# The singular set of 1-1 integral currents 

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#### Abstract

We prove that 2 dimensional integer multiplicity 2 dimensional rectifiable currents which are almost complex cycles in an almost complex manifold admitting locally a compatible positive symplectic form are smooth surfaces aside from isolated points and therefore are $J$-holomorphic curves.


## I. Introduction

Let $\left(M^{2 p}, J\right)$ be an almost complex manifold. Let $k \in \mathbb{N}, k \leq p$. We shall adopt classical notation from Geometric Measure Theory [Fe]. We say that a $2 k$-current $C$ in $\left(M^{2 p}, J\right)$ is an almost complex integral cycle whenever it fulfills the following three conditions
i) Rectifiability: There exists an at most countable union of disjoint oriented $C^{1} 2 k$-submanifolds $\mathcal{C}=\cup_{i} N_{i}$ and an integer multiplicity $\theta \in L_{\text {loc }}^{1}(\mathcal{C})$ such that for any smooth compactly supported in $M 2 k$-form $\psi$ one has

$$
C(\psi)=\sum_{i} \int_{N_{i}} \theta \psi .
$$

ii) Closedness: $C$ is a cycle,

$$
\partial C=0 ; \quad \text { i.e., } \forall \alpha \in \mathcal{D}^{2 k-1}(M), \quad C(d \alpha)=0 .
$$

iii) Almost complex: For $\mathcal{H}^{2 k}$ and almost every point $x$ in $\mathcal{C}$, the approximate tangent plane $T_{x}$ to the rectifiable set $\mathcal{C}$ is invariant under the almost complex structure $J$; i.e.,

$$
J\left(T_{x}\right)=T_{x} .
$$

In this work we address the question of the regularity of such a cycle: Does there exist a smooth almost complex manifold $\left(\Sigma^{2 k}, j\right)$ without boundary and a smooth $j-J$-holomorphic map $u\left(\forall x \in \Sigma\right.$ and $\left.\forall X \in T_{x} \Sigma d u_{x} j \cdot X=J \cdot d u_{x} X\right)$ such that $u$ would realize an embedding in $M^{2 p}$ aside from a locally finite $2 k-2$
measure closed subset of $M$ and such that $C=u_{*}\left[\Sigma^{2 k}\right]$; i.e., $\forall \psi \in C_{0}^{\infty}\left(\wedge^{2 k} M\right)$

$$
C(\psi)=\int_{\Sigma} u^{*} \psi ?
$$

In the very particular case where the almost complex structure $J$ is integrable, this regularity result is optimal ( $C$ is the integral over multiples of algebraic subvarieties of $M$ ) and was established in [HS] and [Ale]. There are numerous motivations for studying the general case of arbitrary, almost complex structures $J$. First, as explained in [RT], the above regularity question for rectifiable almost complex cycles is directly connected to the regularity question of $J$-holomorphic maps into complex projective spaces. It is conjectured, for instance, that the singular set of $W^{1,2}\left(M^{2 p}, N\right) J$-holomorphic maps between almost complex manifolds $M$ and $N$ should be of finite ( $2 p-4$ )-Hausdorff measure. The resolution of that question leads, for instance, to the characterization of stable-bundle, almost complex structures over almost Kähler manifolds via Hermite-Einstein Structures and extends Donaldson, Uhlenbeck-Yau characterization in the integrable case (see [Do], [UY]) to the nonintegrable one. Another motivation for studying the regularity of almost complex rectifiable cycles is the following. In [Li] and [Ti] it is explained how the loss of compactness of solutions to geometric PDEs having a given conformal invariant dimension $q$ (a dimension at which the PDE is invariant under conformal transformations - $q=2$ for harmonic maps, $q=4$ for Yang-Mills Fields...etc) arises along $m-q$ rectifiable cycles (if $m$ denotes the dimension of the domain). These cycles happen sometimes to be almost complex (see more details in [Ri1]).

By trying to produce in $\left(\mathbb{R}^{2 p}, J\right)$ an almost complex graph of real dimension $2 k$ in a neighborhood of a point $x_{0} \in \mathbb{R}^{2 p}$ as a perturbation of a complex one ( $J_{x_{0}}$-holomorphic), one realizes easily that, for generic almost complex structures $J$, the problem is overdetermined whenever $k>1$ and well posed for $k=1$. Therefore the case of 2 -dimensional integer rectifiable almost complex cycles is the generic one from the existence point of view. We shall restrict to that important case in the present paper. After complexification of the tangent bundle to $M^{2 p}$ a classical result asserts that a 2-plane is invariant under $J$ if and only if it has a 1-1 tangent 2 -vector. Therefore we shall also speak about 1-1 integral cycles for the almost complex 2-dimensional integral cycles. In the present work we consider the locally symplectic case: We say that $\left(M^{2 p}, J\right)$ has the locally symplectic property if at a neighborhood of each point $x_{0}$ in $M^{2 p}$ there exist a positive symplectic structure compatible with $J$ and a neighborhood $U$ of $x_{0}$ and a smooth closed 2 -form $\omega$ such that $\omega(\cdot, J \cdot)$ defines a scalar product. It was proved in $[\mathrm{RT}]$ that arbitrary, 4 -dimensional, almost complex manifolds satisfy the locally symplectic property. This is no more the case in larger dimension: one can find an almost complex structure in $S^{6}$ which admits no compatible positive symplectic form even locally; see $[\mathrm{Br}]$.

Our main result is the following.
Theorem I.1. Let $\left(M^{2 p}, J\right)$ be an almost complex manifold satisfying the locally symplectic property above. Let $C$ be an integral 2 dimensional almost complex cycle. Then, there exist a J-holomorphic curve $\Sigma$ in $M$, smooth aside from isolated points, and a smooth integer-valued function $\theta$ on $\Sigma$ such that, for any 2 form $\psi \in C_{0}^{\infty}(M)$,

$$
C(\psi)=\int_{\Sigma} \theta \psi
$$

In the "locally symplectic case" being an almost-complex 2 cycle is equivalent for a 2 -cycle to being calibrated by the local symplectic form $\omega$ for the local metric $\omega(\cdot, J \cdot)$. Therefore the regularity question for almost complex cycles is embedded into the problem of calibrated current and hence the theory of area-minimizing rectifiable 2 -cycles. Therefore our result appears to be a consequence of the "Big Regularity Paper" of F. Almgren [Alm] combined with the PhD thesis of his student S. Chang [Ch]. Our attempt here is to present an alternative proof independent of Almgren's monumental work and adapted to the case we are interested in. The motivation is to give a proof that could be modified in order to solve the general case (non locally symplectic one) which cannot be "embedded" in the theory of area-minimizing cycles anymore.

A proof for the regularity of almost complex cycle in the locally symplectic, $p=2$ case, independent of the regularity theory for area-minimizing surfaces, was also one of the results of the work $\mathrm{Gr} \Longrightarrow \mathrm{SW}$ of C . Taubes [Ta] for $p=2$. In particular, [Ta] presents a proof of Theorem I. 1 when $p=2$. In [RT], we gave an alternative proof for this special case. Theorem I. 1 can be seen as the generalization to higher dimension $(p>2)$ of these works.

One of the main difficulties arising in dimension $p>2$ is the nonnecessary existence of $J$-holomorphic foliations transverse to our almost complex current $C$ in a neighborhood of a point. This then prevents describing the current as a $Q$-multivalued graph from $D^{2}$ into $\mathbb{C}^{p-1},\left\{\left(a_{i}^{k}(z)\right)_{k=1^{p-1}}\right\}_{i=1 \cdots Q}$ in a neighborhood of a point of density $N$ solving locally an equation of the form

$$
\begin{equation*}
\partial_{\bar{z}} a_{i}^{k}=\sum_{l=1}^{p-1} A\left(z, a_{i}\right)_{l}^{k} \cdot \nabla a_{i}^{l}+\alpha^{k}\left(a_{i}, z\right), \tag{I.1}
\end{equation*}
$$

where $A$ and $\alpha$ are small in $C^{2}$ norm, as for $p=2$ in [RT]. What we can only ensure instead is to describe the current $C$, in a neighborhood of a point of multiplicity $Q$, as an "algebraic $Q$-valued graph" from $D^{2}$ into $\mathbb{C}^{p-1}$; that is, a family of points in $\mathbb{C}^{p-1},\left\{a_{1}(z), \ldots, a_{P}(z), b_{1}(z), \ldots, b_{N}(z)\right\}$ where only $P-N=Q$ is independent on $z$ (neither $P$ nor $N$ is a priori independent on $z), a_{i}$ are the positive intersection points and $b_{j}$ are the negative ones. This "algebraic $Q$-valued graph" solves locally a much less attractive equation
than (I.1),

$$
\begin{equation*}
\partial_{\bar{z}} a_{i}^{k}=\sum_{l=1}^{p-1} A_{k}^{l}\left(z, a_{i}, \nabla a_{i}\right) \cdot \nabla a_{i}^{l}+\sum_{l=1}^{p-1} B_{k}^{l}\left(z, a_{i}\right) \cdot \nabla a_{i}^{l}+C^{k}\left(z, a_{i}\right) \tag{I.2}
\end{equation*}
$$

where $A(z, a, p), B(z, a)$ and $C(z, a)$ are also small in $C^{2}$ norm but the dependence on $p$ in $A(z, a, p)$ is linear and therefore as $\nabla a_{i}$ gets bigger, which can happen, the right-hand side of (I.2) cannot be handled as a perturbation of the left-hand one in steps such as the "unique continuation argument". This was used in [RT] for proving that singularities of multiplicity $Q$ cannot have an accumulation point in the carrier $\mathcal{C}$ of $C$.

The strategy of the proof goes as follows. A classical blow-up analysis tells us that, for an arbitrary point $x_{0}$ of the manifold $M^{2 p}$, the limiting density $\theta\left(x_{0}\right)=\lim _{r \rightarrow 0} r^{-2} \mathrm{M}\left(C\left\llcorner B_{r}\left(x_{0}\right)\right)\right.$. Here $M$ denotes the mass of a current and $L$ is the restriction operator which equals $\pi$ times an integer $Q$. Since the density function $r \rightarrow r^{-2} \mathrm{M}\left(C\left\llcorner B_{r}\left(x_{0}\right)\right)\right.$ at every point is a monotonic increasing function, the complement of the set $\mathcal{C}_{Q}:=\{x \in M ; \theta(x) \leq Q\}$ is closed in $M$ and this permits us to perform an inductive proof of Theorem I. 1 restricting the current to $\mathcal{C}_{Q}$ and considering increasing integers $Q$. A point of multiplicity $Q$ is called a singular point of $C$ if it is in the closure of points of nonzero multiplicity strictly less than $Q$.

The goal of the proof is then to show that singularities of multiplicity less than $Q$ are isolated. We assume this fact for $Q-1$ and the paper is devoted to the proof that this then holds for $Q$ itself. From a now classical result of B. White (see [Wh]), the dilated currents at a point $x_{0}$ of density $Q \neq 0$ converge in flat norm to a sum of $Q$ flat $J_{x_{0}}$-holomorphic disks. Moreover, for any $\varepsilon>0$ and $r$ sufficiently small $C\left\llcorner B_{r}\left(x_{0}\right)\right.$ is supported in the cones whose axes are the limiting disks and angle $\varepsilon$. For $Q>1$, if two of these limiting disks are different it is then easy to observe that $x_{0}$ cannot be an accumulation point of singularities of multiplicity $Q$; this is the so called "easy case". If the limiting disks are all identical, equal to $D_{0}$, then we are in the "difficult case" and much more work has to be done in order to reach the same statement.

Contrary to the special case of dimension $4(p=2)$ considered by the authors in $[R T]$, we could not find nice coordinates that would permit us to write $C$ as a $Q$-valued graph over the limiting disk $D_{0}$. Considering then some $J_{x_{0}}$-complex coordinates $\left(z, w_{1}, \ldots, w_{p-1}\right)$ in a neighborhood $B_{\rho_{x_{0}}}^{2 p}\left(x_{0}\right)$ such that $\cap_{i} w_{i}^{-1}\{0\}$ corresponds to $D_{0}$, by the mean of the "lower-epiperimetric inequality" proved by the first author in [Ri2], one can construct a WhitneyBesicovitch covering, $\left\{B_{\rho_{i}}^{2}\left(z_{i}\right)\right\}_{i \in I}$, of the orthogonal projection on $D_{0}$ of the points in $B_{\rho_{x_{0}}}^{2 p}\left(x_{0}\right)$ having a positive density strictly less than $Q$. This covering is such that for every $i \in I$ there exists $x_{i}=\left(z_{i}, w_{i}\right) \in B_{\rho_{x_{0}}}^{2 p}\left(x_{0}\right)$ verifying that the restriction of $C$ to the tube $B_{\rho_{i}}^{2}\left(z_{i}\right) \times B_{\rho_{x_{0}}}^{2 p-2}(0)$ is in fact supported in the
ball $B_{2 \rho_{i}}^{2 p}\left(x_{i}\right)$ of radius $2 \rho_{i}$, two times the width of the tube. Moreover if one looks inside $B_{\rho_{i}}^{2 p}\left(x_{i}\right), C$ is "split": this last word means that $C$ restricted to $B_{\rho_{i}}^{2 p}\left(x_{i}\right)$ is at a flat distance comparable to $\rho_{i}^{3}$ from the $Q$ multiple of any graph over $B_{\rho_{i}}^{2}\left(z_{i}\right)$. This comes from the fact that the density ratio $\rho_{i}^{-2} \mathrm{M}\left(B_{\rho_{i}}\left(x_{i}\right)\right)$ is strictly less than $\pi Q$ minus a constant $\alpha$ depending only on $p, Q, J$ and $\omega$.

We then construct an average curve for $C$. In the 4 -dimensional case since $C$ was a $Q$-valued graph over $D_{0}$ we simply took the average of the $Q$ points over any point in $D_{0}$. Here, in arbitrary dimension, the construction of the average curve is more delicate and uses the covering. We first approximate $C\left\llcorner B_{\rho_{i}}^{2 p}\left(x_{i}\right)\right.$ by a $J_{x_{i}}$-holomorphic graph $C_{i}$ using a technique introduced in [Ri3], and choosing a $J_{x_{i}}$-holomorphic disk $D_{i}$ approximating $D_{0}$ we can express $C_{i}$ as a $Q$-valued graph over $D_{i}$ for which we take the average $\tilde{C}_{i}$ that happens to be Lipschitz with a uniformly bounded Lipschitz constant.

Therefore the $J_{x_{i}}$-holomorphic curve $\tilde{C}_{i}$ can be viewed as a graph $\tilde{a}_{i}$ over $B_{\rho_{i}}^{2}\left(z_{i}\right)$. Patching the $\tilde{a}_{i}$ together we get a graph $\tilde{a}$ that extends over the whole $B_{\rho_{x_{0}}}^{2}(0)$ as a $C^{1, \alpha}$ graph for any $\alpha<1$ which is almost $J$-holomorphic and which passes through all the $B_{\rho_{i}}^{2 p}\left(x_{i}\right)$. The fact that the average curve is more regular than the $J$-holomorphic cycle $C$ from which it is produced is clear in the integrable case (since it is holomorphic); $(z, \pm \sqrt{z})$ is a $C^{0, \frac{1}{2}} 2$-valued graph whereas its average $(z, 0)$ is smooth. This was extended in the nonintegrable case in the particular case of the 4 dimension in [ST].

The points of multiplicity $Q$ in $C$ are contained in the average curve $\tilde{a}$. We then show, by means of a unique continuation argument in the spirit of the one developed in $[\mathrm{Ta}]$ in 4 dimensions, that the points where $C$ gets to coincide with $\tilde{a}$ are either isolated or coincide with the whole curve $\tilde{a}$. We have then shown that any point $x_{0}$ of multiplicity $Q$ is either surrounded by points of multiplicity $Q$ only, and in $B_{\rho_{x_{0}}}^{2 p}, C$ coincides with $Q$ times a smooth graph over $D_{0}$ or $x_{0}$ is not an accumulation point of points of multiplicity $Q$ and is surrounded in $B_{\rho_{x_{0}}}^{2 p}$ by points of multiplicity strictly less than $Q$. It remains at the end to show that it cannot be an accumulation point of singularities of lower density. This is obtained again using an approximation argument by holomorphic curves introduced in [Ri3].

The paper is organized as follows. In Section II we establish preliminaries, introduce notation and give the main statement, assertion $\mathcal{P}_{Q}$, to be proved by induction in the rest of the paper. In Section III, with the help of the "upper-epiperimetric inequality" of B . White, we establish the uniqueness of the tangent cone and a quantitative version of it; see Lemma III.2. In Section IV we prove the relative Lipschitz estimate together with a tilting control of the tangent cones of density $Q$ points in a neighborhood of a density $Q$ point; see Lemma IV.2. In Section V we proceed to the covering argument, Lemma V.3, which is based on the "splitting before tilting" lemma (see Lemma V.1, proved in [Ri2]). In Section VI we construct the approxi-
mated average curve and prove the $C^{1, \alpha}$ estimate for this curve, Lemma VI.3. In Section VII we perform the unique continuation argument showing that singularities of multiplicity $Q$ cannot be accumulation points of singularities of multiplicity $Q$. In Section VIII we show that singularities of multiplicity $Q$ cannot be accumulation points of singularities of multiplicity less than $Q$ either.

## II. Preliminaries

Notation. We shall adopt standard notation from the Geometric Measure Theory [Fe] such as $\mathrm{M}(A)$ for the Mass of a current $A, \mathcal{F}(A)$ for its flat norm, $A\llcorner E$ for its restriction to a measurable subset $E \ldots$...tc; we refer the reader to [Fe].

Preliminaries. Since our result is a local one we shall work in a neighborhood $U$ of a point $x_{0}$ and use a symplectic form $\omega$ compatible with $J$. We denote by $g$ the metric generated by $J$ and $\omega: g(\cdot, \cdot)=\omega(\cdot, J \cdot)$. We also introduce normal coordinates $\left(x_{1}, x_{2}, \ldots, x_{2 p-1}, x_{2 p}\right)$ about $x_{0}$ in $U$ which can be chosen such that at $x_{0}$,

$$
\begin{equation*}
J_{x_{0}} \cdot \frac{\partial}{\partial x_{2 i+1}}=\frac{\partial}{\partial x_{2 i+2}} \quad \text { for } i=0 \ldots p-1 \tag{II.1}
\end{equation*}
$$

Since $C$ is a calibrated current in $(U, \omega, J)$, it is an area-minimizing current and its generalized mean curvature vanishes (see [All] or [Si]). One may isometrically embed $(U, g)$ into a euclidian space $\mathbb{R}^{2 p+k}$ and the generalised mean curvature of $C$ in $\mathbb{R}^{2 p+k}$ coincides with the mean curvature of the embedding of $(U, g)$ and is therefore a bounded function. Combining this fact together with the monotonicity formula (17.3) of [Si] we get that

$$
\begin{equation*}
\frac{\mathrm{M}\left(C\left\llcorner B_{r}\left(x_{0}\right)\right)\right.}{r^{2}}=f(r)+O(r) \tag{II.2}
\end{equation*}
$$

where $f(r)$ is an increasing function, M denotes the Mass of a current and $C\left\llcorner B_{r}\left(x_{0}\right)\right.$ is the restriction of $C$ to the geodesic ball of center $x_{0}$ and radius $r$. There exists in fact a constant $\alpha$ depending only on $g$ such that $e^{\alpha r} \frac{\mathrm{M}\left(C L B_{r}\left(x_{0}\right)\right)}{r^{2}}$ is an increasing function in $r$ (see [Si]). The factor $e^{\alpha r}$ is a perturbation of an order which will have no influence on the analysis below; therefore, by an abuse of notation we will often omit to write it and consider outright that $\frac{\mathrm{M}\left(C L B_{r}\left(x_{0}\right)\right)}{r^{2}}$ is an increasing function.

By means of the coordinates $\left(x_{1} \ldots x_{2 p}\right)$ we shall identify $U$ with a subdomain in $\mathbb{R}^{2 p}$ and use the same notation $C$ for the push forward of $C$ in $\mathbb{R}^{2 p}$ by this chart. For small radii $r$ we introduce the dilation function $\lambda^{r, x_{0}}(x)=\frac{x-x_{0}}{r}$, and we introduce the following dilation of $C$ about $x_{0}$ with rate $r$ as being the
following current in $\mathbb{R}^{2 p}$,

$$
\begin{equation*}
C_{r, x_{0}}:=\left(\lambda^{r, x_{0}}{ }_{*} C\right)\left\llcorner B_{1}^{2 p}(0) .\right. \tag{II.3}
\end{equation*}
$$

Observe that $r^{2} M_{0}\left(C_{r, x_{0}}\right)=M_{0}\left(C\left\llcorner B_{r}^{2 p}(0)\right)\right.$ where $M_{0}$ denotes the mass in the flat metric $g_{0}$ in $\mathbb{R}^{2 p}$. Since $g=g_{0}+O\left(r^{2}\right)$, we deduce from (II.2) that $M_{0}\left(C_{r, x_{0}}\right)$ is uniformly bounded as $r$ tends to zero. Again since $g$ and $g_{0}$ coincide up to the second order, it does not hurt in the analysis below if one mixes the notations for the two masses $M$ and $M_{0}$ and speaks only about $M$. Since now $C$ is a cycle in $U, \partial C_{r, x_{0}}\left\llcorner B_{1}^{2 p}(0)=0\right.$ and we can apply the FedererFleming compactness theorem to deduce that, from any sequence $r_{i} \rightarrow 0$ one can extract a subsequence $r_{i^{\prime}}$ such that $C_{r_{i^{\prime}}, x_{0}}$ converges in Flat norm to a limiting current $C_{0, x_{0}}$ called a tangent cone of $C$ at $x_{0}$. One of the purposes of the next section will be to establish that $C_{0, x_{0}}$ is independent of the subsequence and that the tangent cone is unique. The lower semi-continuity of the mass under weak convergence implies that

$$
\begin{equation*}
\lim _{r \rightarrow 0} \frac{\mathrm{M}\left(C\left\llcorner B_{r}\left(x_{0}\right)\right)\right.}{r^{2}}=\lim _{r \rightarrow 0} \mathrm{M}\left(C_{r, x_{0}}\right) \geq \mathrm{M}\left(C_{0, x_{0}}\right) . \tag{II.4}
\end{equation*}
$$

Now, from the fact that $C$ is calibrated by $\omega$ we deduce that the inequality (II.4) is an equality. Indeed

$$
\mathrm{M}\left(C_{r, x_{0}}\right)=r^{-2} C\left\llcorner B_{r}^{2 p}(0)(\omega)=C_{r, x_{0}}\left(r^{2}\left(\lambda^{r, x_{0}}\right)^{*} \omega\right) .\right.
$$

It is clear that $\lim _{r \rightarrow 0}\left\|r^{2}\left(\lambda^{r, x_{0}}\right)^{*} \omega-\omega_{0}\right\|_{\infty}$ where $\omega_{0}=\sum_{i=1}^{p} d x_{2 i-1} \wedge d x_{2 i}$. Therefore $C_{r, x_{0}}\left(r^{2}\left(\lambda^{r, x_{0}}\right)^{*} \omega-\omega_{0}\right) \longrightarrow 0$ and we get that

$$
\begin{equation*}
\lim _{r \rightarrow 0} \mathrm{M}\left(C_{r, x_{0}}\right)=\lim _{i^{\prime} \rightarrow+\infty} C_{r_{i^{\prime}}, x_{0}}\left(\omega_{0}\right)=C_{0, x_{0}}\left(\omega_{0}\right) \tag{II.5}
\end{equation*}
$$

Since the comass of $\omega_{0}$ is equal to $1, C_{0, x_{0}}\left(\omega_{0}\right) \leq \mathrm{M}\left(C_{0, x_{0}}\right)$. Combining this last fact with (II.4) and (II.5) we have established that

$$
\begin{equation*}
\lim _{r \rightarrow 0} \mathrm{M}\left(C_{r, x_{0}}\right)=\mathrm{M}\left(C_{0, x_{0}}\right)=C_{0, x_{0}}\left(\omega_{0}\right) \tag{II.6}
\end{equation*}
$$

which means in particular that $C_{0, x_{0}}$ is calibrated by the Kähler form $\omega_{0}$ in $\left(R^{2 p}, J_{0}\right) \simeq \mathbb{C}^{p}$ which is equivalent to the fact that $C_{0, x_{0}}$ is $J_{0}$-holomorphic. Using the explicit form of the monotonicity formula (see [Si] page 202), one observes that for any $s \in \mathbb{R}_{+}^{*}$

$$
C_{0, x_{0}}=\lambda_{*}^{s} C_{0, x_{0}}
$$

which means that for $\mathcal{H}^{2}$ almost everywhere on the carrier $\mathcal{C}_{0, x_{0}}$ of $C_{0, x_{0}}, \frac{\partial}{\partial r}$ is in the approximate tangent plane to $C_{0, x_{0}}$; in other words, $C_{0, x_{0}}$ is a cone. Since it is $J_{0}$-holomorphic, $\mathcal{H}^{2}$-a.e. $x$ in $\mathcal{C}_{0, x_{0}}$, the approximate tangent cone is given by

$$
T_{x} C_{0, x_{0}}=\operatorname{Span}\left\{\frac{\partial}{\partial r}, J_{0} \frac{\partial}{\partial r}\right\}
$$

Integral curves of $J_{0} \frac{\partial}{\partial r}$ are great-circles, fibers of the Hopf fibration

$$
\left(z_{1}=x_{1}+i x_{2}, \ldots, z_{p}\right) \longrightarrow\left[z_{1}, \ldots, z_{p}\right]
$$

Therefore we deduce that $C_{0, x_{0}}$ is the sum of the integrals over radial extensions of such great circles $\Gamma_{1} \ldots \Gamma_{Q}$ in $S^{2 p-1}$ which is the integral over a sum of $Q$ flat holomorphic disks. We adopt the following notation (in fact identical to the one used in [Wh]) for the radial extensions in $B_{1}^{2 p}(0)$ of currents supported in $\partial B_{1}^{2 p}(0)$ :

$$
C_{0, x_{0}}=\oplus_{i=1}^{Q} 0 \sharp \Gamma_{i} .
$$

Then we deduce that

$$
\begin{equation*}
\lim _{r \rightarrow 0} \mathrm{M}\left(C_{r, x_{0}}\right)=\pi Q \in \pi \mathbb{Z} \tag{II.7}
\end{equation*}
$$

For any $x \in U$ one denotes $Q_{x}$ the integer such that

$$
\lim _{r \rightarrow 0} \frac{\mathrm{M}\left(C\left\llcorner B_{r}(x)\right)\right.}{\pi r^{2}}=Q_{x}
$$

Using the monotonicity formula, it is straightforward to deduce that for any $Q \in \mathbb{N}$,

$$
\mathcal{C}_{Q}=\left\{x \in U: 0<Q_{x} \leq Q\right\}
$$

is an open subset of $\mathcal{C}_{*}=\left\{x \in U: 0<Q_{x}\right\}$. For $Q>1$, let us also denote

$$
\operatorname{Sing}^{Q}=\left\{x \in \mathcal{C}_{*}: Q_{x}=Q \quad \text { and } x \text { is an acc. point of } \mathcal{C}_{Q-1}\right\}
$$

Observe that, from Allard's theorem, it is clear that $C\left\llcorner\left(U \backslash \cup_{Q} \operatorname{Sing}^{Q}\right)\right.$ is the integral along a smooth surface with a smooth integer multiplicity. Although we won't make use of Allard's theorem this justifies a priori our notation. The whole purpose of our paper is to show that $\cup_{Q} \operatorname{Sing}^{Q}$ is made of isolated points. As we said, we won't make use of Allard's paper below since the relative Lipschitz estimate we establish in Lemma IV. 2 gives Allard's result in our case which is more specific. Because of this nice stratification of $C\left(\mathcal{C}_{Q}\right.$ is open in $\mathcal{C}_{*}$ ) we can argue by induction on $Q$. Let $\mathcal{P}_{Q}$ be the following assertion

$$
\begin{equation*}
\mathcal{P}_{Q}: \cup_{q \leq Q} \operatorname{Sing}^{q} \text { is made of isolated points } \tag{II.8}
\end{equation*}
$$

From the beginning of Section IV until Section VII we will assume either $Q=2$ or that $\mathcal{P}_{Q-1}$ holds and the goal will be to establish $\mathcal{P}_{Q}$.

## III. The uniqueness of the tangent cone

The uniqueness of the tangent cone means that the limiting cone $C_{0, x_{0}}$, obtained in the previous section while dilating at a point following a subsequence of radii $r_{i^{\prime}}$, is independent of the subsequence and is unique. Since our calibrated two-dimensional rectifiable cycle is area-minimizing, this fact
is a consequence of B . White upper-epiperimetric inequality in [Wh] (see also [Ri2] for the justification of the prefix "upper"). We need, however, a more quantitative version of this uniqueness of the tangent cone and express how far we are from the unique tangent cone in terms of the closedness of the density of area $\mathrm{M}\left(C\left\llcorner B_{r}\left(x_{0}\right)\right) / \pi r^{2}\right.$ to the limiting density $Q$. Precisely the goal of this section is to prove the following lemma:

Lemma III. 1 (Uniqueness of the tangent cone). For any $\varepsilon>0$ and $Q \in \mathbb{N}$ there exists $\delta>0$ and $\rho_{\varepsilon} \leq 1$ such that, for any compatible pair $(J, \omega)$ almost complex structure-symplectic form over $B_{1}^{2 p}(0)$ satisfying $J(0)=J_{0}(0), \omega(0)=$ $\omega_{0}(0)$,

$$
\begin{equation*}
\left\|J-J_{0}\right\|_{C^{2}\left(B_{1}\right)}+\left\|\omega-\omega_{0}\right\|_{C^{2}\left(B_{1}\right)} \leq \delta, \tag{III.1}
\end{equation*}
$$

for any J-holomorphic integral 2-cycle $C$ in $B_{1}(0)$ such that $Q_{0}=Q$, if

$$
\mathrm{M}\left(C\left\llcorner B_{1}^{2 p}(0)\right) \leq \pi Q+\delta\right.
$$

Then, there exist $Q J_{0}$ holomorphic flat discs $D_{1} \ldots D_{Q}$ passing through 0 , intersection of holomorphic lines of $\mathbb{C}^{p}$ with $B_{1}^{2 p}(0)$, such that, for any $\rho \leq \rho_{\varepsilon}$

$$
\begin{equation*}
\mathcal{F}\left(C_{\rho, 0}-\oplus_{i=1}^{Q} D_{i}\right) \leq \varepsilon \tag{III.2}
\end{equation*}
$$

and for any $\psi \in C_{0}^{\infty}\left(B^{1} \backslash\left\{x \in B^{1} ; \operatorname{dist}\left(x, \cup_{i} D_{i}\right) \leq \varepsilon|x|\right\}\right)$,

$$
\begin{equation*}
C_{\rho, 0}(\psi)=0 . \tag{III.3}
\end{equation*}
$$

Before proving Lemma III.1, we first establish the following intermediate result:

Lemma III.2. For any $\varepsilon>0$ and $Q \in \mathbb{N}$ there exists $\delta>0$ such that the following is true. If $(J, \omega)$ is a compatible pair of almost complex structuresymplectic form over $B_{1}^{2 p}(0)$ satisfying $J(0)=J_{0}(0), \omega(0)=\omega_{0}(0)$

$$
\left\|J-J_{0}\right\|_{C^{2}\left(B_{1}\right)}+\left\|\omega-\omega_{0}\right\|_{C^{2}\left(B_{1}\right)} \leq \delta,
$$

for any J-holomorphic integer rectifiable 2-cycle $C$ such that $Q_{x}=Q$, if

$$
\mathrm{M}\left(C\left\llcorner B_{1}^{2 p}(0)\right) \leq \pi Q+\delta\right.
$$

then, there exist $Q J_{0}$ holomorphic flat discs $D_{1} \ldots D_{Q}$ passing through 0 , intersection of holomorphic lines of $\mathbb{C}^{p}$ with $B_{1}^{2 p}(0)$, such that,

$$
\mathcal{F}\left(C\left\llcorner B_{1}(0)-\oplus_{i=1}^{Q} D_{i}\right) \leq \varepsilon\right.
$$

Remark III.1. Lemma III. 2 give much less information than Lemma III.1. Since a priori in Lemma III. 2 the disks $D_{i}$ may vary a lot as one dilates $C$ about 0 , whereas Lemma III. 1 controls such a tilting as one dilates the current further.

Proof of Lemma III.2. We prove Lemma III. 2 by contradiction. Assume there exist $\varepsilon_{0}>0, \delta_{n} \rightarrow 0$, compatible $J_{n}$ and $\omega_{n}$ and $C_{n}$ such that
i) $\quad\left\|J_{n}-J_{0}\right\|_{C^{2}}+\left\|\omega_{n}-\omega_{0}\right\|_{C^{2}} \leq \delta_{n}$,
ii) $\quad \lim _{r \rightarrow 0} \pi^{-1} r^{-2} \mathrm{M}\left(C_{n}\left\llcorner B_{r}(0)\right)=Q\right.$,
iii) $\quad \mathrm{M}\left(C_{n}\left\llcorner B_{1}\right) \leq \pi Q+\delta_{n}\right.$,
iv) $\quad \inf \left\{\mathcal{F}\left(C_{n}\left\llcorner B_{1}-\oplus_{i=1}^{Q} D_{i}\right)\right.\right.$ s.t. $D_{i}$ flat holom discs, $\left.0 \in D_{i}\right\} \geq \varepsilon_{0}$.

Since $\partial C_{n}\left\llcorner B_{1}=0\right.$ and since the mass of $C_{n}$ is uniformly bounded, one may assume, modulo extraction of a subsequence if necessary, that $C_{n}$ converges to a limiting rectifiable cycle $C_{\infty}$. Exactly as in Section III we have the fact that for any $0<r \leq 1$

$$
\begin{equation*}
\lim _{n \rightarrow+\infty} \mathrm{M}\left(C_{n}\left\llcorner B_{r}\right)=\mathrm{M}\left(C_{\infty}\left\llcorner B_{r}\right)=C_{\infty}\left\llcorner B_{r}\left(\omega_{0}\right) .\right.\right.\right. \tag{III.5}
\end{equation*}
$$

We deduce then that $C_{\infty}$ is calibrated by $\omega_{0}$ and is therefore a $J_{0}$-holomorphic cycle. Using ii) we deduce also that

$$
\lim _{r \rightarrow 0} \pi^{-1} r^{-2} \mathrm{M}\left(C_{\infty}\left\llcorner B_{r}(0)\right)=Q\right.
$$

and finally, from iii) and the lower semicontinuity of the mass, we have that $\mathrm{M}\left(C_{\infty}\left\llcorner B_{1}\right)=\pi Q\right.$ Thus, since $\pi^{-1} r^{-2} \mathrm{M}\left(C_{\infty}\left\llcorner B_{r}(0)\right)\right.$ is an increasing function, we have established that on $[0,1]$,

$$
\begin{equation*}
\pi^{-1} r^{-2} \mathrm{M}\left(C_{\infty} L B_{r}(0)\right) \equiv Q \tag{III.6}
\end{equation*}
$$

Let, for almost every $r, S_{\infty}^{r}=\left\langle C_{\infty}, \operatorname{dist}(\cdot, 0), r\right\rangle$ be the slice current obtained by slicing $C_{\infty}$ with $\partial B_{r}(0)$ (see [Fe, 4.2.1]). By Fubini, we have that for a.e. $0<r<1$

$$
\mathrm{M}\left(\left\langle C_{\infty}, \operatorname{dist}(\cdot, 0), r\right\rangle\right) \leq 2 \pi Q r
$$

Let $0 \sharp S_{\infty}^{r}$ be the radial extension of $S_{\infty}^{r}$ in $B_{r}(0)$.

$$
\mathrm{M}\left(0 \sharp S_{\infty}^{r}\right)=\frac{r}{2} \mathrm{M}\left(S_{\infty}^{r}\right)=\pi Q r^{2}=\mathrm{M}\left(C_{\infty}\left\llcorner B_{r}(0)\right) .\right.
$$

Since $\partial\left(C_{\infty}\left\llcorner B_{r}(0)-0 \sharp S_{\infty}^{r}\right)=0\right.$, and since $C_{\infty}\left\llcorner B_{r}(0)\right.$ is area-minimizing we have that $0 \sharp S_{\infty}^{r}$ is also area-minimizing. Let $\alpha$ be such that $d \alpha=\omega$ and

$$
\mathrm{M}\left(0 \sharp S_{\infty}^{r}\right)=\mathrm{M}\left(C_{\infty}\left\llcorner B_{r}(0)\right)=C_{\infty}\left\llcorner B_{r}(0)\left(\omega_{0}\right)=S_{\infty}^{r}(\alpha)=\left(0 \sharp S_{\infty}^{r}\right)\left(\omega_{0}\right) .\right.\right.
$$

Therefore $0 \sharp S_{\infty}^{r}$ is a holomorphic cone which is a cycle. So we deduce as in Section II that $0 \sharp S_{\infty}^{r}$ is a sum of flat holomorphic disks for any $r$. Thus $C_{\infty}$ is also a sum

$$
C_{\infty}\left\llcorner B_{1}(0)=\sum_{i=1}^{Q} D_{i}\right.
$$

where each $D_{i}$ is the intersection of a complex straight line in $\mathbb{C}^{p}$ with $B_{1}^{2 p}$. From the Federer-Fleming compactness theorem we have the fact that the weak convergence of $C_{n}$ to $C_{\infty}$ holds in flat norm

$$
\mathcal{F}\left(C_{n}\left\llcorner B_{1}-\sum_{i=1}^{Q} D_{i}\right) \longrightarrow 0\right.
$$

which contradicts iv), and Lemma III. 2 is proved.
Proof of Lemma. III.1. We first prove assertion (III.2) and recall Brian White's upper-epiperimetric inequality adapted to our present context: White's upper-epiperimetric inequality was proved for area-minimizing surfaces in $\mathbb{R}^{2 p}$. Here, in the present situation, we are dealing with area-minimizing currents which are $J$-holomorphic for a metric $g=\omega(\cdot, J \cdot)$ which gets as close as we want to the standard one because of assumption (III.1). Therefore very minor changes have to be provided to adapt White's theorem to the present context. An adaptation of the epiperimetric inequality for ambient nonflat metric is also given in [Ch, App. A]. So we have the following result.

Given an integer $Q$, there exists a positive number $\varepsilon_{Q}>0$, such that, for any compatible pair $\omega, J$ in $B_{2}^{2 p}(0)$ satisfying $\left\|\omega-\omega_{0}\right\|_{C^{2}\left(B_{2}\right)}+\left\|J-J_{0}\right\|_{C^{2}\left(B_{2}\right)} \leq$ $\varepsilon_{Q}$ and for any $C J$-holomorphic 2-rectifiable integral current in $B_{2}^{2 p}(0)$, satisfying $\partial C\left\llcorner B_{2}^{2 p}=0\right.$, assuming there exist $Q$ flat holomorphic disks $D_{1} \ldots D_{Q}$ in $\left(B_{1}^{2 p}(0), J\right) \simeq \mathbb{C}^{p} \cap B_{1}^{2 p}(0)$ passing through the origin such that

$$
\begin{equation*}
\mathcal{F}\left(C_{2,0}\left\llcorner B_{1}(0)-\sum_{i=1}^{Q} D_{i}\right) \leq \varepsilon_{Q}\right. \tag{III.7}
\end{equation*}
$$

(where we used a common notation for the oriented 2-disks $D_{i}$ and the corresponding 2-currents) we have

$$
\begin{equation*}
\mathrm{M}\left(C\left\llcorner B_{1}^{2 p}\right)-\pi Q \leq\left(1-\varepsilon_{Q}\right)\left(\frac{1}{2} M\left(\partial\left(C\left\llcorner B_{1}^{2 p}\right)\right)-\pi Q\right)\right.\right. \tag{III.8}
\end{equation*}
$$

Remark III.2. Observe that in the statement of the epiperimetric property in Definition 2 of [ Wh ] the epiperimetric constant $\varepsilon_{Q}$ may a priori also depend on the cone $\sum_{i=1}^{Q} D_{i}$. It is however elementary that this space of cones made of the intersection of $Q$ holomorphic straight lines passing through the origin with $B_{1}^{2 p}(0)$ is compact for the flat distance. Now by using a simple finite covering argument for this space of cones by balls (for the flat distance) one may obtain a constant $\varepsilon_{Q}>0$ for which the epiperimetric property holds independently of the cone $\sum_{i=1}^{Q} D_{i}$.

Once again we shall ignore the factor $e^{\alpha r}$ in front of $r^{-2} \mathrm{M}\left(C\left\llcorner B_{r}\right)\right.$ which induces lower order perturbations and argue as if $r^{-2} \mathrm{M}\left(C\left\llcorner B_{r}\right)\right.$ itself would
be an increasing function (observe also that $\alpha$ may be taken arbitrarily small because $J$ and $\omega$ are chosen as close as we want to $J_{0}$ and $\omega_{0}$ ).

Then $\varepsilon>0, \varepsilon<\varepsilon_{Q}$, and $\delta>0$, given by Lemma III. 2 for that $\varepsilon$. Assuming then $\mathrm{M}\left(C\left\llcorner B_{1}\right) \leq \pi Q+\delta\right.$ implies from the monotonicity formula that for any $r<1, r^{-2} \mathrm{M}\left(C\left\llcorner B_{r}(0)\right)=\mathrm{M}\left(C_{r, 0}\right) \leq(\pi Q+\delta)\right.$. Applying Lemma III. 2 to $C_{2 r, 0}$ for $r<1 / 2$ we deduce the existence of $Q$ flat disks $D_{1} \ldots D_{Q}$ such that

$$
\begin{equation*}
\mathcal{F}\left(C_{2 r, 0}-\sum_{i=1}^{Q} D_{i}\right) \leq \varepsilon \tag{III.9}
\end{equation*}
$$

We can then apply the epiperimetric inequality to $C_{r, 0}$ and get, after rescaling,

$$
\begin{equation*}
\mathrm{M}\left(C\left\llcorner B_{r}(0)\right)-\pi Q r^{2} \leq\left(1-\varepsilon_{Q}\right)\left(\frac{r}{2} \mathrm{M}\left(\partial\left(C\left\llcorner B_{r}(0)\right)\right)-\pi Q r^{2}\right)\right.\right. \tag{III.10}
\end{equation*}
$$

Denote $f(r)=\mathrm{M}\left(C\left\llcorner B_{r}(0)\right)-\pi Q r^{2}, f^{\prime}(r) \geq \mathrm{M}\left(\partial\left(C\left\llcorner B_{r}(0)\right)\right)-2 \pi Q r\right.\right.$. Therefore (III.10) implies

$$
\frac{1-\varepsilon_{Q}}{2} r f^{\prime}(r) \geq f(r)
$$

Integrating this differential inequality between $s$ and $\sigma(1 / 2>s>\sigma)$, we see that $f(s) \geq\left(\frac{s}{\sigma}\right)^{\frac{2}{1-\varepsilon_{Q}}} f(\sigma)$. When $\nu=\frac{2}{1-\varepsilon}-2>0$,

$$
\begin{equation*}
\frac{f(s)}{s^{2}} \geq\left(\frac{s}{\sigma}\right)^{\nu} \frac{f(\sigma)}{\sigma^{2}} \tag{III.11}
\end{equation*}
$$

Let $F(x)=\frac{x}{|x|}$. Then,

$$
\mathrm{M}\left(F_{*}\left(C\left\llcorner B_{s}(0) \backslash B_{\sigma}(0)\right)\right)=\int_{B_{S}(0) \backslash B_{\sigma}(0)} \frac{1}{|x|^{3}}\left|\tau \wedge \frac{x}{|x|}\right||\theta| d \mathcal{H}^{2}\llcorner\mathcal{C}\right.
$$

where $\tau$ denotes the unit 2-vector associated to the oriented approximate tangent plane to $C$ and is defined $\mathcal{H}^{2}$-a.e. along the carrier $\mathcal{C}$ of the rectifiable current, $\theta$ is the $L^{1}(\mathcal{C})$ integer-valued multiplicity of $C$ (i.e using classical GMT notations: $C=\langle\mathcal{C}, \theta, \tau\rangle)$ and $d \mathcal{H}^{2}\llcorner\mathcal{C}$ is the restriction to $\mathcal{C}$ of the 2-dimensional Hausdorff measure. Using Cauchy-Schwarz and 5.4.3 (2) of [Fe] (the explicit formulation of the monotonicity formula) and (III.11), we have

$$
\begin{align*}
& \mathrm{M}\left(F_{*}\left(C\left\llcorner B_{s}(0) \backslash B_{\sigma}(0)\right)\right)\right.  \tag{III.12}\\
& \leq\left[\frac{\mathrm{M}\left(C\left\llcorner B_{s}(0)\right)\right.}{s^{2}}-\frac{\mathrm{M}\left(C\left\llcorner B_{\sigma}(0)\right)\right.}{\sigma^{2}}\right]^{\frac{1}{2}}\left[\frac{\mathrm{M}\left(C\left\llcorner B_{s}(0)\right)\right.}{\sigma^{2}}\right]^{\frac{1}{2}} \\
& \leq\left[\frac{\mathrm{M}\left(C\left\llcorner B_{s}(0)\right)\right.}{s^{2}}-\pi Q\right]^{\frac{1}{2}}\left[\frac{\mathrm{M}\left(C\left\llcorner B_{s}(0)\right)\right.}{\sigma^{2}}\right]^{\frac{1}{2}} \\
& \leq\left[\frac{f(s)}{s^{2}}\right]^{\frac{1}{2}}\left[\frac{s^{2}}{\sigma^{2}} \frac{\mathrm{M}\left(C\left\llcorner B_{s}(0)\right)\right.}{s^{2}}\right]^{\frac{1}{2}} \leq K s^{\frac{\nu}{2}} \frac{s}{\sigma}
\end{align*}
$$

With $r<\rho<1 / 2$, applying (III.12) for $s=2^{-k} \rho$ and $\sigma=2^{-k-1} \rho$ for $k \leq \log _{2} \frac{\rho}{r}$ and summing over $k$ we get

$$
\begin{equation*}
\mathrm{M}\left(F_{*}\left(C\left\llcorner B_{\rho}(0) \backslash B_{r}(0)\right)\right) \leq C \rho^{\frac{\nu}{2}}\right. \tag{III.13}
\end{equation*}
$$

Observe that $\partial\left(F_{*}\left(C\left\llcorner B_{\rho}(0) \backslash B_{r}(0)\right)\right)=\partial C_{\rho, 0}-\partial C_{r, 0}\right.$. Therefore we deduce

$$
\begin{equation*}
\mathcal{F}\left(\left(C_{\rho, 0}-C_{r, 0}\right)\left\llcorner B_{1}(0) \backslash B_{\frac{1}{2}}(0)\right) \leq C \rho^{\frac{\nu}{2}}\right. \tag{III.14}
\end{equation*}
$$

Since
$\mathcal{F}\left(\left(C_{\rho, 0}-C_{r, 0}\right)\left\llcorner B_{\frac{1}{2}}(0) \backslash B_{\frac{1}{4}}(0)\right) \leq\left(\frac{1}{3}\right)^{\frac{1}{3}} \mathcal{F}\left(\left(C_{\frac{\rho}{2}, 0}-C_{\frac{r}{2}, 0}\right)\left\llcorner B_{1}(0) \backslash B_{\frac{1}{2}}(0)\right)\right.\right.$,
applying (III.14) for $\rho, r$ replaced by $2^{-k} \rho, 2^{-k} r$ and summing over $k=$ $1, \ldots, \infty$ we finally obtain

$$
\begin{equation*}
\mathcal{F}\left(\left(C_{\rho, 0}-C_{r, 0}\right)\left\llcorner B_{1}(0)\right) \leq C \rho^{\frac{\nu}{2}}\right. \tag{III.15}
\end{equation*}
$$

which is the desired inequality (III.2).
It remains to show (III.3) in order to finish the proof of Lemma III.1. We argue by contradiction. Assume there exists $\varepsilon_{0}>0, \rho_{n} \rightarrow 0$ and $\psi_{n} \in$ $C_{0}^{\infty}\left(\wedge^{2} B_{1}\right)$ such that

$$
\operatorname{supp} \psi_{n} \subset E_{0}=\left\{x \in B^{1} ; \operatorname{dist}\left(x, \cup_{i} D_{i}\right) \leq \varepsilon_{0}|x|\right\}
$$

where $C_{0,0}=\oplus_{i=1}^{Q} D_{i}$, and

$$
C_{\rho_{n}, 0}\left(\psi_{n}\right) \neq 0
$$

This latter fact implies in particular that there exists $x^{n} \in E_{0}$ such that $\lim _{r \rightarrow 0} \mathrm{M}\left(C_{r, x_{n}}\right) \neq 0$. Using the monotonicity formula we deduce then that

$$
\mathrm{M}\left(C_{\rho_{n}\left|x_{n}\right|, 0}\left\llcorner B_{\varepsilon_{0} / 2}\left(\frac{x_{n}}{\left|x_{n}\right|}\right)\right) \geq \frac{\pi}{4} \varepsilon_{0}^{2} .\right.
$$

We may then extract a subsequence such that $\frac{x_{n}}{\left|x_{n}\right|} \rightarrow x_{\infty}$. Thus,

$$
\mathrm{M}\left(C_{\rho_{n}\left|x_{n}\right|, 0}\left\llcorner B_{3 \varepsilon_{0} / 4}\left(x_{\infty}\right)\right) \geq \frac{\pi}{4} \varepsilon_{0}^{2}\right.
$$

Now,

$$
\mathrm{M}\left(C_{\rho_{n}\left|x_{n}\right|, 0}\left\llcorner B_{3 \varepsilon_{0} / 4}\left(x_{\infty}\right)\right)=C_{\rho_{n}\left|x_{n}\right|, 0}\left\llcorner B_{3 \varepsilon_{0} / 4}\left(x_{\infty}\right)\left({\frac{x}{\rho_{n}\left|x_{n}\right|}}^{*} \omega_{n}\right) .\right.\right.
$$

Since $\left\|\omega_{n}-\omega_{0}\right\|_{C^{2}} \rightarrow 0$ and since $\omega_{n}(0)=\omega_{0}(0)$ we clearly have that

$$
\left\|{\frac{x}{\rho_{n}\left|x_{n}\right|}}^{*} \omega_{n}-\omega_{0}\right\|_{\infty} \longrightarrow 0
$$

Therefore

$$
\begin{aligned}
& \left\lvert\, C_{\rho_{n}\left|x_{n}\right|, 0}\left\llcorner\left. B_{3 \varepsilon_{0} / 4}\left(x_{\infty}\right)\left({\frac{x}{\rho_{n}\left|x_{n}\right|}}^{*} \omega_{n}-\omega_{0}\right) \right\rvert\,\right.\right. \\
& \quad \leq \mathrm{M}\left(C_{\rho_{n}\left|x_{n}\right|, 0}\left\llcorner B_{3 \varepsilon_{0} / 4}\left(x_{\infty}\right)\right)\left\|\frac{x}{\rho_{n}\left|x_{n}\right|} \omega_{n}-\omega_{0}\right\|_{\infty} \longrightarrow 0\right.
\end{aligned}
$$

Thus

$$
C_{0,0}\left\llcorner B_{3 \varepsilon_{0} / 4}\left(x_{\infty}\right)\left(\omega_{0}\right)=\lim _{n \rightarrow+\infty} C_{\rho_{n}\left|x_{n}\right|, 0}\left\llcorner B_{3 \varepsilon_{0} / 4}\left(x_{\infty}\right)\left(\omega_{0}\right) \geq \frac{\pi}{4} \varepsilon_{0}^{2}\right.\right.
$$

which contradicts the fact that $B_{3 \varepsilon_{0} / 4}\left(x_{\infty}\right) \subset E_{0}$. Therefore (III.3) holds and Lemma III. 1 is proved.

## IV. Consequences of Lemma III.1. No accumulation of points in Sing ${ }^{Q}$ in the easy case. The relative Lipschitz estimate in the difficult case

In this section we expose two important consequences of Lemma III.1. Before explaining them we first observe that proving the implication $\mathcal{P}_{Q-1} \Longrightarrow$ $\mathcal{P}_{Q}$, will require considering two cases separately. The first case (the easy one) is the case where the tangent cone at the point $x_{0}$ of multiplicity $Q$ (i.e. $\left.\pi^{-1} r^{-2} \mathrm{M}\left(B_{r}\left(x_{0}\right)\right) \rightarrow Q\right)$ is not made of $Q$ times the same disk. The second case is the case where the tangent case is made of $Q$ times the same disk. In the first case we will deduce almost straight from Lemma III. 1 that such an $x_{0}$ cannot be an accumulation point of points of multiplicity $Q$ also; see Lemma IV. 1 below. In the second case, much more analysis will be needed to reach the same statement and this is the purpose of Sections IV through IX.

We can nevertheless deduce in this section an important consequence of our quantitative version of the uniqueness of the tangent cone (Lemma III.1) for the difficult case: this is the so called "relative Lipshitz estimate" (see Lemma IV. 2 below). This property says that, given a point $x_{0}$ of multiplicity $Q$ whose tangent cone is $Q$ times a flat disk and given an $\varepsilon>0$, there exists a radius $r_{\varepsilon, x_{0}}>0$ such that given any two points of $\mathcal{C}_{*} \cap B_{r_{\varepsilon, x_{0}}}\left(x_{0}\right)$, one of the two being also of multiplicity $Q$, the slope they realize relative to the tangent cone of $x_{0}$ is less than $\varepsilon$.

The condition that one of the two points has multiplicity $Q$ (this could be $x_{0}$ itself for instance) is a crucial assumption. It is indeed straightforward to find counterexamples to any Lipschitz estimates of multivalued graphs of holomorphic curves. Take for instance $w^{2}=z$ in $\mathbb{C}^{2} \simeq\{(z, w) z, w \in \mathbb{C}\}$ viewed as a 2 -valued graph over the line $\{w=0\}$, all points having multiplicity $1,(0,0)$ included of course, but the best possible estimate is a Hölder one $C^{0, \frac{1}{2}}$. We cannot exclude that such a configuration exists as we dilate at a point $x_{0}$ of mupltiplicity $Q>1$.

We first prove the following consequence of Lemma III. 1
Lemma IV. 1 (no accumulation - the easy case). Let $Q \in \mathbb{N}, Q \geq 2$. Let $x_{0}$ be a point in $\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\left(\right.$ i.e., $\pi^{-1} r^{-2} \mathrm{M}\left(B_{r}\left(x_{0}\right)\right) \rightarrow Q$ as $\left.r \rightarrow 0\right)$. Assume that the tangent cone at $x_{0}, C_{0, x_{0}}$, contains at least two different flat $J_{x_{0}}$-holomorphic disks (i.e., $C_{0, x_{0}} \neq Q D$ for the single flat $J_{x_{0}-h o l o m o r p h i c ~}$
disk $D$ ). Then, there exists $r>0$ such that

$$
B_{r}\left(x_{0}\right) \cap\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right)=\left\{x_{0}\right\}
$$

Proof of Lemma IV.1. Let $x_{0}$ be as in the statement of the lemma: $x_{0} \in$ $\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}$ and $C_{0, x_{0}}=\oplus_{i=1}^{K} Q_{i} D_{i}\left(\right.$ where $D_{i} \neq D_{j}$ for $i \neq j$ and $K>1$ ). Let $\varepsilon>0$ be a positive number smaller than $1 / 4 \mathrm{Q}$ times the maximal angle $\alpha$ between the various disks $D_{1}, \ldots, D_{K}$ in the tangent cones in such a way that there exists $i \neq j$ such that

$$
E_{\varepsilon}\left(D_{i}\right) \cap E_{\varepsilon}\left(D_{j}\right)=\left\{x_{0}\right\},
$$

where we use the following notation

$$
\left.E_{\varepsilon}\left(D_{i}\right)\left\{x \in \mathbb{R}^{2 p} ; \operatorname{dist}\left(x, D_{i}\right) \leq \varepsilon\left|x-x_{0}\right|\right\}\right) .
$$

By taking $\varepsilon$ as small as above, we even have ensured that $\cup E_{\varepsilon}\left(D_{i}\right) \backslash x_{0}$ has at least two connected components whose intersections with $\partial B_{1}\left(x_{0}\right)$ are at a distance larger than $\alpha / 2$. We now prove Lemma IV. 1 by contradiction. Assume there exists $x_{n} \in \mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}$ such that $x_{n} \rightarrow x_{0}$ and $x_{n} \neq x_{0}$. Let $\delta>0$ be given by Lemma III. 1 for $\varepsilon$ chosen as above. Let $\rho>0$ be such that $\rho^{-2} \mathrm{M}\left(C\left\llcorner B_{\rho}\left(x_{0}\right)\right) \leq \pi Q+\delta / 2\right.$. For any $x \in B_{\rho_{\frac{\delta}{4 \pi Q}}}\left(x_{0}\right)$ we have

$$
\begin{align*}
\mathrm{M}\left(C\left\llcorner B_{\rho\left(1-\frac{\delta}{4 \pi Q}\right)}(x)\right)\right. & \leq \mathrm{M}\left(C\left\llcorner B_{\rho}\left(x_{0}\right)\right) \leq \rho^{2}\left(\pi Q+\frac{\delta}{2}\right)\right.  \tag{IV.1}\\
& \leq\left(\rho\left(1-\frac{\delta}{4 \pi Q}\right)\right)^{2}\left(1-\frac{\delta}{4 \pi Q}\right)^{-2}\left(\pi Q+\frac{\delta}{2}\right) \\
& \leq\left(\rho\left(1-\frac{\delta}{4 \pi Q}\right)\right)^{2}(\pi Q+\delta) .
\end{align*}
$$

Choose then $x_{n} \in B_{\rho_{\frac{\delta}{8 \pi Q}}}\left(x_{0}\right)$. Applying (III.3) for $x_{0}$ we know that $x_{n}$ is contained in one of the $E_{\varepsilon}\left(D_{i}\right)$, say $E_{\varepsilon}\left(D_{1}\right)$. Denote $E_{1}$ the connected component of $\cup E_{\varepsilon}\left(D_{i}\right) \backslash\left\{x_{0}\right\}$ that contains $E_{\varepsilon}\left(D_{1}\right)$. We have chosen $\varepsilon$ small enough such that $\cup E_{\varepsilon}\left(D_{i}\right)$ has at least two connected components. Therefore we can chose $D_{j}$ such that $E_{\varepsilon}\left(D_{j}\right)$ is disjoint from the component containing $E_{\varepsilon}\left(D_{1}\right) \backslash\left\{x_{0}\right\}$. Let $\alpha$ be the angular distance, relative to $x_{0}$, from $E_{\varepsilon}\left(D_{j}\right)$ and the component containing $E_{\varepsilon}\left(D_{1}\right) \backslash\left\{x_{0}\right\}$. Clearly, al is bounded from below by a positive number as one chooses $\varepsilon$ smaller and smaller. Applying Lemma III. 1 this time to $x_{n}$, we know that in $B_{4\left|x_{n}\right|}\left(x_{n}\right) \backslash B_{\left|x_{n}\right|}\left(x_{n}\right)$ the support of $C$ is at the $\varepsilon\left|x_{n}\right|$ distance from a union of flat disks passing through $x_{n}$ (the tangent cone at $\left.x_{n}\right)$. This implies that the angular distance between the tangent cone at $x_{n}$ and $D_{1}$ is less than $C_{Q} \varepsilon$, where $C_{Q}$ depends on $Q$ only. Therefore
$($ IV.2 $) \operatorname{supp}\left(C\left\llcorner B_{2\left|x_{n}\right|}\left(x_{n}\right) \backslash B_{\left|x_{n}\right|}\left(x_{n}\right)\right) \subset \tilde{E}_{1}=\left\{x ; \operatorname{dist}\left(x, D_{1}\right) \leq C_{Q} \varepsilon\left|x_{n}\right|\right\}\right.$.
Observe that $\operatorname{dist}\left\{E_{\varepsilon}\left(D_{j}\right) \cap\left(B_{\left|x_{n}\right|}\left(x_{0}\right) \backslash B_{\left|x_{n}\right| / 2}\left(x_{0}\right)\right) ; \tilde{E}_{1}\right\} \geq \alpha / 4$. This later fact combined with (IV.2) contradicts (III.2). Lemma IV. 1 is then proved.

From now on until the beginning of Section X we will be dealing with the difficult case only: the case where the point $x_{0}$ of multiplicity $Q$ has a tangent cone which is made of $Q$ times the same disk. As we have been doing since Section II, we will work in a neighborhood of $x_{0}$ where a compatible simplectic form $\omega$ for $J$ exists, and we shall use normal coordinates for $g(\cdot, c d o t)=\omega(\cdot, J \cdot)$ about $x_{0}$, compatible with $J_{x_{0}}$ at $x_{0}$, satisfying (II.1); we can also assume that the tangent cone at $x_{0}$ is

$$
\begin{equation*}
C_{0, x_{0}}\left\llcorner B_{1}(0)=Q\left[D_{0}\right]\right. \tag{IV.3}
\end{equation*}
$$

where $D_{0}$ is the flat, oriented disk whose tangent 2 -vector is $\frac{\partial}{\partial x_{1}} \wedge \frac{\partial}{\partial x_{2}}$. From now on we also use the following notation for complex coordinates about $x_{0}$ :

$$
\begin{equation*}
z=x_{1}+i x_{2} \quad \text { and } \quad w_{i}=x_{2 k+1}+i x_{2 k+2} \text { for } k=1 \ldots p-1 \tag{IV.4}
\end{equation*}
$$

We will also denote $w=\left(w_{1}, \ldots, w_{p-1}\right)$. A second consequence to Lemma III. 1 is the following result:

Lemma IV. 2 (the relative Lipschitz estimate). Let $x_{0}$ be a point of multiplicity $Q$ (i.e. $x_{0} \in \mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}$ ), assume the tangent cone $C_{0, x_{0}}\left\llcorner B_{1}(0)\right.$ at $x_{0}$ is $Q$ times a flat disk (i.e. of the form (IV.3)). Let $\varepsilon>0$; then there exists $r_{\varepsilon, x_{0}}$ such that for any $r \leq r_{\varepsilon, x_{0}}$

$$
\begin{equation*}
\forall x \in B_{r}\left(x_{0}\right) \cap\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right), \quad \mathcal{F}\left(\left(C_{0, x}-C_{0, x_{0}}\right)\left\llcorner B_{1}(0)\right) \leq \varepsilon\right. \tag{IV.5}
\end{equation*}
$$

Also, for all $x=(z, w) \in B_{r}\left(x_{0}\right) \cap\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right)$ and $x^{\prime}=\left(z^{\prime}, w^{\prime}\right) \in B_{r}\left(x_{0}\right) \cap \mathcal{C}_{*}$ we have

$$
\begin{equation*}
\left|w-w^{\prime}\right| \leq \varepsilon\left|z-z^{\prime}\right| \tag{IV.6}
\end{equation*}
$$

Proof of Lemma IV.2. Let $\varepsilon>0$ and $\delta>0$ be given by Lemma III.1. Choose $r_{1}$ such that

$$
\mathrm{M}\left(C\left\llcorner B_{r_{1}}\left(x_{0}\right)\right) \leq r^{2}\left(\pi Q+\frac{\delta}{2}\right)\right.
$$

This implies in particular that for any $r<r_{1}$,

$$
\begin{equation*}
\mathcal{F}\left(\left(C_{r, x_{0}}-C_{0, x_{0}}\right)\left\llcorner B_{1}(0)\right) \leq \varepsilon .\right. \tag{IV.7}
\end{equation*}
$$

As in the proof of Lemma IV.1, see (IV.1), we have that for any $x \in$ $B_{r_{1} \frac{\delta}{4 \pi Q}}\left(x_{0}\right)$ and $r<r_{1}\left(1-\frac{\delta}{4 \pi Q}\right)$

$$
\mathrm{M}\left(C\left\llcorner B_{r}(x)\right) \leq r^{2}(\pi Q+\delta)\right.
$$

Then, letting $x \in B_{r_{1} \frac{\delta}{4 \pi Q Q}}\left(x_{0}\right) \cap\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right)$ and applying Lemma III.1, we have for $r<r_{1}\left(1-\frac{\delta}{4 \pi Q}\right)=r_{2}$,

$$
\begin{equation*}
\mathcal{F}\left(\left(C_{r, x}-C_{0, x}\right)\left\llcorner B_{1}(0)\right) \leq \varepsilon .\right. \tag{IV.8}
\end{equation*}
$$

Choose now $x \in B_{\varepsilon r_{1}}\left(x_{0}\right) \cap\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right)$. $C_{r_{2}, x}$ is fast eine $\varepsilon$-translation from $C_{r_{1}, x_{0}}$; therefore, since $\mathrm{M}\left(C_{r, x_{0}}\right) \leq 2 \pi Q$,

$$
\begin{equation*}
\mathcal{F}\left(\left(C_{r_{2}, x_{0}}-C_{r_{2}, x}\right)\left\llcorner B_{1}(0)\right) \leq 2 \pi Q .\right. \tag{IV.9}
\end{equation*}
$$

Taking $r_{\varepsilon, x_{0}}=\min \left\{\varepsilon r_{1}, \frac{\delta}{4 \pi Q} r_{1}\right\}$ and combining (IV.7), (IV.8) and (IV.9), we deduce that
(IV.10)

$$
\forall x \in B_{r_{\varepsilon, x_{0}}}\left(x_{0}\right) \cap\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right), \quad \mathcal{F}\left(\left(C_{0, x}-C_{0, x_{0}}\right)\left\llcorner B_{1}(0)\right) \leq(2+2 \pi Q) \varepsilon\right.
$$

It remains to check (IV.6) which is in fact an almost direct consequence of (III.3) and (IV.5). Lemma IV. 2 is then proved.

## V. The covering argument

Let $x_{0}$ be a point of multiplicity $Q>1$ whose tangent cone $C_{0, x_{0}}$ is $Q$ times the integral over the flat disk $D_{0}$ given by $w_{i}=0$ for $i \ldots Q-1$ (we use the system of coordinates introduced in the beginning of $\S I V$ ). The purpose of this section is to construct a Whitney-Besicovitch covering $B_{r_{i}}^{2}\left(z_{i}\right)$ of $\Pi\left(\mathcal{C}_{Q-1}\right) \cap B_{\rho}^{2}\left(x_{0}\right)$ where $\Pi$ is the projection on $D_{0}$ which gives the first complex coordinate of each point $\left(\Pi\left(z, w_{1} \ldots w_{p-1}\right)=z\right)$, for some radius $\rho$, small enough depending on $x_{0}$. This covering will be chosen in such a way that the following striking facts hold: first, $C\left\llcorner\Pi^{-1}\left(B_{r_{i}}^{2}\left(z_{i}\right)\right)\right.$ is in fact supported in a ball of radius $2 r_{i}, B_{2 r_{i}}^{2 p}\left(x_{i}\right)$, moreover $C\left\llcorner B_{2 r_{i}}^{2 p}\left(x_{i}\right)\right.$ is "split". This last word means that the flat distance between $C\left\llcorner\Pi^{-1}\left(B_{r_{i}}^{2}\left(z_{i}\right)\right)\right.$ and the $Q$ multiple of any single-valued graph over $D_{0}$ is larger than $K r_{i}^{3}$ where $K$ only depends on $p, J$ and $\omega$. This will come from the fact that $r_{i}$ may be chosen in such a way that $r_{i}^{-2} \mathrm{M}\left(B_{r_{i}}^{2 p}\left(x_{i}\right) \leq \pi-K^{\prime}\right.$ where again $K^{\prime}>0$ only depends on $p, Q$, $J$ and $\omega$. The existence of such a covering is a consequence of the "splitting before tilting" lemma proved in [Ri2].

Let $\alpha$ be given by Lemma V. 1 and let $\varepsilon>0$ be chosen small enough; compare to $\alpha$ later. Let $r_{\varepsilon, x_{0}}$ be the radius given by Lemma IV.2. We may choose also $r_{\varepsilon, x_{0}}$ small enough in such a way that

$$
\begin{equation*}
\forall r \leq r_{\varepsilon, x_{0}}, \quad \mathrm{M}\left(C_{r, x_{0}}\left\llcorner B_{1}(0)\right) \leq \pi Q+\varepsilon^{2} .\right. \tag{V.1}
\end{equation*}
$$

Using the proof of Lemma III. 1 (from (III.12) until the end of the proof), we deduce that

$$
\begin{equation*}
\forall r \leq r_{\varepsilon, x_{0}}, \quad \mathcal{F}\left(\left(C_{r, x_{0}}-C_{0, x_{0}}\right)\left\llcorner B_{1}(0)\right) \leq K \varepsilon\right. \tag{V.2}
\end{equation*}
$$

(In fact $\delta=O\left(\varepsilon^{2}\right)$ works in the statement of Lemma III.1.) On one hand, as in the proof of Lemma IV.1, we have that for any $x \in B_{\varepsilon^{2} r_{\varepsilon, x_{0}}}\left(x_{0}\right)$ and $r \leq r_{\varepsilon, x_{0}}\left(1-\varepsilon^{2}\right)$,

$$
\begin{equation*}
\mathrm{M}\left(C\left\llcorner B_{r}(x)\right) \leq r^{2}\left(\pi Q+\varepsilon^{2}\right)\right. \tag{V.3}
\end{equation*}
$$

On the other hand, arguing as in the proof of Lemma IV.2, between (IV.8) and (IV.10), we have, using also (V.23),

$$
\begin{equation*}
\forall x \in B_{\varepsilon^{2}}\left(x_{0}\right), \quad \mathcal{F}\left(C_{\frac{r_{\varepsilon}, x_{0}}{2}, x_{0}}\left\llcorner B_{1}(0)-\pi Q\left[D_{0}\right]\right) \leq K \varepsilon .\right. \tag{V.4}
\end{equation*}
$$

Having chosen $K \varepsilon<\alpha$ we are in a position to apply the "Splitting before tilting" lemma of [Ri2] which is a key step in our proof of the regularity of 1-1 rectifiable cycles.

Lemma V. 1 (splitting before tilting, [Ri2]). There exists $\alpha>0$ such that for any $x_{0} \in \mathcal{C}_{Q-1}$ and for any radius $0<\rho<\alpha$,

$$
\begin{equation*}
\mathcal{F}\left(C_{2 \rho, x_{0}}\left\llcorner B_{1}(0)-Q\left[D_{0}\right]\right) \leq \alpha\right. \tag{V.5}
\end{equation*}
$$

where $D_{0}$ is a flat $J_{x_{0}}$-holomorphic disk passing through $x_{0}$. Then, for any $r<\rho$ and any $J_{x_{0}}$-holomorphic flat disk, $D_{1}$, passing through $x_{0}$ and satisfying

$$
\begin{equation*}
\mathcal{F}\left(\left[D_{0}\right]-\left[D_{1}\right]\right) \geq \frac{1}{4} \tag{V.6}
\end{equation*}
$$

we have

$$
\begin{equation*}
\mathcal{F}\left(C_{r, x_{0}}\left\llcorner B_{1}(0)-Q\left[D_{1}\right]\right) \geq \alpha\right. \tag{V.7}
\end{equation*}
$$

Moreover, there exist $r_{0}<\rho$ and $K_{0}$ a constant depending only on $\|\omega\|_{C^{1}}$ and $\varepsilon_{Q}^{ \pm}$, the epiperimetric constants, such that

$$
\begin{align*}
& \mathrm{M}\left(C_{r_{0}, x_{0}}\left\llcorner B_{1}(0)\right)\right.=\pi Q-K_{0} \alpha,  \tag{V.8}\\
& \mathcal{F}\left(C_{r_{0}, x_{0}}\left\llcorner B_{1}(0)-Q\left[D_{0}\right]\right) \leq K \sqrt{\alpha}\right. \tag{V.9}
\end{align*}
$$

for some constant $K$ depending also only on $\|\omega\|_{C^{1}}$ and $\varepsilon_{Q}^{ \pm}$. Finally, for any $J_{x_{0}}$-holomorphic disk $D$ passing through 0 ,

$$
\begin{equation*}
\forall r \leq r_{0}, \quad \mathcal{F}\left(C_{r, x_{0}}\left\llcorner B_{1}(0)-Q[D]\right) \geq \alpha\right. \tag{V.10}
\end{equation*}
$$

For any $x \in \mathcal{C}_{p-1} \cap B_{\varepsilon^{2} r_{\varepsilon, x_{0}}}\left(x_{0}\right)$ we denote by $r_{x}$ the radius $r_{0}$ given by the lemma. We then have

$$
\begin{gather*}
\mathrm{M}\left(C_{r_{x}, x}\left\llcorner B_{1}(0)\right)=\pi Q-K_{0} \alpha,\right.  \tag{V.11}\\
\mathcal{F}\left(C_{r_{x}, x}\left\llcorner B_{1}(0)-\pi Q\left[D_{0}\right]\right) \leq K \sqrt{\alpha}\right. \tag{V.12}
\end{gather*}
$$

for some constant $K$ depending only on $\|\omega\|_{C^{1}}$ and $\varepsilon_{Q}^{ \pm}$. Moreover, for $\alpha$ chosen small enough and $\varepsilon$ small enough compared to $\alpha$, the following lemma holds

Lemma V.2. Under the above notation, for any $x \in \mathcal{C}_{p-1} \cap B_{\varepsilon r_{\varepsilon, x_{0}}}\left(x_{0}\right)$,

$$
\begin{equation*}
\operatorname{supp}\left(C \left\llcorner\Pi^{-1}\left(B_{r_{x}}^{2}(\Pi(x)) \cap B_{r_{\varepsilon, x_{0}}}\left(x_{0}\right)\right) \subset B_{r_{x}}^{2}(\Pi(x)) \times B_{r_{x}}^{2 p-2}(0),\right.\right. \tag{V.13}
\end{equation*}
$$

and

$$
\begin{equation*}
\operatorname{supp}\left(C \left\llcorner\Pi^{-1}\left(B_{r_{x}}^{2}(\Pi(x)) \cap B_{r_{\varepsilon, x_{0}}}\left(x_{0}\right)\right) \subset \mathcal{C}_{p-1}\right.\right. \tag{V.14}
\end{equation*}
$$

Proof of Lemma V.2. We claim that for any $r$ between $\frac{r_{\varepsilon, x_{0}}}{2}$ and $r_{x}$ one has

$$
\begin{equation*}
\operatorname{supp}\left(C_{r, x}\left\llcorner B_{1}(0)\right) \subset E\left(\alpha^{\frac{1}{16}}\right)\right. \tag{V.15}
\end{equation*}
$$

where we use the notation

$$
\begin{equation*}
E(\lambda)=\left\{y=(z, w) \in B_{1}(0): \quad|w| \leq \lambda\right\} . \tag{V.16}
\end{equation*}
$$

We show (V.15) arguing by contradiction. First of all from the proof of Lemma V. 1 applied to $x$ we have the fact that for any $r \in\left[r_{x}, \frac{r_{e, x_{0}}}{2}\right]$

$$
\begin{equation*}
\mathcal{F}\left(C_{r, x}\left\llcorner B_{1}(0)-\pi Q\left[D_{0}\right]\right) \leq K \sqrt{\alpha} .\right. \tag{V.17}
\end{equation*}
$$

Let $\omega_{\alpha}=\chi_{\alpha} \omega=\chi\left(\frac{|w|}{\alpha^{\frac{1}{4}}}\right) \omega$ where $\chi$ is a smooth cut-off function on $\mathbb{R}_{+}$ satisfying $\chi \equiv 1$ on $[0,1 / 2]$ and $\chi \equiv 0$ in $[1,+\infty)$. Let $S$ and $R$ be a 3 and a 2-current satisfying $\left(C_{r, x}\left\llcorner B_{1}(0)-\pi Q\left[D_{0}\right]\right)=\partial S+R\right.$ with $\mathrm{M}(S)+\mathrm{M}(R) \leq$ $2 K \sqrt{\alpha}$. Now,

$$
\begin{aligned}
\mid\left(C_{r, x}\left\llcorner B_{1}(0)-\pi Q\left[D_{0}\right]\right)\left(\omega_{\alpha}\right) \mid\right. & =\left|S\left(\omega \wedge d \chi_{\alpha}\right)+R\left(\omega_{\alpha}\right)\right| \\
& \leq\left\|\nabla \chi_{\alpha}\right\|_{\infty}\left\|\omega_{\infty}\right\| \mathrm{M}(S)+\left\|\omega_{\alpha}\right\|_{\infty} \mathrm{M}(R) \leq K \alpha^{\frac{1}{4}}
\end{aligned}
$$

Thus we get in particular

$$
\begin{align*}
\pi Q-K \alpha^{\frac{1}{4}} \leq \mid C_{r, x}\left\llcorner B_{1}(0)\left(\omega_{\alpha}\right) \mid\right. & \leq \mathrm{M}\left(C_{r, x}\left\llcorner B_{1}(0) \cap E\left(\alpha^{\frac{1}{4}}\right)\right)\left\|\omega_{\alpha}\right\|_{\infty}\right.  \tag{V.18}\\
& \leq \mathrm{M}\left(C_{r, x}\left\llcorner B_{1}(0) \cap E\left(\alpha^{\frac{1}{4}}\right)\right) .\right.
\end{align*}
$$

Assume now that there exists $y \in\left(\mathcal{C}_{r, x}\right)_{*} \cap B_{1}(0) \cap\left(\mathbb{R}^{2 p} \backslash E\left(\alpha^{\frac{1}{16}}\right)\right)$. From the monotonicity formula we deduce that

$$
\begin{equation*}
\mathrm{M}\left(C\left\llcorner B_{\frac{\alpha \frac{1}{16}}{2} r}(y)\right) \geq \frac{\pi}{4} \alpha^{\frac{1}{8}} r^{2} .\right. \tag{V.19}
\end{equation*}
$$

Combining (V.18) and (V.19), we obtain

$$
\begin{equation*}
\mathrm{M}\left(C\left\llcorner B_{r}(x)\right) \geq r^{2}\left(\pi Q-K \alpha^{\frac{1}{4}}+\frac{\pi}{4} \alpha^{\frac{1}{8}}\right) .\right. \tag{V.20}
\end{equation*}
$$

For $\alpha$ small enough (V.20) contradicts (V.11) and (V.15) holds true for any $r \in\left[r_{x}, \frac{r_{\varepsilon, x_{0}}}{2}\right]$. From this latter fact one deduces (V.13).

It remains to prove (V.14). Again we argue by contradiction. Assume there exists $y \in\left(\mathcal{C}_{p} \backslash \mathcal{C}_{p-1}\right) \cap \Pi^{-1}\left(B_{r_{x}}^{2}(\Pi(x)) \cap B_{r_{e, x_{0}}}\left(x_{0}\right)\right.$. Because of (V.3) and since $y \in B_{r_{x_{0}, \varepsilon}}\left(x_{0}\right)$ we can apply Lemma IV. 2 to $y$ in order to deduce that $\mathcal{C}_{\star} \cap B_{r_{x}}(x)$ is included in a cone of center $y$, axis parallel to $D_{0}$ and angle $\varepsilon$. This cone of course contains $x$ and then we can deduce that

$$
\begin{equation*}
\operatorname{supp}\left(C_{r_{x}, x}\left\llcorner B_{1}(0)\right) \subset E(4 \varepsilon)\right. \tag{V.21}
\end{equation*}
$$

(The notation $E(\lambda)$ is introduced in (V.16)). We have $\partial C_{r, x}\left\llcorner B_{1}(0)=0\right.$; moreover, because of (V.17); for $\alpha$ small enough we deduce that the intersection
number of $C\left\llcorner B_{r_{x}}^{2}(\Pi(x)) \times B_{r_{x}}^{2 p-2}(0)\right.$ with any vertical current $\Pi^{-1}(z)$ for $z \in$ $B_{r_{x}}^{2}$ is $Q$. Combining this fact with (V.21), using Fubini, one deduces that

$$
\begin{equation*}
\mathrm{M}\left(C_{r_{x}, x}\left\llcorner B_{1}(0)\right) \geq \pi Q-O\left(\varepsilon^{2}\right) .\right. \tag{V.22}
\end{equation*}
$$

For $\varepsilon$ small enough compared to $\alpha$ we get a contradiction while comparing (V.22) and (V.11) and (V.14) is proved. This concludes the proof of Lemma V.2.

In the following second lemma of this section, we show that the covering $\left(B_{r_{x}}^{2}(\Pi(x))\right)_{x \in \mathcal{C}_{p-1} \cap B_{\varepsilon^{2} r_{\varepsilon, x_{0}}}^{2}\left(x_{0}\right)}$ of $\Pi\left(\mathcal{C}_{p-1} \cap B_{\varepsilon^{2} r_{\varepsilon, x_{0}}}^{2 p}\left(x_{0}\right)\right.$ has the "Whitney" property: two balls intersecting each-other have comparable size. From now on we adopt the following notation granting the fact that $\alpha$ and $\varepsilon$ are fixed small enough for the constraint mentionend above to be fulfilled:

$$
\begin{equation*}
\rho_{x_{0}}:=\varepsilon^{2} r_{\varepsilon, x_{0}} . \tag{V.23}
\end{equation*}
$$

Precisely we have:
Lemma V. 3 (Whitney property of the covering). There exists $\gamma>0$ depending only on $Q$ such that, given $x_{0} \in \mathcal{C}_{p} \backslash \mathcal{C}_{p-1}$ whose tangent cone is $Q\left[D_{0}\right]$ and letting $\left(B_{r_{x}}^{2}(\Pi(x))\right)$ for $x \in \mathcal{C}_{p-1} \cap B_{\rho_{x_{0}}}^{2}\left(x_{0}\right)$ the covering of $\Pi\left(\mathcal{C}_{p-1} \cap\right.$ $B_{\rho_{x_{0}}}^{2 p}\left(x_{0}\right)$ described above, assuming for some $x, y \in \mathcal{C}_{p-1} \cap B_{\rho_{x_{0}}}^{2}\left(x_{0}\right)$

$$
B_{r_{x}}^{2}(x) \cap B_{r_{y}}^{2}(y) \neq \emptyset,
$$

then

$$
\begin{equation*}
r_{x} \geq \alpha^{\gamma} r_{y} . \tag{V.24}
\end{equation*}
$$

Proof of Lemma V.3. This lemma is again a consequence of the upper and lower-epiperimetric inequalities. Assume for instance that $r_{x} \leq r_{y}$. From (V.10) we have

$$
\begin{equation*}
\mathcal{F}\left(C_{r_{y}, y}\left\llcorner B_{1}(0)-Q\left[D_{0}\right]\right) \geq \alpha\right. \tag{V.25}
\end{equation*}
$$

Which implies that for all $r \leq r_{y}$

$$
\begin{equation*}
\mathcal{F}\left(C\left\llcorner B_{r}^{2}\left(z_{y}\right) \times B_{\rho_{x_{0}}}^{2 p}(0)-Q\left[B_{r}^{2}\left(z_{y}\right) \times\{0\}\right]\right) \geq \alpha r^{3}\right. \tag{V.26}
\end{equation*}
$$

where $y=\left(z_{y}, w_{y}\right)$. Since $\left|z_{x}-z_{y}\right| \leq 2 \max \left\{r_{x}, r_{y}\right\}=2 r_{y}$ and $B_{r_{y}}\left(z_{y}\right) \subset$ $B_{3 r_{y}}\left(z_{x}\right)$, (V.26) implies that

$$
\begin{equation*}
\mathcal{F}\left(C\left\llcorner B_{4 r_{y}}^{2}\left(z_{x}\right) \times B_{\rho_{x_{0}}}^{2 p}(0)-Q\left[B_{4 r_{y}}^{2}\left(z_{x}\right) \times\left\{w_{y}\right\}\right]\right) \geq \frac{\alpha}{3} r_{y}^{3} .\right. \tag{V.27}
\end{equation*}
$$

This passage from (V.26) to (V.27) is obtained by applying a Fubini type argument. Indeed, let $A=C\left\llcorner B_{4 r_{y}}^{2}\left(z_{x}\right) \times B_{\rho_{x_{0}}}^{2 p}(0)-Q\left[B_{4 r_{y}}^{2}\left(z_{x}\right) \times\left\{w_{y}\right\}\right]\right.$ and let
$S$ and $R$ be such that $A=\partial S+R$ and $\mathrm{M}(S)+\mathrm{M}(R)^{\frac{3}{2}} \leq 2 \mathcal{F}(A)$. For almost every $r$ in $\left[r_{y} / 2, r_{y}\right]$ we have

$$
\partial\left(S\left\llcorner B_{r}^{2}\left(z_{y}\right) \times B_{\rho_{x_{0}}}^{2 p}(0)\right)=\partial S\left\llcorner B_{r}^{2}\left(z_{y}\right) \times B_{\rho_{x_{0}}}^{2 p}(0)\right)+\left\langle S, \operatorname{dist}\left(\cdot,\left\{z=z_{y}\right\}, r\right\rangle\right.\right.
$$

where $\left\langle S\right.$, $\left.\operatorname{dist}\left(\cdot,\left\{z=z_{y}\right\}\right), r\right\rangle$ is the slice current between $S$ and the boundary of the cylinder $B_{r}^{2}\left(z_{y}\right) \times B_{\rho_{x_{0}}}^{2 p}(0)$ and $\operatorname{dist}\left(\cdot,\left\{z=z_{y}\right\}\right)$ denotes the distance function to the axis of this cylinder (see [Fe, 4.2.1, pp. 395...]). Thus

$$
\begin{align*}
A\left\llcorner B_{r}^{2}\left(z_{y}\right) \times B_{\rho_{x_{0}}}^{2 p}(0)=\right. & p\left(S\left\llcorner B_{r}^{2}\left(z_{y}\right) \times B_{\rho_{x_{0}}}^{2 p}(0)\right)\right.  \tag{V.28}\\
& -\left\langle S, \operatorname{dist}\left(\cdot,\left\{z=z_{y}\right\}, r\right\rangle+R\left\llcorner B_{r}^{2}\left(z_{y}\right) \times B_{\rho_{x_{0}}}^{2 p}(0) .\right.\right.
\end{align*}
$$

We have, see [Fe, 4.2.1, p. 395],

$$
\begin{equation*}
\int_{\frac{r_{y}}{2}}^{r_{y}} M\left(\left\langle S, \operatorname{dist}\left(\cdot,\left\{z=z_{y}\right\}, r\right\rangle\right) \leq M\left(S\left\llcorner\left(B_{r_{y}}^{2} \backslash B_{\frac{r_{y}}{2}}^{2}\right) \times B_{\rho_{x_{0}}}^{2 p}(0)\right) .\right.\right. \tag{V.29}
\end{equation*}
$$

Using Fubini's theorem we may then find $r=r_{1} \in\left[r_{y} / 2, r_{y}\right]$ such that

$$
\begin{equation*}
M\left(\left\langle S, \operatorname{dist}\left(\cdot,\left\{z=z_{y}\right\}, r_{1}\right\rangle\right) \leq \frac{2}{r_{y}} M\left(S\left\llcorner\left(B_{r_{y}}^{2} \backslash B_{\frac{r_{y}}{2}}^{2}\right) \times B_{\rho_{x_{0}}}^{2 p}(0)\right) \leq \frac{2}{r_{y}} \mathrm{M}(S) .\right.\right. \tag{V.30}
\end{equation*}
$$

Combining (V.28), (V.29) and (V.30) we deduce that

$$
\begin{align*}
& \mathcal{F}\left(C\left\llcorner B_{r_{1}}^{2}\left(z_{y}\right) \times B_{\rho_{x_{0}}}^{2 p}(0)-Q\left[B_{r_{1}}^{2}\left(z_{y}\right) \times\{0\}\right]\right)\right.  \tag{V.31}\\
& \quad=\mathcal{F}\left(A\left\llcorner B_{r}^{2}\left(z_{y}\right) \times B_{\rho_{x_{0}}}^{2 p}(0)\right) \leq \mathrm{M}(S)+\left(\mathrm{M}(R)+\frac{2}{r_{y}} \mathrm{M}(S)\right)^{\frac{3}{2}} .\right.
\end{align*}
$$

Thus, combining (V.26) for $r=r_{1}$ and (V.31), we have

$$
\mathrm{M}(S)+\left(\mathrm{M}(R)+\frac{2}{r_{y}} \mathrm{M}(S)\right)^{\frac{3}{2}} \geq \alpha r_{y}^{3}
$$

and since

$$
\mathcal{F}\left(C\left\llcorner B_{4 r_{y}}^{2}\left(z_{x}\right) \times B_{\rho_{x_{0}}}^{2 p}(0)-Q\left[B_{4 r_{y}}^{2}\left(z_{x}\right) \times\left\{w_{y}\right\}\right]\right) \geq \frac{1}{2}\left[\mathrm{M}(S)+\mathrm{M}(R)^{\frac{3}{2}}\right]\right.
$$

we obtain (V.27). Therefore we deduce that

$$
\begin{equation*}
\mathcal{F}\left(C_{4 r_{y}, x}\left\llcorner B_{1}(0)-Q\left[D_{0}\right]\right) \geq \frac{1}{3 \times 4^{3}} \alpha .\right. \tag{V.32}
\end{equation*}
$$

Let $\frac{\rho_{x_{0}}}{\varepsilon^{2}}>s_{x}>r_{x}$ be such that

$$
\begin{equation*}
\mathrm{M}\left(C_{s_{x}, x}\left\llcorner B_{1}(0)\right)=\pi Q\right. \tag{V.33}
\end{equation*}
$$

Because of (V.3), arguing as in the proof of Lemma IV.2, between (IV.8) and (IV.10), we have

$$
\begin{equation*}
\forall s_{x} \leq r \leq \frac{\rho_{x_{0}}}{\varepsilon^{2}}, \quad \mathcal{F}\left(C_{r, x}\left\llcorner B_{1}(0)-Q\left[D_{0}\right]\right) \leq K \varepsilon\right. \tag{V.34}
\end{equation*}
$$

Assuming, as we did above that $\alpha \gg \varepsilon$, comparing (V.32) and (V.26) we deduce that $s_{x}>4 r_{y}$. Let $\lambda=4 \frac{r_{x}}{r_{y}}$. From the proof of Lemma V.1, in fact from (V.9) precisely, for any $r \in\left[r_{x}, s_{x}\right]$ we have

$$
\begin{equation*}
\mathcal{F}\left(C_{2 r, x}\left\llcorner B_{1}(0)-\pi Q\left[D_{0}\right]\right) \leq K \sqrt{\alpha} \leq \varepsilon_{Q}^{-},\right. \tag{V.35}
\end{equation*}
$$

which means in particular that we are in the position to apply the lowerepiperimetric inequality and the differential inequality (III.27) in [Ri2] deduced from it. Integrating then this inequality between $r_{x}$ and $4 r_{y}$ we have

$$
\begin{align*}
\pi Q-\mathrm{M}\left(C_{4 r_{y}, x}\left\llcorner B_{1}(0)\right)\right. & \leq\left(\frac{1}{2^{2 \varepsilon_{Q}^{-}}}\right)^{\log _{2} \lambda}\left[\pi Q-\mathrm{M}\left(C_{r_{x}, x}\left\llcorner B_{1}(0)\right)\right]\right.  \tag{V.36}\\
& =\left(\frac{1}{2^{2 \varepsilon_{Q}^{-}}}\right)^{\log _{2} \lambda} \alpha .
\end{align*}
$$

Using now (III.28) from [Ri2], we deduce from (V.36) that

$$
\begin{align*}
\mathcal{F}\left(\left(C_{s_{x}, x}-C_{4 r_{y}, x}\right)\left\llcorner B_{1}(0)\right)\right. & \leq \sqrt{\pi Q-\mathrm{M}\left(C_{4 r_{y}, x}\left\llcorner B_{1}(0)\right)\right.}  \tag{V.37}\\
& =\left(\frac{1}{2^{\varepsilon_{Q}^{-}}}\right)^{\log _{2} \lambda} \alpha^{\frac{1}{2}} .
\end{align*}
$$

Combining (V.34) and (V.37) one gets that

$$
\begin{equation*}
\mathcal{F}\left(C_{4 r_{y}, x}\left\llcorner B_{1}(0)-Q\left[D_{0}\right]\right) \leq\left(\frac{1}{2^{\varepsilon_{Q}^{-}}}\right)^{\log _{2} \lambda} \alpha^{\frac{1}{2}}+K \varepsilon .\right. \tag{V.38}
\end{equation*}
$$

Comparing (V.32) and (V.38) we obtain

$$
\begin{equation*}
\frac{1}{3 \times 4^{3}} \alpha \leq\left(\frac{1}{2^{\varepsilon_{Q}^{-}}}\right)^{\log _{2} \lambda} \alpha^{\frac{1}{2}} . \tag{V.39}
\end{equation*}
$$

Since $K \varepsilon \leq \frac{1}{6 \times 4^{3}} \alpha$ we have $\frac{1}{6 \times 4^{3}} \alpha \leq\left(\frac{1}{2^{\varepsilon^{-}}}\right)^{\log _{2} \lambda} \alpha^{\frac{1}{2}}$. Taking the log of this last inequality we obtain

$$
\frac{\log \lambda}{\log 2} \varepsilon_{Q}^{-} \log \frac{1}{2}+\frac{1}{2} \log \alpha \geq \log \alpha-\log \left(6 \times 4^{3}\right) .
$$

Thus

$$
\frac{1}{2} \log \frac{1}{\alpha}+\log \left(6 \times 4^{3}\right) \geq \varepsilon_{Q}^{-} \log \lambda
$$

By taking $\alpha$ small enough, we may always assume that $\log \frac{1}{\alpha} \geq 4 \times \log \left(6 \times 4^{3}\right)$ and we finally get that

$$
\frac{1}{\varepsilon_{Q}^{-}} \log \frac{1}{\alpha} \geq \log \lambda
$$

which leads to the desired inequality (V.24) and, now, Lemma V. 3 is proved. $\square$

Constructing a partition of unity adapted to the covering. From the covering $\left(B_{r_{x}}^{2}(x)\right)$ for $x \in \mathcal{C}_{p-1} \cap B_{\rho_{x_{0}}}\left(x_{0}\right)$ of $\Pi\left(\mathcal{C}_{p-1} \cap B_{\rho_{x_{0}}}\left(x_{0}\right)\right)$ we extract a Besicovitch covering $\left(B_{r_{x_{i}}}^{2}\left(x_{i}\right)\right)$ for $i \in I$ ( $I$ is a countable set) of $\Pi\left(\mathcal{C}_{p-1} \cap\right.$ $\left.B_{\rho_{x_{0}}}\left(x_{0}\right)\right)$ that is a covering such that

$$
\begin{equation*}
\forall z \in B_{\rho_{x_{0}}}^{2}\left(x_{0}\right), \quad \operatorname{Card}\left\{i \in I: z \in B_{r_{x_{i}}}^{2}\left(x_{i}\right)\right\} \leq n \tag{V.40}
\end{equation*}
$$

where $N$ is some universal number (see [Fe]). To simplify the notation we will simply write $r_{i}$ for $r_{x_{i}}$. Note that since balls intersecting each other have comparable size (see Lemma V.3), each ball $B_{r_{i}}^{2}\left(z_{i}\right)$ intersects a uniformly bounded number of other balls: there exists $N_{Q, \alpha}$ such that

$$
\begin{equation*}
\forall i \in I, \quad \operatorname{Card}\left\{j \in I: B_{r_{j}}^{2}\left(z_{j}\right) \cap B_{r_{i}}^{2}\left(z_{i}\right) \neq \emptyset\right\} \leq N_{Q, \alpha} . \tag{V.41}
\end{equation*}
$$

We now construct a partition of unity adapted to a slightly modified covering. Considering the covering $\left(B_{r_{z_{i}}}^{2}\left(z_{i}\right)\right)$ for $i \in I$ ( $I$ is a countable set) of $\Pi\left(\mathcal{C}_{p-1} \cap\right.$ $\left.B_{\rho_{x_{0}}}^{2 p}\left(x_{0}\right)\right)$, we can apply Lemma A. 1 and obtain $\delta$ depending on $\alpha$ and $Q$ such that (A.3) holds true for some $P \in \mathbb{N}$. Letting $i \in I$, we can deduce from (A.3) and (V.24) that the radii $r_{j}$ of balls $B_{r_{j}}^{2}\left(z_{j}\right)$ intersecting $B_{r_{i}(1+\delta)}^{2}\left(z_{i}\right)$ satisfy $\alpha^{\gamma P} r_{i} \leq r_{j} \leq \alpha^{-\gamma P} r_{i}$. From this latter fact we deduce that there exists a number $M \in \mathbb{N}$ depending only on $\alpha$ and $Q$ such that

$$
\begin{equation*}
\operatorname{Card}\left\{j \in I: B_{r_{i}}\left(z_{i}\right) \cap B_{(1+\delta) r_{j}}\left(z_{j}\right) \neq \emptyset\right\} \leq M . \tag{V.42}
\end{equation*}
$$

Indeed, assuming $B_{r_{i}}\left(z_{i}\right) \cap B_{r_{j}}\left(z_{j}\right)=\emptyset$, if $B_{r_{i}}\left(z_{i}\right) \cap B_{(1+\delta) r_{j}}\left(z_{j}\right)$ we have just seen that $\alpha^{\gamma P} r_{j} \leq r_{i} \leq \alpha^{-\gamma P} r_{j}$ : the two radii have comparable size which is of course also comparable with the distance $\left|z_{i}-z_{j}\right|$. From (V.40) it is then clear that the number of such ball $B_{r_{j}}\left(z_{j}\right)$ is bounded by a constant depending only on the variables $\alpha$ and $Q$. It is now not difficult to deduce that $\left(B_{r_{i}\left(1+\frac{\delta}{2}\right)}^{2}\right)_{i \in I}$ realizes a locally finite covering of $\Pi\left(\mathcal{C}_{p-1} \cap B_{\rho_{x_{0}}}^{2 p}\left(x_{0}\right)\right)$ satisfying
(V.43) $\forall i \in I$,

$$
\operatorname{Card}\left\{j \in I: B_{\left(1+\frac{\delta}{2}\right) r_{i}}\left(z_{i}\right) \cap B_{\left(1+\frac{\delta}{2}\right) r_{j}}\left(z_{j}\right) \neq \emptyset\right\} \leq M+P .
$$

Indeed, when we assume for instance that $r_{i} \geq r_{j}$, then

$$
B_{\left(1+\frac{\delta}{2}\right) r_{i}}\left(z_{i}\right) \cap B_{\left(1+\frac{\delta}{2}\right) r_{j}}\left(z_{j}\right) \neq \emptyset
$$

implies clearly that $B_{(1+\delta) r_{i}}\left(z_{i}\right) \cap B_{r_{j}}\left(z_{j}\right) \neq \emptyset$ and the number of such a $j$ is controlled by $P$ (see A.3), whereas if $r_{i} \leq r_{j}, B_{\left(1+\frac{\delta}{2}\right) r_{i}}\left(z_{i}\right) \cap B_{\left(1+\frac{\delta}{2}\right) r_{j}}\left(z_{j}\right) \neq \emptyset$ implies clearly that $B_{r_{i}}\left(z_{i}\right) \cap B_{r_{j}(1+\delta)}\left(z_{j}\right) \neq \emptyset$ and the number of such a $j$ is controlled by $M$ (see V.42). Thus (V.43) holds true. And we shall use from now on the notation

$$
\begin{equation*}
\forall i \in I, \quad \rho_{i}:=r_{i}\left(1+\frac{\delta}{2}\right) . \tag{V.44}
\end{equation*}
$$

For any $i \in I$ we define $\chi_{i}$ to be a smooth nonnegative function satisfying
i) $\quad \chi_{i} \equiv 1 \quad$ in $B_{r_{i}}^{2}\left(z_{i}\right)$.
ii) $\quad \chi_{i} \equiv 0 \quad$ in $\mathbb{R}^{2} \backslash B_{\left(1+\frac{\delta}{2} r_{i}\right)}^{2}\left(z_{i}\right)$
iii) $\quad \forall k \in \mathbb{N},\left\|\nabla^{k} \chi_{i}\right\|_{\infty} \leq \frac{K_{k}}{r_{i}^{k}}$,
where $K_{k}$ depends only on $k$ and $Q$.
We define now

$$
\begin{equation*}
\varphi_{i}:=\frac{\chi_{i}}{\sum_{i \in I} \chi_{i}} . \tag{V.48}
\end{equation*}
$$

It is clear that $\left(\varphi_{i}\right)$ defines a partition of unity adapted to $B_{\left(1+\frac{\delta}{2} r_{i}\right)}^{2}\left(z_{i}\right)$ and satisfying the following estimates

$$
\begin{equation*}
\forall k \in \mathbb{N}, \quad\left\|\nabla^{k} \varphi_{i}\right\|_{\infty} \leq \frac{K_{k}}{r_{i}^{k}} \tag{V.49}
\end{equation*}
$$

where $K_{k}$ depends only on $k$ and $Q$.

## VI. The approximated average curve

This section is another step towards the proof that $\mathcal{P}_{Q-1} \Longrightarrow \mathcal{P}_{Q}$ which continues until Section VIII. We thus assume that $\mathcal{P}_{Q-1}$ holds (or that $Q=1$ ). Again in this part we consider the difficult case which is the case where we are blowing-up the current at a point $x_{0}$ of multiplicity $Q>1$ whose tangent cone $C_{0, x_{0}}$ is $Q$ times the integral over the flat disk $D_{0}$ given by $w_{i}=0$ for $i \ldots Q-1$ (we use the system of coordinates introduced at the beginning of $\S$ II) and where $x_{0}$ belongs to the closure of $\mathcal{C}_{Q-1}$. The purpose of this section is to approximate first our current over each ball of the covering introduced in the previous section $C\left\llcorner\Pi^{-1}\left(B_{r_{i}}^{2}\left(z_{i}\right)\right)\right.$ by a $Q$-valued graph $\left\{a_{i}^{k}\right\}_{k=1 \ldots Q}$ over $B_{r_{i}}^{2}\left(z_{i}\right)$ which is almost $J$-holomorphic ( $J_{x_{i}}$-holomorphic in fact where $x_{i} \in \mathcal{C}_{*}$ and $\left.\Pi\left(x_{i}\right)=z_{i}\right)$. Gluing the average curves $\tilde{a}_{i}=\frac{1}{Q} \sum_{k=1}^{Q} a_{i}^{k}$ of each of these $J_{x_{i}}-$ holomorphic $Q$-valued graphs together we shall produce a single-valued graph $\tilde{a}$ over $B_{\rho_{x_{0}}}^{2}\left(x_{0}\right)$ which approximates $C$ and for which we will study regularity properties that will be used in the following Section VII devoted to the unique continuation argument. Finally in the second subsection of this section we construct new coordinates adapted to the average curve.
VI.1. Constructing the average curve. Let $\rho_{x_{0}}$ be given by (V.23) and let $\left(B_{\rho_{i}}^{2}\left(z_{i}\right)\right)_{i \in I}$ be the Besicovitch-Whitney covering of $\Pi\left(\mathcal{C}_{Q-1} \cap B_{\rho_{x_{0}}}\left(x_{0}\right)\right)$ obtained at the end of the previous section. As we have seen above, for any
$i \in I, \mathcal{C}_{*} \cap \Pi^{-1}\left(B_{\rho_{i}}^{2}\left(z_{i}\right)\right) \subset \mathcal{C}_{Q-1} \cap B_{\rho_{i}}^{2}\left(z_{i}\right) \times B_{2 \rho_{i}}^{2 p-2}\left(w_{i}\right)$ where $x_{i}=\left(z_{i}, w_{i}\right)$ is in $\mathcal{C}_{*}$ (see Lemma V.2). For convenience we shall adopt the following notation:

$$
\begin{equation*}
N_{r}^{i}:=B_{r}^{2}\left(z_{i}\right) \times B_{2 \rho_{i}}^{2 p-2}\left(w_{i}\right) . \tag{VI.1}
\end{equation*}
$$

Assuming $\mathcal{P}_{Q-1}, C\left\llcorner N_{2 \rho_{i}}\right.$ is a $J$-holomorphic curve: there exist a smooth Riemannian surface and a smooth $J$-holomorphic map

$$
\begin{align*}
\Psi_{i}: \Sigma_{2, i} & \longrightarrow N_{2 \rho_{i}}^{i}  \tag{VI.2}\\
\xi & \longrightarrow \Psi_{i}(\xi)
\end{align*}
$$

such that $\Psi_{*}\left[\Sigma_{2, i}\right]=C\left\llcorner N_{2 \rho_{i}}\right.$. Let $H_{ \pm}^{0}\left(\Sigma_{2, i}\right)$ be the sets respectively of holomorphic and antiholomorphic functions on $\Sigma_{2, i}$. We introduce now $\eta_{i}$ the map from $\Sigma_{2, i}$ into $\mathbb{R}^{2 p-2}$ chosen such that the perturbation $\Psi_{i}+\eta_{i}$ is $J_{x_{i}}$ holomorphic; precisely, $\eta_{i}$ is given by

$$
\left\{\begin{array}{l}
\frac{\partial}{\partial \xi_{1}}\left(\Psi_{i}+\eta_{i}\right)+J_{x_{i}} \frac{\partial}{\partial \xi_{2}}\left(\Psi_{i}+\eta_{i}\right)=0 \quad \text { in } \Sigma_{2, i}  \tag{VI.3}\\
\forall h \in H\left(\Sigma_{2, i}\right) \quad \int_{\partial \Sigma_{2, i}} h d \eta_{i}=0
\end{array}\right.
$$

where $\left(\xi_{1}, \xi_{2}\right)$ are local coordinates on $\Sigma_{2, i}$ compatible with the complex structure. The existence of $\eta_{i}$ is justified in a few lines below. To this aim we shall make use of the following notation. Since $J$ is smooth the inverse function theorem gives the existence of a smooth map $\Lambda: B_{\rho_{x_{0}}}^{2 p}\left(x_{0}\right) \times \mathbb{R}^{2 p} \longrightarrow \mathbb{R}^{2 p}$ for $\rho_{x_{0}}$ chosen small enough such that
iii) $\quad \forall x \in B_{\rho_{x_{0}}}^{2 p}\left(x_{0}\right) \quad J_{x_{0}}=\left(\begin{array}{ccccc}0 & -1 & 0 \ldots & 0 & 0 \\ 1 & 0 & 0 \ldots & 0 & 0 \\ \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & . & . & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot \\ 0 & 0 & \ldots & 0 & -1 \\ 0 & 0 & \ldots & 1 & 0\end{array}\right)=\Lambda_{x} J_{x} \Lambda_{x}^{-1}$.

We shall denote by $\left(z^{i}=x_{1}^{i}+i x_{2}^{i}, w_{1}^{i}=x_{3}^{i}+i x_{4}^{i}, \ldots, w_{p}^{i}=x_{2 p-1}^{i}+i x_{2 p}^{i}\right)$ the following complex coordinates in $N_{2 \rho_{i}}$

$$
\left(\begin{array}{c}
x_{1}^{i}  \tag{VI.7}\\
x_{2}^{i} \\
\cdot \\
\cdot \\
\cdot \\
x_{2 p-1}^{i} \\
x_{2 p}^{i}
\end{array}\right)=\Lambda_{x_{i}} \cdot\left(\begin{array}{c}
x_{1} \\
x_{2} \\
\cdot \\
\cdot \\
\cdot \\
x_{2 p-1} \\
x_{2 p}
\end{array}\right)-x_{i} .
$$

We will also denote by $\Pi_{i}$ the map that assign to any point $x$ in $N_{\rho_{i}}^{i}$ the complex coordinate $z^{i}$ and by $D^{i}$ we denote the $J_{x_{i}}$-holomorphic 2 -disk

$$
\begin{equation*}
D^{i}:=\left\{x ; \forall k=1 \ldots p-1 \quad w_{k}^{i}=0\right\}=\Lambda_{x_{i}}^{-1} D_{0} . \tag{VI.8}
\end{equation*}
$$

Using these complex coordinates in $N_{2 \rho_{i}}^{i}$, we see that (VI.3) means

$$
\left\{\begin{array}{l}
\bar{\partial} \eta_{i}=-\bar{\partial} \Psi_{i} \quad \text { in } \Sigma_{2, i}  \tag{VI.9}\\
\forall h \in H\left(\Sigma_{2, i}\right), \quad \int_{\partial \Sigma_{2, i}} h d \eta_{i}=0 .
\end{array}\right.
$$

The existence and uniqueness of $\eta_{i}$ are given by Proposition A. 3 of [Ri3]. Since $\Psi_{i}$ is $J$-holomorphic we have $\partial_{\xi_{1}} \Psi_{i}+J\left(\Psi_{i}(\xi)\right) \partial_{\xi_{2}} \Psi_{i}=0$; thus

$$
\left|\partial_{\xi_{1}} \Psi_{i}+J\left(x_{i}\right) \partial_{\xi_{2}} \Psi_{i}\right| \leq\left|J\left(\Psi_{i}(\xi)\right)-J\left(x_{i}\right)\right||\nabla \psi|
$$

Combining this fact with the second part of Proposition A. 3 (i.e. estimate (A.13) of [Ri3]) we obtain

$$
\begin{equation*}
\int_{\Sigma_{2, i}}\left|\nabla \eta_{i}\right|^{2} \leq K r_{i}^{2} \int_{\Sigma_{2, i}}\left|\nabla \Psi_{i}\right|^{2} \leq K r_{i}^{4} . \tag{VI.10}
\end{equation*}
$$

For $\lambda \leq 2$, we denote by $\Sigma_{\lambda, i}$ the surface $\Sigma_{\lambda, i}=\Sigma_{2, i} \cap \Psi_{i}^{-1}\left(N_{\lambda \rho_{i}}^{i}\right)$ such that

$$
\begin{equation*}
\Psi_{i *}\left[\Sigma_{i}\right]=C\left\llcorner B_{\rho_{i}}^{2}\left(z_{i}\right) \times B_{\rho_{x_{0}}}^{n-2}(0)\right. \tag{VI.11}
\end{equation*}
$$

We then prove in [Ri3] the following lemma:
Lemma VI.1. Under the notation above,

$$
\begin{equation*}
\left\|\eta_{i}\right\|_{L^{\infty}\left(\Sigma_{\frac{3}{2}, i}\right)} \leq K r_{i}^{2} \tag{VI.12}
\end{equation*}
$$

where $K$ is a constant depending only on $\|\nabla J\|_{\infty}$ and the choice of $\alpha$ made in the previous section.

Consider now the $J_{x_{i}}$-holomorphic curve $C^{i}$ given by the image by $\Psi_{i}+\eta_{i}$ of $\Sigma_{\frac{3}{2}, i}$. Since $\partial \Psi_{i *}\left[\Sigma_{\frac{3}{2}, i}\right]$ is supported in $\Pi^{-1}\left(\partial B_{\frac{3}{2} \rho_{i}}^{2}\right)$, we know from Lemma VI. 1 that $\left|\eta_{i}\right|_{\infty} \leq C r_{i}^{2}$. Therefore $\partial \Psi_{i}+\eta_{i_{*}}\left[\Sigma_{\frac{3}{2}, i}\right]$ is supported in an $r_{i}^{2}$ neighborhood of $\Pi^{-1}\left(\partial B_{\frac{3}{2} \rho_{i}}^{2}(0)\right)$ (for $r_{i}$ small enough : that holds if $\varepsilon$ has been chosen small enough in Section VI). Thus we have that $\left.\Psi_{i *}\left[\Sigma_{\frac{3}{2}}, i\right]\right)$ is a cycle in $\Pi_{i}^{-1}\left(B_{\frac{5}{4} \rho_{i}}^{2}(0)\right)$ and the part of the image of $\Sigma_{\frac{3}{2}, i}$ included in $\Pi_{i}^{-1}\left(B_{\frac{5}{4} \rho_{i}}^{2}(0)\right)$ by $\Psi_{i}+\eta_{i}$ is a $J_{x_{i}}$-holomorphic cycle and therefore it is a $Q$-valued graph over $D^{i}$ for the complex coordinates given by $\left(z^{i}, w^{i}\right)$. We denote by $\left\{a_{k}^{i}\right\}_{k=1 \ldots Q}$ this $Q$-valued graph (i.e. $a_{k}^{i}\left(z_{0}^{i}\right)$ are the $w^{i}$ coordinates, in the chart $\left(z^{i}, w^{i}\right)$, of the $Q$ intersection points between the $J_{x_{i}}$-holomorphic curve $\Psi_{i}+\eta_{i}\left(\sum_{\frac{3}{2}, i}\right)$ and the $J_{x_{i}}$-holomorphic submanifold given by $z^{i}=z_{0}^{i}$ ). We now define $\tilde{C}^{i}$ to be the $J_{x_{i}}$ holomorphic curve in $\Pi_{i}^{-1}\left(B_{\rho_{i}}^{2}\right)$ given by

$$
\begin{equation*}
\tilde{C}^{i}:=\left\{x=\Lambda_{x_{i}}^{-1}\left(\left(z^{i}, \tilde{a}_{i}^{i}\left(z_{i}\right)=\frac{1}{Q} \sum_{k=1}^{Q} a_{k}^{i}\left(z^{i}\right)\right)+x_{i}\right) \quad \forall z^{i} \in B_{\frac{5}{4} \rho_{i}}^{2}(0)\right\} . \tag{VI.13}
\end{equation*}
$$

Observe that

$$
\begin{equation*}
\frac{\partial}{\partial z^{i}} \tilde{a}_{i}^{i}=0 \quad \text { in } \mathcal{D}^{\prime}\left(B_{\frac{5}{4} \rho_{i}}^{2}(0)\right) \tag{VI.14}
\end{equation*}
$$

Moreover, the conformal invariance of the Dirichlet energy gives

$$
\begin{align*}
& \int_{B_{\frac{5}{4} \rho_{i}}^{2}} \sum_{k=1} \sum_{k=1}^{Q}\left|\nabla a_{k}^{i}\right|^{2}\left(z^{i}\right) d z^{i} \wedge d \bar{z}^{i}  \tag{VI.15}\\
&=\int_{\left(\Psi_{i}+\eta_{i}\right)^{-1}\left(C_{i} \cap \Pi_{i}^{-1}\left(B_{\frac{5}{4} \rho_{i}}^{2}(0)\right)\right)}\left|\nabla\left(\Psi_{i}+\eta_{i}\right)\right|^{2}(\xi) d \xi \leq K r_{i}^{2} .
\end{align*}
$$

We then deduce that

$$
\begin{equation*}
\int_{B_{\frac{5}{4} \rho_{i}}^{2}(0)}\left|\nabla \tilde{a}_{i}^{i}\right|^{2}\left(z^{i}\right) d z^{i} \leq K r_{i}^{2} . \tag{VI.16}
\end{equation*}
$$

Combining (VI.15) and (VI.16) and using standard elliptic estimates we get that for any $l \in \mathbb{N}$

$$
\begin{equation*}
\left\|\nabla^{l} \tilde{a}_{i}^{i}\right\|_{L^{\infty}\left(B_{\frac{\rho_{5} p_{i}}{2}}^{2}(0)\right)} \leq K_{l} r_{i}^{-l+1} \tag{VI.17}
\end{equation*}
$$

The subscript $i$ in the notation $\tilde{a}_{i}^{i}$ is here to recall that we express $\tilde{C}^{i}$ as a graph in the $\left(z^{i}, w^{i}\right)$ coordinates. The same $J_{x_{i}}$-holomorphic curve $\tilde{C}^{i}$ can also, due to (VI.17), be expressed as a graph in a neighboring system of coordinate $\left(z^{j}, w^{j}\right)$
where $B_{\rho_{i}}\left(z_{i}\right) \cap B_{\rho_{j}}\left(z_{j}\right) \neq \emptyset$ (indeed the passage from $\left(z^{i}, w^{i}\right)$ to $\left(z^{j}, w^{j}\right)$ is given by a transformation matrix in $\mathbb{R}^{2 p}$ close to the identity at a distance of the order $\left.r_{i}\right)$. In such system of coordinates $\left(z^{j}, w^{j}\right)$, we shall denote $\tilde{a}_{j}^{i}\left(z^{j}\right)$ the graph corresponding to $\tilde{C}^{i}$.

Since $\tilde{C}^{i}$ is a graph over $w^{i}=0$ given by $\left(z^{i}, \tilde{a}_{i}^{i}\left(z^{i}\right)\right)$ whose gradient is bounded (see (VI.17)), and since the passage from the $(z, w)$ coordinates to $\left(z^{i}, w^{i}\right)$ coordinates is given by a transformation $\Lambda_{x_{i}}$ whose distance to the identity is bounded by $\left|x_{i}\right|$ that tends to zero, $\tilde{C}^{i}$ is then also realized by a graph over $B_{\frac{7}{6}}^{2} \rho_{i}\left(\Pi\left(x_{i}\right)\right)$ that we shall now denote $\left(z, \tilde{a}_{i}(z)\right)$ :

$$
\begin{equation*}
\tilde{C}^{i}\left\llcorner\Pi^{-1}\left(B_{\frac{7}{6} \rho_{i}}^{2}\left(\Pi\left(x_{i}\right)\right)\right)=\left(z, \tilde{a}_{i}(z)\right)_{*}\left[B_{\frac{7}{6} \rho_{i}}^{2}\left(\Pi\left(x_{i}\right)\right)\right]\right. \tag{VI.18}
\end{equation*}
$$

Consider now $i$ and $j$ such that $B_{\rho_{i}}\left(z_{i}\right) \cap B_{\rho_{j}}\left(z_{j}\right) \neq \emptyset$. We shall compare $\tilde{a}_{i}$ and $\tilde{a}_{j}$ in $\Pi^{-1}\left(B_{\rho_{i}}\left(z_{i}\right) \cap B_{\rho_{j}}\left(z_{j}\right)\right)$. Precisely, we have the following lemma:

Lemma VI.2. Under the notation above,

$$
\begin{equation*}
\forall l \in \mathbb{N}, \quad\left\|\nabla^{l}\left(\tilde{a}_{i}-\tilde{a}_{j}\right)\right\|_{L^{\infty}\left(B_{\rho_{i}}^{2}\left(z_{i}\right) \cap B_{\rho_{j}}^{2}\left(z_{j}\right)\right)} \leq K_{l} \rho_{i}^{2-l} \tag{VI.19}
\end{equation*}
$$

Proof of Lemma VI.2. First of all we compare $C_{i}$ and $C_{j}$ in $\Pi^{-1}\left(B_{\rho_{i}}\left(z_{i}\right) \cap\right.$ $\left.B_{\rho_{j}}\left(z_{j}\right)\right)$. We can always assume that $\Sigma_{2, i}$ and $\Sigma_{2, j}$ are part of a same Riemannian surface $\Sigma$ with a joint parametrization $\Psi=\Psi_{i}$ on $\Sigma_{2, i}$ and $\Psi=\Psi_{j}$ on $\Sigma_{2, j}$ and such that $\Psi_{*}[\Sigma]=C\left\llcorner N_{2 \rho_{i}}^{i} \cup N_{2 \rho_{j}}^{j}\right.$. Let $\Sigma^{i j}:=\Psi^{-1}(\operatorname{supp}(C)) \cap N_{2 \rho_{i}}^{i} \cap N_{2 \rho_{j}}^{j}$. We consider the following mapping

$$
\begin{align*}
\Xi^{i j} \Sigma^{i j} \times[0,1] & \longrightarrow N_{3 \rho_{i}}^{i} \cap N_{3 \rho_{j}}^{j}  \tag{VI.20}\\
(\xi, t) & \longrightarrow \Psi(\xi)+t \eta_{j}(\xi)+(1-t) \eta_{i}(\xi)
\end{align*}
$$

Clearly for any $\lambda<\frac{3}{2}$

$$
\begin{equation*}
\partial \Xi^{i j}\left[\Sigma^{i j}\right] \times[0,1]\left\llcorner N_{\lambda \rho_{i}}^{i} \cap N_{\lambda \rho_{j}}^{j}=C^{j}-C^{i}\left\llcorner N^{i} \lambda \rho_{i} \cap N_{\lambda \rho_{j}}^{j}\right.\right. \tag{VI.21}
\end{equation*}
$$

We have

$$
\begin{equation*}
\mathrm{M}\left(\Xi^{i j}{ }_{*}\left[\Sigma^{i j}\right] \times[0,1]\right)=\int_{0}^{1} \int_{\Sigma^{i j}} J_{3} \Xi^{i j} \tag{VI.22}
\end{equation*}
$$

where $\left(J_{3} \Xi^{i j}\right)^{2}$ is the sum of the squares of the determinants of the $3 \times 3$ submatrices of $D \Xi^{i j}$. Clearly
(VI.23)

$$
\left|J_{3} \Xi^{i j}\right|(\xi, t) \leq K\left[\left\|\eta_{i}\right\|_{\infty}+\left\|\eta_{j}\right\|_{\infty}\right]\left[|\nabla \Psi|^{2}(\xi)+\left|\nabla \eta_{i}\right|^{2}(\xi)+\left|\nabla \eta_{j}\right|^{2}(\xi)\right]
$$

Combining Lemma V.3, Lemma VI.1, (VI.22) and (VI.23), we get that for any $\lambda<\frac{3}{2}$

$$
\begin{equation*}
\mathrm{M}\left(\Xi^{i j}\left[\Sigma^{i j}\right] \times[0,1]\left\llcorner N^{i} \lambda \rho_{i} \cap N_{\lambda \rho_{j}}^{j}\right) \leq K r_{i}^{4}\right. \tag{VI.24}
\end{equation*}
$$

Therefore, combining (VI.21) and (VI.24), using a standard slicing and Fubini type argument, we may find $\lambda \in\left(\frac{5}{4}, \frac{3}{2}\right)$ such that

$$
\begin{equation*}
\mathcal{F}\left(\left(C^{i}-C^{j}\right)\left\llcorner N_{\lambda \rho_{i}}^{i} \cap N_{\lambda \rho_{j}}^{j}\right) \leq K r_{i}^{4} .\right. \tag{VI.25}
\end{equation*}
$$

We shall now compare $\tilde{C}^{i}$ and $\tilde{C}^{j}$. Denote $\left(x_{1}^{t}, x_{2}^{t} \ldots x_{2 p}^{t}\right)$ the coordinates given by $\left(x_{1}^{t}, x_{2}^{t} \ldots x_{2 p}^{t}\right)^{T}=\Lambda_{x_{t}} \cdot\left[\left(x_{1}, x_{2} \ldots x_{2 p}\right)^{T}-x_{i}\right]$ where we keep denoting ( $x_{1}, x_{2} \ldots x_{2 p}$ ) our original normal coordinates vanishing at the center $x_{0}$ introduced in (II.1) and $\Lambda_{x_{t}}$ is the transformation matrix introduced in (VI.4). Observe that with this notation $\left(x_{1}^{0}, x_{2}^{0} \ldots x_{2 p}^{0}\right)=\left(x_{1}^{i}, x_{2}^{i} \ldots x_{2 p}^{i}\right)$, that $\left(x_{1}^{1}, x_{2}^{1} \ldots x_{2 p}^{1}\right)=\left(x_{1}^{j}, x_{2}^{j} \ldots x_{2 p}^{j}\right)+\Lambda_{x^{j}} \cdot\left(x_{i}-x_{j}\right)$ and that $\left(x_{1}^{t}, x_{2}^{t} \ldots x_{2 p}^{t}\right)$ has been chosen in order to vanish at a fixed point $x_{i}$. Observe also that

$$
\begin{equation*}
\left|\frac{d}{d t} \Lambda_{x^{t}}\right| \leq K r_{i}, \quad\left\|\frac{d}{d t} x^{t}\right\|_{L^{\infty}\left(B_{4 r_{i}}^{2 p}\left(x_{i}\right)\right)}+\left\|\frac{d}{d t} y^{t}\right\|_{L^{\infty}\left(B_{4 r_{i}}^{2 p}\left(x_{i}\right)\right)} \leq K r_{i}^{2} . \tag{VI.26}
\end{equation*}
$$

We also adopt the notation $z^{t}:=x_{1}^{t}+i x_{2}^{t}$ and for $k=1 \ldots p-1 w^{t}:=$ $x_{2 k+1}^{t}+i x_{2 k+2}^{t}$. Observe then that $z^{t}=$ constant or $w_{k}^{t}=$ constant are $J_{x^{t-}}$ holomorphic $2 p-2$ submanifolds, or simply complex varieties in $\left(\mathbb{R}^{2 p}, J_{x^{t}}\right)$. In order to compare $\tilde{C}^{i}$ and $\tilde{C}^{j}$ we shall perturb $\Xi^{i j}$ in the following way: denote first $\Psi^{t}, \eta_{i}^{t}$ and $\eta_{j}^{t}$ the maps $\Psi, \eta_{i}$ and $\eta_{j}$ expressed in the coordinates $\left(z^{t}, w^{t}\right)$, and consider the map $s^{t}:\left(\Sigma^{\prime}\right)^{i j} \rightarrow \mathbb{C}^{p}$ solving

$$
\begin{gather*}
\partial_{\bar{\xi}} s^{t}=\partial_{\bar{\xi}}\left(\Psi^{t}+t \eta_{j}^{t}+(1-t) \eta_{i}^{t}\right) \quad \text { in }\left(\Sigma^{\prime}\right)^{i j}  \tag{VI.27}\\
\forall h \in H\left(\Sigma^{i j}\right), \quad \int_{\partial \Sigma^{i j}} s^{t} d h=0,
\end{gather*}
$$

where $\Sigma^{i j}:=\Psi^{-1}(\operatorname{supp}(C)) \cap N_{\frac{3}{2} \rho_{i}}^{i} \cap N_{\frac{3}{2} \rho_{j}}^{j}$. The existence and uniqueness of $s^{t}$ are given by Proposition A. 3 of [Ri3]. We shall now replace the map $\Xi^{i j}$ on $\left(\Sigma^{\prime}\right)^{i j}$ by the map

$$
\left.\begin{array}{rl}
\left(\Xi^{\prime}\right)^{i j}:\left(\Sigma^{\prime}\right)^{i j} \times & {[0,1]} \tag{VI.28}
\end{array}\right) N_{3 \rho_{i}}^{i} \cap N_{3 \rho_{j}}^{j} .
$$

Observe that for each $t \in[0,1]$ the map $\Xi^{\prime i j}(\cdot, t)$ is a $J_{x^{t}}$-holomorphic curve. Observe also that, for $t=0, s^{t}=0$ and for $t=1, \partial_{\bar{\xi}}\left(\Psi^{1}+\eta_{j}^{1}\right)=0$, since $\Psi+\eta_{j}$ is $J_{x^{j}}$-holomorphic and $\left(z^{1}, w^{1}\right)$ are $J_{x^{j}}$ coordinates, thus we have also $s^{1}=0$. One can easily verify, as in proving (VI.10) that for all $t \in[0,1]$

$$
\int_{\left(\Sigma^{\prime}\right)^{i j}}\left|\partial_{\bar{\xi}}\left(\Psi^{t}+t \eta_{j}^{t}+(1-t) \eta_{i}^{t}\right)\right|^{2} \leq K r_{i}^{4}
$$

Now, using Lemma VI. 1 we have $\left\|s^{t}\right\|_{\infty}\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right) \leq K r_{i}^{2}$ where

$$
\left(\Sigma^{\prime \prime}\right)^{i j}:=\Psi^{-1}(\operatorname{supp}(C)) \cap N_{\frac{5}{4} \rho_{i}}^{i} \cap N_{\frac{5}{4} \rho_{j}}^{j} .
$$

Therefore for any $\lambda<\frac{6}{5}$ we have

$$
\begin{equation*}
\partial\left(\Xi^{\prime}\right)^{i j}{ }_{*}\left[\left(\Sigma^{\prime}\right)^{i j}\right] \times[0,1]\left\llcorner N^{i} \lambda \rho_{i} \cap N_{\lambda \rho_{j}}^{j}=C^{j}-C^{i}\left\llcorner N^{i} \lambda \rho_{i} \cap N_{\lambda \rho_{j}}^{j} .\right.\right. \tag{VI.29}
\end{equation*}
$$

We consider now the following interpolation between $\tilde{C}^{i}$ and $\tilde{C}^{j}$ : let $\tilde{\Xi}^{i j}$ be the following map

$$
\begin{align*}
\tilde{\Xi}^{i j} \Pi\left(\Psi\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right)\right) \times & {[0,1] \longrightarrow N_{3 \rho_{i}}^{i} \cap N_{3 \rho_{j}}^{j} }  \tag{VI.30}\\
& (z, t) \longrightarrow \Lambda_{x^{t}}^{-1} \cdot\left[\left(z, \tilde{a}^{t}(z)\right]+x_{i}\right.
\end{align*}
$$

where the $p-1$ complex components are given by the slices of

$$
C^{t}:=\left(\Xi^{\prime}\right)^{i j}{ }_{*}\left[\left(\Sigma^{\prime}\right)^{i j}\right] \times\{t\}
$$

by $z^{t}=\xi$ evaluated on the functions $w_{k}^{t}$. Precisely, using the notation of $[\mathrm{Fe}$, 4.3], we have

$$
\begin{equation*}
\tilde{a}_{k}^{t}(z):=\left\langle C^{t}, z^{t}, z\right\rangle\left(w_{k}^{t}\right) . \tag{VI.31}
\end{equation*}
$$

It is clear that for any $\lambda<\frac{6}{5}$

$$
\begin{equation*}
\tilde{\Xi}_{*}^{i j}\left[\Pi_{i}\left(\Psi\left(\left(\Sigma^{\prime}\right)^{i j}\right)\right)\right] \times[0,1]\left\llcorner N^{i} \lambda \rho_{i} \cap N_{\lambda \rho_{j}}^{j}=\tilde{C}^{j}-\tilde{C}^{i}\left\llcorner N^{i} \lambda \rho_{i} \cap N_{\lambda \rho_{j}}^{j} .\right.\right. \tag{VI.32}
\end{equation*}
$$

In order to get a bound for $\mathcal{F}\left(\tilde{C}^{j}-\tilde{C}^{i}\left\llcorner N^{i} \lambda \rho_{i} \cap N_{\lambda \rho_{j}}^{j}\right)\right.$, it remains to evaluate the mass of $\tilde{\Xi}_{*}^{i j}\left[\Pi_{i}\left(\Psi\left(\left(\Sigma^{\prime}\right)^{i j}\right)\right)\right] \times[0,1]\left\llcorner N^{i} \lambda \rho_{i} \cap N_{\lambda \rho_{j}}^{j}\right.$ for $\lambda=\frac{6}{5}$ for instance. We have

$$
\begin{equation*}
\left|J_{3} \tilde{\Xi}^{i j}\right|(z, t) \leq \left\lvert\, \frac{\partial}{\partial t} \Lambda_{x^{t}}^{-1} \cdot\left[\left(z, \tilde{a}^{t}(z)\right] \mid(z, t)\left[1+\left|\nabla_{z} \tilde{a}^{t}(z)\right|^{2}\right] .\right.\right. \tag{VI.33}
\end{equation*}
$$

Because of the same arguments developed to prove (VI.17), since $\tilde{a}^{t}(\xi)$ is holomorphic, we have

$$
\begin{equation*}
\left\|\nabla_{z} \tilde{a}^{t}(z)\right\|_{L^{\infty}\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right)} \leq K \tag{VI.34}
\end{equation*}
$$

Thus

$$
\begin{align*}
\mathrm{M} & \left(\tilde{\Xi}_{*}^{i j}\left[\Pi \Psi\left(\left(\Sigma^{\prime}\right)^{i j}\right)\right] \times[0,1]\left\llcorner N^{i} \lambda \rho_{i} \cap N_{\lambda \rho_{j}}^{j}\right) \leq \int_{0}^{1} \int_{\Pi\left(\Psi\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right)\right)}\left|J_{3} \tilde{\Xi}^{i j}\right|\right.  \tag{VI.35}\\
& \leq K \int_{0}^{1} \int_{\Pi\left(\Psi\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right)\right)} \left\lvert\, \frac{\partial}{\partial t} \Lambda_{x^{t}}^{-1} \cdot\left[\left(z, \tilde{a}^{t}(z)\right] \mid d z \wedge d \bar{z} \wedge d t\right.\right. \\
& \leq K \int_{0}^{1} \int_{\Pi\left(\Psi\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right)\right)}\left[r_{i}\left|\left(z, \tilde{a}^{t}(z)\right)\right|+\left|\frac{\partial \tilde{a}^{t}}{\partial t}\right|\right] .
\end{align*}
$$

On the one hand, since $\Pi_{i}\left(\Psi\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right)\right)\left|\left(z, \tilde{a}^{t}(z)\right)\right| \leq K r_{i}$, we have

$$
\begin{equation*}
\int_{0}^{1} \int_{\Pi\left(\Psi\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right)\right)} r_{i}\left|\left(z, \tilde{a}^{t}(z)\right)\right| \leq K r_{i}^{4} \tag{VI.36}
\end{equation*}
$$

On the other hand
(VI.37)
$\int_{0}^{1} \int_{\Pi\left(\Psi\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right)\right)}\left|\frac{\partial \tilde{a}^{t}}{\partial t}\right|=\lim _{N \rightarrow+\infty} \frac{1}{N} \sum_{l=1}^{N-1} \int_{\Pi\left(\Psi\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right)\right)} N\left|\tilde{a}^{\frac{l}{N}}(z)-\tilde{a}^{\frac{l+1}{N}}(z)\right| d z \wedge d \bar{z}$.
We have

$$
\begin{align*}
\left|\tilde{a}_{k}^{\frac{l}{N}}(z)-\tilde{a}_{k}^{\frac{l+1}{N}}(z)\right|= & \left|\left\langle C^{\frac{l}{N}}, z^{\frac{l}{N}}, z\right\rangle\left(w_{k}^{\frac{l}{N}}\right)-\left\langle C^{\frac{l+1}{N}}, z^{\frac{l+1}{N}}, z\right\rangle\left(w_{k}^{\frac{l+1}{N}}\right)\right|  \tag{VI.38}\\
\leq & \left|\left\langle C^{\frac{l}{N}}-C^{\frac{l+1}{N}}, z^{\frac{l}{N}}, z\right\rangle\left(w_{k}^{\frac{l}{N}}\right)\right| \\
& +\left|\left\langle C^{\frac{l+1}{N}}, z^{\frac{l}{N}}, z\right\rangle\left(w_{k}^{\frac{l}{N}}\right)-\left\langle C^{\frac{l+1}{N}}, z^{\frac{l+1}{N}}, z\right\rangle\left(w_{k}^{\frac{l}{N}}\right)\right| \\
& +\left|\left\langle C^{\frac{l+1}{N}}, z^{\frac{l+1}{N}}, z\right\rangle\left(w_{k}^{\frac{l}{N}}-w_{k}^{\frac{l+1}{N}}\right)\right| .
\end{align*}
$$

We have to control the sum over $l$ of the integral over $\Pi\left(\Psi\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right)\right)$ of the three absolute values on the right-hand side of (VI.38) one by one. For the first term of the right-hand side of (VI.38) we have, using [Fe, 4.3.1], since $\left\|w_{k}^{\frac{L}{N}}\right\|_{\infty}+\left\|d w_{k}^{\frac{l}{N}}\right\|_{\infty} \leq 1$,

$$
\begin{align*}
& \int_{\Pi\left(\Psi\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right)\right)}\left|\left\langle C^{\frac{l}{N}}-C^{\frac{l+1}{N}}, z^{\frac{l}{N}}, z\right\rangle\left(w_{k}^{\frac{l}{N}}\right)\right| d z \wedge d \bar{z}  \tag{VI.39}\\
& \leq \operatorname{Lip}\left(z^{\frac{k}{N}}\right) \mathcal{F}_{N_{\lambda_{\rho_{i}}}^{i} \cap N_{\lambda_{\rho_{j}}}^{j}}\left(C^{\frac{k}{N}}-C^{\frac{k+1}{N}}\right) \leq K \mathcal{F}_{N_{\lambda_{\rho_{i}}}^{i} \cap N_{\lambda_{\rho_{j}}}^{j}}\left(C^{\frac{k}{N}}-C^{\frac{k+1}{N}}\right) .
\end{align*}
$$

Similarly, as we established estimate (VI.25) we can show that

$$
\begin{equation*}
\mathcal{F}_{N_{\lambda_{\rho_{i}}}^{i} \cap N_{\lambda \rho_{j}}^{j}}\left(C^{\frac{k}{N}}-C^{\frac{k+1}{N}}\right) \leq K \frac{1}{N} r_{i}^{4} . \tag{VI.40}
\end{equation*}
$$

Thus

$$
\begin{equation*}
\lim _{N \rightarrow+\infty} \sum_{l=1}^{N-1} \int_{\Pi\left(\Psi\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right)\right)}\left|\left\langle C^{\frac{l}{N}}-C^{\frac{l+1}{N}}, z^{\frac{l}{N}}, z\right\rangle\left(w_{k}^{\frac{l}{N}}\right)\right| \leq K r_{i}^{4} \tag{VI.41}
\end{equation*}
$$

For the second term on the right-hand side of (VI.38) we use 4.3.9 (3) of [Fe] and get

$$
\begin{align*}
& \int_{\Pi\left(\Psi\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right)\right)}\left|\left\langle C^{\frac{l+1}{N}}, z^{\frac{l}{N}}, z\right\rangle\left(w_{k}^{\frac{l}{N}}\right)-\left\langle C^{\frac{l+1}{N}}, z^{\frac{l+1}{N}}, z\right\rangle\left(w_{k}^{\frac{l}{N}}\right)\right|  \tag{VI.42}\\
& K \int_{\frac{l}{N}}^{\frac{l+1}{N}} d t \int_{\left(z^{t}\right)^{-1}\left(\Pi\left(\Psi\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right)\right)\right)}\left|z^{\frac{l}{N}}-z^{\frac{l+1}{N}}\right| d\left\|C^{\frac{l+1}{N}}\right\| \\
& K \frac{r_{i}^{2}}{N} \mathrm{M}\left(C^{\frac{l+1}{N}}\left\llcorner N^{i} \lambda \rho_{i} \cap N_{\lambda \rho_{j}}^{j}\right)\right. \\
& \leq K \frac{r_{i}^{4}}{N}
\end{align*}
$$

where we have used (VI.26). Therefore we obtain
(VI.43)

$$
\lim _{l \rightarrow+\infty} \sum_{l}^{N-1} \int_{\Pi\left(\Psi\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right)\right)}\left|\left\langle C^{\frac{l+1}{N}}, z^{\frac{l}{N}}, z\right\rangle\left(w_{k}^{\frac{l}{N}}\right)-\left\langle C^{\frac{l+1}{N}}, z^{\frac{l+1}{N}}, z\right\rangle\left(w_{k}^{\frac{l}{N}}\right)\right| \leq K r_{i}^{4}
$$

Finally for the second term of the right-hand side of (VI.38), we use again (VI.26) and 4.3.2 (2) of [Fe] to obtain that

$$
\begin{align*}
\int_{\Pi\left(\Psi\left(\left(\Sigma^{\prime \prime}\right)^{i j}\right)\right)} \left\lvert\,\left\langle C^{\frac{l+1}{N}}, z^{\frac{l+1}{N}}\right.\right. & , z\rangle \left.\left(w_{k}^{\frac{l}{N}}-w_{k}^{\frac{l+1}{N}}\right) \right\rvert\,  \tag{VI.44}\\
& \leq \mathrm{M}\left(C^{\frac{l+1}{N}}\left\llcorner N^{i} \lambda \rho_{i} \cap N_{\lambda \rho_{j}}^{j}\right) \frac{r_{i}^{2}}{N} \leq K \frac{r_{i}^{4}}{N}\right.
\end{align*}
$$

Combining (VI.35), (VI.36), (VI.37), (VI.38), (VI.41), (VI.43) and (VI.44), we obtain that

$$
\begin{equation*}
\mathrm{M}\left(\tilde{\Xi}_{*}^{i j}\left[\Pi \Psi\left(\left(\Sigma^{\prime}\right)^{i j}\right)\right] \times[0,1]\left\llcorner N_{\lambda \rho_{i}}^{i} \cap N_{\lambda \rho_{j}}^{j}\right) \leq r_{i}^{4}\right. \tag{VI.45}
\end{equation*}
$$

Combining this last inequality with (VI.32) and a Fubini type argument we obtain that there exists $\lambda \in\left[\frac{7}{6}, \frac{6}{5}\right]$ such that

$$
\begin{equation*}
\mathcal{F}\left(\left(\tilde{C}_{i}-\tilde{C}_{j}\right)\left\llcorner N^{i} \lambda \rho_{i} \cap N_{\lambda \rho_{j}}^{j}\right) \leq r_{i}^{4}\right. \tag{VI.46}
\end{equation*}
$$

From this fact we then deduce, since $\tilde{C}_{i}$ and $\tilde{C}_{j}$ are single valued graphs with uniformly bounded gradients

$$
\begin{equation*}
\int_{\Pi_{i}\left(N^{i} \frac{7}{6} \rho_{i} \cap N_{\frac{7}{6} \rho_{j}}^{j}\right)}\left|\tilde{a}_{j}^{i}\left(z^{i}\right)-\tilde{a}_{j}^{j}\right| \leq K r_{i}^{4} \tag{VI.47}
\end{equation*}
$$

Using the notation introduced in (VI.4) and (VI.13), we have that for any $z$ there exists $\xi$ such that

$$
\begin{equation*}
\left(z-z_{i}, \tilde{a}_{i}(z)-w_{i}\right)=\Lambda_{x_{i}}^{-1}\left(\xi, \tilde{a}_{i}^{i}(\xi)\right) \tag{VI.48}
\end{equation*}
$$

where $\left|\Lambda_{x_{i}}-\mathrm{id}\right| \leq K\left|x_{i}\right| \leq K \rho_{x_{0}}$. Let $z^{\prime}:=\rho_{i}^{-1}\left(z-z_{i}\right)$ and $\hat{a}_{i}\left(z^{\prime}\right):=$ $\rho_{i}^{-1}\left(\tilde{a}_{i}(z)-w_{i}\right)$. Let also $\xi^{\prime}:=\rho_{i}^{-1} \xi$ and $\hat{a}_{i}^{i}\left(\xi^{\prime}\right):=\tilde{a}_{i}^{i}(\xi)$. Since $\tilde{a}_{i}^{i}$ is holomorphic $\left(\right.$ see $(\mathrm{VI} .14), \hat{a}_{i}^{i}$ is also clearly holomorphic and since $\left\|\hat{a}_{i}^{i}\right\|_{L^{\infty}\left(B_{\frac{3}{2}}(0)\right)} \leq K$, we have that for any $l \in \mathbb{N}$

$$
\begin{equation*}
\left\|\nabla^{l} \hat{a}_{i}^{i}\right\|_{L^{\infty}\left(B_{\frac{5}{4}}(0)\right)} \leq K_{l} \tag{VI.49}
\end{equation*}
$$

Using the above notation we have

$$
\begin{equation*}
\left(z^{\prime}, \hat{a}_{i}\left(z^{\prime}\right)\right)=\Lambda_{x_{i}}^{-1}\left(\xi^{\prime}, \hat{a}_{i}^{i}\left(\xi^{\prime}\right)\right) \tag{VI.50}
\end{equation*}
$$

From the inverse function theorem, since $\left|\Lambda_{x_{i}}-i d\right| \leq K\left|x_{i}\right| \leq K, \rho_{x_{0}}$ can be taken as small as we want by taking $\rho_{x_{0}}$ small enough, we have that for all $l \in \mathbb{N}$ there exists $K_{l}$ such that

$$
\begin{equation*}
\left\|\nabla_{z^{\prime}}^{l} \xi^{\prime}\right\|_{\infty} \leq K_{l} \tag{VI.51}
\end{equation*}
$$

Therefore, combining (VI.49), (VI. 50 and (VI.51, we obtain

$$
\begin{equation*}
\left\|\nabla_{z^{\prime}}^{l} \hat{a}_{i}\left(z^{\prime}\right)\right\|_{\infty} \leq K_{l} \tag{VI.52}
\end{equation*}
$$

From that estimate we then deduce

$$
\begin{equation*}
\left\|\nabla_{z}^{l} \tilde{a}_{i}\right\|_{L^{\infty}\left(B_{\frac{f}{6} \rho_{i}}^{2}\left(z_{i}\right)\right)} \leq K_{l} r_{i}^{l-1} \tag{VI.53}
\end{equation*}
$$

Since $\tilde{C}^{i}$ is $J_{x_{i}}$-holomorphic, we have the existence of $\lambda_{1}^{i}, \mu_{1}^{i}, \lambda_{2}^{i}, \mu_{2}^{i}$ such that

$$
\left\{\begin{array}{c}
J_{x_{i}} \cdot\left(\begin{array}{c}
1 \\
0 \\
\frac{\partial \tilde{a}_{i}}{\partial x}
\end{array}\right)=\lambda_{1}^{i}\left(\begin{array}{c}
1 \\
0 \\
\frac{\partial \tilde{a}_{i}}{\partial x}
\end{array}\right)+\mu_{1}^{i}\left(\begin{array}{c}
0 \\
1 \\
\frac{\partial \tilde{a}_{i}}{\partial y}
\end{array}\right)  \tag{VI.54}\\
J_{x_{i}} \cdot\left(\begin{array}{c}
0 \\
1 \\
\frac{\partial \tilde{a}_{i}}{\partial y}
\end{array}\right)=\lambda_{2}^{i}\left(\begin{array}{c}
1 \\
0 \\
\frac{\partial \tilde{a}_{i}}{\partial x}
\end{array}\right)+\mu_{2}^{i}\left(\begin{array}{c}
0 \\
1 \\
\frac{\partial \tilde{a}_{i}}{\partial y}
\end{array}\right)
\end{array}\right.
$$

Writing $J_{x_{i}}=J_{0}+\delta\left(x_{i}\right)$, we first observe that

$$
\begin{equation*}
\left|\delta\left(x_{i}\right)\right| \leq\|J\|_{C^{1}}\left|x_{i}\right| \leq K \rho_{x_{0}} . \tag{VI.55}
\end{equation*}
$$

From this notation we deduce using (VI.54),

$$
\left\{\begin{array}{l}
\lambda_{1}^{i}=\delta_{1,1}\left(x_{i}\right)+\sum_{l=3}^{2 p} \delta_{1, l}\left(x_{i}\right) \frac{\partial \tilde{a}_{i}^{l}}{\partial x} \\
\mu_{1}^{i}=1+\delta_{2,1}\left(x_{i}\right)+\sum_{l=3}^{2 p} \delta_{2, l}\left(x_{i}\right) \frac{\partial \tilde{a}_{i}^{l}}{\partial x}  \tag{VI.56}\\
\lambda_{2}^{i}=-1+\delta_{1,2}\left(x_{i}\right)+\sum_{l=3}^{2 p} \delta_{1, l}\left(x_{i}\right) \frac{\partial \tilde{a}_{i}^{l}}{\partial y} \\
\mu_{2}^{i}=\delta_{2,2}\left(x_{i}\right)+\sum_{l=3}^{2 p} \delta_{2, l}\left(x_{i}\right) \frac{\partial \tilde{a}_{i}^{l}}{\partial y}
\end{array}\right.
$$

Therefore the equation solved by $\tilde{a}_{i}$ is for any $k=1 \ldots p-1$,

$$
\left\{\begin{align*}
\frac{\partial \tilde{a}_{i}^{2 k+1}}{\partial x}-\frac{\partial \tilde{a}_{i}^{2 k+2}}{\partial y} & =\left[\delta_{1,1}\left(x_{i}\right)+\sum_{l=3}^{2 p} \delta_{1, l}\left(x_{i}\right) \frac{\partial \tilde{a}_{i}^{l}}{\partial x}\right] \frac{\partial \tilde{a}_{i}^{2 k+2}}{\partial x} \\
& +\left[\delta_{2,1}\left(x_{i}\right)+\sum_{l=3}^{2 p} \delta_{2, l}\left(x_{i}\right) \frac{\partial \tilde{a}_{i}^{l}}{\partial x}\right] \frac{\partial \tilde{a}_{i}^{2 k+2}}{\partial y} \\
& -\delta_{2 k+2,1}\left(x_{i}\right)-\sum_{l=3}^{2 p} \delta_{2 k+2, l} \frac{\partial \tilde{a}_{i}^{l}}{\partial x}  \tag{VI.57}\\
\frac{\partial \tilde{a}_{i}^{2 k+1}}{\partial y}+\frac{\partial \tilde{a}_{i}^{2 k+2}}{\partial x} & =\left[\delta_{1,2}\left(x_{i}\right)+\sum_{l=3}^{2 p} \delta_{1, l}\left(x_{i}\right) \frac{\partial \tilde{a}_{i}^{l}}{\partial y}\right] \frac{\partial \tilde{a}_{i}^{2 k+2}}{\partial x} \\
& +\left[\delta_{2,2}\left(x_{i}\right)+\sum_{l=3}^{2 p} \delta_{2, l}\left(x_{i}\right) \frac{\partial \tilde{a}_{i}^{l}}{\partial y}\right] \frac{\partial \tilde{a}_{i}^{2 k+2}}{\partial y} \\
& -\delta_{2 k+2,1}\left(x_{i}\right)-\sum_{l=3}^{2 p} \delta_{2 k+2, l}\left(x_{i}\right) \frac{\partial \tilde{a}_{i}^{l}}{\partial y}
\end{align*}\right.
$$

Then we deduce that there exist a linear map

$$
A\left(x_{i}, \cdot\right): \mathbb{R}^{2} \otimes \mathbb{R}^{2 p-2} \longrightarrow \mathbb{C}^{p-1} \otimes_{\mathbb{R}}\left(\mathbb{R}^{2} \otimes \mathbb{R}^{2 p-2}\right)^{*}
$$

and an element

$$
B\left(x_{i}, \cdot\right) \in \mathbb{C}^{p-1} \otimes_{\mathbb{R}}\left(\mathbb{R}^{2} \otimes \mathbb{R}^{2 p-2}\right)^{*}
$$

such that $\tilde{a}_{i}$ solves

$$
\begin{equation*}
\frac{\partial \tilde{a}_{i}}{\partial \bar{z}}=A\left(x_{i}, \nabla \tilde{a}_{i}\right) \cdot \nabla \tilde{a}_{i}+B\left(x_{i}, \nabla \tilde{a}_{i}\right)+D\left(x_{i}, \tilde{a}_{i}\right) . \tag{VI.58}
\end{equation*}
$$

Observe also that the dependence of $A$ and $B$ in $B_{\rho_{x_{0}}}\left(x_{0}\right)$ is smooth and that $A\left(x_{0}, \cdot\right)=0, B\left(x_{0}, \cdot\right)=0, D\left(x_{0}\right)=0$ and because of (VI.55) we have an estimate of the sort

$$
\begin{equation*}
\forall p \in \mathbb{R}^{2} \otimes \mathbb{R}^{2 p} \quad\left|A\left(x_{i}, p\right)\right|+\left|B\left(x_{i}, p\right)\right| \leq K\left|x_{i}\right|(1+|p|) . \tag{VI.59}
\end{equation*}
$$

Consider now $i$ and $j$ such that $B_{\rho_{i}}^{2}\left(z_{i}\right) \cap B_{\rho_{j}}^{2}\left(z_{j}\right) \neq \emptyset$. On $B_{\frac{7}{6} \rho_{i}}^{2}\left(z_{i}\right) \cap$ $B_{\frac{7}{6} \rho_{j}}^{2}\left(z_{j}\right) \tilde{a}_{i}-\tilde{a}_{j}$ solves the following equation

$$
\begin{align*}
\partial_{\bar{z}}\left(\tilde{a}_{i}-\tilde{a}_{j}\right)= & A\left(x_{i}, \nabla \tilde{a}_{i}\right) \cdot \nabla \tilde{a}_{i}+B\left(x_{i}, \nabla \tilde{a}_{i}\right)  \tag{VI.60}\\
& -A\left(x_{j}, \nabla \tilde{a}_{j}\right) \cdot \nabla \tilde{a}_{j}-B\left(x_{j}, \nabla \tilde{a}_{j}\right) \\
= & C\left(x_{i}, \nabla \tilde{a}_{i}, \nabla \tilde{a}_{j}\right) \cdot \nabla\left(\tilde{a}_{i}-\tilde{a}_{j}\right) \\
& +E\left(x_{i}, \nabla \tilde{a}_{j}\right)-E\left(x_{j}, \nabla \tilde{a}_{j}\right),
\end{align*}
$$

where
$\left(\mathrm{VI.61)} C\left(x_{i}, \nabla \tilde{a}_{i}, \nabla \tilde{a}_{j}\right) \cdot \nabla\left(\tilde{a}_{i}-\tilde{a}_{j}\right):=A\left(x_{i}, \nabla \tilde{a}_{i}\right) \cdot \nabla \tilde{a}_{i}-A\left(x_{i}, \nabla \tilde{a}_{j}\right) \cdot \nabla \tilde{a}_{j}\right.$

$$
+B\left(x_{i}, \nabla \tilde{a}_{i}\right)-B\left(x_{i}, \nabla \tilde{a}_{j}\right),
$$

(where we have used the linear dependence in p of $A\left(x_{i}, p\right)$ and $B\left(x_{i}, p\right)$ ), and where

$$
\begin{equation*}
E(x, p):=A(x, p) \cdot p+B(x, p)+D(x) . \tag{VI.62}
\end{equation*}
$$

Observe, on the one hand, that $C(x, p, q)$ has a linear dependence in $p$ and $q$ in $\mathbb{R}^{2} \otimes \mathbb{R}^{2 p-2}$, that

$$
\begin{equation*}
|C(x, p, q)| \leq K|x|(1+|p|+|q|), \tag{VI.63}
\end{equation*}
$$

and that the following estimates hold for $D(x, p)$, for all $l \in \mathbb{N}$ :

$$
\begin{equation*}
\left|\nabla_{x}^{l} E(x, p)\right| \leq K_{l}\left(1+|p|^{2}\right) . \tag{VI.64}
\end{equation*}
$$

On $B_{\frac{7}{6}}^{2}\left(\rho_{i}^{-1} z_{i}\right) \cap B_{\frac{7}{6}}^{2}\left(\rho_{i}^{-1} z_{j}\right)$ the function $f\left(z^{\prime}\right):=a_{i}\left(\rho_{i} z^{\prime}\right)-a_{j}\left(\rho_{i} z^{\prime}\right)$ solves

$$
\begin{equation*}
\partial_{\bar{z}^{\prime}} f-C\left(z^{\prime}\right) \cdot \nabla f=g\left(z^{\prime}\right), \tag{VI.65}
\end{equation*}
$$

where

$$
C\left(z^{\prime}\right):=C\left(x_{i}, \nabla \tilde{a}_{i}, \nabla \tilde{a}_{j}\right)\left(\rho_{i} z^{\prime}\right),
$$

and

$$
g\left(z^{\prime}\right):=\rho_{i}\left[D\left(x_{i}, \nabla \tilde{a}_{j}\left(\rho_{i} z^{\prime}\right)\right)-D\left(x_{j}, \nabla \tilde{a}_{j}\left(\rho_{i} z^{\prime}\right)\right)\right] .
$$

Using (VI.53), observe that for any $l \in \mathbb{N}$,

$$
\begin{align*}
& \left\|\nabla_{z^{\prime}}^{l}\left(C\left(x_{i},\left(\nabla_{z} \tilde{a}_{i}\right)\left(\rho_{i} z^{\prime}\right),\left(\nabla_{z} \tilde{a}_{j}\right)\left(\rho_{i} z^{\prime}\right)\right)\right)\right\|_{\infty}  \tag{VI.66}\\
& \quad \leq K\left|x_{i}\right| \rho_{i}^{l}\left[\left\|\nabla_{z}^{l+1} \tilde{a}_{i}\right\|_{\infty}+\left\|\nabla_{z}^{l+1} \tilde{a}_{j}\right\|_{\infty}\right] \leq K_{l} \rho_{x_{0}}
\end{align*}
$$

Therefore, for $\rho_{x_{0}}$ small enough, $L:=\partial_{\overline{z^{\prime}}}-C\left(z^{\prime}\right) \cdot \nabla_{z^{\prime}}$ is an elliptic coercive first order operator with smooth coefficients whose derivatives are uniformly bounded. Observe also that, using again (VI.53),

$$
\begin{align*}
& \left\|\nabla_{z^{\prime}}^{l} \rho_{i}\left[D\left(x_{i}, \nabla \tilde{a}_{j}\left(\rho_{i} z^{\prime}\right)\right)-D\left(x_{j}, \nabla \tilde{a}_{j}\left(\rho_{i} z^{\prime}\right)\right)\right]\right\|_{\infty}  \tag{VI.67}\\
& \quad \leq K \rho_{i}^{2}\left[\rho_{i}^{l} \sum_{s=0}^{[l / 2]}\left\|\nabla_{z}^{s+1} \tilde{a}_{j}\right\|_{\infty}\left\|\nabla_{z}^{l-s+1} \tilde{a}_{j}\right\|_{\infty}+\rho_{i}^{l}\left\|\nabla_{z^{\prime}}^{l+1} \tilde{a}_{j}\right\|_{\infty}\right] .
\end{align*}
$$

Then we have

$$
\begin{equation*}
\left\|\nabla_{z^{\prime}}^{l} g\right\|_{\infty} \leq K_{l} \rho_{i}^{2} \tag{VI.68}
\end{equation*}
$$

From (VI.46) we deduce that

$$
\begin{equation*}
\int_{B_{\frac{7}{6}}^{2}\left(\rho_{i}^{-1} z_{i}\right) \cap B_{\frac{7}{6}}^{2}\left(\rho_{i}^{-1} z_{j}\right)}|f| \leq K \rho_{i}^{2} . \tag{VI.69}
\end{equation*}
$$

Thus combining (VI.65)...(VI.69) and using standard elliptic estimates we obtain that for any $l \in \mathbb{N}$

$$
\begin{equation*}
\left\|\nabla^{l} f\right\|_{L^{\infty}\left(B_{1}^{2}\left(\rho_{i}^{-1} z_{i}\right) \cap B_{1}^{2}\left(\rho_{i}^{-1} z_{j}\right)\right)} \leq K_{l} \rho_{i}, \tag{VI.70}
\end{equation*}
$$

which yields, going back to the original scale, the estimate (VI.19) and Lemma VI. 2 is proved.

Definition of the approximated average curve. On $B_{\rho_{x_{0}}}^{2}(0)=\Pi\left(B_{\rho_{x_{0}}}\left(x_{0}\right)\right)$ we define the approximated average curve as follows. Let $\varphi$ be the partition of unity of $\Pi\left(\mathcal{C}_{Q-1} \cap B_{\rho_{x_{0}}}\left(x_{0}\right)\right)$ defined in (V.48). We denote

$$
\left\{\begin{array}{ll}
\tilde{a}\left(z_{0}\right):=\sum_{i \in I} \varphi\left(z_{0}\right) \tilde{a}_{i}\left(z_{0}\right) & \forall z_{0} \in \Pi\left(\mathcal{C}_{Q-1} \cap B_{\rho_{x_{0}}}\left(x_{0}\right)\right),  \tag{VI.71}\\
\tilde{a}\left(z_{0}\right):=\left\langle C, z, z_{0}\right\rangle(w) &
\end{array} \forall z_{0} \in \Pi\left(\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right) \cap B_{\rho_{x_{0}}}\left(x_{0}\right)\right) .\right.
$$

Observe that, because of Lemma IV.2, for any $z_{0} \in \Pi\left(\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right) \cap B_{\rho_{x_{0}}}\left(x_{0}\right)\right)$ the slice $\left\langle C, z, z_{0}\right\rangle$ consists of exactly one point and $\tilde{a}\left(z_{0}\right)$ is simply the $w$ coordinates of that point. The following estimates for $\tilde{a}$ holds :

Lemma VI.3. Under the above notation, for any $q<+\infty$ there exists a constant $K_{q}$ independent of $\rho_{x_{0}}$ such that

$$
\begin{equation*}
\int_{\frac{B_{\frac{\rho_{x_{0}}^{2}}{2}}^{2}(0)}{}\left|\nabla^{2} \tilde{a}\right|^{q} \leq K_{q} \rho_{x_{0}}^{2} . . . . . . . .} . \tag{VI.72}
\end{equation*}
$$

Proof of Lemma VI.3. We claim first that $\tilde{a}$ is a Lipschitz map over $B_{\rho_{x_{0}}}^{2}(0)$. In $\Pi\left(\mathcal{C}_{Q-1} \cap B_{\rho_{x_{0}}}\left(x_{0}\right)\right)$, by the assumptions of the inductive procedure, we know that $\tilde{a}$ is smooth and using both (VI.53) and (VI.19), because also of (V.49), we have

$$
\begin{equation*}
\|\nabla \tilde{a}\|_{L^{\infty}\left(\Pi\left(\mathcal{C}_{Q-1} \cap B_{\rho_{x_{0}}}\left(x_{0}\right)\right)\right)} \leq K . \tag{VI.73}
\end{equation*}
$$

Consider now two arbitrary points $x$ and $y$ of $B_{\rho_{x_{0}}}^{2}(0)$. Either the segment $[x, y]$ in $B_{\rho_{x_{0}}}^{2}(0)$ is included in $\Pi\left(\mathcal{C}_{Q-1} \cap B_{\rho_{x_{0}}}\left(x_{0}\right)\right)$ and then we can integrate (VI.73) all along that segment to get

$$
\begin{equation*}
|\tilde{a}(x)-\tilde{a}(y)| \leq K|x-y|, \tag{VI.74}
\end{equation*}
$$

or there exists $z \in[x, y] \cap \Pi\left(\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right) \cap B_{\rho_{x_{0}}}\left(x_{0}\right)\right)$. Using this time (IV.6) we also get (VI.74), which proves the desired claim. Using the equation (VI.58) solved by the $\tilde{a}_{i}$ s we obtain that, in $\Pi\left(\mathcal{C}_{Q-1} \cap B_{\rho_{x_{0}}}\left(x_{0}\right)\right)$, $\tilde{a}$ is a solution of

$$
\begin{equation*}
\partial_{\tilde{z}} \tilde{a}-A((z, \tilde{a}(z)), \nabla \tilde{a}) \cdot \nabla \tilde{a}-B((z, \tilde{a}(z)), \nabla \tilde{a})-D(z, \tilde{a}(z))=\zeta(z), \tag{VI.75}
\end{equation*}
$$

where

$$
\begin{align*}
\zeta(z):= & \sum_{i \in I} \varphi_{i}\left[A\left(x_{i}, \nabla \tilde{a}_{i}\right) \cdot \nabla \tilde{a}_{i}-A((z, \tilde{a}(z)), \nabla \tilde{a}) \cdot \nabla \tilde{a}\right]  \tag{VI.76}\\
& +\sum_{i \in I}\left[B\left(x_{i}, \nabla \tilde{a}_{i}\right)-B((z, \tilde{a}(z)), \nabla \tilde{a})\right] \\
& +\sum_{i \in I} \partial_{\bar{z}} \varphi_{i}\left[\tilde{a}_{i}-\tilde{a}\right] .
\end{align*}
$$

Observe that $\sum_{i \in I} \partial_{\bar{z}} \varphi_{i} \tilde{a}$ was added at the end because this quantity vanishes. Since for any $x_{i}$ and $r_{i}$, (V.10) holds, since also for any $z \in \Pi\left(\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right) \cap\right.$ $\left.B_{\rho_{x_{0}}}\left(x_{0}\right)\right)$ we have for an $\varepsilon$ as small as we want (recall $\alpha$ was fixed independently of $\varepsilon$ ), the relative Lipschitz estimate (IV.6) holds and granting the fact that $\Pi\left(\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right) \cap B_{\rho_{x_{0}}}\left(x_{0}\right)\right)$ is a compact subset of $B_{\rho_{x_{0}}}\left(x_{0}\right)$, it is clear that

$$
\begin{align*}
& \forall \eta>0 \quad \exists \delta>0: \forall i \in I  \tag{VI.77}\\
& \operatorname{dist}\left(B_{\rho_{i}}\left(z_{i}\right), \Pi\left(\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right)\right)\right) \leq \delta \quad \Longrightarrow \quad\left|\rho_{i}\right| \leq \eta .
\end{align*}
$$

Observe that

$$
\begin{align*}
&|\zeta(z)| \leq K  \tag{VI.78}\\
& \sum_{i, j \in I} \varphi_{i}(z) \varphi_{j}(z)\left|\tilde{a}_{j}-w_{i}\right| \\
& \sum_{i, j \in I} \varphi_{i}(z) \varphi_{j}(z)\left|\nabla\left(\tilde{a}_{i}-\tilde{a}_{j}\right)\right| \\
& \sum_{i, j \in I}\left|\nabla \varphi_{i}(z)\right| \varphi_{j}(z)\left|\tilde{a}_{i}(z)-\tilde{a}_{j}(z)\right| .
\end{align*}
$$

Using (VI.19), we have that

$$
\begin{equation*}
|\zeta(z)| \leq K \max \left\{r_{i} i \in I ; z \in B_{\rho_{i}}^{2}\left(z_{i}\right)\right\} \leq K \operatorname{dist}\left(z, \Pi\left(\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right)\right)\right) \tag{VI.79}
\end{equation*}
$$

Combining (VI.77) and (VI.79), we have that $\zeta(z)$ converges uniformly to 0 as $z$ tends to $\Pi\left(\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right)\right)$. We then extend $\zeta$ by 0 in $\Pi\left(\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right)\right)$. Now $\zeta$ is a continuous function in $B_{\rho_{x_{0}}}^{2}\left(x_{0}\right)$ and we claim that (VI.80)

$$
\partial_{\tilde{z}} \tilde{a}-A((z, \tilde{a}(z)), \nabla \tilde{a}) \cdot \nabla \tilde{a}-B((z, \tilde{a}(z)), \nabla \tilde{a})=\zeta(z) \quad \text { in } \quad \mathcal{D}^{\prime}\left(B_{\rho_{x_{0}}}^{2}\left(x_{0}\right)\right) .
$$

Since $\tilde{a}$ is a Lipschitz function in $B_{\rho_{x_{0}}}^{2}\left(x_{0}\right), \partial_{\tilde{z}} \tilde{a}-A((z, \tilde{a}(z)), \nabla \tilde{a}) \cdot \nabla \tilde{a}-$ $B((z, \tilde{a}(z)), \nabla \tilde{a})$ is a bounded function in $B_{\rho_{x_{0}}}^{2}\left(x_{0}\right)$ and therefore, in order to prove (VI.80), it suffices to prove that for $\mathcal{H}^{2}$ almost every $z$ in $\Pi\left(\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right)\right) \cap$ $B_{\rho_{x_{0}}}^{2}\left(x_{0}\right)$,

$$
\begin{equation*}
\partial_{\tilde{z}} \tilde{a}-A((z, \tilde{a}(z)), \nabla \tilde{a}) \cdot \nabla \tilde{a}-B((z, \tilde{a}(z)), \nabla \tilde{a})-D(z, \tilde{a}(z))=0 . \tag{VI.81}
\end{equation*}
$$

This latter equality, because of the computations between (VI.54) ...(VI.58), is just equivalent to the fact that the tangent plane of the graph $(z, \tilde{a}(z))$ at that point is $J$-holomorphic, which is the case at every point of $\Pi\left(\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right)\right) \cap$
$B_{\rho_{x_{0}}}^{2}\left(x_{0}\right)$ due to Lemma IV.2. Thus (VI.80) is proved. Using again (VI.19) we observe that

$$
\begin{align*}
|\nabla \zeta(z)| \leq & K \sum_{i, j \in I} \varphi_{i}(z) \varphi_{j}(z)\left[\left|\nabla \tilde{a}_{j}\right|+\left|\nabla \tilde{a}_{j}\right|^{3}\right]  \tag{VI.82}\\
& \sum_{i, j \in I} \varphi_{i}(z) \varphi_{j}(z)\left|\nabla \tilde{a}_{i}\right|(z)\left|\nabla^{2}\left(\tilde{a}_{i}-\tilde{a}_{j}\right)\right|(z) \\
& \sum_{i, j \in I}\left[\left|\nabla^{2} \varphi_{i}(z)\right| \varphi_{j}(z)+\left|\nabla \varphi_{i}\right|(z)\left|\nabla \varphi_{j}(z)\right|\right]\left|\tilde{a}_{i}(z)-\tilde{a}_{j}(z)\right|+\left|\tilde{a}_{j}-w_{i}\right| \\
\leq & K
\end{align*}
$$

We claim now that $\zeta$ is Lipschitz in $B_{\rho_{x_{0}}}^{2}\left(x_{0}\right)$. Indeed, arguing as for $\tilde{a}$, given $x$ and $y$ in $B_{\rho_{x_{0}}}^{2}\left(x_{0}\right)$, if the segment $[x, y]$ has no intersection with $\Pi\left(\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right)\right)$, then, integrating (VI.78) on that segment gives

$$
\begin{equation*}
|\zeta(x)-\zeta(y)| \leq K|x-y| \tag{VI.83}
\end{equation*}
$$

Otherwise, if there exists $z \in[x, y] \cap \Pi\left(\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right)\right)$, then, (VI.79) gives

$$
|\zeta(x)|+|\zeta(y)| \leq K|x-z|+|y-z|
$$

which gives (VI.83) and the claim is proved. We claim now that $\tilde{a} \in$ $W^{2, q}\left(B_{\rho_{x_{0} / 2}}^{2}(0)\right)$ for any $q<+\infty$. Let $e$ be a unit vector in $B_{\rho_{x_{0}} / 2}^{2}(0)$. For small $h$ we denote $\tilde{a}_{h}(z):=\tilde{a}(z+h)$ and $\zeta_{h}(z):=\zeta(z+h)$. We have, using the linear dependencies of

$$
\begin{align*}
\partial_{\bar{z}}\left(\tilde{a}-\tilde{a}_{h}\right) & -A((z, \tilde{a}), \nabla \tilde{a}) \cdot \nabla\left(\tilde{a}-\tilde{a}_{h}\right)-A\left((z, \tilde{a}), \nabla\left(\tilde{a}-\tilde{a}_{h}\right)\right) \cdot \nabla \tilde{a}_{h}  \tag{VI.84}\\
& -B\left((z, \tilde{a}), \nabla\left(\tilde{a}-\tilde{a}_{h}\right)\right) \\
= & \zeta-\zeta_{h}+A\left((z, \tilde{a}), \nabla \tilde{a}_{h}\right) \cdot \nabla \tilde{a}_{h}-A\left(\left(z+h, \tilde{a}_{h}\right), \nabla \tilde{a}_{h}\right) \cdot \nabla \tilde{a}_{h} \\
& B\left((z, \tilde{a}), \nabla \tilde{a}_{h}\right)-B\left(\left(z+h, \tilde{a}_{h}\right), \nabla \tilde{a}_{h}\right)+D(z, \tilde{a})-D\left(z+h, \tilde{a}_{h}\right) .
\end{align*}
$$

Denote $\Delta_{h}$ the right-hand side of (VI.84) and observe that there exists a constant $K$ such that

$$
\begin{equation*}
\left|\Delta_{h}\right| \leq K h \tag{VI.85}
\end{equation*}
$$

Let $\chi_{\rho_{x_{0}}}$ be a cut-off function such that $\chi_{\rho_{x_{0}}} \equiv 1$ in $B_{\frac{\rho_{x_{0}}}{2}}^{2}(0)$ and $\chi_{\rho_{x_{0}}} \equiv 0$ in $\mathbb{R}^{2} \backslash B_{\rho_{x_{0}}}^{2}(0)$. Let $f_{h}:=\chi_{\rho_{x_{0}}}\left(\tilde{a}-\tilde{a}_{h}\right)$, and let $L_{h}$ be the operator such that (VI.86) $L_{h} f:=\partial_{\bar{z}} f-A((z, \tilde{a}), \nabla \tilde{a}) \cdot \nabla f-A((z, \tilde{a}), \nabla f) \cdot \nabla \tilde{a}_{h}-B((z, \tilde{a}), \nabla f)$.

Now,

$$
\begin{align*}
L_{h} f_{h}= & \chi_{\rho_{x_{0}}} \Delta_{h}+\left(\tilde{a}-\tilde{a}_{h}\right) \partial_{\tilde{z}} \chi_{\rho_{x_{0}}}-A((z, \tilde{a}), \nabla \tilde{a}) \cdot\left(\tilde{a}-\tilde{a}_{h}\right) \nabla \chi_{\rho_{x_{0}}}  \tag{VI.87}\\
& -A\left((z, \tilde{a}),\left(\tilde{a}-\tilde{a}_{h}\right) \nabla \chi_{\rho_{x_{0}}}\right) \cdot \nabla \tilde{a}_{h}-B\left((z, \tilde{a}),\left(\tilde{a}-\tilde{a}_{h}\right) \nabla \chi_{\rho_{x_{0}}}\right) .
\end{align*}
$$

Since $\tilde{a}$ is Lipschitz, using (VI.85) and (VI.59), we have that

$$
\begin{equation*}
\left|L_{h} f_{h}\right| \leq K h \tag{VI.88}
\end{equation*}
$$

Observe that $\left|L_{h} f_{h}\right| \geq\left|\partial_{\bar{z}} f_{h}\right|-K \rho_{x_{0}}\left|\nabla f_{h}\right|$. Since $f_{h}=0$ on $\partial B_{\rho_{x_{0}}}^{2}(0)$, for any $p<+\infty$ we have that
(VI.89)

$$
\int_{B_{\rho_{x_{0}}}^{2}(0)}\left|\nabla f_{h}\right|^{q} \leq K_{q} \int_{B_{\rho_{x_{0}}}^{2}(0)}\left|\partial_{z} f_{h}\right|^{q} \leq K_{q} \int_{B_{\rho_{x_{0}}}^{2}(0)}\left|L_{h} f_{h}\right|^{q}+K \rho_{x_{0}}^{q}\left|\nabla f_{h}\right|^{q}
$$

Dividing by $h^{q}$ and making $h$ tend to zero, we get that for $\rho_{x_{0}}$ small enough inequality (VI.72) holds and Lemma VI. 3 is proved.
VI.2. Constructing adapted coordinates to $C$ in a neighborhood of $x_{0} \in$ $\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}$. We consider a point $x_{0}$ in the support of our $J$-holomorphic current $C$ and we assume, as above, that the multiplicity at $x_{0}$ is $Q$ and that the tangent cone $C_{0, x_{0}}$ is $Q$ times a $J_{x_{0}}$-holomorphic disk $D$. We start with the coordinates $\left(z, w_{1}, \ldots, w_{p-1}\right)$ chosen in (II.1) such that $C_{0, x_{0}}=Q[D]$ is $Q$ times the "horizontal" disk given by $w_{i}=0$ for $i=1, \ldots, p-1$ and we work in the ball $B_{\rho_{x_{0}}}^{2 p}\left(x_{0}\right)$ whose radius $\rho_{x_{0}}$ is given by (V.23). The purpose of this subsection is to construct new coordinates $\left(\xi, \lambda_{1}, \ldots, \lambda_{p-1}\right)$ in $B_{\rho_{x_{0}}}^{2 p}\left(x_{0}\right)$ such that the set $\lambda_{i}=0$ for $i=1, \ldots, p-1$ coincides with the graph of the average map $\tilde{a}$ constructed in the previous subsection.

On the graph $\tilde{A}(z):=(z, \tilde{a}(z))$ we consider the complex structure $j$ given by the metric induced by $g:=\omega(J \cdot, \cdot)$. Let $X$ be a vector tangent to $\tilde{A}$ at $(z, \tilde{a}(z))$. We compare $j X$ and $J X$ in $\mathbb{R}^{2 p}$. Let $n(z, \tilde{a}(z))$ be the $2 p-2$-unit vector normal to $T_{(z, \tilde{a}(z))} \tilde{A}$, making the identification between $2 p-1$-vectors and the vector given by the ambient metric,

$$
j X:=n \wedge X .
$$

We have seen that

$$
|J n-n|(z, \tilde{a}(z)) \leq r_{(z, \tilde{a}(z))} .
$$

Therefore

$$
|j J X-J j X| \leq|n \wedge J X-J(n \wedge X)| \leq|n \wedge J X-J n \wedge J X| \leq r_{(z, \tilde{a}(z))}|X|
$$

Thus, $|(J-j)(J+j) X| \leq r_{(z, \tilde{a}(z))}|X|$ where we extend $j$ to the normal bundle to $\tilde{A}$ again by means of the induced metric. Since $|(J+j) X|$ and $|X|$ are comparable independent of $X$, we have

$$
\begin{equation*}
\forall X \in T_{(z, \tilde{a}(z))} \tilde{A} \quad|(J-j) X| \leq K r_{(z, \tilde{a}(z))}|X| \tag{VI.90}
\end{equation*}
$$

We choose now coordinates $\xi=\left(\xi_{1}, \xi_{2}\right)$ on $\tilde{A}$ compatible with $j$ (i.e. $j \frac{\partial}{\partial \xi_{1}}=$ $\left.\frac{\partial}{\partial \xi_{2}}\right)$. Let $\left(z^{\prime}, \hat{a}\left(z^{\prime}\right)\right):=\left(\rho_{x_{0}}^{-1}, \rho_{x_{0}}^{-1} \tilde{a}\left(\rho_{x_{0}} z^{\prime}\right)\right)$ and in $B_{1}^{2 p}(0)$ consider the metric $\hat{g}\left(z^{\prime}, w^{\prime}\right):=\rho_{x_{0}}^{-2}\left(\rho_{x_{0}} z^{\prime}, \rho_{x_{0}} w^{\prime}\right)^{*} g$ where $g$ in $B_{\rho_{x_{0}}}^{2 p}(0)$ is the original metric $g(\cdot, \cdot)=$ $\omega(J \cdot, \cdot)$. After this scaling we have, using (VI.72),

$$
\begin{equation*}
\int_{B_{1}^{2}}\left|\nabla^{2} \hat{a}\right|^{p} d z^{\prime} \leq K_{q} \rho_{x_{0}}^{q} \tag{VI.91}
\end{equation*}
$$

and

$$
\begin{equation*}
\hat{g}^{i j}=\delta^{i j}+h^{i j} \quad \text { where } h^{i j}(0,0)=0 \quad \text { and }\left\|\nabla h^{i j}\right\|_{\infty} \leq K \rho_{x_{0}} \tag{VI.92}
\end{equation*}
$$

We look for isothermal coordinates $\left(\xi_{1}^{\prime}, \xi_{2}^{\prime}\right)$ in $\hat{A}=\left\{\left(z^{\prime}, \hat{a}\left(z^{\prime}\right)\right), z^{\prime} \in B_{1}^{2}(0)\right\}$ of the form $\xi^{\prime}=z^{\prime}+\delta\left(z^{\prime}\right)$ where $\delta$ will be small in $W^{2, p}$. On $B_{1}^{2}(0)$ we consider the metric $\hat{k}=\left(z^{\prime}, \hat{a}\left(z^{\prime}\right)\right)^{*} \hat{g}=\left(1+k_{11}\right)\left(d x_{1}^{\prime}\right)^{2}+2 k_{12} d x_{1}^{\prime} d x_{2}^{\prime}+\left(1+k_{22}\right)\left(d x_{2}^{\prime}\right)^{2}$. From the estimates above we have for any $q>0$, (since $\nabla \hat{a}(0,0)=0$ and $\left\|\nabla^{2} \hat{a}\right\|_{q} \leq \rho_{x_{0}}$ that for $\left.\|\nabla \hat{a}\|_{\infty} \leq \rho_{x_{0}}\right)$,

$$
\begin{equation*}
\int_{B_{1}^{2}}|\nabla k|^{q} \leq K_{q} \rho_{x_{0}}^{2 q} \tag{VI.93}
\end{equation*}
$$

Following [DNF, pp. 110-111], it suffices to find $\delta_{1}$ solving (VI.94)

$$
-\frac{\partial}{\partial x_{1}^{\prime}}\left[\frac{\left(1+k_{11}\right) \frac{\partial \delta_{1}}{\partial x_{1}^{\prime}}-k_{12} \frac{\partial \delta_{1}}{\partial x_{2}^{\prime}}}{\sqrt{\left(1+k_{11}\right)\left(1+k_{22}\right)-k_{12}^{2}}}\right]-\frac{\partial}{\partial x_{2}^{\prime}}\left[\frac{\left(1+k_{22}\right) \frac{\partial \delta_{1}}{\partial x_{2}^{\prime}}-k_{12} \frac{\partial \delta_{1}}{\partial x_{1}^{\prime}}}{\sqrt{\left(1+k_{11}\right)\left(1+k_{22}\right)-k_{12}^{2}}}\right]=0 .
$$

Taking $\delta_{1}=0$ on $\partial B_{1}^{2}(0)$ we get a well-posed elliptic problem and obtain the existence of $\delta_{1}$ satisfying

$$
\sum_{i=1}^{2} a_{i j} \frac{\partial^{2} \delta_{1}}{\partial x_{i}^{\prime} \partial x_{j}^{\prime}}=F \cdot \nabla \delta_{1}
$$

where $a_{i j}$ are Hölder continuous, $\left\|a_{i j}-\delta_{i j}\right\|_{C^{0, \alpha}\left(B_{1}^{2}\right)} \leq K_{\alpha} \rho^{2}$ and $F \in L^{q}$ with $\int|F|^{q} \leq K_{p} \rho_{x_{0}}^{2 q}$. Standard elliptic estimates give then

$$
\begin{equation*}
\left\|\delta_{1}\right\|_{W^{2, q}\left(B_{1}^{2}\right)} \leq K \rho_{x_{0}}^{2} \tag{VI.95}
\end{equation*}
$$

Therefore, going back to the original scale, we have found coordinates $\xi_{i}=$ $x_{i}+\rho_{x_{0}} \delta_{i}\left(\rho_{x_{0}}^{-1} z\right)=x_{i}+\alpha_{i}(z)$ such that

$$
\begin{equation*}
\|\nabla \alpha\|_{\infty} \leq K \rho_{x_{0}}^{2} \quad \text { and } \quad j \frac{\partial}{\partial \xi_{1}}=\frac{\partial}{\partial \xi_{2}} \tag{VI.96}
\end{equation*}
$$

We translate these coordinates in such a way that $\alpha(0,0)=(0,0)$.
Inside $G l\left(\mathbb{R}^{2 p}\right)$, the space of invertible $2 p \times 2 p$ matrices with real coefficients, we denote $U(p)$ the subspace of matrices $M$ which commute with $J_{0}$. $U(p)$ is a compact submanifold of $G l\left(\mathbb{R}^{2 p}\right)$ and for some metric in $G l\left(\mathbb{R}^{2 p}\right)$ we
denote $\pi_{U(p)}$ the orthogonal projection from a neighborhood of $U(p)$ onto $U(p)$. We consider $\mathrm{M}(z)$ the matrix which is given by

$$
\begin{equation*}
\mathrm{M}(z):=\Lambda_{(z, \tilde{a}(z))}^{-1} \pi_{U(p)}\left(\Lambda_{(z, \tilde{a}(z))}\right) \tag{VI.97}
\end{equation*}
$$

where we recall that $\Lambda_{x}$ is given by (VI.6). We have clearly, since $\|\nabla \tilde{a}\| \leq K$ and $\int_{B_{\rho_{x_{0}}}^{2}(0)}\left|\nabla^{2} \tilde{a}\right|^{q} \leq K_{q} \rho_{x_{0}}^{2}$ for any $p<+\infty$,

$$
\begin{equation*}
\|\nabla \mathrm{M}(z)\|_{L^{\infty}\left(B_{\rho x_{0}}^{2}(0)\right)} \leq K \quad \text { and } \quad\left\|\nabla^{2} \mathrm{M}(z)\right\|_{L^{q}\left(B_{\rho_{x_{0}}}^{2}(0)\right)} \leq K_{q} \rho_{x_{0}}^{\frac{2}{\bar{q}}} \tag{VI.98}
\end{equation*}
$$

We continue denoting $e_{1}, e_{2}, \ldots, e_{2 p}$ the canonical basis of $\mathbb{R}^{2 p}$. Let

$$
\begin{equation*}
\varepsilon_{k}(z):=\mathrm{M}(z) \cdot e_{i} \tag{VI.99}
\end{equation*}
$$

We have for all $i=1 \ldots p$ that
(VI.100) $J\left((z, \tilde{a}(z)) \cdot \varepsilon_{2 i-1}(z)=J\left((z, \tilde{a}(z)) \cdot \Lambda_{(z, \tilde{a}(z))}^{-1} \pi_{U(p)}\left(\Lambda_{(z, \tilde{a}(z))}\right) \cdot e_{2 i-1}\right.\right.$

$$
\begin{aligned}
& =\Lambda_{(z, \tilde{a}(z))}^{-1} J_{0} \pi_{U(p)}\left(\Lambda_{(z, \tilde{a}(z))}\right) \cdot e_{2 i-1} \\
& =\Lambda_{(z, \tilde{a}(z))}^{-1} \pi_{U(p)}\left(\Lambda_{(z, \tilde{a}(z))}\right) J_{0} \cdot e_{2 i-1} \\
& =\Lambda_{(z, \tilde{a}(z))}^{-1} \pi_{U(p)}\left(\Lambda_{(z, \tilde{a}(z))}\right) \cdot e_{2 i} \\
& =\varepsilon_{2 i}(z) .
\end{aligned}
$$

In $B_{\rho_{x_{0}}}^{2 p}\left(x_{0}\right)$ we consider the new coordinates $(\xi, \lambda)$ given by

$$
\begin{equation*}
\Psi:(\xi, \lambda) \longrightarrow \Psi(\xi, \lambda):=(z(\xi), \tilde{a}(z(\xi)))+\sum_{l=1}^{2 p} \lambda_{l} \varepsilon_{l+2}(z(\xi)) \tag{VI.101}
\end{equation*}
$$

Letting $\tilde{J}$ be the expression of the almost complex structure in these coordinates (i.e. $\tilde{J}_{(\xi, \lambda)} \cdot X:=d \Psi^{-1} J_{\Psi(\xi, \lambda)} \cdot d \Psi \cdot X$ ), we shall now estimate $\left|\tilde{J}-J_{0}\right|$ for points satisfying $|\lambda| \leq r_{\Psi(\xi, \lambda)}$; recall that $r_{x}$ was defined in the beginning of Section VI (see (V.11) and (V.12)), which corresponds in $B_{\rho_{x_{0}}}^{2 p}\left(x_{0}\right)$ to a neighborhood of $\tilde{A}(z)=(z, \tilde{a}(z))$ containing the support of $C$. We have first for $i=1,2$, using (VI.90),

$$
\begin{align*}
d \Psi \tilde{J}_{(\xi, 0)} e_{i}=J_{\Psi(\xi, 0)} \cdot \frac{\partial}{\partial \xi_{i}} & =j_{\Psi(\xi, 0)} \cdot \frac{\partial}{\partial \xi_{i}}+\left(J_{\Psi(\xi, 0)}-j_{\Psi(\xi, 0)}\right) \cdot \frac{\partial}{\partial \xi_{i}}  \tag{VI.102}\\
& =(-1)^{i+1} \frac{\partial}{\partial \xi_{i+1}}+O\left(r_{\Psi}(\xi, 0)\right)
\end{align*}
$$

where we are using the convention $\frac{\partial}{\partial \xi_{i+1}}=\frac{\partial}{\partial \xi_{i-1}}$. We have then, for $1<l \leq p$,

$$
\begin{equation*}
d \Psi \tilde{J}_{(\xi, 0)} e_{2 l}=J_{\Psi(\xi, 0)} \frac{\partial}{\partial \lambda_{2 l}}=J_{\Psi(\xi, 0)} \varepsilon_{2 l}=-\varepsilon_{2 l-1}=\frac{\partial}{\partial \xi_{2 l+1}} \tag{VI.103}
\end{equation*}
$$

For $|\lambda| \leq r_{\Psi(\xi, \lambda)}$ we have for $i=1,2$,

$$
\begin{align*}
d \Psi \tilde{J}_{(\xi, \lambda)} e_{i}= & J_{\Psi(\xi, \lambda)} \cdot d \Psi_{(\xi, \lambda)} e_{i}=J_{\Psi(\xi, 0)} \cdot d \Psi_{(\xi, 0)} e_{i}  \tag{VI.104}\\
& +J_{\Psi(\xi, 0)} \cdot\left[d \Psi_{(\xi, \lambda)} e_{i}-d \Psi_{(\xi, 0)} e_{i}\right] \\
& +\left[J_{\Psi(\xi, \lambda)}-J_{\Psi(\xi, 0)}\right] \cdot d \Psi_{(\xi, \lambda)} e_{i} .
\end{align*}
$$

Using the facts that $\left|J_{\Psi(\xi, \lambda)}-J_{\Psi(\xi, 0)}\right| \leq\|J\|_{C^{1}}|\Psi(\xi, \lambda)-\Psi(\xi, 0)| \leq K r_{\Psi(\xi, \lambda)}$ and that $d \Psi_{(\xi, \lambda)} e_{i}-d \Psi_{(\xi, 0)} e_{i}=\frac{\partial \Psi}{\partial \xi_{i}}(\xi, \lambda)-\frac{\partial \Psi}{\partial \xi_{i}}(\xi, 0)=\sum_{l=1} \lambda_{l} \partial_{\xi_{i}} \varepsilon_{l+2}$, we have then

$$
\begin{equation*}
\left|d \Psi \tilde{J}_{(\xi, \lambda)} e_{1}-J_{\Psi(\xi, 0)} \cdot d \Psi_{(\xi, 0)} e_{1}\right| \leq O\left(r_{\Psi(\xi, \lambda)}\right) . \tag{VI.105}
\end{equation*}
$$

Using (VI.90) again, we have $\left|\left[J_{\Psi(\xi, 0)}-j(\Psi(\xi, 0))\right] \cdot d \Psi_{(\xi, 0)} e_{i}\right| \leq O\left(r_{\Psi(\xi, 0)}\right)$, thus, since $\Psi$ is Lipschitz, $r_{\Psi(\xi, \lambda)}$ and $r_{\Psi(\xi, 0)}$ from Lemma V. 3 are comparable and (VI.105) implies

$$
\begin{align*}
& \mid d \Psi \tilde{J}_{(\xi, \lambda)} e_{1}-d \Psi_{(\xi, \lambda)} \cdot e_{2} \mid  \tag{VI.106}\\
& \leq\left|d \Psi \tilde{J}_{(\xi, \lambda)} e_{1}-d \Psi_{(\xi, 0)} \cdot e_{2}\right|+\left|d \Psi_{(\xi, 0)} \cdot e_{2}-d \Psi_{(\xi, \lambda)} \cdot e_{2}\right|
\end{align*}
$$

Using again the fact that $\left|d \Psi_{(\xi, 0)} \cdot e_{2}-d \Psi_{(\xi, \lambda)} \cdot e_{2}\right|=\left|\frac{\partial \Psi}{\partial \xi_{2}}(\xi, \lambda)-\frac{\partial \Psi}{\partial \xi_{2}}(\xi, 0)\right|=$ $\left|\sum_{l=1} \lambda_{l} \partial_{\xi_{2}} \varepsilon_{l+2}\right| \leq O\left(r_{\Psi(\xi, \lambda)}\right)$, we have finally that

$$
\begin{equation*}
\left|d \Psi \tilde{J}_{(\xi, \lambda)} e_{1}-d \Psi_{(\xi, \lambda)} \cdot e_{2}\right| \leq O\left(r_{\Psi(\xi, \lambda)}\right) \tag{VI.107}
\end{equation*}
$$

Finally, for $1<l \leq p$,

$$
\begin{align*}
d \Psi \tilde{J}_{(\xi, \lambda)} e_{2 l}= & J_{\Psi(\xi, \lambda)} \cdot d \Psi_{(\xi, \lambda)} e_{2 l}=-d \Psi_{(\xi, \lambda)} e_{2 l-1}  \tag{VI.108}\\
& +\left[d \Psi_{(\xi, \lambda)}-d \Psi_{(\xi, 0)}\right] e_{2 l-1}+J_{\Psi(\xi, 0)} \cdot\left[d \Psi_{(\xi, \lambda)}-d \Psi_{(\xi, 0)} e_{2 l}\right] \\
& +\left[J_{\Psi(\xi, \lambda)}-J_{\Psi(\xi, 0)}\right] \cdot d \Psi_{(\xi, \lambda)}
\end{align*}
$$

Using the estimates from the above lines, (VI.108) becomes for $1<l \leq p$,

$$
\begin{equation*}
\left|d \Psi \tilde{J}_{(\xi, \lambda)} e_{2 l}+d \Psi_{(\xi, \lambda)} e_{2 l-1}\right| \leq O\left(r_{\Psi(\xi, \lambda)}\right) . \tag{VI.109}
\end{equation*}
$$

Thus combining (VI.107) and (VI.109), we obtain

$$
\begin{equation*}
\forall(\xi, \lambda), \text { such that }|\lambda| \leq r_{\Psi(\xi, \lambda)}, \quad\left|\tilde{J}_{(\xi, \lambda)}-J_{0}\right| \leq K r_{\Psi(\xi, \lambda)} . \tag{VI.110}
\end{equation*}
$$

## VII. The unique continuation argument

In this part we show that, assuming $\mathcal{P}_{Q-1}$ (recall that the definition is given by (II.8)), a point $x_{0}$ in $\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}$ is isolated in $\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}$ unless all points in $\mathcal{C}_{*}$ in a neighborhood from $x_{0}$ are of multiplicity $Q$. This fact has already been proved in Section IV in the case where $C_{0, x_{0}}$ was not $Q$ times the same flat holomorphic disk (the easy case). Here we assume that we are in the difficult case $C_{0, x_{0}}=Q\left[D_{0}\right]$ where $D_{0}$ is the horizontal unit disk as before.

We adopt the coordinate system about $x_{0}$ constructed in Section VIII.2. We denote by $\Pi$ the map that assigns the first complex coordinate $\xi=\xi_{1}+i \xi_{2}$. Assuming there exists a sequence of points $x_{n} \in \mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}$ different from $x_{0}$ and converging to $x_{0}$, the goal of this section is to show that $C$ in a neighborhood is a $Q$ times the same graph. The strategy is inspired by [Ta, Ch. 1]; in our coordinates, the points in $\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}$ are contained in the disk $\lambda_{i}=0$ for $i=1, \ldots, 2 p-2$ and we shall use a unique continuation argument based on the proof of a Carleman estimate to show that our assumption implies that the whole cycle in the neighborhood of $x_{0}$ is included in that disk. Let $\left(\xi_{n}, 0\right)$ be the coordinates of $x_{n} \rightarrow x_{0}$. We can always extract a subsequence such that $\left|\xi_{n+1}\right| \leq\left|\xi_{n}\right|^{2}$. We then introduce the function $g_{N}(\xi):=\prod_{j=1}^{N}\left(\xi-\xi_{n}\right)$. Because of the speed of convergence of our sequence $\xi_{n}$ to zero it is not difficult to check that there exists a constant $K$ independent of $N$ such that for any $\xi \in B_{\rho_{x_{0}}}^{2}$ the following holds:

$$
\begin{equation*}
\frac{K^{-1}}{|\xi|^{N-N_{\xi}}} \frac{1}{\left|\xi-\xi_{N_{\xi}}\right|} \prod_{j=1}^{N_{\xi}-1} \frac{1}{\left|\xi_{j}\right|} \leq\left|g_{N}^{-1}\right|(\xi) \leq \frac{K}{|\xi|^{N-N_{\xi}}} \frac{1}{\left|\xi-\xi_{N_{\xi}}\right|} \prod_{j=1}^{N_{\xi}-1} \frac{1}{\left|\xi_{j}\right|} \tag{VII.1}
\end{equation*}
$$

where $N_{\xi}$ is the index less than $N$ such that $\left|\xi-\xi_{N_{\xi}}\right|$ is minimal among the $\left|\xi-\xi_{n}\right|$. It is also straightforward to check that

$$
\begin{align*}
& \left|\nabla g_{N}^{-1}\right|(\xi) \leq \frac{K\left(N-N_{\xi}\right)}{|\xi|^{N-N_{\xi}+1}} \frac{1}{\left|\xi-\xi_{N_{\xi}}\right|} \prod_{j=1}^{N_{\xi}-1} \frac{1}{\left|\xi_{j}\right|}  \tag{VII.2}\\
& +\frac{K}{|\xi|^{N-N_{\xi}}} \frac{1}{\mid \xi-\xi_{\left.N_{\xi}\right|^{2}}} \prod_{j=1}^{N_{\xi}-1} \frac{1}{\left|\xi_{j}\right|}+\frac{K}{|\xi|^{N-N_{\xi}}} \frac{1}{\left|\xi-\xi_{N_{\xi}}\right|} \sum_{l=1}^{N_{\xi}-1} \frac{1}{\left|\xi_{l}\right|} \prod_{j=1}^{N_{\xi}-1} \frac{1}{\left|\xi_{j}\right|}
\end{align*}
$$

and we have a corresponding estimate for $\left|\nabla^{k} g_{N}^{-1}\right|(\xi)$ for arbitrary $k$ in general. Let $\xi \in \Pi\left(\mathcal{C}_{Q-1}\right)$ with $\xi$ belonging to some $B_{\rho_{i}}^{2}\left(\xi_{i}\right)$ of the covering constructed in Section VI. To every such $i$ we assign $k_{i}$, an index such that $\left|\xi_{k_{i}}\right|+\rho_{k_{i}} \leq\left|\xi_{i}\right|-\rho_{i}$ and such that $\left|\xi_{k_{i}}-\xi_{i}\right| \leq K \rho_{i}$ and such that $\rho_{i}$ and $\rho_{k_{i}}$ are comparable:

$$
\begin{equation*}
K^{-1} \rho_{i} \leq \rho_{k_{i}} \leq K \rho_{i} . \tag{VII.3}
\end{equation*}
$$

This is always possible due to the Whitney-Besicovitch nature of our covering; moreover for every $k$ there exists a uniformly bounded number of $i$ such that $k_{i}=k$. Observe also, because of the relative Lipschitz estimate (IV.6) with constant $\varepsilon$ and because of the "splitting stage" of $C_{\xi_{i}, \rho_{i}}$ characterized by (V.10) we have that for any $\delta>0$ one may choose $\varepsilon$ small enough compared to $\alpha$ defined in Section V such that for any $\xi \in \Pi\left(\mathcal{C}_{Q-1}\right)$

$$
\begin{equation*}
\operatorname{dist}\left(\xi, \Pi\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right)\right) \geq \delta^{-1} \rho_{i} \tag{VII.4}
\end{equation*}
$$

where $\xi \in B_{\rho_{i}}^{2}\left(\xi_{i}\right)$. Combining (VII.3) and (VII.4) we get that $\forall i \in I, \forall \xi \in$ $B_{\rho_{i}}^{2}\left(\xi_{i}\right), \forall \zeta \in B_{\rho_{k_{i}}}^{2}\left(\xi_{k_{i}}\right)$,
(VII.5)
$\frac{1}{2} \operatorname{dist}\left(\xi, \Pi\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right)\right) \leq \operatorname{dist}\left(\zeta, \Pi\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right)\right) \leq 2 \operatorname{dist}\left(\xi, \Pi\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right)\right)$.
From (VII.1), (VII.2) and (VII.5) we get that for any $N \in \mathbb{N}$, for any $\xi \in B_{\rho_{i}}^{2}\left(\xi_{i}\right)$ and for any $\zeta \in B_{\rho_{k_{i}}}^{2}\left(\xi_{k_{i}}\right)$

$$
\begin{equation*}
\left|g_{N}^{-1}\right|(\xi) \leq K\left|g_{N}^{-1}\right|(\zeta) \tag{VII.6}
\end{equation*}
$$

and

$$
\begin{equation*}
\rho_{i} \frac{\left|\nabla g_{N}\right|}{\left|g_{N}\right|^{2}}(\xi)+\rho_{i}^{2} \frac{\left|\nabla^{2} g_{N}\right|}{\left|g_{N}\right|^{2}}(\xi) \leq K \frac{1}{\left|g_{N}\right|}(\zeta) \tag{VII.7}
\end{equation*}
$$

Let $\chi_{\rho_{x_{0}}}$ be a cut-off function identically equal to 1 in $B_{\rho_{x_{0}} / 2}^{2}(0)$ and equal to 0 outside $B_{\rho_{x_{0}}}^{2}$. In $B_{\rho_{x_{0}}}^{2}\left(z_{0}\right) \times \mathbb{R}^{2 p-2}$ we introduce the cycle $C^{g_{N}}$ which is given by

$$
\begin{equation*}
\forall \xi \in B_{\rho_{x_{0}}}^{2} \quad\left\langle C^{g_{N}}, \Pi, \xi\right\rangle:=g_{N}^{-1}(\xi)_{*}\langle C, \Pi, \xi\rangle \tag{VII.8}
\end{equation*}
$$

In other words if $\Psi: \Sigma \rightarrow B_{\rho_{x_{0}}}^{2}(0) \times \mathbb{R}^{2 p-2}$ is a parametrization of a piece of $C$, a parametrization of the corresponding piece in $C^{g_{N}}$ is given by $\left(\Psi_{\xi}, g_{N}^{-1} \circ\right.$ $\Pi \circ \Psi \Psi_{\lambda}$ ) where $\left(\Psi_{\xi}, \Psi_{\lambda}\right)$ are the coordinates of $\Psi$. Since $C^{g_{N}}$ is a cycle in $B_{\rho_{x_{0}}}^{2}\left(z_{0}\right) \times \mathbb{R}^{2 p-2}$, we have, denoting $\left.\Omega:=\sum_{l=1}^{p-1} d \lambda_{2 l-1} \wedge d \lambda_{2 l}\right)$,

$$
\begin{equation*}
C^{g_{N}}\left(\chi_{\rho_{x_{0}}} \circ \Pi \Omega\right)=-C^{g_{N}}\left(d \chi_{\rho_{x_{0}}} \circ \Pi \wedge \sum_{l=1}^{p-1} \lambda_{2 l-1} d \lambda_{2 l}\right) \tag{VII.9}
\end{equation*}
$$

Splitting $C^{g_{N}}\left(\chi_{\rho_{x_{0}}} \circ \Pi \Omega\right)=C^{g_{N}}\left\llcorner B_{\rho_{x_{0}} / 2}^{2} \times \mathbb{R}^{2 p-2}(\Omega)+C^{g_{N}}\left\llcorner\left(B_{\rho_{x_{0}}}^{2} \backslash B_{\rho_{x_{0}} / 2}^{2}\right) \times\right.\right.$ $\mathbb{R}^{2 p-2}\left(\chi_{\rho_{x_{0}}} \circ \Pi \Omega\right)$, we have

$$
\begin{equation*}
C^{g_{N}}\left\llcorner B_{\rho_{x_{0}} / 2}^{2} \times \mathbb{R}^{2 p-2}(\Omega)=\sum_{i \in I} C^{g_{N}}\left\llcorner B_{\rho_{x_{0}} / 2}^{2} \times \mathbb{R}^{2 p-2}\left(\varphi_{i} \circ \Pi \sum_{l=1}^{p-1} \Omega\right)\right.\right. \tag{VII.10}
\end{equation*}
$$

where we recall that the partition of unity was constructed in (V.48) adapted to the covering $B_{\rho_{i}}^{2}\left(\xi_{i}\right)$. Let $\Psi_{i}: \Sigma_{i} \rightarrow B_{2 \rho_{i}}^{2}\left(\xi_{i}\right) \times B_{2 \rho_{i}}^{2 p-2}(0)$ be a smooth parametrization of $C\left\llcorner B_{2 \rho_{i}}^{2}\left(\xi_{i}\right) \times B_{\rho_{x_{0}}}^{2 p-2}(0)\right.$ and denote by $\eta_{i}$ the map from $\Sigma_{i}$ into $\mathbb{R}^{2 p}$ given by [Ri3, Prop. A.3], such that $\Psi_{i}+\eta_{i}$ is $J_{0}$-holomorphic. Since $J$ in $B_{2 \rho_{i}}^{2 p}\left(\left(\xi_{i}, 0\right)\right)$ is closed to $J_{0}$ at a distance comparable to $\rho_{i}-$ see (VI.110) - we have

$$
\begin{equation*}
\left\|\nabla \eta_{i}\right\|_{L^{2}\left(\Sigma_{\frac{3}{2}, i}\right)} \leq \rho_{i}^{2} \tag{VII.11}
\end{equation*}
$$

where we recall that $\Sigma_{\frac{3}{2}, i}=\Sigma_{i} \cap \Psi_{i}^{-1}\left(B_{\frac{3 \rho_{i}}{2}}^{2}\left(\xi_{i}\right) \times \mathbb{R}^{2}\right)$. Using now Lemma II. 2 of [Ri3], which does not require $J$ to be $C^{1}$ in these coordinates but just the
metric $g$ to be close to the flat one, one has

$$
\begin{equation*}
\left\|\eta_{i}\right\|_{L^{2}\left(\Sigma_{\frac{4}{3}, i}\right)} \leq \rho_{i}^{3} \tag{VII.12}
\end{equation*}
$$

From the parametrization $\Psi_{i}=\left(\Psi_{i, \xi}, \Psi_{i, \lambda}\right)$,
(VII.13)

$$
\begin{aligned}
C^{g_{N}}\left\llcorner B_{\rho_{x_{0}} / 2}^{2} \times \mathbb{R}^{2 p-2}\right. & \left(\varphi_{i} \circ \Pi \Omega\right) \\
& =\int_{\Sigma_{i}} \varphi\left(\Psi_{i, \xi}\right) \sum_{l=1}^{p-1} d\left[\frac{\Psi_{i, \lambda}^{2 l-1}}{g\left(\Psi_{i, \xi}\right)}\right] \wedge d\left[\frac{\Psi_{i, \lambda}^{2 l}}{g\left(\Psi_{i, \xi}\right)}\right] .
\end{aligned}
$$

We compare this quantity with
(VII.14)

$$
\begin{aligned}
& C_{i}^{g_{N}}\left\llcorner\left(\varphi_{i} \circ \Pi \Omega\right)\right. \\
& \quad:=\int_{\Sigma_{i}} \varphi\left(\Psi_{i, \xi}\right) \sum_{l=1}^{p-1} d\left[\frac{\Psi_{i, \lambda}^{2 l-1}+\eta_{i, \lambda}^{2 l-1}}{g\left(\Psi_{i, \xi}+\eta_{i, \xi}\right)}\right] \wedge d\left[\frac{\Psi_{i, \lambda}^{2 l}+\eta_{i, \lambda}^{2 l}}{g\left(\Psi_{i, \xi}+\eta_{i, \xi}\right)}\right] .
\end{aligned}
$$

Now,

$$
\begin{align*}
& \left|\left(C_{i}^{g_{N}}-C^{g_{N}}\right)\left(\varphi_{i} \circ \Pi \sum_{l=1}^{p-1} d \lambda_{2 l-1} \wedge d \lambda_{2 l}\right)\right|  \tag{VII.15}\\
& \quad \leq K \int_{\Sigma_{i}}\left|\nabla\left[\frac{\Psi_{i, \lambda}+\eta_{i, \lambda}}{g\left(\Psi_{i, \xi}+\eta_{i, \xi}\right)}-\frac{\Psi_{i, \lambda}}{g\left(\Psi_{i, \xi}\right)}\right]\right|\left|\nabla\left[\frac{\Psi_{i, \lambda}+\eta_{i, \lambda}}{g\left(\Psi_{i, \xi}+\eta_{i, \xi}\right)}\right]\right| \\
& \quad \leq \delta K \int_{\Sigma_{i}} \varphi_{k_{i}}\left|\nabla\left[\frac{\Psi_{i, \lambda}+\eta_{i, \lambda}}{g\left(\Psi_{i, \xi}+\eta_{i, \xi}\right)}\right]\right|^{2} \\
& \quad+\frac{K}{\delta} \int_{\Sigma_{i}}\left|\nabla\left[\frac{\Psi_{i, \lambda}+\eta_{i, \lambda}}{g\left(\Psi_{i, \xi}+\eta_{i, \xi}\right)}-\frac{\Psi_{i, \lambda}}{g\left(\Psi_{i, \xi}\right)}\right]\right|^{2} .
\end{align*}
$$

Next,
(VII.16) $\int_{\Sigma_{i}}\left|\nabla\left[\frac{\Psi_{i, \lambda}+\eta_{i, \lambda}}{g\left(\Psi_{i, \xi}+\eta_{i, \xi}\right)}-\frac{\Psi_{i, \lambda}}{g\left(\Psi_{i, \xi}\right)}\right]\right|^{2}$

$$
\begin{aligned}
\leq & \int_{\Sigma_{i}}\left|\nabla\left[\Psi_{i}\left(\frac{1}{g\left(\Psi_{i, \xi}+\eta_{i, \xi}\right)}-\frac{1}{g\left(\Psi_{i, \xi}\right)}\right)\right]\right|^{2} \\
& +\int_{\Sigma_{i}}\left|\nabla \eta_{i}\right| \sup _{\xi \in B_{p_{i}}^{2}\left(\xi_{i}\right)} \frac{1}{|g(\xi)|^{2}}+\int_{\Sigma_{i}}\left|\eta_{i}\right|^{2} \sup _{\xi \in B_{p_{i}}^{2}\left(\xi_{i}\right)} \frac{|\nabla g|^{2}}{|g|^{4}}(\xi) .
\end{aligned}
$$

Let $f_{0}(\xi)$ be the flat norm of the slice of $C^{g_{N}}$ minus the average curve, $\Psi_{\lambda}=0$, by $\Pi^{-1}(\xi)$,

$$
f_{N}(\xi)=\mathcal{F}\left(\left\langle C^{g_{N}}, \Pi, \xi\right\rangle-Q \delta_{0}\right) .
$$

Using (V.8) (observe that the difference of the densities for the metrics $\omega(\cdot, J \cdot)$ and $\omega_{0}\left(\cdot, J_{0} \cdot\right)$ is as small as we want for $\rho_{x_{0}}$ chosen small enough) and using also (VII.11), (VII.12), (VII.6) and (VII.7), we have

$$
\begin{align*}
& \int_{\Sigma_{i}}\left|\nabla \eta_{i}\right|^{2} \sup _{\xi \in B_{\rho_{i}}^{2}\left(\xi_{i}\right)} \frac{1}{|g(\xi)|^{2}}+\int_{\Sigma_{i}}|\eta|^{2} \sup _{\xi \in B_{\rho_{i}}^{2}\left(\xi_{i}\right)} \frac{\left|\nabla g_{N}\right|^{2}}{\left|g_{N}\right|^{4}}(\xi)  \tag{VII.17}\\
& \leq K \int_{B_{\rho_{k_{i}}}^{2}\left(\xi_{k_{i}}\right)}\left|f_{N}\right|^{2}
\end{align*}
$$

where we have also used the fact that $\int_{\Sigma_{k_{i}}} \varphi_{k_{i}}\left|\Psi_{k_{i}, \lambda}\right|^{2} \geq K \rho_{i}^{4}$. This lower bound, a crucial point in our paper, comes from the fact that $C$ restricted to $B_{\rho_{k_{i}}}\left(x_{k_{i}}\right)$ is split; $\rho_{k_{i}}^{-2} \mathrm{M}\left(C\left\llcorner B_{\rho_{k_{i}}}\left(x_{k_{i}}\right)\right)\right.$ is less that $\pi Q-K_{0} \alpha$ (see (V.11)) where $K_{0}$ and $\alpha$ only depend on $p, Q, J$ and $\omega$. If $\Psi_{k_{i}, \lambda}$ had been too close to 0 in the $L^{2}$ norm, since the intersection number between $\left(\Psi_{k_{i}}\right)_{*}\left[\Sigma_{k_{i}}\right]$ and the $2 p-2$-planes $\Pi^{-1}(\xi)$ for $\xi \in B_{\rho_{k_{i}}}^{2}\left(\xi_{k_{i}}\right)$ is $Q, \rho_{k_{i}}^{-2} \mathrm{M}\left(C\left\llcorner B_{\rho_{k_{i}}}\left(x_{k_{i}}\right)\right)\right.$ would have been too large which contradicts the upper bound (V.11). The first term on the right-hand side of (VII.16) can be bounded as follows

$$
\begin{align*}
\int_{\Sigma_{i}} & \left|\nabla\left[\Psi_{i}\left(\frac{1}{g_{N}\left(\Psi_{i, \xi}+\eta_{i, \xi}\right)}-\frac{1}{g_{N}\left(\Psi_{i, \xi}\right)}\right)\right]\right|^{2}  \tag{VII.18}\\
\leq & \int_{\Sigma_{i}}\left|\frac{1}{g_{N}\left(\Psi_{i, \xi}+\eta_{i, \xi}\right)}-\frac{1}{g_{N}\left(\Psi_{i, \xi}\right)}\right|^{2} \\
& +\int_{\Sigma_{i}} \rho_{i}^{2}\left|\nabla\left(\frac{1}{g_{N}\left(\Psi_{i, \xi}+\eta_{i, \xi}\right)}-\frac{1}{g_{N}\left(\Psi_{i, \xi}\right)}\right)\right|^{2} \\
\leq & K \int_{\Sigma_{i}} \rho_{i}^{4}\left|\sup _{\xi \in B_{p_{i}}^{2}\left(\xi_{i}\right)} \frac{\left|\nabla g_{N}\right|}{\left|g_{N}\right|^{2}}\right|^{2}(\xi) \\
& +K \int_{\Sigma_{i}} \rho_{i}^{2}\left|\nabla \eta_{i}\right|^{2}\left|\sup _{\xi \in B_{\rho_{i}}^{2}}\left(\xi_{i}\right) \frac{\left|\nabla g_{N}\right|}{\left|g_{N}\right|^{2}}\right|^{2}(\xi) \\
& +K \int_{\Sigma_{i}} \rho_{i}^{2}\left|\eta_{i}\right|^{2} \sup _{\xi \in B_{\rho_{i}}^{2}\left(\xi_{i}\right)}^{\left|\frac{\left|\nabla g_{N}\right|^{2}}{\left|g_{N}\right|^{3}}\right|^{2}(\xi)} \\
& +K \int_{\Sigma_{i}} \rho_{i}^{2}\left|\eta_{i}\right|^{2} \sup _{\xi \in B_{\rho_{i}}^{2}\left(\xi_{i}\right)}\left|\frac{\left|\nabla^{2} g_{N}\right|}{\left|g_{N}\right|^{2}}\right|^{2}(\xi) .
\end{align*}
$$

Using (VII.6) and (VII.7) as above and combining (VII.15)-(VII.18), we have finally for every $\delta>0$

$$
\begin{align*}
&\left|\left(C_{i}^{g_{N}}-C^{g_{N}}\right)\left(\varphi_{i} \circ \Pi \sum_{l=1}^{p-1} d \lambda_{2 l-1} \wedge d \lambda_{2 l}\right)\right|  \tag{VII.19}\\
& \leq \delta K \int_{\Sigma_{i}} \varphi_{i}\left|\nabla\left[\frac{\Psi_{i, \lambda}+\eta_{i, \lambda}}{g\left(\Psi_{i, \xi}+\eta_{i, \xi}\right)}\right]\right|^{2}+\frac{K}{\delta} K \int_{B_{\rho_{k_{i}}\left(\xi_{k_{i}}\right)}\left|f_{N}\right|^{2} .} .
\end{align*}
$$

Since $\Psi_{i}+\eta_{i} / g \circ \Pi \circ \Psi_{i}+\eta_{i}$ is a holomorphic map into $\mathbb{C}^{p}$, the $\lambda$-coordinate of it, $\Psi_{i, \lambda}+\eta_{i, \lambda} / g \circ \Pi \circ \Psi_{i}+\eta_{i}$, is also a holomorphic map but into $\mathbb{C}^{p-1}$. Now,
(VII.20)

$$
\begin{aligned}
\frac{\Psi_{i, \lambda}+\eta_{i, \lambda}}{g\left(\Psi_{i, \xi}+\eta_{i, \xi}\right)} *\left(\sum_{l=1}^{p-1} d \lambda_{2 l-1} \wedge d \lambda_{2 l}\right) & =\frac{1}{2}\left(\Psi_{i, \lambda}+\eta_{i, \lambda}\right)^{*}\left(\sum_{l=1}^{p-1} d \Lambda_{l} \wedge d \bar{\Lambda}_{l}\right) \\
& =\left|\nabla \frac{\Psi_{i, \lambda}+\eta_{i, \lambda}}{g\left(\Psi_{i, \xi}+\eta_{i, \xi}\right)}\right|^{2} d \zeta \wedge d \bar{\zeta}
\end{aligned}
$$

where $\zeta$ denotes local complex coordinates on $\Sigma_{i}$, and $\Lambda_{l}$ is the complex coordinate $\Lambda_{l}=\lambda_{2 l-1}+i \lambda_{2 l}$. Therefore, combining (VII.19) and (VII.20) we have for $\delta$ chosen such that $\delta K<\frac{1}{2}$

$$
\begin{align*}
& C^{g_{N}}\left(\varphi_{i} \circ \Pi \sum_{l=1}^{p-1} d \lambda_{2 l-1} \wedge d \lambda_{2 l}\right)  \tag{VII.21}\\
& \geq \frac{1}{2} \int_{\Sigma_{i}} \varphi_{i}\left|\nabla \frac{\Psi_{i, \lambda}+\eta_{i, \lambda}}{g\left(\Psi_{i, \xi}+\eta_{i, \xi}\right)}\right|^{2}-\frac{K}{\delta} \int_{B_{\rho_{k_{i}}}^{2}\left(\xi_{k_{i}}\right)}\left|f_{N}\right|^{2} .
\end{align*}
$$

Letting $\left\{a_{i}^{l}(\xi)\right\}_{l=1 \ldots Q}$ be the holomorphic $Q$-valued graph realized by

$$
\left(\Psi_{i, \xi}+\eta_{i, \xi}, g_{N}^{-1}\left(\Psi_{i, \xi}+\eta_{i, \xi}\right) \Psi_{i, \lambda}+\eta_{i, \lambda}\right)
$$

we have

$$
\begin{equation*}
\int_{\Sigma_{i}} \varphi_{i}\left|\nabla \frac{\Psi_{i, \lambda}+\eta_{i, \lambda}}{g\left(\Psi_{i, \xi}+\eta_{i, \xi}\right)}\right|^{2}=\int_{B_{\rho_{i}}^{2}} \varphi_{i} \sum_{l=1}^{p-1}\left|\nabla a_{i}^{l}\right|^{2}(\xi) d \xi \wedge d \bar{\xi} \tag{VII.22}
\end{equation*}
$$

Clearly this quantity is larger than $\int_{B_{p_{i}}^{2}} \varphi_{i}\left|\nabla \tilde{f}_{N}\right|^{2}$ where $\tilde{f}_{N}(\xi)$ is the Flat norm of the slice by $\Pi^{-1}(\xi)$ of the difference between $C_{i}^{g_{N}}$ and the average curve. Replacing $\tilde{f}_{N}$ by $f_{N}$ itself, we see the slice by $\Pi^{-1}(\xi)$, of $C^{g_{N}}$ minus the average curve, in the integral $\int_{B_{\rho_{i}}^{2}} \varphi_{i}\left|\nabla \tilde{f}_{N}\right|^{2}$, induces error terms which can be controlled by $\int_{B_{\rho_{k_{i}}}^{2}\left(\xi_{k_{i}}\right)}\left|f_{N}\right|^{2}$ as in the computation of the error between $C^{g_{N}}$ and $C_{i}^{g_{N}}$ above. Therefore we have

$$
\begin{equation*}
C^{g_{N}}\left(\varphi_{i} \circ \Pi \Omega\right) \geq \frac{1}{2} \int_{\Sigma_{i}} \varphi_{i}\left|\nabla f_{N}\right|^{2}-\frac{K}{\delta} \int_{B_{\rho_{k_{i}}}^{2}\left(\xi_{k_{i}}\right)}\left|f_{N}\right|^{2} \tag{VII.23}
\end{equation*}
$$

Because of the relative Lipschitz estimate, $f_{N}$ extends as a $W^{1,2}$ function on all of $B_{\rho_{x_{0}}}^{2}(0)$. Standard Poincaré estimates yield

$$
\begin{equation*}
\int_{B_{\rho_{x_{0}}}^{2}}\left|\chi_{\rho_{x_{0}}} f_{N}\right|^{2} \leq K \rho_{x_{0}}^{2} \int_{B_{\rho_{x_{0}}}^{2}}\left|\nabla\left(\chi_{\rho_{x_{0}}} f_{N}\right)\right|^{2} \tag{VII.24}
\end{equation*}
$$

Taking $\rho_{x_{0}}$ small enough we can ensure that $K \rho_{x_{0}}^{2}>O(\delta)$ and combining (VII.9), (VII.23) and (VII.24) we finally get that

$$
\begin{align*}
\int_{B_{\rho_{x_{0}} / 2}^{2}}\left|f_{N}\right|^{2} \leq & K_{1} \int_{B_{\rho_{x_{0}}}^{2} \backslash B_{\rho_{x_{0}} / 2}^{2}}\left|f_{N}\right|^{2}+\left|\nabla f_{N}\right|^{2}  \tag{VII.25}\\
& +C^{g_{N}}\left(d \chi_{\rho_{x_{0}}} \circ \Pi \wedge \sum_{l=1}^{p-1} \lambda_{2 l-1} d \lambda_{2 l}\right)
\end{align*}
$$

where $K_{1}$ is a constant independent of $N$. By taking the sequence $\xi_{n}$ such that the largest $\left|\xi_{1}\right|$ satisfies $\left|\xi_{1}\right| \leq \frac{\rho_{x_{0}}}{4}$, if $C$ does not coincide with the average curve ( $f_{N}$ is not identically zero) near the origin we would have $\left(\frac{\rho_{x_{0}}}{4}\right)^{2 N} \int_{B_{\rho_{x_{0}} / 2}^{2}}\left|f_{N}\right|^{2}$ tending to infinity, whereas, it is not difficult to check that the right-hand side of (VII.25) which involves quantities supported in $B_{\rho_{x_{0}}}^{2} \backslash B_{\rho_{x_{0}} / 2}^{2}$ is bounded by $K_{\rho_{x_{0}}} N^{2}\left(\frac{\rho_{x_{0}}}{2}\right)^{-2 N}$. The multiplication of it by $\left(\frac{\rho_{x_{0}}}{4}\right)^{2 N}$ tends clearly to zero as $N$ tend to infinity. We have then obtained a contradiction and have proved that any point inside $\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}$ is surrounded in $\mathcal{C}_{*}$ by points which are all in $\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}$ or by points which are all in $\mathcal{C}_{Q-1}$. It remains to show that a point in $\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}$ is not an accumulation point of $\cup_{q \leq Q-1} \operatorname{Sing}^{q}$. This is the purpose of the next section.

## VIII. Points in $\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}$ are not accumulation points of $\cup_{q \leq Q-1} \operatorname{Sing}^{q}$

In this section we prove, assuming $\mathcal{P}_{Q-1}$, that points in $\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}$ are not accumulation points of $\cup_{q \leq Q-1} \operatorname{Sing}^{q}$ and combining this fact with the result in the previous section we will have proved $\mathcal{P}_{Q}$.

Let then $x_{0} \in \mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}$, and assume that $x_{0}$ is an accumulation point of $\mathcal{C}_{Q-1}$, which means, by the monotonicity formula, Lemma IV. 1 together with the result obtained in the previous section, that there exists a radius $\rho$ such that $\mathcal{C}_{*} \cap B_{\rho}\left(x_{0}\right) \subset \mathcal{C}_{Q}$ and that $\left(\mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}\right) \cap B_{\rho}\left(x_{0}\right)=\left\{x_{0}\right\}$. From the assumed hypothesis $\mathcal{P}_{Q-1}$, we have then that there exists a Riemann surface $\Sigma$ and a smooth $J$-holomorphic map $\Psi$ such that $C\left\llcorner B_{r_{0}}\left(x_{0}\right)=\Psi_{*}[\Sigma]\right.$. The goal is to show that $\Sigma$ has a finite topology and that it is a closed Riemann surface. The idea is to perturb $\Psi$ by finding $\eta \in L^{\infty}(\Sigma)$ such that $\Psi+\eta$ is $J_{0}$-holomorphic and $(\Psi+\eta)_{*}[\Sigma]$ is a cycle.

For any $r<r_{0}$, we denote $\Sigma_{r}$ the finite Riemann surface obtained by taking $\Sigma \cap \Psi^{-1}\left(B_{\rho}\left(x_{0}\right) \backslash B_{r}\left(x_{0}\right)\right)$ and we shall denote $\Gamma_{r}$ the part of the boundary of $\Sigma_{r}$ which is disjoint from $\partial \Sigma \subset\left(\left|\Psi-x_{0}\right|\right)^{-1}\left(r_{0}\right)$. On $\Sigma_{r}$ we consider $\eta_{r}$ the map which is given by Proposition A. 3 in [Ri3]. It satisfies in particular, when
we use the complex coordinates induced by $J_{0}$,

$$
\begin{gather*}
\bar{\partial}\left(\Psi+\eta_{r}\right)=0 \quad \text { in } \Sigma_{r}  \tag{VIII.1}\\
\forall r \leq r_{0}, \quad \int_{\Sigma_{r}}\left|\nabla \eta_{r}\right|^{2} \leq \int_{\Sigma_{r}}\left|J(\Psi)-J_{0}\right|^{2}|\nabla \Psi|^{2} \leq K r_{0}^{4},
\end{gather*}
$$

where for the induced metric by $\Psi$ on $\Sigma$ ( $\Psi$ is an isometry), we have $\int_{\Sigma}|\nabla \Psi|^{2}=$ $\mathrm{M}\left(C\left\llcorner B_{r_{0}}\left(x_{0}\right)\right) \leq K r_{0}^{2}\right.$. Using local $\xi_{1}, \xi_{2}$ coordinates in $\Sigma_{r}$, we have for all $k=1 \ldots 2 p$,

$$
\frac{\partial \Psi_{i}^{k}}{\partial \xi_{1}}=-\sum_{l=1}^{2 p} J_{l}^{k}\left(\Psi_{i}\right) \frac{\partial \Psi^{l}}{\partial \xi_{2}} \quad \text { and } \quad \frac{\partial \Psi_{i}^{k}}{\partial \xi_{2}}=\sum_{l=1}^{2 p} J_{l}^{k}\left(\Psi_{i}\right) \frac{\partial \Psi^{l}}{\partial \xi_{1}}
$$

Taking respectively the $\xi_{1}$ derivative and the $\xi_{2}$ derivative of these two equations we obtain

$$
\begin{equation*}
\forall k=1 \ldots 2 p, \quad \Delta_{\Sigma_{r}} \Psi_{i}^{k}=*\left(\sum_{l=1}^{2 p} d\left(J_{l}^{k}\left(\Psi_{i}\right)\right) \wedge d \Psi_{i}^{l}\right) \tag{VIII.2}
\end{equation*}
$$

From (VIII.1) we deduce that $\Delta_{\Sigma_{r}}\left(\Psi+\eta_{r}\right)=0$; therefore this yields

$$
\begin{equation*}
\forall k=1 \ldots 2 p, \quad \Delta_{\Sigma_{r}} \eta_{r}^{k}=-*\left(\sum_{l=1}^{2 p} d\left(J_{l}^{k}\left(\Psi_{i}\right)\right) \wedge d \Psi_{i}^{l}\right) \tag{VIII.3}
\end{equation*}
$$

Let $\delta_{r}^{k}$ be given by

$$
\left\{\begin{align*}
\Delta_{\Sigma_{r}} \delta_{r}^{k} & =*\left(\sum_{l=1}^{2 p} d\left(J_{l}^{k}\left(\Psi_{i}\right)\right) \wedge d \Psi_{i}^{l}\right) & & \text { in } \Sigma_{r}  \tag{VIII.4}\\
\delta_{r}^{k} & =0 & & \text { on } \partial \Sigma_{r}
\end{align*}\right.
$$

From [Ge] and $[\mathrm{To}]$ there exists a universal constant $K$ such that

$$
\begin{equation*}
\left\|\delta_{r}\right\|_{L^{\infty}\left(\Sigma_{r}\right)}+\|\nabla \delta\|_{L^{2}\left(\Sigma_{r}\right)} \leq K\|J\|_{C^{1}} \int_{\Sigma_{r}}|\nabla \Psi|^{2} \leq K r_{0}^{2} . \tag{VIII.5}
\end{equation*}
$$

Because of the above estimates, taking some sequence $r_{n} \rightarrow 0$, one can always extract a subsequence $r_{n^{\prime}}$ such that $\eta_{r_{n^{\prime}}}$ and $\delta_{r_{n^{\prime}}}$ converge to limits $\eta_{0}$ and $\delta_{0}$ that satisfy in particular

$$
\begin{align*}
& \bar{\partial}\left(\Psi+\eta_{0}\right)=0  \tag{VIII.6}\\
& \text { in } \Sigma, \\
& \Delta_{\Sigma}\left(\eta_{0}+\delta_{0}\right)=0 \\
&\left\|\nabla \delta_{0}\right\|_{L^{2}(\Sigma)}+\left\|\delta_{0}\right\|_{L^{\infty}(\Sigma)} \leq K r_{0}^{2}
\end{align*}
$$

For any $k=1, \ldots, 2 p$ we consider the harmonic function $u^{k}:=\eta^{k}+\delta^{k}$. Using the coarea formula we have, for any $r<r_{0}$,

$$
\begin{equation*}
\int_{0}^{r} d s \int_{\Gamma_{s}}\left|\nabla u^{k}\right|=\int_{\Sigma \backslash \Sigma_{r}}\left|\nabla u^{k}\right||\nabla| \Psi| | \leq r\left(\int_{\Sigma \backslash \Sigma_{r}}\left|\nabla u^{k}\right|^{2}\right)^{\frac{1}{2}} \tag{VIII.7}
\end{equation*}
$$

Therefore, by use of a mean formula, for any $\varepsilon>0$, there exists $s>0$ such that

$$
\begin{equation*}
\int_{\Gamma_{s}}\left|\nabla u^{k}\right| \leq \varepsilon . \tag{VIII.8}
\end{equation*}
$$

Now,

$$
\begin{equation*}
0=\int_{\Sigma_{s}} \Delta_{\Sigma} u^{k}=\int_{\Gamma_{s}} \frac{\partial u^{k}}{\partial \nu}+\int_{\partial \Sigma} \frac{\partial u^{k}}{\partial \nu} . \tag{VIII.9}
\end{equation*}
$$

By choosing $\varepsilon$ smaller and smaller and taking the corresponding $s$ given by (VIII.8), one gets

$$
\begin{equation*}
\int_{\partial \Sigma} \frac{\partial u^{k}}{\partial \nu} \tag{VIII.10}
\end{equation*}
$$

Let $m<M$ be two values such that $\sup _{\partial \Sigma} u^{k}<m$ and consider the truncation $T_{m}^{M} u^{k}$ equal to $m$ if $u^{k} \leq m$ equal to $M$ if $u^{k} \geq M$ and equal to $u^{k}$ otherwise.
We have
(VIII.11)

$$
0=\int_{\Sigma_{s}} T_{m}^{M} u^{k} \Delta_{\Sigma} u^{k}=-\int_{\Sigma_{s}}\left|\nabla T_{m}^{M} u^{k}\right|^{2}+\int_{\Gamma_{s}} T_{m}^{M} u^{k} \frac{\partial u^{k}}{\partial \nu}+m \int_{\partial \Sigma} \frac{\partial u^{k}}{\partial \nu} .
$$

Therefore

$$
\begin{equation*}
\int_{\Sigma_{s}}\left|\nabla T_{m}^{M} u^{k}\right|^{2} \leq M \int_{\Gamma_{s}}\left|\nabla u^{k}\right| \tag{VIII.12}
\end{equation*}
$$

and by choosing again $s$ tending to zero according to (VIII.8), one gets that $T_{m}^{M} u^{k}$ is identically equal to $m$ and we deduce that $u^{k} \leq m$ in $\Sigma$. Similarly one gets that $u^{k}$ is bounded from below and then we have proved that $\|u\|_{L^{\infty}(\Sigma)}<$ $+\infty$. Combining this fact with (VIII.6) we have that

$$
\begin{equation*}
\left\|\eta_{0}\right\|_{L^{\infty}(\Sigma)}<+\infty \tag{VIII.13}
\end{equation*}
$$

Being more careful above by taking eventually $\Sigma_{r}$ instead of $\Sigma$ for some $r \in$ $\left[r_{0} / 2, r_{0}\right]$, and using [Ri3] we could have shown that $\left\|\eta_{0}\right\|_{L^{\infty}(\Sigma)}<K r_{0}^{2}$. We claim now that $\partial\left(\Psi+\eta_{0}\right)_{*}[\Sigma]=\left(\Psi+\eta_{0}\right)_{*}[\partial \Sigma]$; that is, for any smooth 1-form $\phi$ equal to zero in a neighborhood of $\left(\Psi+\eta_{0}\right)(\partial \Sigma)$, one has

$$
\begin{equation*}
\int_{\Sigma}\left(\Psi+\eta_{0}\right)^{*} d \phi=0 \tag{VIII.14}
\end{equation*}
$$

Now,

$$
\begin{equation*}
\left|\int_{\Sigma_{s}}\left(\Psi+\eta_{0}\right)^{*} d \phi\right|=\left|\int_{\Gamma_{s}}\left(\Psi+\eta_{0}\right)^{*} \phi\right| \leq K_{\phi} \int_{\Gamma_{s}}\left|\nabla \Psi+\eta_{0}\right| . \tag{VIII.15}
\end{equation*}
$$

Arguing as in the proof of (VIII.8), for any $\varepsilon$, we can find $s$ such that $\int_{\Gamma_{s}}\left|\nabla \Psi+\eta_{0}\right| \leq \varepsilon$ and we then deduce (VIII.14). Thus in $B_{r_{0}}^{2 p}\left(x_{0}\right) \backslash \Psi+\eta_{0}(\partial \Sigma)$, $\left(\Psi+\eta_{0}\right)_{*}[\Sigma]$ is an integer-multiplicity, rectifiable, holomorphic cycle. Using the
results of Harvey-Shiffman and King ([HS] and [Ki]) we have that there exists a compact Riemann surface with boundary $\Sigma^{\prime}$ and a holomorphic map $\Psi^{\prime}$ such that $\left(\Psi+\eta_{0}\right)_{*}[\Sigma]=\Psi_{*}^{\prime}\left[\Sigma^{\prime}\right]$. Therefore $(\phi+\eta)(\Sigma)$ is a holomorphic curve with boundary - in $\mathbb{C}^{2 p}$. We claim that $\Psi+\eta_{0}$ is a holomorphic simple covering of $\Sigma^{\prime}$. Indeed letting $\omega_{\Sigma}$ be the pull-back by $\Psi$ of the symplectic form $\omega$ in $\mathbb{R}^{2 p}$ we have $\int_{\Sigma} \omega_{\Sigma}=\int_{\Sigma} \Psi^{*} \omega=\int_{\Sigma}|\nabla \Psi|^{2} \geq \pi Q r_{0}^{2}$, because of the monotonicity formula ( $x_{0} \in \mathcal{C}_{Q} \backslash \mathcal{C}_{Q-1}$ ). Let $\omega_{\Sigma^{\prime}}$ be the restriction of $\omega_{0}$ to $\Sigma^{\prime}$. We have $\int_{\Sigma}\left(\Psi+\eta_{0}\right)^{*} \omega_{\Sigma^{\prime}}=\int_{\Sigma}\left(\Psi+\eta_{0}\right)^{*} \omega_{0}=\int_{\Sigma}\left|\nabla\left(\Psi+\eta_{0}\right)\right|^{2}$. Because of (VIII.1), the holomorphic covering $\Psi+\eta_{0}$ from $\Sigma$ onto $\Sigma^{\prime}$ satisfies

$$
\begin{equation*}
\left|\int_{\Sigma} \omega_{\Sigma}-\int_{\Sigma}\left(\Psi+\eta_{0}\right) \omega_{\Sigma^{\prime}}\right|=o_{r_{0}}\left(\int_{\Sigma} \omega_{\Sigma}\right) . \tag{VIII.16}
\end{equation*}
$$

Therefore, for $r_{0}$ small enough this covering has to be a simple one and $\Sigma$ is a compact Riemann surface. Now $\Psi$ is a $J$-holomorphic map from a compact Riemann surface $\Sigma$ into $\left(B^{2 p}, J\right)$; it is then smooth and $C\left\llcorner B_{r_{0}}^{2 p}\right.$ is a $J$-holomorphic curve.

## A. Appendix

Lemma A.1. Let $U$ be an open subset of $\mathbb{R}^{2}$, let $0<\lambda<1$ and let $\left(B_{r_{i}}^{2}\left(z_{i}\right)\right)_{i \in I}$ a covering of $U$ which is locally finite. There exists $n \in \mathbb{N}$ such that

$$
\begin{equation*}
\forall z \in U, \quad \operatorname{Card}\left\{i \in I: z \in B_{r_{z_{i}}}^{2}\left(z_{i}\right)\right\} \leq N \tag{A.1}
\end{equation*}
$$

Moreover one assumes that

$$
\begin{equation*}
\forall i, j \in I, \quad B_{r_{i}}\left(z_{i}\right) \cap B_{r_{j}}\left(z_{j}\right) \neq \emptyset \quad \Longrightarrow \quad r_{i} \geq \lambda r_{j} \tag{A.2}
\end{equation*}
$$

Then there exist $\delta$ and $P \in \mathbb{N}$ depending on $\lambda$ only such that

$$
\begin{equation*}
\forall i \in I, \quad \operatorname{Card}\left\{j \in I: B_{r_{j}}\left(z_{j}\right) \cap B_{(1+\delta) r_{i}}\left(z_{i}\right) \neq \emptyset\right\} \leq P . \tag{A.3}
\end{equation*}
$$

Proof of Lemma A.1. We argue by contradiction. Assume there exist $\delta_{n} \rightarrow 0$, a sequences of coverings of $U,\left(B_{r_{n, i}}^{2}\left(z_{n, i}\right)\right)$ for $i \in I$ satisfying (A.1) and (A.2) and a sequence of indices $i_{n}$ such that

$$
\begin{equation*}
\operatorname{Card}\left\{j \in I: B_{r_{j, n}}\left(z_{j, n}\right) \cap B_{\left(1+\delta_{n}\right) r_{i_{n}, n}}\left(z_{i_{n}, n}\right) \neq \emptyset\right\} \rightarrow+\infty \quad \text { as } n \rightarrow+\infty \tag{A.4}
\end{equation*}
$$

After a possible rescaling of the whole covering and a translation we can assume that $r_{i_{n}, n}=1$ and $z_{i_{n}, n}=0$. Also, after extraction of a subsequence, we can ensure that there exists $A \in \partial B_{1}(0)$ such that for any $r>0$

$$
\begin{equation*}
\operatorname{Card}\left\{j \in I: B_{r_{j, n}}\left(z_{j, n}\right) \cap B_{r}(A) \neq \emptyset\right\} \rightarrow+\infty \quad \text { as } n \rightarrow+\infty \tag{A.5}
\end{equation*}
$$

For a given $r$ and $n$ we take the longest sequence of distinct balls of our covering $B_{r_{j_{p}, n}}\left(z_{j_{p}, n}\right)$ for $p=0 \ldots P_{n}$ satisfying

$$
B_{r_{j_{0}, n}}\left(z_{j_{0}, n}\right)=B_{1}(0),
$$

ii)

$$
\forall p \leq P_{n}-1, \quad B_{r_{j_{p}, n}}\left(z_{j_{p}, n}\right) \cap B_{r_{j_{p+1}, n}}\left(z_{j_{p+1}, n}\right) \neq \emptyset
$$

iii) $\quad \forall p \leq P_{n}, \quad B_{r_{j_{p}, n}}\left(z_{j_{p}, n}\right) \cap B_{r}(A) \neq \emptyset$.

It is clear that for a given $r, P_{n} \rightarrow+\infty$; indeed if it were not the case, i.e. $P_{n} \leq P_{*}<+\infty$, this would imply that the minimal radius for the balls of the covering intersecting $B_{r}(A)$ is $\lambda^{P_{*}}$. Combining this fact with (A.5) would contradict (A.1). Therefore we can find $r_{m} \rightarrow+\infty$ and $n_{m} \rightarrow+\infty$ as $m \rightarrow+\infty$ and sequences $\left(B_{r_{j_{p}, n_{m}}}\right)$ for $1 \leq p \leq Q_{m}$ and for $m=0 \cdots+\infty$ such that

$$
\begin{align*}
\text { i) } & B_{r_{j_{0}, n_{m}}}\left(z_{j_{0}, n_{m}}\right)=B_{1}(0) \\
\text { ii) } \quad \forall p \leq Q_{m-1}, & B_{r_{j_{p}, n_{m}}}\left(z_{j_{p}, n_{m}}\right) \cap B_{r_{j_{p+1}, n_{m}}}\left(z_{j_{p+1}, n_{m}}\right) \neq \emptyset  \tag{A.6}\\
\text { iii }) \quad \forall p \leq Q_{m}, & B_{r_{j_{p}, n_{m}}}\left(z_{j_{p}, n_{m}}\right) \cap B_{r_{m}}(A) \neq \emptyset  \tag{А.7}\\
\text { iv) } & Q_{m} \rightarrow+\infty
\end{align*}
$$

Since $\lambda \leq r_{j_{1}, n_{m}} \leq \lambda^{-1}$ and since the distance $\left|z_{j_{1}, n_{m}}\right|$ is bounded, we can extract from $n_{m}$ a subsequence that we still denote $n_{m}$ such that $B_{r_{j_{1}, n_{m}}}\left(z_{j_{1}, n_{m}}\right)$ converges to a limiting ball $B_{r_{1, \infty}}\left(z_{1, \infty}\right)$ with $\lambda \leq r_{1, \infty} \leq \lambda^{-1}, z_{1, \infty} \leq 2$ and $A \in$ $\overline{B_{r_{1, \infty}}\left(z_{1, \infty}\right)}$. This procedure can be iterated and using a diagonal argument we can assume that

$$
\forall p \in \mathbb{N}, \quad r_{j_{p}, n_{m}} \rightarrow r_{p, \infty}, \quad z_{j_{p}, n_{m}} \rightarrow z_{p, \infty}
$$

such that

$$
\forall p \in \mathbb{N}, \quad \lambda^{p} \leq r_{p, \infty} \leq \lambda^{-p}, \quad\left|z_{p, \infty}\right| \leq 2 \quad \text { and } A \in \overline{B_{r_{p, \infty}}\left(z_{p, \infty}\right)}
$$

Moreover because of (A.1) we have that

$$
\begin{equation*}
\forall z \in \mathbb{R}^{2}, \quad \text { Card }\left\{p \in \mathbb{N}: z \in B_{r_{p, \infty}}^{2}\left(z_{p, \infty}\right)\right\} \leq N \tag{A.8}
\end{equation*}
$$

Because of this latter fact, since $A \in \overline{B_{r_{p, \infty}}\left(z_{p, \infty}\right)}$ for all $p$, it is clear that $r_{p, \infty} \rightarrow+\infty$ as $p \rightarrow+\infty$. Because of (A.8) again, the number of open balls $B_{r_{p, \infty}}\left(z_{p, \infty}\right)$ containing $A$ is bounded by $N$ and we can therefore forget them while considering the sequence and assume that

$$
\forall p, \quad A \in \partial B_{r_{p, \infty}}^{2}\left(z_{p, \infty}\right)
$$

Let $\vec{t}_{p}(A) \in S^{1}$ be the unit exterior normal to $\partial B_{r_{p, \infty}}^{2}\left(z_{p, \infty}\right)$ at $A$. Let $\vec{t}_{\infty}$ be an accumulation unit vector of the sequence $\overrightarrow{t_{p}}(A)$. Given a direction $\vec{t}$ and an open disk containing $A$ in its boundary and whose exterior normal at $A$ is given by $\vec{t}$ any other open disk containing $A$ in its boundary and whose exterior unit at $A$ is not $-\vec{t}$ as a nonempty intersection with that disk. When $B_{r_{p_{0}, \infty}}^{2}\left(z_{p_{0}, \infty}\right)$, such that $\vec{t}_{p_{0}}(A) \neq-\vec{t}_{\infty}$, there exist infinitely many disks $B_{r_{p}, \infty}^{2}\left(z_{p, \infty}\right)$ having nonempty intersections with $B_{r_{p_{0}, \infty}}^{2}\left(z_{p_{0}, \infty}\right)$. But then, because of (A.2), that passes to the limit and all these infinitely many disks have radii which are bounded from below by positive numbers and this contradicts the fact that $r_{p, \infty} \rightarrow+\infty$ as implied by (A.8). Thus Lemma A. 1 is proved.

Acknowledgements. The authors are very grateful to Robert Bryant for having pointed out to them almost complex structures admitting no locally compatible positive symplectic forms.

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(Received October 21, 2003)
(Revised September 29, 2007)

