# Embedded self-similar shrinkers of genus 0 

By Simon Brendle<br>Dedicated to Professor Leon Simon on the occasion of his seventieth birthday


#### Abstract

We confirm a well-known conjecture that the round sphere is the only compact, embedded self-similar shrinking solution of mean curvature flow in $\mathbb{R}^{3}$ with genus 0 . More generally, we show that the only properly embedded self-similar shrinkers in $\mathbb{R}^{3}$ with vanishing intersection form are the sphere, the cylinder, and the plane. This answers two questions posed by T. Ilmanen.


## 1. Introduction

This paper is concerned with self-similar shrinking solutions to the mean curvature flow in $\mathbb{R}^{3}$. A surface $M \subset \mathbb{R}^{3}$ is called a self-similar shrinker if it satisfies the equation $H=\frac{1}{2}\langle x, \nu\rangle$, where $\nu$ and $H$ denote the unit normal vector and the mean curvature, respectively. This condition guarantees that the surface $M$ moves by homotheties when evolved by the mean curvature flow.

The classification of self-similar solutions to geometric flows is a central problem with important implications for the analysis of singularities. Indeed, Huisken's montonicity formula [15] implies that any tangent flow to a compact solution of mean curvature flow is a self-similar shrinker (see also [10] and [12]). The simplest example of a compact self-similar shrinker in $\mathbb{R}^{3}$ is the round sphere of radius 2 centered at the origin. G. Drugan [11] has recently constructed an example of a self-similar shrinker of genus 0 which is immersed but fails to be embedded. Angenent [1] has constructed an example of an embedded self-similar shrinker of genus 1. Moreover, N. Kapouleas, S. Kleene, and N.M. Møller [18] have constructed new examples of noncompact self-similar shrinkers using gluing techniques. These examples are embedded and have high genus.

A well-known conjecture asserts that the round sphere of radius 2 should be the only embedded self-similar shrinker of genus 0 . Our main result confirms this conjecture:

[^0]Theorem 1. Let $M$ be a compact, embedded self-similar shrinker in $\mathbb{R}^{3}$ of genus 0 . Then $M$ is a round sphere.

In view of the examples constructed by Drugan and Angenent, the assumptions that $M$ is embedded and has genus 0 are both necessary. In that respect, Theorem 1 shares some common features with Lawson's Conjecture on embedded minimal tori in $S^{3}$ (cf. [19]). This conjecture was recently confirmed in [3]; see [6] for a survey.

In the noncompact case, our arguments imply the following:
Theorem 2. Suppose that $M$ is a properly embedded self-similar shrinker in $\mathbb{R}^{3}$ with the property that any two loops in $M$ have vanishing intersection number mod 2. Then $M$ is a round sphere or a cylinder or a plane.

Theorem 2 confirms two conjectures of T. Ilmanen, the Wiggly Plane Conjecture and the Planar Domain Conjecture (cf. [17]). We note that the topological assumption in Theorem 2 is equivalent to the condition that $M$ is homeomorphic to an open subset of $S^{2}$; this follows, e.g., from the simple exhaustion theorem in Section 4 in [14].

We next discuss some related results. In 1990, G. Huisken [15] proved that the round sphere is the only compact self-similar shrinking solution with positive mean curvature. Using a similar argument, Huisken was able to show that a noncompact self-similar shrinker which has bounded curvature and positive mean curvature must be a cylinder (cf. [16]). Moreover, K. Ecker and G. Huisken proved that a self-similar shrinker which can be written as an entire graph must be a plane (cf. [13, p. 471]). In a remarkable recent work, T. Colding and W. Minicozzi [9] proved that a self-similar shrinker which is a stable critical point of a certain entropy functional must be a sphere or a cylinder or a plane. Furthermore, T. Colding, T. Ilmanen, W. Minicozzi, and B. White recently showed that the round sphere has smallest entropy among all compact self-similar shrinkers (see [8]). We note that L. Wang [23] has obtained a classification of self-similar shrinkers which are asymptotic to cones. X. Wang [24] has proved a uniqueness result for convex translating solutions to the mean curvature flow which can be expressed as graphs over $\mathbb{R}^{3}$. Furthermore, we recently obtained a classification of steady gradient Ricci solitons in dimension 3 and 4 under a noncollapsing assumption (cf. [4], [5]).

We now sketch the main ideas involved in the proof of Theorem 1. Suppose that $M$ is a compact, embedded self-similar shrinker in $\mathbb{R}^{3}$ of genus 0 . In the first step, we show that, for any plane $P \subset \mathbb{R}^{3}$ which passes through the origin, the intersection $M \cap P$ consists of a single Jordan curve which is piecewise $C^{1}$. This argument is inspired in part by the two-piece property for embedded minimal surfaces in $S^{3}$ (cf. Ros [21]). We next prove that $M$ is star-shaped.

Indeed, if $\langle\bar{x}, \nu(\bar{x})\rangle=0$ for some point $\bar{x} \in M$, then we consider the tangent plane $P$ to $M$ at $\bar{x}$. Clearly, $P$ passes through the origin, so the intersection $M \cap P$ consists of a single Jordan curve. On the other hand, $M \cap P$ contains at least two arcs which intersect transversally at $\bar{x}$. This gives a contradiction. Having established that $M$ is star-shaped, it follows that the mean curvature of $M$ does not change sign. Huisken's theorem then implies that $M$ is a round sphere, thereby completing the proof of Theorem 1.

The proof of Theorem 2 uses similar techniques; this is discussed in Section 4.

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## 2. The key estimate

We begin by collecting some basic identities for self-similar shrinkers in $\mathbb{R}^{3}$.
Proposition 3. Let $\Sigma$ be a self-similar shrinker in $\mathbb{R}^{3}$. Moreover, suppose that $\Xi$ is a smooth vector field on $\mathbb{R}^{3}$, and let $\xi$ denote the projection of $\Xi$ to the tangent plane of $\Sigma$. Then

$$
\operatorname{div}_{\Sigma} \xi-\frac{1}{2}\langle x, \xi\rangle=\sum_{i=1}^{2}\left\langle\bar{D}_{e_{i}} \Xi, e_{i}\right\rangle-\frac{1}{2}\langle x, \Xi\rangle
$$

Here, $\bar{D}$ denotes the Levi-Civita connection on the ambient space $\mathbb{R}^{3}$, and $\left\{e_{1}, e_{2}\right\}$ is a local orthonormal frame on $\Sigma$.

Proof. Since $\Sigma$ is a self-similar shrinker, we have $H=\frac{1}{2}\langle x, \nu\rangle$. This implies

$$
\begin{aligned}
\operatorname{div}_{\Sigma} \xi-\frac{1}{2}\langle x, \xi\rangle & =\sum_{i=1}^{2}\left\langle\bar{D}_{e_{i}} \Xi, e_{i}\right\rangle-H\langle\Xi, \nu\rangle-\frac{1}{2}\langle x, \xi\rangle \\
& =\sum_{i=1}^{2}\left\langle\bar{D}_{e_{i}} \Xi, e_{i}\right\rangle-\frac{1}{2}\langle x, \nu\rangle\langle\Xi, \nu\rangle-\frac{1}{2} \sum_{i=1}^{2}\left\langle x, e_{i}\right\rangle\left\langle\Xi, e_{i}\right\rangle \\
& =\sum_{i=1}^{2}\left\langle\bar{D}_{e_{i}} \Xi, e_{i}\right\rangle-\frac{1}{2}\langle x, \Xi\rangle
\end{aligned}
$$

This proves the assertion.
Corollary 4. Let $\Sigma$ be a self-similar shrinker in $\mathbb{R}^{3}$. Suppose that $F$ : $\mathbb{R}^{3} \rightarrow \mathbb{R}$ is a smooth function, and let $f: \Sigma \rightarrow \mathbb{R}$ denote the restriction of $F$ to $\Sigma$. Then

$$
\Delta_{\Sigma} f-\frac{1}{2}\left\langle x, \nabla^{\Sigma} f\right\rangle=\sum_{i=1}^{2}\left(\bar{D}^{2} F\right)\left(e_{i}, e_{i}\right)-\frac{1}{2}\langle x, \bar{\nabla} F\rangle .
$$

Here, $\bar{\nabla} F$ and $\bar{D}^{2} F$ denote gradient and Hessian of $F$ with respect to the Euclidean metric, and $\left\{e_{1}, e_{2}\right\}$ is a local orthonormal frame on $\Sigma$.

Proof. Apply Proposition 3 to the gradient vector field $\Xi=\bar{\nabla} F$.
It is well known that self-similar shrinkers can be characterized as critical points of a functional. More precisely, $\Sigma$ is a self-similar shrinker if and only if $\Sigma$ is a critical point of the functional

$$
\mathscr{F}(\Sigma)=\int_{\Sigma} e^{-\frac{|x|^{2}}{4}} .
$$

Following Colding and Minicozzi, we define a differential operator $L$ on $\Sigma$ by

$$
L f=\Delta_{\Sigma} f+|A|^{2} f+\frac{1}{2} f-\frac{1}{2}\left\langle x, \nabla^{\Sigma} f\right\rangle
$$

(cf. [9, eq. (4.13)]). The second variation of $\mathscr{F}$ is given by

$$
-\int_{\Sigma} e^{-\frac{|x|^{2}}{4}} f L f=\int_{\Sigma} e^{-\frac{|x|^{2}}{4}}\left(\left|\nabla^{\Sigma} f\right|^{2}-|A|^{2} f^{2}-\frac{1}{2} f^{2}\right)
$$

where $f: \bar{\Sigma} \rightarrow \mathbb{R}$ is a test function which has compact support and vanishes along the boundary of $\Sigma$ (see [9, Th. 4.14]).

We next consider a self-similar shrinker whose boundary is contained in a plane. In this case, we can use the height function as a test function in the stability inequality. This leads to the following result:

Proposition 5. Let $\Sigma$ be a smooth surface in $\mathbb{R}^{3}$ with boundary $\partial \Sigma=\Gamma$, and let $k \geq 4$. Suppose that $H=\frac{1}{2}\langle x, \nu\rangle$ on $\Sigma \cap\{|x| \leq k\}$. Moreover, suppose that the stability inequality

$$
0 \leq-\int_{\Sigma} e^{-\frac{|x|^{2}}{4}} f L f
$$

holds for every smooth function $f: \bar{\Sigma} \rightarrow \mathbb{R}$ which vanishes on the set $\Gamma \cup(\bar{\Sigma} \cap$ $\{|x| \geq k\}$ ). Finally, we assume that $\Gamma \cap\{|x| \leq k\} \subset\left\{x \in \mathbb{R}^{3}:\langle a, x\rangle=0\right\}$ for some unit vector $a \in \mathbb{R}^{3}$. Then

$$
\int_{\Sigma \cap\{|x| \leq \sqrt{k}\}}|A|^{2} e^{-\frac{|x|^{2}}{4}}\langle a, x\rangle^{2} \leq \frac{C}{\log k} \int_{\Sigma \cap\{\sqrt{k} \leq|x| \leq k\}} e^{-\frac{|x|^{2}}{4}},
$$

where $C$ is a positive constant independent of $k$.
Proof. Let us fix a smooth cutoff function $\eta:(-\infty, \infty) \rightarrow[0,1]$ satisfying $\eta=1$ on $\left(-\infty, \frac{1}{2}\right], \eta=0$ on $[1, \infty)$, and $\eta^{\prime} \leq 0$ on $(-\infty, \infty)$. We define a smooth function $F: \mathbb{R}^{3} \rightarrow \mathbb{R}$ by

$$
F(x)=\langle a, x\rangle \eta\left(\frac{\log |x|}{\log k}\right) .
$$

Note that

$$
\langle x, \bar{\nabla} F\rangle=F+\frac{1}{\log k}\langle a, x\rangle \eta^{\prime}\left(\frac{\log |x|}{\log k}\right) .
$$

Moreover, we have

$$
\left|\bar{D}^{2} F\right| \leq \frac{C}{\log k} \frac{1}{|x|} 1_{\{\sqrt{k} \leq|x| \leq k\}},
$$

where $C$ is a positive constant independent of $k$. This implies

$$
|F|\left|\bar{D}^{2} F\right| \leq \frac{C}{\log k} 1_{\{\sqrt{k} \leq|x| \leq k\}} .
$$

Let $f: \bar{\Sigma} \rightarrow \mathbb{R}$ denote the restriction of $F$ to $\bar{\Sigma}$. Using Corollary 4, we obtain

$$
\begin{aligned}
& -f\left(\Delta_{\Sigma} f-\frac{1}{2}\left\langle x, \nabla^{\Sigma} f\right\rangle\right) \\
& =-F\left(\sum_{i=1}^{2}\left(\bar{D}^{2} F\right)\left(e_{i}, e_{i}\right)-\frac{1}{2}\langle x, \bar{\nabla} F\rangle\right) \\
& =-F \sum_{i=1}^{2}\left(\bar{D}^{2} F\right)\left(e_{i}, e_{i}\right)+\frac{1}{2} F^{2}+\frac{1}{2 \log k}\langle a, x\rangle^{2} \eta\left(\frac{\log |x|}{\log k}\right) \eta^{\prime}\left(\frac{\log |x|}{\log k}\right) \\
& \leq \frac{C}{\log k} 1_{\{\sqrt{k} \leq|x| \leq k\}}+\frac{1}{2} f^{2} .
\end{aligned}
$$

In the last step, we have used the inequality $\eta^{\prime} \leq 0$. Consequently,

$$
-f L f \leq \frac{C}{\log k} 1_{\{\sqrt{k} \leq|x| \leq k\}}-|A|^{2} f^{2}
$$

Note that $f$ vanishes on the set $\Gamma \cup(\bar{\Sigma} \cap\{|x| \geq k\})$. Using $f$ as a test function in the stability inequality gives

$$
\begin{aligned}
0 & \leq-\int_{\Sigma} e^{-\frac{|x|^{2}}{4}} f L f \\
& \leq \frac{C}{\log k} \int_{\Sigma \cap\{\sqrt{k} \leq|x| \leq k\}} e^{-\frac{|x|^{2}}{4}}-\int_{\Sigma}|A|^{2} e^{-\frac{|x|^{2}}{4}} f^{2} \\
& \leq \frac{C}{\log k} \int_{\Sigma \cap\{\sqrt{k} \leq|x| \leq k\}} e^{-\frac{|x|^{2}}{4}}-\int_{\Sigma \cap\{|x| \leq \sqrt{k}\}}|A|^{2} e^{-\frac{|x|^{2}}{4}}\langle a, x\rangle^{2} .
\end{aligned}
$$

This proves the assertion.

## 3. Proof of Theorem 1

In this section, we describe the proof of Theorem 1. Let $M$ be a compact, embedded self-similar shrinker in $\mathbb{R}^{3}$ of genus 0 . Moreover, suppose that $M$ is not a round sphere. By Theorem 4.1 in [15], the mean curvature $H$ must change sign. In particular, we can find a point $\bar{x} \in M$ such that $H(\bar{x})=0$. Using the shrinker equation, we conclude that $\langle\bar{x}, \nu(\bar{x})\rangle=0$. For abbreviation, let $a:=\nu(\bar{x})$ and $Z:=\{x \in M:\langle a, x\rangle=0\}$. Clearly, $\bar{x} \in Z$. The structure of the set $Z$ is described in the following lemma.

Lemma 6. The set $Z=\{x \in M:\langle a, x\rangle=0\}$ is a union of finitely many $C^{1}$-arcs which meet at isolated points. More precisely, for each point $x_{0} \in Z$, there exists an open neighborhood $U \subset M$ of $x_{0}$ such that $Z \cap U$ is a union of $m C^{1}$-arcs which intersect transversally at $x_{0}$. Here, $m$ can be characterized as the order of vanishing of the function $x \mapsto\langle a, x\rangle$ at $x_{0}$.

Proof. The set $Z$ can be viewed as the nodal set of a solution of an elliptic equation. Indeed, it follows from Corollary 4 that the function $f(x):=\langle a, x\rangle$ satisfies the equation

$$
\Delta_{M} f-\frac{1}{2}\left\langle x, \nabla^{M} f\right\rangle=-\frac{1}{2} f
$$

(see also [9, Lemma 3.20]). This identity can be rewritten as

$$
\Delta_{M}\left(e^{-\frac{|x|^{2}}{8}} f\right)=h e^{-\frac{|x|^{2}}{8}} f,
$$

where $h:=e^{\frac{|x|^{2}}{8}} \Delta_{M}\left(e^{-\frac{|x|^{2}}{8}}\right)-\frac{1}{2}$. If we apply Lemma 2.4 and Theorem 2.5 in [7] to the function $e^{-\frac{|x|^{2}}{8}} f$, the assertion follows.

We now continue with the proof of Theorem 1. In view of our choice of $\bar{x}$ and $a$, the function $x \mapsto\langle a, x\rangle$ vanishes to order $m \geq 2$ at the point $\bar{x}$. Consequently, there exists an open neighborhood $U \subset M$ of $\bar{x}$ such that $Z \cap U$ is a union of at least two $C^{1}$-arcs which intersect transversally at $x_{0}$. In particular, $Z$ cannot be a Jordan curve. Hence, we can find a closed Jordan curve $\Gamma$ with the property that $\Gamma$ is piecewise $C^{1}$ and $\Gamma \subsetneq Z$. Since $M$ has genus $0, \Gamma$ bounds a disk in $M$.

The complement $\mathbb{R}^{3} \backslash M$ has two connected components which we denote by $\Omega$ and $\tilde{\Omega}$. To fix notation, let us assume that $\Omega$ is unbounded and $\tilde{\Omega}$ is bounded.

Proposition 7. There exists a smooth surface $\Sigma \subset \Omega$ such that $\bar{\Sigma} \backslash \Sigma=\Gamma$ and $|A|^{2}=\langle x, \nu\rangle=0$ at each point on $\Sigma$.

Proof. For $k$ sufficiently large, we denote by $\mathscr{C}_{k}$ the set of all embedded disks $S \subset \bar{\Omega} \cap\{|x| \leq 2 k\}$ with the property that $\partial S=\Gamma$. The fact that $\Gamma$ bounds a disk in $M$ implies that $\mathscr{C}_{k}$ is nonempty if $k$ is sufficiently large. Moreover, we choose a smooth cutoff function $\psi_{k}:[0, \infty) \rightarrow[0,1]$ satisfying $\psi_{k}=0$ on $[0, k]$ and $\psi_{k}^{\prime}(2 k)>k$. We now consider the functional

$$
\mathscr{F}_{k}(S)=\int_{S} e^{-\frac{|x|^{2}}{4}+\psi_{k}(|x|)}
$$

for $S \in \mathscr{C}_{k}$. We can interpret $\mathscr{F}_{k}$ as the area functional for the conformal metric $e^{-\frac{|x|^{2}}{4}+\psi_{k}(|x|)} \delta_{i j}$. For $k$ sufficiently large, the region $\bar{\Omega} \cap\{|x| \leq 2 k\}$ is a mean convex domain with respect to this conformal metric. Therefore, general results from [20] guarantee that there exists a smooth embedded surface
$\Sigma_{k} \in \mathscr{C}_{k}$ which minimizes the functional $\mathscr{F}_{k}$. Since $\Sigma_{k}$ is a global minimizer for the functional $\mathscr{F}_{k}$, it is easy to see that

$$
\sup _{k} \mathscr{F}_{k}\left(\Sigma_{k}\right)<\infty .
$$

This implies

$$
\sup _{k} \int_{\Sigma_{k}} e^{-\frac{|x|^{2}}{4}}<\infty
$$

Using the first variation formula, we deduce that $H=\frac{1}{2}\langle x, \nu\rangle$ on $\Sigma_{k} \cap\{|x| \leq k\}$. Finally, the stability inequality implies that

$$
0 \leq-\int_{\Sigma_{k}} e^{-\frac{|x|^{2}}{4}} f L f
$$

for every smooth function $f: \bar{\Sigma}_{k} \rightarrow \mathbb{R}$ which vanishes on the set

$$
\Gamma \cup\left(\bar{\Sigma}_{k} \cap\{|x| \geq k\}\right) .
$$

Using Proposition 5, we obtain

$$
\begin{align*}
& \limsup _{k \rightarrow \infty} \int_{\Sigma_{k} \cap\{|x| \leq \sqrt{k}\}}|A|^{2} e^{-\frac{|x|^{2}}{4}}\langle a, x\rangle^{2}  \tag{1}\\
& \quad \leq \limsup _{k \rightarrow \infty} \frac{C}{\log k} \int_{\Sigma_{k} \cap\{\sqrt{k} \leq|x| \leq k\}} e^{-\frac{|x|^{2}}{4}}=0
\end{align*}
$$

Finally, it follows from Theorem 3 in [22] that

$$
\limsup _{k \rightarrow \infty} \sup _{\Sigma_{k} \cap W}|A|^{2}<\infty
$$

for every compact set $W \subset \mathbb{R}^{3} \backslash \Gamma$. Hence, after passing to a subsequence if necessary, the surfaces $\Sigma_{k}$ converge in $C_{l o c}^{\infty}\left(\mathbb{R}^{3} \backslash \Gamma\right)$ to a smooth surface $\Sigma \subset \mathbb{R}^{3} \backslash \Gamma$ which satisfies the shrinker equation $H=\frac{1}{2}\langle x, \nu\rangle$. Using (1), we conclude that $\Sigma$ is totally geodesic. In particular, $\langle x, \nu\rangle=0$ at each point on $\Sigma$. Moreover, it is easy to see that $\Sigma \subset \bar{\Omega}$. Since $M$ is not totally geodesic, the strict maximum principle implies that $\Sigma$ cannot touch $M$. Consequently, $\Sigma \subset \Omega$.

We next show that $\Gamma \subset \bar{\Sigma}$. If $\Gamma \backslash \bar{\Sigma} \neq \emptyset$, we can construct a one-form $\alpha$ on $\mathbb{R}^{3}$ such that $\alpha$ has compact support, $\alpha=0$ in an open neighborhood of $\bar{\Sigma}$, $d \alpha=0$ in an open neighborhood of $\Gamma$, and $\int_{\Gamma} \alpha \neq 0$. This implies $\int_{\Sigma_{k}} d \alpha=$ $\int_{\Gamma} \alpha \neq 0$ for each $k$, and $\int_{\Sigma_{k}} d \alpha \rightarrow 0$ as $k \rightarrow \infty$. This is a contradiction. Thus, $\Gamma \subset \bar{\Sigma}$. Since $\bar{\Sigma} \backslash \Sigma \subset \Gamma$, we conclude that $\bar{\Sigma} \backslash \Sigma=\Gamma$. This completes the proof of Proposition 7.

Proposition 8. There exists a smooth surface $\tilde{\Sigma} \subset \tilde{\Omega}$ such that $\overline{\tilde{\Sigma}} \backslash \tilde{\Sigma}=\Gamma$ and $|A|^{2}=\langle x, \nu\rangle=0$ at each point on $\tilde{\Sigma}$.

Proof. We consider the set $\tilde{\mathscr{C}}$ of all embedded disks $S \subset \overline{\tilde{\Omega}}$ with boundary $\partial S=\Gamma$. As above, $\tilde{\mathscr{C}}$ is nonempty since $\Gamma$ bounds a disk in $M$. We now consider the functional $\mathscr{F}$ defined in Section 2. The functional $\mathscr{F}$ can be viewed as the area functional for the conformal metric $e^{-\frac{|x|^{2}}{4}} \delta_{i j}$. Clearly, $\overline{\tilde{\Omega}}$
is a mean convex domain with respect to this conformal metric. Using results in [20], we can find a smooth embedded surface $\tilde{\Sigma} \in \tilde{\mathscr{C}}$ which minimizes the functional $\mathscr{F}$. The first variation formula implies that the surface $\tilde{\Sigma}$ satisfies $H=\frac{1}{2}\langle x, \nu\rangle$. Moreover, the stability inequality gives

$$
0 \leq-\int_{\tilde{\Sigma}} e^{-\frac{|x|^{2}}{4}} f L f
$$

for every smooth function $f: \overline{\tilde{\Sigma}} \rightarrow \mathbb{R}$ which vanishes on the boundary $\Gamma$. Using Proposition 5 with $k$ sufficiently large, we obtain

$$
\int_{\tilde{\Sigma}}|A|^{2} e^{-\frac{\mid x x^{2}}{4}}\langle a, x\rangle^{2}=0 .
$$

Consequently, $\tilde{\Sigma}$ is totally geodesic. This implies $\langle x, \nu\rangle=0$ at each point on $\tilde{\Sigma}$. Finally, we clearly have $\tilde{\Sigma} \subset \tilde{\Omega}$. Since $M$ is not totally geodesic, the strict maximum principle implies that $\tilde{\Sigma}$ cannot touch $M$. Therefore, $\tilde{\Sigma} \subset \tilde{\Omega}$, as claimed.

Proposition 9. The unit normal vectors to $\Sigma$ and $\tilde{\Sigma}$ are parallel to a at all points.

Proof. Suppose that there exists a point $x \in \Sigma$ such that $\nu(x)=b$, where $a$ and $b$ are linearly independent. Let us define

$$
\Sigma^{\prime}=\{x \in \Sigma: \nu(x)=b\} \neq \emptyset .
$$

By Proposition 7, $\Sigma^{\prime}$ is a subset of $\left\{x \in \mathbb{R}^{3}:\langle b, x\rangle=0\right\} \backslash \Gamma$. Moreover, Proposition 7 implies that $\Sigma^{\prime}$, viewed as a subset of $\left\{x \in \mathbb{R}^{3}:\langle b, x\rangle=0\right\} \backslash \Gamma$, is open and closed. On the other hand, we have

$$
\left\{x \in \mathbb{R}^{3}:\langle b, x\rangle=0\right\} \cap \Gamma \subset\left\{x \in \mathbb{R}^{3}:\langle a, x\rangle=\langle b, x\rangle=0\right\}=: L .
$$

Hence, the closure of $\Sigma^{\prime}$ is either an entire plane or a halfplane with boundary $L$. In the latter case, we have $L \subset \Gamma \subset M$, but this is impossible since $M$ is compact. Consequently, the closure of $\Sigma^{\prime}$ is the entire plane $\left\{x \in \mathbb{R}^{3}:\langle b, x\rangle=0\right\}$. Since $\Sigma^{\prime} \subset \Omega$, it follows that the plane $\left\{x \in \mathbb{R}^{3}:\langle b, x\rangle=0\right\}$ is contained in $\bar{\Omega}$, and $M$ lies on one side of this plane. This contradicts the fact that $\int_{M} e^{-\frac{|x|^{2}}{4}}\langle b, x\rangle=0$. Consequently, the normal vector to $\Sigma$ is parallel to $a$ at each point on $\Sigma$. An analogous argument shows that the normal vector to $\tilde{\Sigma}$ is parallel to $a$ at each point on $\tilde{\Sigma}$. This completes the proof of Proposition 9 .

Combining Propositions 7, 8, and 9, we conclude that the surfaces $\Sigma$ and $\tilde{\Sigma}$ are contained in the plane $\left\{x \in \mathbb{R}^{3}:\langle a, x\rangle=0\right\}$. Moreover, we have $\Sigma \subset \Omega$ and $\tilde{\Sigma} \subset \tilde{\Omega}$; in particular, $\Sigma$ and $\tilde{\Sigma}$ are disjoint. Finally, $\Sigma$ and $\tilde{\Sigma}$ have the same boundary $\Gamma$. Therefore, the union $\Sigma \cup \tilde{\Sigma} \cup \Gamma$, viewed as a subset of
$\left\{x \in \mathbb{R}^{3}:\langle a, x\rangle=0\right\}$, is open and closed. This implies

$$
\left\{x \in \mathbb{R}^{3}:\langle a, x\rangle=0\right\}=\Sigma \cup \tilde{\Sigma} \cup \Gamma \subset \Omega \cup \tilde{\Omega} \cup \Gamma=\left(\mathbb{R}^{3} \backslash M\right) \cup \Gamma .
$$

Consequently,

$$
\{x \in M:\langle a, x\rangle=0\} \subset \Gamma
$$

In other words, the set $Z$ coincides with $\Gamma$. This contradicts our choice of $\Gamma$. This completes the proof of Theorem 1.

## 4. Proof of Theorem 2

In this final section, we discuss the proof of Theorem 2. Throughout this section, we assume that $M$ is a properly embedded self-similar shrinker in $\mathbb{R}^{3}$. We first recall a well-known result, which is an immediate consequence of Brakke's local area bound (see [2] or [12, Prop. 4.9]):

Proposition 10. For $k$ large, the area of $M \cap\{|x| \leq k\}$ is at most $O\left(k^{2}\right)$.
Proof. We sketch the proof for the convenience of the reader. By assumption, the surfaces $M_{t}=\sqrt{-t} M$ form a solution of mean curvature flow. Applying Proposition 4.9 in [12] (with $t_{0}=0$ and $\rho=4$ ) gives

$$
\operatorname{area}\left(M_{t} \cap\{|x| \leq 2\}\right) \leq 8 \operatorname{area}\left(M_{-1} \cap\{|x| \leq 4\}\right)
$$

for all $t \in[-1,0)$. This implies

$$
\operatorname{area}(M \cap\{|x| \leq 2 k\}) \leq 8 k^{2} \operatorname{area}(M \cap\{|x| \leq 4\})
$$

for $k \geq 1$. From this, the assertion follows.
We will also need the following result, which is a special case of a much more general theorem of Brian White (see [25] for an announcement):

Theorem 11 (B. White [25]). Suppose that $M$ contains the line $\left\{x \in \mathbb{R}^{3}\right.$ : $\left.x_{1}=x_{2}=0\right\}$, and $M$ is disjoint from the halfplane $\left\{x \in \mathbb{R}^{3}: x_{1}<0, x_{2}=0\right\}$. Then $M$ is a plane.

Proof. We again sketch an argument for the convenience of the reader. We define a vector field $\Xi$ on $\mathbb{R}^{3}$ by $\Xi\left(x_{1}, x_{2}, x_{3}\right)=\left(-x_{2}, x_{1}, 0\right)$, and let $\xi$ denote the projection of $\Xi$ to the tangent plane of $M$. By assumption, $M$ is disjoint from the halfplane $\left\{x \in \mathbb{R}^{3}: x_{1}<0, x_{2}=0\right\}$. Hence, every point $x \in M \backslash\left\{x \in \mathbb{R}^{3}: x_{1}=x_{2}=0\right\}$ can be uniquely written in the form $x=\left(\sqrt{x_{1}^{2}+x_{2}^{2}} \cos \theta, \sqrt{x_{1}^{2}+x_{2}^{2}} \sin \theta, x_{3}\right)$ for some $\theta \in(-\pi, \pi)$. This defines a smooth function $\theta: M \backslash\left\{x \in \mathbb{R}^{3}: x_{1}=x_{2}=0\right\} \rightarrow(-\pi, \pi)$ satisfying $\left(x_{1}^{2}+x_{2}^{2}\right) \nabla^{M} \theta=\xi$.

Using Proposition 3, we obtain

$$
\operatorname{div}_{M} \xi-\frac{1}{2}\langle x, \xi\rangle=\sum_{i=1}^{2}\left\langle\bar{D}_{e_{i}} \Xi, e_{i}\right\rangle-\frac{1}{2}\langle x, \Xi\rangle=0 .
$$

In the last step, we have used that $\Xi$ is a Killing vector field in ambient space. This implies

$$
\begin{aligned}
\operatorname{div}_{M}\left(e^{-\frac{|x|^{2}}{4}} \theta \xi\right) & =e^{-\frac{|x|^{2}}{4}} \theta\left(\operatorname{div}_{M} \xi-\frac{1}{2}\langle x, \xi\rangle\right)+e^{-\frac{|x|^{2}}{4}}\left\langle\nabla^{M} \theta, \xi\right\rangle \\
& =e^{-\frac{\mid x x^{2}}{4}}\left(x_{1}^{2}+x_{2}^{2}\right)\left|\nabla^{M} \theta\right|^{2}
\end{aligned}
$$

on $M \backslash\left\{x \in \mathbb{R}^{3}: x_{1}=x_{2}=0\right\}$. Integrating over $M \backslash\left\{x \in \mathbb{R}^{3}: x_{1}=x_{2}=0\right\}$ gives

$$
\begin{aligned}
0 & =\int_{M \backslash\left\{x \in \mathbb{R}^{3}: x_{1}=x_{2}=0\right\}} \operatorname{div}_{M}\left(e^{-\frac{|x|^{2}}{4}} \theta \xi\right) \\
& =\int_{M \backslash\left\{x \in \mathbb{R}^{3}: x_{1}=x_{2}=0\right\}} e^{-\frac{|x|^{2}}{4}}\left(x_{1}^{2}+x_{2}^{2}\right)\left|\nabla^{M} \theta\right|^{2} .
\end{aligned}
$$

Consequently, $\nabla^{M} \theta=0$ at each point $x \in M \backslash\left\{x \in \mathbb{R}^{3}: x_{1}=x_{2}=0\right\}$. Thus, $M$ is a plane. This completes the proof of Theorem 11.

We now continue with the proof of Theorem 2. Let $M$ be a properly embedded self-similar shrinker in $\mathbb{R}^{3}$ with the property that any two loops in $M$ have vanishing intersection number mod 2 . Moreover, we assume that $M$ is neither a round sphere, nor a cylinder, nor a plane. By Proposition 10, $M$ has polynomial area growth. By a theorem of Colding and Minicozzi, the mean curvature must change sign (see [9, Th. 10.1]). In particular, we can find a point $\bar{x} \in M$ such that $H(\bar{x})=0$, and hence $\langle\bar{x}, \nu(\bar{x})\rangle=0$. As above, we put $a:=\nu(\bar{x})$. We now consider two cases.

Case 1: Suppose that the sets $\{x \in M:\langle a, x\rangle>0\}$ and $\{x \in M:\langle a, x\rangle<0\}$ are both connected. In this case, we can construct two loops with the property that the first loop is contained in $\{x \in M:\langle a, x\rangle>0\} \cup\{\bar{x}\}$, the second loop is contained in $\{x \in M:\langle a, x\rangle<0\} \cup\{\bar{x}\}$, and the two loops intersect transversally at $\bar{x}$. This contradicts our assumption that any two loops in $M$ have vanishing intersection number mod 2.

Case 2: For the remainder of this section, we will assume that one of the sets $\{x \in M:\langle a, x\rangle>0\}$ and $\{x \in M:\langle a, x\rangle<0\}$ is not connected. Without loss of generality, we may assume that $\{x \in M:\langle a, x\rangle>0\}$ is disconnected. Let $D$ be an arbitrary connected component of $\{x \in M:\langle a, x\rangle>0\}$. Moreover, let $\Omega$ and $\tilde{\Omega}$ denote the connected components of $\mathbb{R}^{3} \backslash M$.

Proposition 12. There exists a smooth surface $\Sigma \subset \Omega$ such that $\bar{\Sigma} \backslash \Sigma=$ $\partial D$ and $|A|^{2}=\langle x, \nu\rangle=0$ at each point on $\Sigma$.

Proof. By Sard's Lemma, we can find a sequence of numbers $r_{k} \in(2 k, 3 k)$ such that the sphere $\left\{|x|=r_{k}\right\}$ intersects $M$ transversally. By smoothing out the domain $\Omega \cap\left\{|x|<r_{k}\right\}$, we can construct an open domain $\Omega_{k}$ with smooth boundary such that $\Omega_{k} \cap\{|x| \leq 2 k\}=\Omega \cap\{|x| \leq 2 k\}$ and $\Omega_{k} \subset\{|x| \leq 3 k\}$. Moreover, we can find a smooth function $\chi_{k}: \bar{\Omega}_{k} \rightarrow[0,1]$ such that $\chi_{k}=0$ on the set $\bar{\Omega}_{k} \cap\{|x| \leq k\}$ and $\bar{\Omega}_{k}$ is a mean convex domain with respect to the conformal metric $e^{-\frac{|x|^{2}}{4}+\chi_{k}(x)} \delta_{i j}$.

By Sard's Lemma, we can find a real number $\rho_{k} \in(k, 2 k)$ such that the sphere $\left\{|x|=\rho_{k}\right\}$ intersects $M$ and $\partial D$ transversally. Clearly, the curve $\Gamma_{k}=\partial\left(D \cap\left\{|x|<\rho_{k}\right\}\right)$ satisfies $\Gamma_{k} \cap\{|x| \leq k\} \subset \partial D \subset\left\{x \in \mathbb{R}^{3}:\langle a, x\rangle=0\right\}$. Let $\Sigma_{k}$ be a surface which minimizes the modified area functional

$$
\int_{S} e^{-\frac{|x|^{2}}{4}+\chi_{k}(x)}
$$

among all embedded, orientable surfaces $S \subset \bar{\Omega}_{k}$ with boundary $\partial S=\Gamma_{k}$. Clearly,

$$
\sup _{k} \int_{\Sigma_{k}} e^{-\frac{|x|^{2}}{4}}<\infty
$$

Moreover, the first variation formula implies that $H=\frac{1}{2}\langle x, \nu\rangle$ at each point on $\Sigma_{k} \cap\{|x| \leq k\}$. Using the stability inequality together with Proposition 5, we conclude that

$$
\begin{align*}
& \limsup _{k \rightarrow \infty} \int_{\Sigma_{k} \cap\{|x| \leq \sqrt{k}\}}|A|^{2} e^{-\frac{|x|^{2}}{4}}\langle a, x\rangle^{2}  \tag{2}\\
& \quad \leq \limsup _{k \rightarrow \infty} \frac{C}{\log k} \int_{\Sigma_{k} \cap\{\sqrt{k} \leq|x| \leq k\}} e^{-\frac{|x|^{2}}{4}}=0
\end{align*}
$$

Finally, it follows from results in [22] that

$$
\limsup _{k \rightarrow \infty} \sup _{\Sigma_{k} \cap W}|A|^{2}<\infty
$$

for every compact set $W \subset \mathbb{R}^{3} \backslash \partial D$. Hence, after passing to a subsequence, the surfaces $\Sigma_{k}$ converge in $C_{\text {loc }}^{\infty}\left(\mathbb{R}^{3} \backslash \partial D\right)$ to a smooth surface $\Sigma \subset \mathbb{R}^{3} \backslash \partial D$ which satisfies the shrinker equation $H=\frac{1}{2}\langle x, \nu\rangle$. Using (2), we conclude that $\Sigma$ is totally geodesic. In particular, $\langle x, \nu\rangle=0$ at each point on $\Sigma$. Moreover, it is easy to see that $\Sigma \subset \bar{\Omega}$. Since $M$ is not totally geodesic, the strict maximum principle implies that $\Sigma$ cannot touch $M$. Consequently, $\Sigma \subset \Omega$. Arguing as above, we obtain $\bar{\Sigma} \backslash \Sigma=\partial D$.

Proposition 13. There exists a smooth surface $\tilde{\Sigma} \subset \Omega$ such that $\overline{\tilde{\Sigma}} \backslash \tilde{\Sigma}=$ $\partial D$ and $|A|^{2}=\langle x, \nu\rangle=0$ at each point on $\tilde{\Sigma}$.

Proof. Analogous to Proposition 12.
Proposition 14. The unit normal vectors to $\Sigma$ and $\tilde{\Sigma}$ are parallel to a at all points.

Proof. Suppose that there exists a point on $\Sigma$ or $\tilde{\Sigma}$ where the unit normal vector is not parallel to $a$. Without loss of generality, we may assume that there exists a point $x \in \Sigma$ such that $\nu(x)=b$, where $a$ and $b$ are linearly independent. Let

$$
\Sigma^{\prime}=\{x \in \Sigma: \nu(x)=b\} \neq \emptyset .
$$

As above, $\Sigma^{\prime}$ is a subset of $\left\{x \in \mathbb{R}^{3}:\langle b, x\rangle=0\right\} \backslash \partial D$, which is both open and closed. Since

$$
\left\{x \in \mathbb{R}^{3}:\langle b, x\rangle=0\right\} \cap \partial D \subset\left\{x \in \mathbb{R}^{3}:\langle a, x\rangle=\langle b, x\rangle=0\right\}=: L,
$$

it follows that the closure of $\Sigma^{\prime}$ is either an entire plane or a halfplane with boundary $L$. If the closure of $\Sigma^{\prime}$ is a halfplane with boundary $L$, then we have $L \subset \partial D \subset M$ and $\Sigma^{\prime} \subset \Omega \subset \mathbb{R}^{3} \backslash M$. In other words, $M$ contains the line $L$, and $M$ is disjoint from the halfplane $\Sigma^{\prime}$. Hence, it follows from Theorem 11 that $M$ is a plane, contrary to our assumption. Consequently, the closure of $\Sigma^{\prime}$ is the entire plane $\left\{x \in \mathbb{R}^{3}:\langle b, x\rangle=0\right\}$. Since $\Sigma^{\prime} \subset \Omega$, it follows that the plane $\left\{x \in \mathbb{R}^{3}:\langle b, x\rangle=0\right\}$ is contained in $\bar{\Omega}$, and $M$ lies on one side of this plane. As above, this contradicts the fact that $\int_{M} e^{-\frac{|x|^{2}}{4}}\langle b, x\rangle=0$. This completes the proof.

It follows from Propositions 12,13 , and 14 that $\Sigma$ and $\tilde{\Sigma}$ are contained in the plane $\left\{x \in \mathbb{R}^{3}:\langle a, x\rangle=0\right\}$. Moreover, $\Sigma$ and $\tilde{\Sigma}$ are disjoint, and have the same boundary $\partial D$. Putting these facts together, we conclude that

$$
\left\{x \in \mathbb{R}^{3}:\langle a, x\rangle=0\right\}=\Sigma \cup \tilde{\Sigma} \cup \partial D \subset \Omega \cup \tilde{\Omega} \cup \partial D=\left(\mathbb{R}^{3} \backslash M\right) \cup \partial D .
$$

Thus, $\{x \in M:\langle a, x\rangle=0\}=\partial D$. This implies $\{x \in M:\langle a, x\rangle>0\}=D$. In particular, the set $\{x \in M:\langle a, x\rangle>0\}$ is connected, contrary to our assumption. This completes the proof of Theorem 2.

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