0(Y)} \rangle \\ &= \langle \text{inv}(\delta_0(Y), \delta_0(Y)(ww_0)), \kappa_{\delta_0(Y)} \rangle \\ &= \langle \text{inv}(\delta_0(Y'), \delta_0(Y')(w')), \kappa_{\delta_0(Y')} \rangle. \end{aligned}$$

Proof. The representatives w', w were defined in the paragraph before Lemma 7.2, and w_0 lies in $G_{\delta_0}^\theta$. Write the three inv terms in the statement as $\text{inv}(w), \text{inv}(ww_0), \text{inv}(w')$. To define $\text{inv}(w)$ we start with the Galois cocycle $\sigma(w)w^{-1}$ in the maximal torus $A^{\delta_0} = \text{Cent}(T^{\delta_0}, G)$ of G . (Earlier we used the notation T^\dagger for A^{δ_0} .) Notice that A^{δ_0} is preserved by $\theta_0 = \text{Int}(\delta_0) \circ \theta$ and $\text{Int}(\delta_0(Y)) \circ \theta$ acts as θ_0 on A^{δ_0} . Let $A_{\text{sc}}^{\delta_0}$ be the corresponding torus in G_{sc} . Then we factor w in the usual manner, as the product of the image of an element w_{sc} of G_{sc} and a central element z . The pair $(\sigma(w_{\text{sc}})w_{\text{sc}}^{-1}, (\theta_0 - 1)z)$ represents $\text{inv}(w)$, an element of $H^1(\Gamma, A_{\text{sc}}^{\delta_0} \xrightarrow{\varphi} B^{\delta_0})$. Here B^{δ_0} is the image of A^{δ_0} under $\theta_0 - 1$ and φ is the composition of $\theta_0 - 1$ with the projection $A_{\text{sc}}^{\delta_0} \rightarrow A^{\delta_0}$. We have arranged that $\sigma(ww_0)(ww_0)^{-1}$ coincides with $\sigma(w)w^{-1} \cdot (-1)^{\alpha_0^\vee}$ up to coboundaries in $A^{\delta_0} \cap (G_{\delta_0}^\theta)_{\text{der}} = T^{\delta_0} \cap (G_{\delta_0}^\theta)_{\text{der}}$. Thus we can factor the corresponding hypercocycle as $(\sigma(w_{\text{sc}})w_{\text{sc}}^{-1}, (\theta_0 - 1)z) \cdot ((-1)^{\alpha_0^\vee}, 1)$. The usual argument (see the proof of Lemma 9.3) shows that the second term in the statement of the present lemma is $\alpha_0^\vee(\mathfrak{s}_T)$ times the first. Since $\alpha_0^\vee(\mathfrak{s}_T) = 1$ (see the proof of Lemma 9.3 again), we are done with the first equality.

Our choices ensure that $\text{inv}(w)$ is represented by a hypercocycle $(a_{\text{sc}}(\sigma), (\theta_0 - 1)z)$ and $\text{inv}(w')$ is represented by $(s(a_{\text{sc}}(\sigma)), (\theta_0 - 1)z)$. Here, recall that

s is a Cayley transform in $(G_{\delta_0}^\theta)_{\text{sc}}$. Since we have also to analyze the dual data, we use our chosen toral data to pass from G to G^* . Then in place of $H^1(\Gamma, A_{\text{sc}}^{\delta_0} \xrightarrow{\varphi} B^{\delta_0})$ we consider $H^1(\Gamma, T_{\text{sc}} \rightarrow (\theta^* - 1)T)$, etc., and we identify T' with T over \mathbb{C} using the inverse Cayley transform $t^* = (s^*)^{-1}$. Consider the pair $(\text{inv}(w)^{-1}, \text{inv}(w'))$ in

$$H^1(\Gamma, T_{\text{sc}} \times T'_{\text{sc}} \rightarrow (\theta^* - 1)(T \times T')).$$

It is represented by

$$((t_{\text{sc}}(\sigma), (\theta^* - 1)z)^{-1}, (t_{\text{sc}}(\sigma), (\theta^* - 1)z)),$$

where $t_{\text{sc}}(\sigma)$ is the image of $a_{\text{sc}}(\sigma)$ under our identification of A^{δ_0} with T . To prove that the (equal) first and second terms in the statement of the lemma coincide with the third, we show that $(\text{inv}(w)^{-1}, \text{inv}(w'))$ pairs trivially with the class in

$$H^1(W_{\mathbb{R}}, [(\theta^* - 1)(T \times T')]^\vee \rightarrow T_{\text{ad}}^\vee \times (T'^\vee)_{\text{ad}})$$

represented by

$$((b_T(w)^{-1}, \mathfrak{s}_{\text{ad}}), (b_{T'}(w)^{-1}, \mathfrak{s}_{\text{ad}})),$$

where $b_T, b_{T'}$ are as constructed on p. 55 of [KS99]. (We will describe them in detail shortly.) Recall

$$S = S(T, T') = T \times T' / \{(z^{-1}, z) : z \in Z(G^*)\}.$$

The projection $(\theta^* - 1)(T \times T') \rightarrow (\theta^* - 1)S$ determines a map on hypercohomology groups under which the image of $(\text{inv}(w)^{-1}, \text{inv}(w'))$ is represented (in the obvious manner) by $((t_{\text{sc}}(\sigma)^{-1}, 1), (t_{\text{sc}}(\sigma), 1))$. Thus by functoriality of the pairing, it is enough to show that $((b_T(w)^{-1}, \mathfrak{s}_{\text{ad}}), (b_{T'}(w)^{-1}, \mathfrak{s}_{\text{ad}}))$ represents a class in the image of

$$H^1(W_{\mathbb{R}}, [(\theta^* - 1)S]^\vee \rightarrow T_{\text{ad}}^\vee \times (T'^\vee)_{\text{ad}})$$

under the (dual) map on dual hypercohomology groups. Thus it is enough to show that the cocycle $(b_T(w), b_{T'}(w))$ in $[(\theta^* - 1)(T \times T')]^\vee$ lies in the subtorus $[(\theta^* - 1)S]^\vee$.

Recall the cocycle $(a_{T_2}(w), a_{T'_2}(w))$ of $W_{\mathbb{R}}$ in $T_2^\vee \times T_2^\vee$ from the construction of Δ_{III} ; see the proof of Lemma 9.3 and p. 45 of [KS99]. Also, the torus $[(\theta^* - 1)(T \times T')]^\vee$ may be identified with $T_2^\vee \times T_2^\vee / T_1^\vee \times T_1^\vee$ (see p. 55 of [KS99]). Then $(b_T(w), b_{T'}(w))$ is, by definition, the image of $(a_{T_2}(w), a_{T'_2}(w))$ under the natural projection

$$\text{proj} : T_2^\vee \times T_2^\vee \rightarrow T_2^\vee \times T_2^\vee / T_1^\vee \times T_1^\vee.$$

By construction, $(a_{T_2}(w), a_{T'_2}(w))$ lies in S_1^\vee (identified as a subtorus of $T_2^\vee \times T_2^\vee$). We denote by θ_2 the extension of θ^* to T_2, T'_2 (p. 42 of [KS99]). The torus $(\theta^* - 1)T$ may be identified as the (isomorphic) image of $(\theta_2 - 1)T_2$ under the projection $T_2 \rightarrow T$, and so $(\theta^* - 1)S$ may be identified with $(\theta_2 - 1)S_1$ under $S_1 \rightarrow S$. Since $[(\theta_2 - 1)S_1]^\vee$ coincides with the image of S_1^\vee under proj , we are done. □

10. Proof of Theorem 5.1 and extension to derivatives

To complete the proof of Theorem 5.1 we return to the formulas of Sections 7 and 8, and we combine them with the results of Section 9. We have only to consider the case that γ_0 is both a $T_1^{s_1}$ -norm and a T_1 -norm, and maintain the toral descent data attached to γ_0 , along with the s -compatible data sets, in Section 7. Write $\Phi_1(\gamma_\nu)$ as

$$\begin{aligned} & \left| \det(\text{Ad}(\gamma_\nu) - I)_{\mathfrak{h}_1/\mathfrak{t}_1} \right|^{1/2} \sum_w \Delta(\gamma_\nu, \delta_\nu(w)) O^{\theta, \varpi}(\delta_\nu(w), f) \\ &= \Delta_I(\gamma_\nu) \Delta_{II}(\gamma_\nu) \Delta_{III}(\gamma_\nu, \delta_\nu; \bar{\gamma}, \bar{\delta}) \Delta_{IV, \text{num}}(\delta_\nu) \\ & \quad \times \sum_w \langle \text{inv}(\delta_\nu, \delta_\nu(w)), \kappa_{\delta_\nu} \rangle O^{\theta, \varpi}(\delta_\nu(w), f). \end{aligned}$$

To pass to the transform $\Psi_{a, \chi}(\gamma_\nu)$, we simply replace $\Delta_{II}(\gamma_\nu)$ by $\Delta_{II, \text{num}}(\gamma_\nu)$. Without changing notation, we drop the terms for those classes in $\mathfrak{D}_\theta(T^{\delta_0})$ with no representative w for which $w\alpha_0 = \pm\alpha_0$. Since

$$\langle \text{inv}(\delta_\nu, \delta_\nu(w)), \kappa_{\delta_\nu} \rangle = \langle \text{inv}(\delta_\nu, \delta_\nu(ww_0)), \kappa_{\delta_\nu} \rangle$$

(Lemma 9.6), we may then replace the sum by a sum over representatives w for $\mathfrak{D}_\theta(\alpha_0)$ and examine

$$\begin{aligned} & \Delta_I(\gamma_\nu) \cdot \Delta_{II, \text{num}}(\gamma_\nu) \Delta_{IV, \text{num}}(\delta_\nu) \Delta_{\alpha_0}(\delta_\nu)^{-1} \cdot \Delta_{III}(\gamma_\nu, \delta_\nu; \bar{\gamma}, \bar{\delta}) \\ & \quad \times (2/d(\alpha_0)) \sum_w \langle \text{inv}(\delta_\nu, \delta_\nu(w)), \kappa_{\delta_\nu} \rangle \Delta_{\alpha_0}(\delta_\nu) O^{\theta, \varpi}(\delta_\nu(w), f). \end{aligned}$$

On the other hand,

$$\begin{aligned} \Psi_{a', \chi'}(\gamma'_\nu) &= \Delta_I(\gamma'_\nu) \cdot \Delta_{II, \text{num}}(\gamma'_\nu) \Delta_{IV, \text{num}}(\delta'_\nu) \Delta_{\pm\alpha'_0}(\delta'_\nu)^{-1} \cdot \Delta_{III}(\gamma'_\nu, \delta'_\nu; \bar{\gamma}, \bar{\delta}) \\ & \quad \times \sum_{w'} \langle \text{inv}(\delta'_\nu, \delta'_\nu(w')), \kappa_{\delta'_\nu} \rangle \Delta_{\pm\alpha'_0}(\delta'_\nu) O^{\theta, \varpi}(\delta'_\nu(w'), f), \end{aligned}$$

where the summation is over representatives w' for the elements of $\mathfrak{D}_\theta(T'^{\delta_0})$. From Lemma 8.1, Corollaries 9.2, 9.4, and Lemmas 9.5, 9.6, we conclude that

$$\lim_{\nu \rightarrow 0^+} \Psi_{a, \chi}(\gamma_\nu) - \lim_{\nu \rightarrow 0^-} \Psi_{a, \chi}(\gamma_\nu) = 2 \lim_{\nu \rightarrow 0} \Psi_{a', \chi'}(\gamma'_\nu).$$

For the final step in the proof of Theorem 5.1, we notice that the Weyl reflection w_1 for α_1 provides a stable conjugation of $\gamma_{-\nu}$ with γ_ν . Since Φ_1 is invariant under stable conjugacy, it is enough to examine the factor

$$\prod_{\mathcal{O}_1} \chi_{\beta_1} \left(\frac{(\beta_1(\gamma_1) - 1)}{a_{\beta_1}} \right),$$

where γ_1 is regular in $T_1(\mathbb{R})$. There is no change in the total contribution from the orbits $\mathcal{O}_1 \neq \{\pm\alpha_1\}$ when γ_1 is replaced by $\gamma_1^{w_1}$ since the definition of s -compatible data ensures that $\chi_{w_1\beta_1} = \chi_{\beta_1}$ and $a_{w_1\beta_1} = a_{\beta_1}$; the contribution

from the orbit of β_1 is interchanged with that from the orbit of $w_1\beta_1$. For the case $\mathcal{O}_1 = \{\pm\alpha_1\}$, we have

$$\chi_{\alpha_1} \left(\frac{(\alpha_1(\gamma_1^{w_1}) - 1)}{a_{\alpha_1}} \right) = -\chi_{\alpha_1}(\alpha_1(\gamma_1))^{-1} \chi_{\alpha_1} \left(\frac{(\alpha_1(\gamma_1) - 1)}{a_{\alpha_1}} \right).$$

Since $\chi_{\alpha_1}(\alpha_1(\gamma_0)) = \chi_{\alpha_1}(1) = 1$,

$$\lim_{\nu \rightarrow 0^-} \Psi_{a,\chi}(\gamma_\nu) = - \lim_{\nu \rightarrow 0^+} \Psi_{a,\chi}(\gamma_\nu),$$

and so we are done with the proof of Theorem 5.1. □

To consider limit formulas for derivatives, let $\mathcal{S}(t_1)$ denote the symmetric algebra of t_1 . Denote by $D \rightarrow \widehat{D}$ the automorphism of $\mathcal{S}(t_1)$ determined by the map $Y_1 \rightarrow Y_1 - n_1\alpha_1(Y_1)I$ of t_1 into $\mathcal{S}(t_1)$, where $2n_1$ is the *odd* integer given by $\chi_{\alpha_1}(z) = (z/\bar{z})^{n_1} = (z/|z|)^{2n_1}$.

For $\gamma_1 = \gamma_0 \exp Y_1$ near γ_0 , define $\chi_{\alpha_1}(\alpha_1(\gamma_1))^{1/2}$ to be $\chi_{\alpha_1}(\exp \frac{1}{2}\alpha_1(Y_1))$. Then the function (germ)

$$\widehat{\Psi}_{a,\chi}(\gamma_1) = \chi_{\alpha_1}(\alpha_1(\gamma_1))^{-1/2} \Psi_{a,\chi}(\gamma_1)$$

is defined for G -regular γ_1 near γ_0 in $T_1(\mathbb{R})$. (Recall the smooth extension from strongly G -regular elements to all G -regular elements in Section 7.) Moreover, this function is odd:

$$\widehat{\Psi}_{a,\chi}(\gamma_1^{w_1}) = -\widehat{\Psi}_{a,\chi}(\gamma_1).$$

LEMMA 10.1. *For all $D \in \mathcal{S}(t_1)$, we have that both $\lim_{\nu \rightarrow 0^-} \widehat{D}\Psi_{a,\chi}(\gamma_\nu)$ and $\lim_{\nu \rightarrow 0^+} \widehat{D}\Psi_{a,\chi}(\gamma_\nu)$ exist. If $D^{w_1} = -D$, then*

$$\lim_{\nu \rightarrow 0^-} \widehat{D}\Psi_{a,\chi}(\gamma_\nu) = \lim_{\nu \rightarrow 0^+} \widehat{D}\Psi_{a,\chi}(\gamma_\nu),$$

and if $D^{w_1} = D$, then

$$\lim_{\nu \rightarrow 0^-} \widehat{D}\Psi_{a,\chi}(\gamma_\nu) = - \lim_{\nu \rightarrow 0^+} \widehat{D}\Psi_{a,\chi}(\gamma_\nu).$$

Proof. Existence of each limit in the statement follows from the basic estimates (see the appendix). The twist \widehat{D} of D was defined expressly to obtain the property

$$\widehat{D}\Psi_{a,\chi}(\gamma_0 \exp Y_1) = \chi_{\alpha_1} \left(\exp \frac{1}{2}\alpha_1(Y_1) \right) \cdot D\widehat{\Psi}_{a,\chi}(\gamma_0 \exp Y_1).$$

Thus we have

$$\lim_{\nu \rightarrow 0^\pm} \widehat{D}\Psi_{a,\chi}(\gamma_\nu) = \lim_{\nu \rightarrow 0^\pm} D\widehat{\Psi}_{a,\chi}(\gamma_\nu).$$

The desired equations are then immediate from the oddness of $\widehat{\Psi}_{a,\chi}$. □

We may choose the χ' -datum $\chi_{\alpha'_1}$ nontrivial. Because α'_1 is real, we define $\widehat{\Psi}_{a',\chi'}$ by

$$\widehat{\Psi}_{a',\chi'}(\gamma'_1) = \chi_{\alpha'_1}(\alpha'_1(\gamma'_1))^{-1} \Psi_{a',\chi'}(\gamma'_1)$$

for G -regular γ'_1 near γ_0 in $T_1(\mathbb{R})$. The Cayley transform s_1 provides us with an isomorphism $D \rightarrow D'$ from $\mathcal{S}(\mathfrak{t}_1)$ to $\mathcal{S}(\mathfrak{t}'_1)$. We write \widehat{D}' for the image of D' under the automorphism given by $Y'_1 \rightarrow Y'_1 - z\alpha'_1(Y'_1)I$, where z is the complex number given by $\chi_{\alpha'_1}(x) = (\text{sgn } x)^\varepsilon |x|^z$. Then for all $D \in \mathcal{S}(\mathfrak{t}_1)$, we have that

$$\lim_{\nu \rightarrow 0} \widehat{D}'\Psi_{a',\chi'}(\gamma'_\nu) = \lim_{\nu \rightarrow 0} D'\widehat{\Psi}_{a',\chi'}(\gamma'_\nu)$$

exists.

LEMMA 10.2. *If $D^{w_1} = D$ then, for any s -compatible data set,*

$$\lim_{\nu \rightarrow 0^+} \widehat{D}\Psi_{a,\chi}(\gamma_\nu) = \lim_{\nu \rightarrow 0} \widehat{D}'\Psi_{a',\chi'}(\gamma'_\nu).$$

Proof. For this we return to the formulas obtained by descent in Section 8 and use Harish-Chandra descent for operators in the center of the universal enveloping algebra of the complex Lie algebra of G as well, extending the arguments for Proposition 4.5 of [She79a] via results of Bouaziz (see Theorem 2.4.1 of [Bou87]). The formula then follows by repeating the steps at the start of this section. \square

This concludes then our extension of Theorem 5.1 to derivatives of $\Psi_{a,\chi}$. The extension applies, in particular, to the setting of Theorem 4.2.

11. Completion of proof of the main theorem

We recall once again that if

$$S(\gamma_1) = \sum_{\delta, \theta\text{-conj}} \Delta(\gamma_1, \delta) O^{\theta, \varpi}(\delta, f),$$

then we have the normalized integral

$$\Phi_1(\gamma_1) = \left| \det(\text{Ad}(\gamma_1) - I)_{\mathfrak{h}_1/\mathfrak{t}_1} \right|^{1/2} S(\gamma_1)$$

and the transform

$$\Psi_{a,\chi}(\gamma_1) = \Delta_{a,\chi}(\gamma_1) S(\gamma_1).$$

Recall also that $S(\gamma_1)$, defined initially for strongly G -regular elements γ_1 , was extended smoothly to all G -regular elements. Next, we extend S to a smooth function around all regular elements in $T_1(\mathbb{R})$. Since $\Delta_{a,\chi}$ is nonvanishing and thus smooth on the regular set in $T_1(\mathbb{R})$, we may replace S by $\Psi_{a,\chi}$ for this extension.

Assume that γ_0 is regular in $T_1(\mathbb{R})$, so that $(H_1)_{\gamma_0} = T_1$. If γ_0 is not a norm, then $\Psi_{a,\chi}(\gamma_1) = 0$ for G -regular γ_1 near γ_0 in $T_1(\mathbb{R})$ by Lemma 6.1, and so S extends trivially. Suppose now that γ_0 is a T_1 -norm of $\delta_0 \in G(\mathbb{R})$. We consider the case that δ_0 is θ -semiregular, by which we mean that $G_{\delta_0}^\theta$ is of type A_1 . As before, we denote by $\pm\alpha_0$ the roots of T^{δ_0} in $G_{\delta_0}^\theta$. If the root α_0 is real or totally compact, then we follow our earlier descent arguments (and include

derivatives) to see that $\Psi_{a,\chi}$ extends smoothly around γ_0 . Suppose then that α_0 is imaginary and not totally compact. By passage to a stable θ -conjugate of δ_0 , we may assume that α_0 itself is noncompact. Again we rely on the earlier descent arguments, except that Lemma 9.6 is replaced by the following.

LEMMA 11.1. *In the present setting, we have $\alpha_0^\vee(\mathfrak{s}_T) = -1$, and then*

$$\langle \text{inv}(\delta_0(Y), \delta_0(Y)(w)), \kappa_{\delta_0(Y)} \rangle = - \langle \text{inv}(\delta_0(Y), \delta_0(Y)(ww_0)), \kappa_{\delta_0(Y)} \rangle.$$

Proof. Since $\sigma\alpha_0 = -\alpha_0$, we also have that $\sigma(N\alpha^\vee) = -N\alpha^\vee$, and then $N\alpha^\vee(\mathfrak{s}_T)^2 = 1$ since \mathfrak{s}_T is Γ -invariant. Suppose α is of type R_2 . If $N\alpha^\vee(\mathfrak{s}_T) = 1$, then α_1 is a root of H_1 and $\alpha_1(\gamma_0) = N\alpha(\delta_0)^2 = 1$ contradicting the regularity of γ_0 . Thus $N\alpha^\vee(\mathfrak{s}_T) = -1$ is the only possibility. In fact, then the coroot β_1 of $2(\alpha^\vee)_{\text{res}}$ is a root of H_1 and $\beta_1(\gamma_0) = N\beta(\delta_0) = N\alpha(\delta_0) = \pm 1$. Since $\beta_1(\gamma_0) \neq 1$, we must have $N\alpha(\delta_0) = -1$, a contradiction since α is of type R_2 . We conclude that α cannot be of type R_2 . Suppose α is of type R_3 . If $N\alpha^\vee(\mathfrak{s}_T) = 1$, then β_1 is a root of H_1 , where now β_1 denotes the coroot of $\frac{1}{2}(\alpha^\vee)_{\text{res}}$. This implies $\beta_1(\gamma_0) = N\beta(\delta_0)^2 = N\alpha(\delta_0)^2 = 1$, which contradicts the regularity of γ_0 . Thus $N\alpha^\vee(\mathfrak{s}_T) = -1 = \alpha_0^\vee(\mathfrak{s}_T)$. Finally if α is of type R_1 , then $\alpha_0^\vee(\mathfrak{s}_T) = N\alpha^\vee(\mathfrak{s}_T) = -1$ since α_1 cannot be a root of H_1 . \square

The argument of Section 10.1 now shows that

$$\lim_{\nu \rightarrow 0^-} \widehat{D}\Psi_{a,\chi}(\gamma_\nu) = + \lim_{\nu \rightarrow 0^+} \widehat{D}\Psi_{a,\chi}(\gamma_\nu)$$

for all $D \in \mathcal{S}(\mathfrak{t}_1)$. Thus $\Psi_{a,\chi}$ extends smoothly around γ_0 .

We then have that $\Psi_{a,\chi}$ extends smoothly around all regular elements γ_0 in $T_1(\mathbb{R})$ that are norms of θ -semiregular elements in $G(\mathbb{R})$. Next, $\Psi_{a,\chi}$ extends to a smooth function around all regular elements γ_0 in $T_1(\mathbb{R})$ that are norms of θ -semisimple elements in $G(\mathbb{R})$. For this, Lemma 6.2 implies immediately that we may apply a familiar principle of Harish-Chandra, which we call *semiregular is sufficient*; see, for example, [Var77, §6 of Part I, §13 of Part II] and also [War72, Lemma 8.4.4.6 and §8.5]. We conclude then that $\Psi_{a,\chi}$, and thus S itself, extends to a smooth function on the full regular set of $T_1(\mathbb{R})$.

To finish the proof of the main theorem, Theorem 2.1, we check that S satisfies all requirements of our characterization theorem for stable orbital integrals on $H_1(\mathbb{R})$, i.e., Properties I–IV of Theorem 12.1 with $G = H_1$, $g_0 = 1$. Recall that we use Haar measures attached to invariant differential forms of highest degree defined over \mathbb{R} , as in [She79a, §4] and [LS87, §1.4], and use provided inner twists or \mathbb{R} -isomorphisms to transport forms when needed (for example, in the formulation of Property I). We assume that the forms on $\mathfrak{g}, \mathfrak{h}_1$ are products corresponding to the Lie algebra decompositions $\mathfrak{g} = \mathfrak{z}^\theta + (1 - \theta)\mathfrak{z} + \mathfrak{g}_{\text{der}}$, $\mathfrak{h}_1 = \mathfrak{z}_1 + \mathfrak{h}$. Suppose strongly G -regular γ_1 is a norm of strongly θ -regular δ . We require that Haar measures on $G_\delta^\theta(\mathbb{R})$ and $T_1(\mathbb{R})$ be compatible

in the following sense. First the underlying forms are to respect $\mathfrak{g}_\delta^\theta = \mathfrak{z}^\theta + (\mathfrak{g}_\delta^\theta \cap \mathfrak{g}_{\text{der}})$, $\mathfrak{t}_1 = \mathfrak{z}_1 + \mathfrak{t}_H$. Because the constant $|\det(\text{Int}(\delta) \circ \theta - I)_{\text{Cent}(\mathfrak{g}_\delta^\theta, \mathfrak{g}) / \mathfrak{g}_\delta^\theta}|$ or, more simply, $|\det(\theta^* - I)_{\mathfrak{t} / \mathfrak{t}^{\theta^*}}|$ was omitted from the normalizing factor Δ_{IV} , we include it now by requiring that the form on \mathfrak{t}_H be obtained by transport of $[\det(\theta^* - I)_{\mathfrak{t} / \mathfrak{t}^{\theta^*}}]^{-1}$ times the form on $\mathfrak{g}_\delta^\theta$. For the Haar measure on $\text{Cent}(\delta, G)(\mathbb{R})$, we extend that on $G_\delta^\theta(\mathbb{R})$.

For Property III, it remains to consider $\Psi_{a,\chi}$ around a T_1 -norm γ_0 annihilated only by real or complex roots. Again we use the *semiregular is sufficient* principle to assume that the root is real and unique up to sign and that both $(H_1)_{\gamma_0}$ and $G_{\delta_0}^\theta$ are of Dynkin type A_1 . Then descent finishes the argument. As in Section 14 of [She08a] for the standard (untwisted) case, an alternative proof that $\Psi_{a,\chi}$ extends to an ϖ_1 -Schwartz function on $T_1(\mathbb{R})_{\text{im-reg}}$ may be given via formulas for parabolic descent (see [Mez12], [She11]).

By Theorem 5.1 and its extension to derivatives, S satisfies Property IV under the additional assumption that γ_0 is G -semiregular. Our (stronger) statements of limit formulas for transfer factors in Section 9 allow us to remove the assumption by an application of the *semiregular is sufficient* principle, and then we are done with the proof of the main theorem. \square

12. The general case: slightly twisted norms

Without the assumption at the beginning of Section 6, the norms of strongly θ -regular elements in $G(\mathbb{R})$ lie in a certain coset of $H_1(\mathbb{R})$ in $H_1(\mathbb{C})$, rather than in $H_1(\mathbb{R})$ itself. This feature requires only a minor modification in the formulation of transfer, as we will recall from [KS99, §5.4]. We consider arbitrary (G, θ, a_ϖ) , endoscopic data \mathfrak{e} , and z -pair (H_1, ξ_1) (see Section 1).

We return to the first paragraph of Section 6. Recall that we work with the variant $m : G \rightarrow G^*$ of the inner twist ψ defined by $m(\delta) = \psi(\delta)g_\theta^{-1}$. Without the assumption of the first paragraph, we have that

$$\sigma(m)(\delta) = z(\sigma)u(\sigma)^{-1}m(\delta)\theta^*(u(\sigma)),$$

where $z(\sigma)$ is a 1-cochain of Γ in Z_{sc}^* . (As usual, we have used the same notation for the image of $z(\sigma)$ in G^* .) The image of $z(\sigma)$ in $(Z_{\text{sc}}^*)_{\theta_{\text{sc}}^*}$ is a 1-cocycle $z_\theta(\sigma)$. As in (5.4) of [KS99], $z_\theta(\sigma)$ determines a 1-cocycle $z_H(\sigma)$ in the center of H which we assume splits in H , since otherwise the transfer statement is empty. Let $z_H(\sigma) = h_0^{-1}\sigma(h_0)$. Then there is a 1-cocycle $z_1(\sigma) = h_1^{-1}\sigma(h_1)$ in the center of H_1 which projects to $z_H(\sigma)$ under $H_1 \rightarrow H$. Write θ_1 for the automorphism $\text{Int}(h_1)$. We replace $H_1(\mathbb{R})$ by the coset $H_1(\mathbb{R})h_1$ in the formulation of transfer.

First we extend the definition of stable orbital integral to this setting and describe a characterization theorem. Until after Theorem 12.1, we return to G as notation for the group on which we consider orbital integrals. Since it is enough for our purposes (i.e., for the case $G = H_1$), we also assume G quasi-split over \mathbb{R} and with simply-connected derived group. Then the complex

points of centralizers of semisimple elements are connected and there are no totally compact imaginary roots.

Fix an element g_0 in $G(\mathbb{C})$ such that $\sigma(g_0)^{-1}g_0$ is central, so that $\theta = \text{Int}(g_0)$ lies in $G_{\text{ad}}(\mathbb{R})$ and $G(\mathbb{R})g_0$ lies in the inverse image of $G_{\text{ad}}(\mathbb{R})$ under the projection $G \rightarrow G_{\text{ad}}$. There will be no harm in assuming that θ preserves the pair $(B_{\text{spl}}, T_{\text{spl}})$, where $B_{\text{spl}}, T_{\text{spl}}$ are from a chosen \mathbb{R} -splitting of G . Then g_0 lies in the maximal torus T_{spl} of the splitting. There is also no harm in assuming g_0 lies in G_{der} . Then $\sigma(g_0)^{-1}g_0 = z(\sigma)$ lies in the center $Z_{\text{der}} = Z_{\text{sc}}$ of $G_{\text{der}} = G_{\text{sc}}$.

Let $\gamma \in G(\mathbb{R})g_0 \subset G(\mathbb{C})$. Then $\text{Cent}(\gamma, G)$ is defined over \mathbb{R} since $\sigma(\gamma)^{-1}\gamma = \sigma(g_0)^{-1}g_0 = z_\sigma$. Suppose γ is regular semisimple, so that $T_\gamma = \text{Cent}(\gamma, G)$ is a maximal torus defined over \mathbb{R} . If $\gamma = \gamma'g_0$, then right translation by g_0 maps bijectively the $\text{Int}(g_0)$ -twisted conjugacy class of γ' to the $G(\mathbb{R})$ -conjugacy class of γ . It also maps the intersection of $G(\mathbb{R})$ with the $\text{Int}(g_0)$ -twisted conjugacy class of γ' in $G(\mathbb{C})$ to the intersection of $G(\mathbb{R})g_0$ with the conjugacy class of γ in $G(\mathbb{C})$. We will call this last set the stable conjugacy class of γ (again since G_{der} is simply-connected). The $G(\mathbb{R})$ -conjugacy classes in the stable conjugacy class of γ are parametrized by untwisted $\mathcal{D}(T_\gamma)$, as for the case $\gamma \in G(\mathbb{R})$.

Let T be a maximal torus over \mathbb{R} in G . Then T contains an element γ in $G(\mathbb{R})g_0$ if and only if z_σ splits in $H(\Gamma, T_{\text{der}}) = H(\Gamma, T_{\text{sc}})$. In that case, $T(\mathbb{R})\gamma$ also lies in $G(\mathbb{R})g_0$ and moreover $T(\mathbb{R})\gamma = T \cap G(\mathbb{R})g_0$. Write $\mathcal{T}(g_0)$ for the collection of all such T . Clearly, $T_{\text{spl}} \in \mathcal{T}(g_0)$ and the set of regular semisimple elements in $G(\mathbb{R})g_0$ is the union over $T \in \mathcal{T}(g_0)$ of the (open, dense) regular set $(T \cap G(\mathbb{R})g_0)_{\text{reg}}$ in $T \cap G(\mathbb{R})g_0$. Suppose $T \in \mathcal{T}(g_0)$, $\gamma_0 \in T \cap G(\mathbb{R})g_0$ is semiregular, and $\alpha(\gamma_0) = 1$, where α is an imaginary root of T . On replacing γ_0 by a stable conjugate we may assume that α is noncompact, i.e., that $\text{Cent}(\gamma_0, G)$ is split modulo center. If T' is a maximally split maximal torus over \mathbb{R} in $\text{Cent}(\gamma_0, G)$, then clearly $T' \in \mathcal{T}(g_0)$. It then follows that if $T \in \mathcal{T}(g_0)$ and s is any Cayley transform relative to an imaginary root α of T , then $T^s \in \mathcal{T}(g_0)$. Also, if $\gamma_0 \in T \cap G(\mathbb{R})g_0$ is semiregular and $\alpha(\gamma_0) = 1$, then $(\gamma_0)^s \in T^s \cap G(\mathbb{R})g_0$. We denote by $(T \cap G(\mathbb{R})g_0)_{\text{im-reg}}$ the set of all elements in $T \cap G(\mathbb{R})g_0$ such that $\alpha(\gamma_0) \neq 1$ for all imaginary roots α of T .

Let $S(\gamma, dt_\gamma, dg)$ be a complex-valued function defined for regular semisimple γ in $G(\mathbb{R})g_0$ and Haar measures dt_γ on $T^\gamma(\mathbb{R}) = \text{Cent}(\gamma, G)(\mathbb{R})$, dg on $G(\mathbb{R})$. Write Φ for the normalized version of S :

$$\Phi(\gamma, dt_\gamma, dg) = \left| \det(\text{Ad}(\gamma) - I)_{\mathfrak{g}/\mathfrak{t}_\gamma} \right|^{1/2} S(\gamma, dt_\gamma, dg).$$

Since it is useful for our application, we assume that there is a central torus Z_1 and character ϖ_1 on $Z_1(\mathbb{R})$ such that

$$S(z_1\gamma, dt_\gamma, dg) = \varpi_1(z_1)^{-1} S(\gamma, dt_\gamma, dg)$$

for all $z_1 \in Z_1(\mathbb{R})$, regular semisimple γ in $G(\mathbb{R})g_0$, and all dt_γ, dg .

Consider the following properties (I)–(IV):

(I) *S is invariant under stable conjugacy.*

This means that if $w \in G(\mathbb{C})$ is such that $\gamma^w = w^{-1}\gamma w$ lies in $G(\mathbb{R})g_0$ and dt_{γ^w} is obtained from dt_γ by transport under w , then

$$S(\gamma^w, dt_{\gamma^w}, dg) = S(\gamma, dt_\gamma, dg).$$

(II) *S transforms under change of measures according to the rule*

$$S(\gamma, \lambda dt_\gamma, \mu dg) = \frac{\mu}{\lambda} S(\gamma, dt_\gamma, dg).$$

Here λ, μ are positive real numbers.

Next, let $T \in \mathcal{T}(g_0)$. For γ in $(T \cap G(\mathbb{R})g_0)_{\text{reg}}$ and fixed Haar measures dt, dg on $T(\mathbb{R}), G(\mathbb{R})$ respectively, set $S^T(\gamma) = S(\gamma, dt, dg)$ and $\Phi^T(\gamma) = \Phi(\gamma, dt, dg)$.

(III) *Φ^T is a ϖ_1 -Schwartz function on $(T \cap G(\mathbb{R})g_0)_{\text{reg}}$ and extends to a ϖ_1 -Schwartz function on $(T \cap G(\mathbb{R})g_0)_{\text{im-reg}}$.*

Here the notion of ϖ_1 -Schwartz function is clear since $T(\mathbb{R})\gamma_0$ lies in the inverse image of $T_{\text{ad}}(\mathbb{R})$ in $T(\mathbb{C})$.

The final property concerns behavior at the imaginary walls. It is simpler to state if we assume (I), (III). Suppose $T \in \mathcal{T}(g_0)$, $\gamma_0 \in T \cap G(\mathbb{R})g_0$ is semiregular and $\alpha(\gamma_0) = 1$, where α is an imaginary root of T . Let s be a Cayley transform for α (in the sense of Section 3), and fix s -compatible a -data, χ -data for $T, T^s = T'$. (We again use $'$ in place of s in notation.) The Haar measure on $T'(\mathbb{R})$ is to be obtained by transport under s from that on $T(\mathbb{R})$ in our earlier sense (Section 8). We have defined the generalized Weyl denominator $\Delta_{a,\chi}(\gamma)$ for $\gamma \in T(\mathbb{R})$. Notice that $\Delta_{a,\chi}(\gamma)$ depends on the image of γ under the natural map $T \rightarrow T_{\text{ad}}$ rather than on γ itself. We may therefore extend the definition of $\Delta_{a,\chi}$ to the inverse image of $T_{\text{ad}}(\mathbb{R})$ in $T(\mathbb{C})$ and so to $T(\mathbb{R})\gamma_0$. We also extend $\Delta_{a',\chi'}$ to $T'(\mathbb{R})\gamma'_0$. Thus we may define the transforms $\Psi_{a,\chi}, \Psi_{a',\chi'}$ on $(T(\mathbb{R})\gamma_0)_{\text{reg}}, (T'(\mathbb{R})\gamma'_0)_{\text{reg}}$ respectively, as before. For ν real and nonzero, set $\gamma_\nu = \exp(\nu a_\alpha \alpha^\vee) \cdot \gamma_0$ and $\gamma'_\nu = \exp(\nu a_{\alpha'} \alpha'^\vee) \cdot \gamma'_0$. Denote by w the Weyl reflection for α . To a differential operator D in $\mathcal{S}(\mathfrak{t})$ attach D' in $\mathcal{S}(\mathfrak{t}')$ and define the twists $\widehat{D}, \widehat{D}'$ as in Section 10.

(IV) *If $D^w = D$, then*

$$\lim_{\nu \rightarrow 0^+} \widehat{D} \Psi_{a,\chi}(\gamma_\nu) = \lim_{\nu \rightarrow 0} \widehat{D}' \Psi_{a',\chi'}(\gamma'_\nu).$$

With the assumption of (I) and (III), there will be no harm in assuming in (IV) that α is noncompact and that the Cayley transform comes from the simply-connected cover $SL(2)$ of $\text{Cent}(\gamma_0, G)$. Then $\gamma'_0 = \gamma_0$. Also, the argument of Section 10, along with (I) and (III), shows that if $D^w = \pm D$, then

$$\lim_{\nu \rightarrow 0^-} \widehat{D} \Psi_{a,\chi}(\gamma_\nu) = \mp \lim_{\nu \rightarrow 0^+} \widehat{D} \Psi_{a,\chi}(\gamma_\nu).$$

Suppose now that f is a ϖ_1 -Schwartz function on $G(\mathbb{R})g_0$. (Again the notion is clear, or see the appendix.) Then the stable orbital integrals

$$\text{SO}(\gamma, f) = \text{SO}(\gamma, f, dt_\gamma, dg) = \sum_{\gamma' \in \mathcal{D}(T_\gamma)} \int_{T\gamma'(\mathbb{R}) \backslash G(\mathbb{R})} f(g^{-1}\gamma'g) \frac{dg}{dt_\gamma}$$

transform by ϖ_1^{-1} under translation by $Z_1(\mathbb{R})$ and satisfy (I)–(IV). A proof of this requires only a very minor variant of the standard argument; see the next proof or the appendix for more general results. Extension of our main theorem to the present setting rests on the converse theorem.

THEOREM 12.1. *Suppose S transforms by ϖ_1^{-1} under translation by $Z_1(\mathbb{R})$ and satisfies (I)–(IV). Then there exists $f \in \mathcal{C}(G(\mathbb{R})g_0, \varpi_1)$ such that*

$$S(\gamma, dt_\gamma, dg) = \text{SO}(\gamma, f, dt_\gamma, dg)$$

for all regular semisimple γ in $G(\mathbb{R})g_0$, and all dt_γ, dg . If also S vanishes off the orbits of some set $Z_1(\mathbb{R})B$, where B is a bounded subset of the regular semisimple set of $G(\mathbb{R})g_0$, then f may be chosen in $C_c^\infty(G(\mathbb{R})g_0, \varpi_1)$.

Proof. To find f in $\mathcal{C}(G(\mathbb{R})g_0, \varpi_1)$ we prove an analog of Lemma 4.8 of [She79a] in which f is constructed satisfying a weaker condition, and we then finish by using the inductive argument for the proof of Theorem 4.7 in [She79a]. Assume $T \in \mathcal{T}(g_0)$. Then an argument shows that we may replace g_0 by an element of $G(\mathbb{R})g_0 \cap T$ if necessary and assume $g_0 \in T$. It is now straightforward to extend the wave packet construction in the proof of Lemma 4.8 to $G(\mathbb{R})g_0$ and thus find the desired f in $\mathcal{C}(G(\mathbb{R})g_0, \varpi_1)$. To pass to a C_c^∞ -function when the support is appropriate, we reduce to Bouaziz’s characterization theorem on $G_{\text{ad}}(\mathbb{R})$. \square

Finally, the extension of Theorem 2.1 requires a recasting of the norm correspondence and transfer factors. This again is straightforward (and described in Section 5.4 of [KS99]). First, for the norm correspondence we consider strongly G -regular elements γ_1 of $H_1(\mathbb{R})h_1$, assuming such elements exist, and strongly θ -regular elements δ of $G(\mathbb{R})$. Then γ_1 is a norm of δ if the θ -conjugacy class of δ in $G(\mathbb{C})$ is the image (under the canonical map) of the conjugacy class of γ_1 in $H_1(\mathbb{C})$. Let $T_1 = \text{Cent}(\gamma_1, H_1)$, a maximal torus over \mathbb{R} in H_1 . Then there are toral data $(T_1 \rightarrow T_{\theta^*}, g)$ as in Section 6 for which $\delta^* = gm(\delta)\theta^*(g)^{-1}$ has the property that $N(\delta^*)$ is the image of γ_1 under $T_1 \rightarrow T_{\theta^*}$. The cochain $v(\sigma)$ in T_{sc} now has the extra term $z(\sigma)$ from Z_{sc} , but that does not affect the assertions of the lemmas in Section 6 when we now take semisimple γ_0 in $H_1(\mathbb{R})h_1$ instead of $H_1(\mathbb{R})$. Nor does it affect the definition of the relative term Δ_{III} in transfer factors, since $(z(\sigma)^{-1}, z(\sigma))$ represents the identity element of the torus U of Section 4.4 of [KS99]. The results of Sections 6, 7 and 9 thus apply. After adjusting the definition of $\text{Trans}(f)$ and $\text{Trans}_c(f)$, we conclude then

THEOREM 12.2. *The assertions of the main theorem (Theorem 2.1) and corollary (Corollary 2.2) remain true in the general setting of Section 6.*

13. Appendix: Harish-Chandra Schwartz functions

We return to the setting of Section 1, where f_θ is a smooth function on $G(\mathbb{R})\theta$. As pointed out by a referee, θ is the product of an inner automorphism defined over \mathbb{R} and an automorphism of finite order also defined over \mathbb{R} ; see the remark near the end of Section 6. We may further assume that the inner automorphism is of the form $\text{Int}(g)$, where $g \in G_{\text{sc}}(\mathbb{R})$. There will be no harm then in assuming that θ itself is of finite order. Following [HC75], let $V = \exp \mathfrak{v}$, where \mathfrak{v} is the maximal \mathbb{R} -split subalgebra of the Lie algebra $\mathfrak{z}(\mathbb{R})$ of $Z(\mathbb{R})$, so that we have $G(\mathbb{R})$ as a direct product $(1 - \theta)V.V^\theta.\circ G(\mathbb{R})$, where θ acts as automorphism of each factor. Then $G(\mathbb{R})\theta$ is a direct product $(1 - \theta)V.G_1$, where $G_1 = V^\theta.\circ G(\mathbb{R})\theta$; G_1 embeds smoothly as an open subset of the Lie group $V^\theta.\circ G(\mathbb{R}) \rtimes \langle \theta \rangle$ to which the results of [Bou87] apply. We will start with the space $\mathcal{C}(G(\mathbb{R})\theta, \varpi)$, where we require f_θ to transform by ϖ^{-1} under the twisted conjugacy action of V since we will need such a space for the twisted orbital integrals. Thus we require $f_\theta(v\theta(v)^{-1}g\theta) = f_\theta(vg\theta v^{-1}) = \varpi(v)^{-1}f_\theta(g\theta)$ for $v \in V, g \in G(\mathbb{R})$. (Since we assume a nonempty norm correspondence, the character ϖ is trivial on V^θ , the kernel of the action.) Call f_θ a ϖ -Schwartz function if the restriction of f_θ to G_1 is Schwartz in the Harish-Chandra sense [HC75]: the functions σ and Ξ appearing in Harish-Chandra's seminorms are well defined on G_1 (see Sections 3.4, 3.5 of [Bou87]). We write $\mathcal{C}(G(\mathbb{R})\theta, \varpi)$ for the Fréchet space of all such functions equipped with the Harish-Chandra seminorms. If \mathcal{O} is open in $G(\mathbb{R})\theta$ and invariant under translation by $(1 - \theta)V$, we define $\mathcal{C}(\mathcal{O}, \varpi)$ analogously. It is clear also how to define the space $\mathcal{C}(G(\mathbb{R})\theta)$ of (purely) Schwartz functions on $G(\mathbb{R})\theta$.

We need a twisted analogue of Theorem 16.1 of [HC75], which asserts that $f \rightarrow {}'F_f$ is a well-defined continuous map on the appropriate Schwartz spaces. To shorten the presentation (but also make it clumsier than necessary), we take our (θ, ϖ) -twisted transform to depend on the endoscopic group also, or more precisely, on endoscopic data and z -pair. To use pieces of the transfer factors in the definition, we will start with familiar constructions on $G(\mathbb{R})$ and then translate to $G(\mathbb{R})\theta$. Fix a strongly θ -regular element δ_0 of $G(\mathbb{R})$ with norm γ_0 in $H_1(\mathbb{R})$, along with toral data, a -data and χ -data associated with the torus $\text{Cent}(\gamma_0, H_1)$. It will be enough for our purposes to define a transform $\Psi_f^{\delta_0}$ on the θ -regular elements δ in $G_{\delta_0}^\theta(\mathbb{R})^0\delta_0$, although extension to a larger set is easy. If $\delta = \exp Y.\delta_0$, then we set $\gamma_1 = \exp Y_H.\gamma_0$ (see Section 6) and define

$$\Psi_f^{\delta_0}(\delta) = \Delta_{III}(\gamma_0, \delta_0; \bar{\gamma}, \bar{\delta}).\Delta_{II}^{\text{num}}(\delta).\Phi^{\theta, \varpi}(\delta, f).$$

We have omitted the term Δ_I since fixed toral data and a -data guarantee that Δ_I is a constant that plays no role here. The term Δ_{II}^{num} (from Section 9) is a

twisted version of the Weyl denominator of Section 3. The presence of the constant $\Delta_{III}(\gamma_0, \delta_0; \bar{\gamma}, \bar{\delta})$ ensures that if $g \in G(\mathbb{R})$, then $\Psi_{f_\theta}^{g^{-1}\delta_0\theta(g)}(g^{-1}\delta\theta(g)) = \Psi_{f_\theta}^{\delta_0}(\delta)$, provided we follow the usual conventions in the choice of Haar measures. If we replace δ_0 by (strongly) θ -regular δ'_0 in $G_{\delta_0}^\theta(\mathbb{R})^0\delta_0$, we obtain a translate of $\Psi_{f_\theta}^{\delta_0}$ which does not matter for the Schwartz properties we seek. (For translation-invariance arguments, see, for example, Section 8.5 of [War72].) To pass to $G(\mathbb{R})\theta$, we set $\Phi^\varpi(\delta\theta, f_\theta) = \Phi^{\theta, \varpi}(\delta, f)$ and $\Psi_{f_\theta}^{\delta_0}(\delta\theta) = \Psi_f^{\delta_0}(\delta)$ for all regular $\delta\theta$ in $\text{Conn}(\delta_0\theta) = G_{\delta_0}^\theta(\mathbb{R})^0\delta_0\theta$, a connected component of $\text{Cent}(\delta_0\theta, G(\mathbb{R})\theta)$. It is now routine to define $\text{Conn}(\delta_0\theta)_{\text{im-reg}}$. Our assertion is that Theorem 16.1 of [HC75] together with the work of Bouaziz already cited implies that $f_\theta \rightarrow \Psi_{f_\theta}^{\delta_0}$ is a well-defined continuous mapping from $\mathcal{C}(G(\mathbb{R})\theta)$ (or from $\mathcal{C}(G(\mathbb{R})\theta, \varpi)$) to $\mathcal{C}(\text{Conn}(\delta_0\theta)_{\text{im-reg}}, \varpi)$. Theorem 16.1 is proved in [Var77] following Harish-Chandra's original argument. (The final steps are in Part II, Section 12.) An alternative argument not dependent on the construction of discrete series characters has been given by Wallach (see Chapter 7 of [Wal88]). Since an analogue for the otherwise needed discrete series result has not yet appeared, we follow step by step the arguments of [Wal88]. In particular, the crucial Lemma 7.4.3 extends to our setting by preparation from Sections 1–3 of [Bou87]. This is enough to finish the argument.

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(Received: October 12, 2010)

(Revised: May 29, 2012)

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