Characterization of Lee-Yang polynomials

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Abstract

The Lee-Yang circle theorem describes complex polynomials of degree n in z with all their zeros on the unit circle $|z| = 1$. These polynomials are obtained by taking $z_1 = \cdots = z_n = z$ in certain multiaffine polynomials $\Psi(z_1, \ldots, z_n)$ which we call Lee-Yang polynomials (they do not vanish when $|z_1|, \ldots, |z_n| < 1$ or $|z_1|, \ldots, |z_n| > 1$). We characterize the Lee-Yang polynomials Ψ in $n+1$ variables in terms of polynomials Φ in n variables (those such that $\Phi(z_1, \ldots, z_n) \neq 0$ when $|z_1|, \ldots, |z_n| < 1$). This characterization gives us a good understanding of Lee-Yang polynomials and allows us to exhibit some new examples. In the physical situation where the Ψ are temperature dependent partition functions, we find that those Ψ which are Lee-Yang polynomials for all temperatures are precisely the polynomials with pair interactions originally considered by Lee and Yang.

Introduction

The Lee-Yang circle theorem [5] states that certain polynomials P of degree n in one complex variable z have all their zeros on the unit circle $|z| = 1$. The polynomials P are obtained by taking $z_1 = \cdots = z_n = z$ in $\Psi(z_1, \ldots, z_n)$ when Ψ is a *Lee-Yang polynomial*, i.e., belongs to a certain class LY_n of complex polynomials separately of degree 1 in n variables (multiaffine polynomials). Specifically, LY_n consists of those Ψ such that $\Psi(z_1, \ldots, z_n) \neq 0$ when $|z_1|, \ldots, |z_n| < 1$ and when $|z_1|, \ldots, |z_n| > 1$ (including $|z_i| = \infty$ in a sense to be made precise later). Our current understanding of Lee-Yang polynomials is based on the concept of *Asano contraction* [1]. We shall define an *inner radius* associated with a multiaffine polynomial Φ, and see that it behaves supermultiplicatively with respect to Asano contraction (Proposition 2). Using the properties of the inner radius, we shall characterize the $\Psi \in \mathrm{LY}_{n+1}$ ($n+1$ variables) in terms of polynomials Φ in n variables such that $\Phi(z_1, \ldots, z_n) \neq 0$ when $|z_1|, \ldots, |z_n| < 1$ (Theorem 3). This characterization will give us a good understanding of LY_{n+1} (Proposition 5), and allow us to exhibit elements of LY_{n+1} outside of the (pair interaction) class originally considered by Lee and Yang (see in particular Example 7(d)). The original

Lee-Yang class is obtained for "temperature dependent polynomials" by imposing the Lee-Yang condition at all temperatures (Theorem 9). For $n > 1$, this class is a lower dimensional set in LY_{n+1}. In Section 10 we discuss the situation where multiaffine polynomials are replaced by polynomials of higher degree in each variable. We conclude by briefly presenting a physical interpretation of the results obtained (§11).

The approach to Lee-Yang theory presented here is based on a small set of ideas, as indicated by the minimal list of references. A broader view of the subject is provided by J. Borcea and P. Brändén [3], [4]. This contribution, while essentially disjoint in content from the present paper, has in particular the interest of giving an extensive bibliography covering the work of C. M. Newman, E. H. Lieb, A. D. Sokal, etc., as well as older classical references.

1. *Definitions.* Let $\mathcal{A}_n \subset \mathbf{C}[z_1, \ldots, z_n]$ consist of the *multiaffine polynomials*, i.e., those which are separately of degree 1 in each variable $z_1, \ldots, z_n$, with $n \geq 1$. A polynomial $\Phi \in \mathcal{A}_n$ is thus of the form

$$\Phi(z_1, \ldots, z_n) = \sum_{X \subset [n]} E_X z^X$$

where we have $[n] = \{1, \ldots, n\}$ and $z^X = \prod_{x \in X} z_x$.

We associate with Φ the polynomial $\Phi^\dagger$ such that

$$\Phi^\dagger(z_1, \ldots, z_n) = \sum_{X \subset [n]} E^*_{[n] \setminus X} z^X$$

and note that

$$\Phi^\dagger(z_1, \ldots, z_n) = z_1 \cdots z_n \Phi(z_1^{*-1}, \ldots, z_n^{*-1})^*$$

so that if $|\alpha_1| = \cdots = |\alpha_n| = 1$ we have $|\Phi^\dagger(\alpha_1, \ldots, \alpha_n)| = |\Phi(\alpha_1, \ldots, \alpha_n)|$.

If $\Phi_1, \Phi_2 \in \mathcal{A}_n$, with

$$\Phi_1(z_1, \ldots, z_n) = \sum_{X \subset [n]} E^1_X z^X , \qquad \Phi_2(z_1, \ldots, z_n) = \sum_{X \subset [n]} E^2_X z^X$$

we define a product,[1] $\Phi_1 * \Phi_2 \in \mathcal{A}_n$, so that

$$\Phi_1 * \Phi_2(z_1, \ldots, z_n) = \sum_{X \subset [n]} E^1_X E^2_X z^X.$$

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[1] As pointed out by the referee, $\Phi_1 * \Phi_2$ is known as the Schur-Hadamard product, and has been denoted by $\Phi_1 \bullet \Phi_2$. If Φ is considered as the Fourier transform of $X \mapsto E_X$, then $\Phi_1 * \Phi_2$ may be viewed as the convolution product of Φ_1 and Φ_2. Convolution should multiply the singularities of Φ_1 and Φ_2. Remarkably, if we think of the zeros (up to sign) as "singularities", then Proposition A2 indeed expresses a kind of multiplicativity of "singularities" (and there is an extension of this fact to the higher degree situation of §10). In view of this we keep the $*$ notation in the multiaffine and also the higher degree situation (where $*$ is no longer the Schur-Hadamard product).

With respect to $*$, $\mathscr{A}_n$ is thus isomorphic to the multiplicative semigroup of complex functions on $\{X : X \subset [n]\}$. Note that $(\Phi_1 * \Phi_2)^\dagger = \Phi_1^\dagger * \Phi_2^\dagger$.

We define the *inner radius* $r(\Phi)$, for $\Phi \in \mathscr{A}_n$, by $r(\Phi) = \infty$ if Φ is a constant $\neq 0$, and by

$$r(\Phi) = \sup\{r \geq 0 : \Phi(z_1 \ldots, z_n) \neq 0 \quad \text{if} \quad |z_1|, \ldots, |z_n| < r\}$$

otherwise. By compactness, it follows that $\Phi(\xi_1, \ldots, \xi_n) = 0$ for some $\xi_1, \ldots, \xi_n$ with $|\xi_1|, \ldots, |\xi_n| \leq r(\Phi)$. If $|\xi_k| < r(\Phi)$ for some k, we may relabel the z_j's so that $\Phi(\xi_1, \ldots, \xi_n) = 0$ with $|\xi_1|, \ldots, |\xi_k| < r(\Phi)$, $|\xi_{k+1}| \ldots, |\xi_n| = r(\Phi)$ and $k < n$. Then, if $k \geq 1$,

$$(z_1, \ldots, z_k) \mapsto \Phi(z_1, \ldots, z_k, \xi_{k+1}, \ldots, \xi_n)$$

must vanish identically. Otherwise for any small change

$$(\xi_{k+1}, \ldots, \xi_n) \to (\eta_{k+1} \ldots, \eta_n)$$

we could find $(\eta_1, \ldots, \eta_k)$ close to $(\xi_1, \ldots, \xi_k)$ so that $\Phi(\eta_1, \ldots, \eta_n) = 0$. In particular we could take $|\eta_1|, \ldots, |\eta_n| < r(\Phi)$, in contradiction to the definition of $r(\Phi)$. A consequence of the above argument is that there are $\zeta_1, \ldots, \zeta_n$ with $\Phi(\zeta_1, \ldots, \zeta_n) = 0$ and $|\zeta_1|, \ldots, |\zeta_n| = r(\Phi)$.

It will be convenient to denote by $[\mathscr{A}_n]_\mathbf{R}, [\mathscr{A}_n]_\mathbf{C}$ the real and complex projective spaces associated with $\mathscr{A}_n$, i.e., the set of nonzero classes $[\Phi] = \{\lambda\Phi : \lambda \neq 0, \lambda \in \mathbf{R} \text{ or } \mathbf{C}\}$. It is readily seen that the function $r(\cdot)$ is well defined on the classes $\in [\mathscr{A}_n]_\mathbf{C}$.

The set LY_n mentioned in the introduction consists of those $\Psi \in \mathscr{A}_n$ such that $r(\Psi) \geq 1$ and $r(\Psi^\dagger) \geq 1$ (and therefore $r(\Psi) = r(\Psi^\dagger) = 1$ by Proposition 2(b) below).

If $\Phi \in \mathscr{A}_n$, we define $\Psi_\Phi \in \mathscr{A}_{n+1}$ by

$$\Psi_\Phi(z_1, \ldots, z_{n+1}) = z_{n+1}\Phi^\dagger(z_1, \ldots, z_n) + \Phi(z_1, \ldots, z_n).$$

It follows that $\Psi_\Phi^\dagger = \Psi_\Phi$ (where $\dagger$ denotes now a map $\mathscr{A}_{n+1} \to \mathscr{A}_{n+1}$), and $\Psi_{\Phi_1 * \Phi_2} = \Psi_{\Phi_1} * \Psi_{\Phi_2}$ (where $*$ in the right-hand side is the product in $\mathscr{A}_{n+1}$).

PROPOSITION 2 (properties of the inner radius $r(\Phi)$). (a) $[\Phi] \mapsto r([\Phi])$ *is continuous* $: [\mathscr{A}_n]_\mathbf{C} \to [0, \infty]$ *(the compactification of the interval* $[0, \infty)$*)*.

(b) $r(\Phi^\dagger) \leq r(\Phi)^{-1}$.

(c) *If* $\Phi_1, \Phi_2 \in \mathscr{A}_n$, *then* $r(\Phi_1 * \Phi_2) \geq r(\Phi_1)r(\Phi_2)$.

(d) $r(\Phi) > 0$ *if and only if* $E_\varnothing \neq 0$ *(in particular,* $r(\Phi) > 0$ *when* Φ *is* $*$-*invertible).*

(e) *If* $r(\Phi) \geq 1$, *then* $|E_X| \leq |E_\varnothing|$ *for all* X.

(f) *For* $\lambda \in \mathbf{C}\backslash\{0\}$ *write* $(\Phi \circ \lambda)(z_1, \ldots, z_n) = \Phi(\lambda z_1, \ldots, \lambda z_n)$, *then* $r(\Phi \circ \lambda) = |\lambda|^{-1}r(\Phi)$.

Proof. (a) Let $\Phi \neq 0$ and $\Phi(z_1 \ldots, z_n) = 0$ with $|z_1|, \ldots, |z_n| \leq r(\Phi)$ finite. Given $\varepsilon > 0$, if $\tilde{\Phi}$ is close to Φ we have $\tilde{\Phi}(\tilde{z}_1, \ldots, \tilde{z}_n) = 0$ for some $(\tilde{z}_1, \ldots, \tilde{z}_n)$ close to $(z_1, \ldots, z_n)$, and therefore $|\tilde{z}_1|, \ldots, |\tilde{z}_n| \leq r(\Phi) + \varepsilon$. Hence $r(\tilde{\Phi}) < r(\Phi) + \varepsilon$. Also, by compactness, we can find $a > 0$ such that $|\Phi(z_1 \ldots, z_n)| \geq 2a$ if $|z_1|, \ldots, |z_n| \leq r(\Phi) - \varepsilon$. Therefore, for $\tilde{\Phi}$ close to Φ, we have $|\tilde{\Phi}(z_1 \ldots, z_n)| \geq a$ if $|z_1|, \ldots, |z_n| \leq r(\Phi) - \varepsilon$, hence $r(\tilde{\Phi}) \geq r(\Phi) - \varepsilon$. Thus, finally, we have $|r(\tilde{\Phi}) - r(\Phi)| \leq \varepsilon$ for $\tilde{\Phi}$ close to Φ if $r(\Phi)$ is finite. We have $r([\Phi]) = \infty$ only when $[\Phi]$ is the class $[1]$ of nonzero constants, and it is readily seen that r is also continuous at $[1]$.

(b) is trivial if $r(\Phi)$ or $r(\Phi^\dagger)$ vanishes. In other cases Φ is nonconstant and $r(\Phi) > 0$. Then we see that there are $\zeta_1, \ldots, \zeta_n$ such that $\Phi(\zeta_1, \ldots, \zeta_n) = 0$ with $|\zeta_1|, \ldots, |\zeta_n| = r(\Phi) > 0$, and therefore $\Phi^\dagger(\zeta_1^{*-1}, \ldots, \zeta_n^{*-1}) = 0$ with $|\zeta_1^{*-1}|, \ldots, |\zeta_n^{*-1}| = r(\Phi)^{-1}$, proving the assertion.

(c) is Corollary A.3 of Appendix A.

(e) Let $E = \max_X |E_X|$ and write $\tilde{\Phi} = \Phi/E$. Then $r(\tilde{\Phi}) \geq 1$ and the coefficients $(E_X/E)^n$ of $\tilde{\Phi}^{*n}$ have a limit over a suitable subsequence $n \to \infty$. Along this subsequence $\tilde{\Phi}^{*n} \to \Phi_0 \in \mathscr{A}_n$, where $\Phi_0 \neq 0$. We have $r(\tilde{\Phi}^{*n}) \geq 1$ by (c); thus $r(\Phi_0) \geq 1$ by (a), and $\Phi_0(0, \ldots, 0) \neq 0$ by (d). Since $\Phi_0(0, \ldots, 0) = \lim(E_\varnothing/E)^n$ we cannot have $|E_\varnothing/E| < 1$. We have thus $|E_\varnothing|/E = 1$. Hence finally $|E_X| \leq E = |E_\varnothing|$.

(d), and (f) are clear from the definitions. □

THEOREM 3 (characterization of LY_{n+1}). *An element $\Psi \in \mathscr{A}_{n+1}$ is in LY_{n+1} if and only if $\Psi = c\Psi_\Phi$ where $|c| = 1$ and $\Phi \in \mathscr{A}_n$ satisfies $r(\Phi) \geq 1$.*

First part of proof ($\Psi \in \text{LY}_{n+1} \Rightarrow \Psi = c\Psi_\Phi$). Let $\Psi \in \text{LY}_{n+1}$, (i.e., $\Psi \in \mathscr{A}_{n+1}$, and $r(\Psi) \geq 1, r(\Psi^\dagger) \geq 1$). We write $\mathbf{T}^n = \{(\alpha_1, \ldots, \alpha_n) : |\alpha_1| = \cdots = |\alpha_n| = 1\}$. When $(\alpha_1, \ldots, \alpha_n) \in \mathbf{T}^n$, the functions

$$z \quad \mapsto \quad \Psi(\alpha_1, \ldots, \alpha_n, z), \quad \Psi^\dagger(\alpha_1, \ldots, \alpha_n, z)$$

are affine on $\mathbf{C}$. If $z \mapsto \Psi(\alpha_1, \ldots, \alpha_n, z)$ is constant $\neq 0$ (thus $\Psi(\alpha_1, \ldots, \alpha_n, 1) \neq 0$) then

$$z \quad \mapsto \quad \Psi^\dagger(\alpha_1, \ldots, \alpha_n, z) = \alpha_1 \cdots \alpha_n z [\Psi(\alpha_1, \ldots, \alpha_n, 1)]^*$$

is nonconstant, vanishing at $z = 0$. By a small change of $(\alpha_1, \ldots, \alpha_n, 0)$ we obtain $(z_1, \ldots, z_n, z_{n+1})$ such that $\Psi^\dagger(z_1, \ldots, z_n, z_{n+1}) = 0$ and $|z_1| < 1, \ldots, |z_{n+1}| < 1$ in contradiction with $r(\Psi^\dagger) \geq 1$. Therefore, if $z \mapsto \Psi(\alpha_1, \ldots, \alpha_n, z)$ is constant, it is identically 0. Similarly, if $z \mapsto \Psi^\dagger(\alpha_1, \ldots, \alpha_n, z)$ is constant, it is identically 0.

If the function $z \mapsto \Psi(\alpha_1, \ldots, \alpha_n, z)$ is not constant, it vanishes at $\beta \in \mathbf{C}$, and $z \mapsto \Psi^\dagger(\alpha_1, \ldots, \alpha_n, z)$ is nonconstant vanishing at β^{*-1}. We cannot have $|\beta| < 1$ because we could, by a small change of $(\alpha_1, \ldots, \alpha_n, \beta)$, obtain $(z_1, \ldots, z_n, z_{n+1})$ such that $\Psi(z_1, \ldots, z_n, z_{n+1}) = 0$ and $|z_1| < 1, \ldots, |z_{n+1}| < 1$ in contradiction with $r(\Psi) \geq 1$. Similarly we cannot have $|\beta| > 1$.

We have thus shown that $z \mapsto \Psi(\alpha_1, \ldots, \alpha_n, z)$, $\Psi^\dagger(\alpha_1, \ldots, \alpha_n, z)$ either are both identically 0, or are both nonconstant and vanish at $z = \beta$ with $|\beta| = 1$. In the latter case, there is $C = C(\alpha_1, \ldots, \alpha_n) \neq 0$ such that $\Psi^\dagger(\alpha_1, \ldots, \alpha_n, z) = C\Psi(\alpha_1, \ldots, \alpha_n, z)$. Since Ψ does not vanish identically and is continuous on $\mathbf{T}^{n+1}$, we can choose a nonempty open set $\mathbb{O} \in \mathbf{T}^{n+1}$ such that $\Psi, \Psi^\dagger$ do not vanish in $\mathbb{O}$, and therefore $\Psi^\dagger(\alpha_1, \ldots, \alpha_n, \alpha_{n+1}) = C(\alpha_1, \ldots, \alpha_n)\Psi(\alpha_1, \ldots, \alpha_n, \alpha_{n+1})$ when $(\alpha_1, \ldots, \alpha_n, \alpha_{n+1}) \in \mathbb{O}$. We have thus $\Psi^\dagger = C\Psi$ in $\mathbb{O}$ with C independent of the coordinate α_{n+1}. But similarly C is independent of $\alpha_1, \ldots, \alpha_n$, and is therefore constant on $\mathbb{O}$. By analyticity, this implies $\Psi^\dagger = C\Psi$ on $\mathbf{C}^{n+1}$.

Note that $\Psi = \Psi^{\dagger\dagger} = (C\Psi)^\dagger = C^*\Psi^\dagger$ so that $C^{*-1} = C$ and $|C| = 1$. Choose c such that $c^{-2} = C$, hence $|c| = 1$. Define $\Psi_c = c^{-1}\Psi$; then $\Psi_c^\dagger = c\Psi^\dagger = cc^{-2}\Psi = c^{-1}\Psi = \Psi_c$. Writing

$$\Psi_c(z_1, \ldots, z_{n+1}) = \sum_{X \subset [n+1]} E_X z^X = \sum_{X \subset [n]} E_X z^X + z_{n+1} \sum_{X \subset [n]} F_X z^X$$

we see that $\Psi_c^\dagger = \Psi_c$ is equivalent to $F_X = E_{[n]\setminus X}^*$, or $\Psi_c = \Psi_\Phi$ with $\Phi(z_1, \ldots, z_n) = \sum_{X \subset [n]} E_X z^X$. We have thus shown that if $\Psi \in LY_{n+1}$, then $\Psi = c\Psi_\Phi$ with $|c| = 1$ and $\Phi \in \mathscr{A}_n$. Furthermore

$$\Phi(z_1, \ldots, z_n) = c^{-1}\Psi(z_1, \ldots, z_n, 0)$$

and since the right-hand side cannot vanish when $|z_1|, \ldots, |z_n| < 1$, we have $r(\Phi) \geq 1$. This concludes the first part of the proof.

Second part of proof ($\Psi \in LY_{n+1} \Leftarrow \Psi = c\Psi_\Phi$). We will now show that $r(\Phi) \geq 1$ implies that $\Psi_\Phi(z_1, \ldots, z_{n+1}) \neq 0$ if $|z_1|, \ldots, |z_{n+1}| < 1$, i.e., $r(\Psi_\Phi) \geq 1$. Since $\Psi_\Phi^\dagger = \Psi_\Phi$, this will imply that $c\Psi_\Phi \in LY_{n+1}$. For $|z_1|, \ldots, |z_n| < 1$, and $z_{n+1} = 0$ we have $\Psi_\Phi(z_1, \ldots, z_n, 0) = \Phi(z_1, \ldots, z_n) \neq 0$. If $z_{n+1} \neq 0$, $\Psi_\Phi(z_1, \ldots, z_n, z_{n+1}) = 0$ is equivalent to

$$z_{n+1}^{-1} = -\frac{\Phi^\dagger(z_1, \ldots, z_n)}{\Phi(z_1, \ldots, z_n)}.$$

We thus have to prove that

$$\left| \frac{\Phi^\dagger(z_1, \ldots, z_n)}{\Phi(z_1, \ldots, z_n)} \right| \leq 1$$

if $|z_1|, \ldots, |z_n| < 1$. Take λ real < 1; then $\Phi \circ \lambda \to \Phi$, $(\Phi \circ \lambda)^\dagger \to \Phi^\dagger$ pointwise when $\lambda \to 1$. Therefore

$$\left| \frac{\Phi^\dagger(z_1, \ldots, z_n)}{\Phi(z_1, \ldots, z_n)} \right| = \lim_{\lambda \to 1} \left| \frac{(\Phi \circ \lambda)^\dagger(z_1, \ldots, z_n)}{(\Phi \circ \lambda)(z_1, \ldots, z_n)} \right|$$

where, by the maximum principle,

$$\left| \frac{(\Phi \circ \lambda)^\dagger(z_1, \ldots, z_n)}{(\Phi \circ \lambda)(z_1, \ldots, z_n)} \right| \leq \max_{|\alpha_1| = \cdots = |\alpha_n| = 1} \left| \frac{(\Phi \circ \lambda)^\dagger(\alpha_1, \ldots, \alpha_n)}{(\Phi \circ \lambda)(\alpha_1, \ldots, \alpha_n)} \right| = 1.$$

This concludes the second part of the proof. □

4. *Remarks.* (a) When P is a complex polynomial of order n, there is a unique $\Psi \in \mathcal{A}_n$ symmetric in its n variables such that $P(z) = \Psi(z, \ldots, z)$. If we write $P \in \mathcal{U}_n$ when all n roots α_i of P satisfy $|\alpha_i| = 1$, then $P \in \mathcal{U}_n$ is equivalent to $\Psi \in \mathrm{LY}_n$ by Grace's Theorem A5. Therefore Theorem 3 shows that:

$P \in \mathcal{U}_{n+1}$ if and only if $P(z) = c(zQ^\dagger(z) + Q(z))$, where $c \neq 0$, $Q(z) = \sum_{\ell=0}^{n} C_\ell z^\ell$ has all its roots β_i in $\{z : |z| \geq 1\} \cup \infty$, and $Q^\dagger(z) = \sum_{\ell=0}^{n} C_{n-\ell}^* z^\ell$.

(b) Let $\phi(z) = (az + b)/(cz + d)$ be a fractional linear transformation of the Riemann sphere (with $ad - bc \neq 0$). If $\Phi \in \mathcal{A}_n$ we may define $\Phi^\phi \in \mathcal{A}_n$ by the replacement $z_i \to (az_i + b)/(cz_i + d)$ in $\Phi(z_1, \ldots, z_n)$, and then chasing the denominators. This yields new versions of Theorem A5 where the unit circle is replaced by the real or the imaginary axis, etc. We shall not discuss here these new versions.

PROPOSITION 5 (the set $[\mathcal{J}_n]$ of Lee-Yang classes and its interior). *Let $\mathcal{H}_n = \{\Phi \in \mathcal{A}_n : \Phi^\dagger = \Phi\}$ and $[\mathcal{H}_n]_{\mathbf{R}}$ be the corresponding real projective space, consisting of classes $[\Phi] = \{\lambda\Phi : \lambda \in \mathbf{R}\backslash\{0\}\}$ of elements $\Phi \in \mathcal{H}_n\backslash\{0\}$. Also let $\mathcal{J}_n = \mathrm{LY}_n \cap \mathcal{H}_n$ and $[\mathcal{J}_n]_{\mathbf{R}} = \{[\Phi] : \Phi \in \mathcal{J}_n\}$. We have then*

(1) $$\mathcal{H}_{n+1} = \{\Psi_\Phi : \Phi \in \mathcal{A}_n\}$$

(2) $$\mathcal{J}_{n+1} = \{\Psi_\Phi : \Phi \in \mathcal{A}_n \quad \text{and} \quad r(\Phi) \geq 1\}.$$

Furthermore, the set $[\mathcal{J}_{n+1}]_{\mathbf{R}}$ is the closure in $[\mathcal{H}_{n+1}]_{\mathbf{R}}$ of its interior $\{[\Psi_\Phi] : r([\Phi]) > 1\} = [\mathcal{J}_{n+1}^\circ]_{\mathbf{R}}$ where $\mathcal{J}_{n+1}^\circ = \{\Psi \in \mathcal{H}_{n+1} : \Psi(z_1, \ldots, z_n, z_{n+1}) \neq 0$ when $|z_1|, \ldots, |z_n| \leq 1$ and $|z_{n+1}| < 1\}$.

We obtain (1) directly from the definitions.

Given $c\Psi_\Phi \in \mathcal{A}_{n+1}$, we see that $c\Phi$ is uniquely determined by taking $z_{n+1} = 0$, and (if $\Phi \neq 0$) $c\Psi_\Phi \in \mathcal{H}_{n+1}$ if and only if c is real. Thus the elements of $\mathrm{LY}_{n+1} \cap \mathcal{H}_{n+1}$ are of the form $\pm\Psi_\Phi = \Psi_{\pm\Phi}$ with $r(\pm\Phi) = r(\Phi) \geq 1$, proving (2).

The class $[\Psi_\Phi]$ is of the form $\{\Psi_{\tilde{\Phi}} : \tilde{\Phi} \in [\Phi]\}$ and $[\Phi] \mapsto [\Psi_\Phi]$ is thus a homeomorphism $[\mathcal{A}_n]_{\mathbf{R}} \to [\mathcal{H}_{n+1}]_{\mathbf{R}}$. The set $[\mathcal{J}_{n+1}]_{\mathbf{R}} = \{[\Psi_\Phi] : r([\Phi]) \geq 1\}$ is homeomorphic to $\{[\Phi] : r([\Phi]) \geq 1\}$, which is closed (hence compact) by continuity of $r(\cdot)$. Similarly, the set $\{[\Psi_\Phi] : r([\Phi]) > 1\}$ is homeomorphic to $\{[\Phi] : r([\Phi]) > 1\}$ which is open by continuity of $r(\cdot)$.

The map $\lambda \mapsto [\Phi \circ \lambda]$ is continuous near $\lambda = 1$ so that, if $r([\Phi]) = 1$, we have $r([\Phi \circ \lambda]) = |\lambda|^{-1}$ which may be > 1 or < 1 for $[\Phi \circ \lambda]$ close to $[\Phi]$. This shows that $\{[\Psi_\Phi] : r([\Phi]) > 1\}$ is the interior of $[\mathcal{J}_{n+1}]_{\mathbf{R}}$, and $[\mathcal{J}_{n+1}]_{\mathbf{R}}$ is the closure of its interior.

Finally we have to show that $\Psi_\Phi \in \mathcal{J}_{n+1}^{\circ} \Leftrightarrow r(\Phi) > 1$. Indeed $\Psi_\Phi \in \mathcal{J}_{n+1}^{\circ} \Rightarrow r(\Phi) > 1$ because $\Phi(z_1, \ldots, z_n) = \Psi(z_1, \ldots, z_n, 0)$. Conversely, if $r(\Phi) > 1$ and $\Psi_\Phi(z_1, \ldots, z_n, z) = 0$ with $|z_1|, \ldots, |z_n| \leq 1$, we must have $z \neq 0$ and

$$z^{-1} = -\frac{\Phi^\dagger(z_1, \ldots, z_n)}{\Phi(z_1, \ldots, z_n)},$$

where

$$\left| -\frac{\Phi^\dagger(z_1, \ldots, z_n)}{\Phi(z_1, \ldots, z_n)} \right| \leq \max_{|\alpha_1| = \cdots = |\alpha_n| = 1} \left| \frac{\Phi^\dagger(\alpha_1, \ldots, \alpha_n)}{\Phi(\alpha_1, \ldots, \alpha_n)} \right| = 1.$$

Hence $|z| \geq 1$, i.e., $\Psi_\Phi \in \mathcal{J}_{n+1}^{\circ}$. $\qquad \square$

6. *Remarks.* (a) Let $\Psi_1, \Psi_2 \in \mathrm{LY}_n$ (or $\mathcal{J}_n$, or $\mathcal{J}_n^{\circ}$); then $\Psi_1 * \Psi_2 \in \mathrm{LY}_n$ (or $\mathcal{J}_n$, or $\mathcal{J}_n^{\circ}$) by Proposition 2(c).

(b) An element Ψ of $\mathcal{H}_{n+1}$ may be written as

$$\Psi(z_1, \ldots, z_{n+1}) = A(z_1, \ldots, z_{n-1}) + B(z_1, \ldots, z_{n-1})z_n$$
$$+ C(z_1, \ldots, z_{n-1})z_{n+1} + D(z_1, \ldots, z_{n-1})z_n z_{n+1}$$

where $A, B, C, D \in \mathcal{A}_{n-1}$ and $D = A^\dagger, C = B^\dagger$ (here $\dagger$ is defined in $\mathcal{A}_{n-1}$). The condition $[\Psi] \in [\mathcal{J}_{n+1}^{\circ}]_{\mathbf{R}}$ is that

$$\Phi(z_1, \ldots, z_n) = A(z_1, \ldots, z_{n-1}) + B(z_1, \ldots, z_{n-1})z_n$$

does not vanish for $|z_1|, \ldots, |z_n| \leq 1$ or equivalently that $A(z_1, \ldots, z_{n-1}) \neq 0$ for $|z_1|, \ldots, |z_{n-1}| \leq 1$ and $|B(\alpha_1, \ldots, \alpha_{n-1})| < |A(\alpha_1, \ldots, \alpha_{n-1})|$ for $|\alpha_1| = \cdots = |\alpha_{n-1}| = 1$. Note that, under this condition, B can equivalently be replaced by C, which corresponds to interchanging z_n and z_{n+1}. This shows that the asymmetric choice of the variable z_{n+1} in our characterization (Theorem 3) of Lee-Yang polynomials is inessential.

(c) Let $E_X^\beta = \exp \beta W_X$ for $X \subset [n]$, with real $\beta > 0$ and $W_X \in \{-\infty\} \cup \mathbf{C}$. Write $\Phi^\beta(z_1, \ldots, z_n) = \sum_{X \subset [n]} E_X^\beta z^X$. If $r(\Phi^1) > 1$ there is an open interval $I \ni 1$ such that $r(\Phi^\beta) > 1$ for $\beta \in I$. For integer m sufficiently large we have $mI \supset (m-1, m+1)$. Hence, by Proposition 2(c), $r(\Phi^\beta) > 1$ for all $\beta \geq m$. Thus, if $\Psi_{\Phi^1} \in \mathcal{J}_n^{\circ}$, then $\Psi_{\Phi^\beta} \in \mathcal{J}_n^{\circ}$ for all sufficiently large β.

(d) *If $\Psi \in \mathcal{J}_{n+1}$ has nontrivial factorization,*

$$\Psi(z_1, \ldots, z_{n+1}) = \Psi'(z_1, \ldots, z_k) \times \Psi''(z_{k+1}, \ldots, z_{n+1}),$$

then $\Psi \notin \mathcal{J}_{n+1}^{\circ}$.

Indeed, we have $\Psi' \in LY_k$, $\Psi'' \in LY_{n-k+1}$, and we may assume $\Psi' \in \mathscr{J}_k$, $\Psi'' \in \mathscr{J}_{n-k+1}$. Therefore we may choose $z_1, \ldots, z_{n+1}$ such that $\Psi'(z_1, \ldots, z_k) = 0$ with $|z_1|, \ldots, |z_k| \le 1$ and $|z_{k+1}|, \ldots, |z_{n+1}| < 1$.

(e) $\mathscr{J}_{n+1}$ has full dimension in $\mathscr{H}_{n+1}$ (the dimension of $\mathscr{H}_{n+1}$ is the real dimension of $\mathscr{A}_n$, i.e., 2^{n+1}). For $n > 1$, this is strictly greater than the dimension $n(n+1)/2 + n + 1 + 1$ of the set of high-temperature polynomials of Theorem 9 below (with b real), which is essentially the class of polynomials originally considered by Lee and Yang [5].

7. *Examples.* In what follows we study
$$\Psi(z_1, \ldots, z_{n+1}) = \sum_{X \subset [n+1]} \left(\prod_U E_{UX} \right) z^X$$
for various choices of the U and E_{UX}.

(a) For $U = \{j, k\} \subset [n+1]$, with $j < k$, write
$$E_{UX} = \begin{cases} a_U & \text{if } j \in X, k \notin X, \; cra_U^* & \text{if } j \notin X, k \in X \\ b_U & \text{if } j, k \notin X \\ b_U^* & \text{if } j, k \in X. \end{cases}$$

Then, if the complex a_U, b_U satisfy $|a_U/b_U| \le 1$ for all U, we have $\Psi \in \mathscr{J}_{n+1}$.

Using Section 6(a), it suffices to consider the case of a single U such that $E_{UX} \ne 1$, and we may thus take $n = 1$, $U = [n+1] = \{1, 2\}$. Then $\Psi = \Psi_\Phi$ with $\Phi(z_1) = b_U + a_U z_1$, hence $r(\Phi) \ge 1$ by assumption, and $\Psi \in \mathscr{J}_{n+1}$ by Proposition 5.   □

(The standard situation [5] corresponds to taking $b_U = 1$ and a_U real. The case of complex a_U was considered by Beauzamy [2]. It is a consequence of Theorem 9 below that for polynomials of the class (a) we may write
$$\Psi(z_1, \ldots, z_{n+1}) = C \tilde{\Psi}(\alpha_1 z_1, \ldots, \alpha_{n+1} z_{n+1})$$
where $C, \alpha_1, \ldots, \alpha_{n+1} \in \mathbf{C}$, $|\alpha_1| = \cdots = |\alpha_{n+1}| = 1$, and $\tilde{\Psi}$ is in the class (a) with real a_U, b_U and $0 \le a_U \le 1, b_U = 1$.)

(b) For $U = \{j, k, l\} \subset [n+1]$, write $E_{UX} = b_U$ if $U \cap X = \varnothing$, $= b_U^*$ if $U \subset X$, $= 1$ otherwise. Then, if b_U is real ≥ 1 for all U, we have $\Psi \in \mathscr{J}_{n+1}$.

It suffices to consider the case of a single U such that $E_{UX} \ne 1$, and we may thus take $n = 2$, $U = [n+1] = \{1, 2, 3\}$. Then $\Psi = \Psi_\Phi$ with
$$\Phi(z_1, z_2) = b_U + z_1 + z_2 + z_1 z_2 = (1 + z_1)(1 + z_2) - (1 - b_U).$$
The set $\{(1 + z_1)(1 + z_2) : |z_1| < 1, |z_2| < 1\}$ is bounded by the cardioid
$$\Gamma = \{\rho e^{i\theta} : \rho = 2(1 + \cos\theta) : \theta \in [-\pi, \pi]\}$$
and does not contain $1 - b_U$, which is real ≤ 0 by assumption. Therefore $r(\Phi) \ge 1$.
                                                                          □

(We note that the situation described in (b), with real $E_{UX} = E_{UX}^{(b)}$, is in fact contained in the situation described in (a), with $E_{UX} = E_{UX}^{(a)}$. Take indeed $b_{\{1,2\}} = b_{\{2,3\}} = b_{\{3,1\}} = b$ in (a). Then if we take $b_{\{1,2,3\}} = b^2$ in (b) we have

$$E_{\{1,2,3\}X}^{(b)} = b^{-1} E_{\{1,2,3\}X}^{(a)}$$

because both sides are equal to b^2 if $|X| = 0$ or 3, and 1 if $|X| = 1$ or 2. Therefore $\Psi^{(b)} = \text{const.} \Psi^{(a)}$. This equivalence fails if b is not real.)

(c) If we take $U = [n+1] = \{1, 2, 3\}$ in (b) and replace the real b_U by $b(\beta) = e^{\beta W}$ with $\text{Re} W > 0$, $\text{Im} W \neq 0$, the set $\{b(\beta) : \beta \text{ real}\}$ is a logarithmic spiral and $\{1 - b(\beta) : \beta > 0\}$ intersects the region $(1 - z_1)(1 - z_2) : |z_1| < 1, |z_2| < 1$ inside the cardioid Γ in a sequence of intervals (the point $1 - b(\beta)$ stays inside (resp. outside) of the cardioid for sufficiently small (resp. large) β). Therefore, there are successive intervals of $\{\beta : \beta > 0\}$ such that $\Psi \notin \mathcal{J}_3$ and $\Psi \in \mathcal{J}_3$.

(d) For $U = \{j, k, l, m\} \subset [n+1]$, and $X \subset [n+1]$, write $E_{UX} = b_U$ if $U \cap X = \varnothing, = b_U^*$ if $U \subset X, = 1$ otherwise. Then, if b_U is real ≥ 2 (or $= 1$) for all U we have $\Psi \in \mathcal{J}_{n+1}$.

It suffices to consider the case of a single U such that $E_{UX} \neq 1$, and we may take $n = 3, U = [n+1] = \{1, 2, 3, 4\}$. Then $\Psi = \Psi_\Phi$ with

$$\Phi(z_1, z_2, z_3) = (1 + z_1)(1 + z_2)(1 + z_3) - (1 - b_U).$$

One checks that $\{(1 + z_1)(1 + z_2)(1 + z_3) : |z_1| < 1, |z_2| < 1, |z_3| < 1\}$ intersects the real interval $(-\infty, 0]$ in $(-1, 0)$, so that for real $b_U \geq 2$ (or $= 1$) we have $r(\Phi) \geq 1$. $\qquad\square$

LEMMA 8. *For real $\beta > 0$, and $W_X \in \mathbf{C}$ if $X \subset [n]$, write*

$$E_X^\beta = \exp \beta W_X, \qquad \Phi^\beta(z_1, \ldots, z_n) = \sum_{X \subset [n]} E_X^\beta z^X.$$

If $r(\Phi^\beta) \geq 1$ for a sequence of β's tending to 0 from above, then

$$W_X = \sum_{j \notin X} \sum_{k \notin X} W_{jk} + \sum_{j \notin X} W_j + W_0$$

with suitable real $W_{jk} = W_{kj} \geq 0, W_{jj} = 0$, and complex W_j, W_0.

For small β we have $E_X^\beta = 1 + \beta W_X + O(\beta^2)$; hence

$$\Phi^\beta(z_1, \ldots, z_n) = \prod_1^n (1 + z_j) + \beta \sum_X W_X z^X + O(\beta^2) \sum_X |z|^X.$$

Writing $\zeta_j = 1 + z_j$ gives (assuming all $\zeta_j \neq 0$, and $\sum |\zeta_j|$ bounded)

$$\Phi^\beta(z_1, \ldots, z_n) = \prod_1^n \zeta_j + \beta \sum_X W'_X \zeta^X + O(\beta^2) \sum_X |\zeta|^X$$

$$= \left(\prod_1^n \zeta_j \right) \left[1 + \beta \sum_X W''_X (\zeta^{-1})^X + O(\beta^2) \sum_X |\zeta^{-1}|^X \right],$$

where the W'_X, W''_X are linear combinations of the W_X with coefficients in $\mathbf{Z}$.

Taking $\gamma = \beta^{1/n}$, $\zeta_1 = \cdots = \zeta_n = \zeta = \gamma\theta^{-1}$, and writing $|X| = \operatorname{card} X$, we have

$$\Phi^\beta(z_1, \ldots, z_n) = \zeta^n \left[1 + W''_{[n]}\theta^n + \sum_{|X|<n} W''_X \gamma^{n-|X|}\theta^{|X|} + O(\gamma^{2n}) \sum_X \gamma^{-|X|}\theta^{|X|} \right]$$

$$= \zeta^n \left[1 + W''_{[n]}\theta^n + O(\gamma) \right]$$

for bounded θ. If $W''_{[n]} \neq 0$ there are thus zeros of Φ^β of the form $z_1 = \cdots = z_n = \zeta - 1$ with

$$\zeta \approx (-W''_{[n]})^{1/n}\gamma \times n\text{-th root of } 1$$

(use the implicit function theorem). For small γ and $n > 2$ this implies that $\Phi^\beta(z_1, \ldots, z_n) = 0$ for some $z_1, \ldots, z_n$ with $|z_1|, \ldots, |z_n| < 1$. If $r(\Phi^\beta) \geq 1$ and $n > 2$ we have thus $W''_{[n]} = 0$.

Put the variables z_j, with $j \notin X$, equal to 0 in Φ^β, and consider the function of the remaining variables z_k, with $k \in X$. Taking now $\gamma = \beta^{1/|X|}$ we see as above that Φ^β vanishes when $z_j = 0$ if $j \notin X, z_k = -1 + \zeta$ if $k \in X$, where

$$\zeta \approx (-W''_X)^{1/|X|}\gamma \times |X|\text{-th root of } 1.$$

Assuming $r(\Phi^\beta) \geq 1$ we have thus $W''_X = 0$ if $|X| > 2$, and W''_X real ≥ 0 if $|X| = 2$ (so that $(-W''_X)^{1/2}$ is pure imaginary).

If $r(\Phi^\beta) \geq 1$ for some arbitrarily small $\beta > 0$, we have thus

$$\sum W_X z^X = \left(\prod_1^n \zeta_j \right) \sum W''_X (\zeta^{-1})^X = \left(\prod_1^n \zeta_j \right) \sum_{X:|X|\leq 2} W''_X (\zeta^{-1})^X$$

$$= \sum_{X:|X|\leq 2} W''_X \zeta^{[n]\setminus X} = \sum_{X:|X|\leq 2} W''_X (1+z)^{[n]\setminus X}$$

$$= \sum_Y z^Y \left[\sum_{\{j,k\}\subset[n]\setminus Y} W''_{\{j,k\}} + \sum_{\{j\}\subset[n]\setminus Y} W''_{\{j\}} + W''_\varnothing \right]$$

and the lemma follows, with $W_{jk} = \frac{1}{2}W''_{\{j,k\}}$, $W_j = W''_{\{j\}}$ and $W_0 = W''_\varnothing$. $\square$

THEOREM 9 (high-temperature Lee-Yang polynomials). *Let $W_X \in \mathbf{C}$, $E_X^\beta = \exp \beta W_X$ and $\Psi^\beta(z_1, \ldots, z_n) = \sum_{X \subset [n]} E_X^\beta z^X$. We say that $(\Psi^\beta)_{\beta > 0}$ is a high-temperature Lee-Yang polynomial if*

(a) $\Psi^\beta \in LY_n$ *for some sequence of real β's tending to zero from above.*
      We claim that (a) *is equivalent to* (b) *and also to* (c):

(b) $\Psi^\beta \in LY_n$ *for all $\beta > 0$.*

(c) *There exist $W_{jk} \in \mathbf{R}, a_j \in \mathbf{R}, b \in \mathbf{C}$ such that $W_{jk} = W_{kj} \geq 0$ and*

$$W_X = -\sum_{j \in X} \sum_{k \notin X} W_{jk} - i \sum_{j \in X} a_j + b.$$

The implication (c)$\Rightarrow$(b) is proved in Section 7(a), and (b)$\Rightarrow$(a) is obvious.

We assume now that (a) holds. Since Ψ^β satisfies the conditions of Lemma 8 we have

$$
\begin{aligned}
W_X &= \sum_{j \notin X} \sum_{k \notin X} W_{jk} + \sum_{j \notin X} W_j + W_0 \\
&= \sum_{j \notin X} \sum_{k} W_{jk} - \sum_{j \notin X} \sum_{k \in X} W_{jk} + \sum_{j \notin X} W_j + W_0 \\
&= -\sum_{j \in X} \sum_{k \notin X} W_{jk} + \sum_{j \notin X} \tilde{W}_j + W_0
\end{aligned}
$$

with real $W_{jk} = W_{kj} \geq 0$, $\tilde{W}_j = W_j + \sum_k W_{jk} \in \mathbf{C}$, $W_0 \in \mathbf{C}$. In view of Theorem 3 the quantity $E_X^\beta / (E_{[n] \setminus X}^\beta)^*$ is independent of X; hence the limit for $\beta \to 0$ of

$$
\begin{aligned}
\frac{d}{d\beta} \exp \beta(W_X - W_{[n] \setminus X}^*) &= \frac{d}{d\beta} \exp \beta \Big(\sum_{j \notin X} \tilde{W}_j - \sum_{j \in X} \tilde{W}_j^* + W_0 - W_0^* \Big) \\
&= \Big(\sum_{j \notin X} \tilde{W}_j - \sum_{j \in X} \tilde{W}_j^* + W_0 - W_0^* \Big) \exp \beta \Big(\sum_{j \notin X} \tilde{W}_j - \sum_{j \in X} \tilde{W}_j^* + W_0 - W_0^* \Big)
\end{aligned}
$$

is also independent of X. Since this limit is $\sum_{j \notin X} \tilde{W}_j - \sum_{j \in X} \tilde{W}_j^* + W_0 - W_0^*$, we have $\tilde{W}_j + \tilde{W}_j^* = 0$, i.e., $\tilde{W}_j = i a_j$ with $a_j \in \mathbf{R}$. Writing $W_0 + i \sum_j a_j = b \in \mathbf{C}$, we have thus proved (a)$\Rightarrow$(c). $\square$

10. *Higher degree.* Given integers $m_1, \ldots, m_n \geq 1$, let $\mathbf{m} = (m_1, \ldots, m_n)$ and $\mathscr{B}_{\mathbf{m}} \subset \mathbf{C}[z_1, \ldots, z_n]$ consist of polynomials separately of degree $m_1, \ldots, m_n$ in $z_1, \ldots, z_n$. A polynomial $\Phi \in \mathscr{B}_{\mathbf{m}}$ is thus of the form

$$\Phi(z_1, \ldots, z_n) = \sum_{k_1=0}^{m_1} \cdots \sum_{k_n=0}^{m_n} E_{k_1 \cdots k_n} \frac{m_1!}{k_1!(m_1 - k_1)!} z_1^{k_1} \cdots \frac{m_n!}{k_n!(m_n - k_n)!} z_n^{k_n}.$$

We define $\Phi^\dagger \in \mathcal{B}_\mathbf{m}$ by

$$\Phi^\dagger(z_1, \ldots, z_n) = \sum_{k_1=0}^{m_1} \cdots \sum_{k_n=0}^{m_n} E^*_{k_1 \cdots k_n} \frac{m_1!}{k_1!(m_1-k_1)!} z_1^{m_1-k_1}$$

$$\cdots \frac{m_n!}{k_n!(m_n-k_n)!} z_n^{m_n-k_n}.$$

We also let $r(\Phi)$ be ∞ if Φ is a constant $\neq 0$, and

$$r(\Phi) = \sup\{r \geq 0 : \Phi(z_1 \ldots, z_n) \neq 0 \quad \text{if} \quad |z_1|, \ldots, |z_n| < r\}$$

otherwise. With this notation we may define a set $\mathrm{LY}_\mathbf{m}$ of Lee-Yang polynomials Ψ in $\mathcal{B}_\mathbf{m}$ by $r(\Psi) = r(\Psi^\dagger) = 1$.

Let now $\mathcal{A}_\mathbf{m} \subset \mathbf{C}[z_{11}, \ldots, z_{1m_1}, \ldots, z_{n1}, \ldots, z_{nm_n}]$ consist of the multiaffine polynomials invariant under permutations of $z_{i1}, \ldots, z_{im_i}$ for $i = 1, \ldots, n$. There is a unique linear isomorphism $\sigma : \mathcal{B}_\mathbf{m} \to \mathcal{A}_\mathbf{m}$ such that $(\sigma^{-1}\Psi)(z_1 \ldots, z_n)$ is obtained if one replaces $z_{i1}, \ldots, z_{im_i}$ by z_i in Ψ. We note that $\sigma\Phi^\dagger = (\sigma\Phi)^\dagger$, and (using Grace's theorem A5) that $r(\Phi) = r(\sigma\Phi)$. If $\Phi_1, \Phi_2 \in \mathcal{B}_\mathbf{m}$ we also define $\Phi_1 * \Phi_2$ by $\sigma(\Phi_1 * \Phi_2) = \sigma\Phi_1 * \sigma\Phi_2$, and we have $r(\Phi_1 * \Phi_2) \geq r(\Phi_1)r(\Phi_2)$.

With the above notation we see that questions about $\mathrm{LY}_\mathbf{m}$ are transformed by use of the isomorphism σ into questions about $\mathrm{LY}_{|\mathbf{m}|}$ where $|\mathbf{m}| = m_1 + \cdots + m_n$. For instance, one finds that if $\Psi \in \mathrm{LY}_\mathbf{m}$, then $\Psi = c\tilde{\Psi}$ where $|c| = 1$ and $\tilde{\Psi} \in \mathrm{LY}_\mathbf{m}$, $\tilde{\Psi}^\dagger = \tilde{\Psi}$.

11. *Physical interpretation.* The polynomials

$$\Psi(z_1, \ldots, z_{n+1}) = \sum_{X \subset [n+1]} E_X z^X$$

relevant to physics are such that $E_X = e^{\beta W_X}$, where $\beta, W_X \in \mathbf{R}$, and $\beta^{-1} > 0$ is interpreted as *temperature*, while $-W_X$ is the *energy* of the *configuration X*. For such Ψ, we have $\Psi \in \mathrm{LY}_{n+1}$ if and only if $\Psi = \Psi_\Phi$ and $r(\Phi) \geq 1$, i.e., $\Psi \in \mathcal{J}_{n+1}$ (Theorem 3 and Proposition 5). Writing $\Psi = \Psi^\beta$, we find (§6(c)) that if $\Psi^\beta \in \mathcal{J}^\circ_{n+1}$ (the interior of $\mathcal{J}_{n+1}$) for some β, then $\Psi^\beta \in \mathcal{J}^\circ_{n+1}$ for all sufficiently large β (i.e., small temperatures). There are cases where $\Psi^\beta \in \mathcal{J}^\circ_{n+1}$ for small temperatures and $\Psi^\beta \notin \mathcal{J}_{n+1}$ for large temperatures (§7(d)). (For complex W_X we may have several successive intervals of temperature where Ψ^β is in and out of $\mathcal{J}_{n+1}$ (§7(c)), I do not know if this can happen for real W_X.) If $\Psi^\beta \in \mathcal{J}_{n+1}$ for a sequence of β's tending to 0 (high temperatures), then $\Psi^\beta \in \mathcal{J}_{n+1}$ for all $\beta > 0$ (all temperatures) and

$$W_X = -\sum_{j \in X} \sum_{k \notin X} W_{jk} + b$$

with $W_{jk}, b \in \mathbf{R}$ such that $W_{jk} = W_{kj} \geq 0$ (Theorem 9). (Since $\Psi^\beta(0, \ldots, 0) = e^{\beta W_\varnothing} = e^{\beta b}$, we have $b \in \mathbf{R}$. Taking only $z_j \neq 0$ in Ψ^β yields $e^{-i\beta a_j} z_j$; hence $a_j = 0$.) This means that those Ψ^β which are Lee-Yang polynomials at high temperatures, hence at all temperatures, are precisely of the form considered by Lee and Yang [5].

The higher degree situation of Section 10 corresponds physically to *higher spins* (the degree m_i corresponds to spin $(m_i + 1)/2$). The $\Psi^\beta \in \mathcal{B}_\mathbf{m}$ relevant to physics which are Lee-Yang polynomials at high temperature, hence at all temperatures, can again be determined. Using the notation of Section 10, they correspond to $E_{k_1 \ldots k_n} = \exp \beta W_{k_1 \ldots k_n}$ (with $\beta > 0$, $W_{k_1 \ldots k_n} \in \mathbf{R}$) such that

$$W_{k_1 \ldots k_n} = -\sum_{i=1}^{n} \sum_{j=1}^{n} W_{ij} k_i (m_j - k_j) + b$$

with $W_{ij}, b \in \mathbf{R}$, $W_{ij} = W_{ji} \geq 0$ (W_{ii} in general does not vanish).

Appendix A. **Some properties of multiaffine polynomials**

See [8] and references quoted there.

LEMMA A.1 ([7]). *Let K_1, K_2 be closed subsets of* $\mathbf{C}$, *with $K_1, K_2 \not\ni 0$. If $\Phi \in \mathcal{A}_2$ and*

$$\Phi(z_1, z_2) \equiv A + B z_1 + C z_2 + D z_1 z_2 \neq 0$$

whenever $z_1 \notin K_1$ and $z_2 \notin K_2$. Then

$$\tilde{\Phi}(z) \equiv A + Dz \neq 0$$

whenever $z \notin -K_1 \cdot K_2$. (We have written $-K_1 \cdot K_2 = \{-uv : u \in K_1, v \in K_2\}$.)

The map $\Phi \mapsto \tilde{\Phi}$ of $\mathcal{A}_2$ to $\mathcal{A}_1$ is called *Asano contraction*. For the convenience of the reader we reproduce here the easy proof of this lemma (which generalizes an earlier result of Asano [1]).

Since $K_1, K_2 \not\ni 0$, we have $A \neq 0$. If $D = 0$, there is nothing to prove. If $D \neq 0$ and $AD - BC = 0$, we have

$$A + B z_1 + C z_2 + D z_1 z_2 = D \left(z_1 + \frac{C}{D} \right) \left(z_2 + \frac{A}{C} \right)$$

which implies $-C/D \in K_1$, $-A/C \in K_2$, hence $A/D \in K_1 \cdot K_2$ and $A + Dz = 0$ only if $z = -A/D \in -K_1 \cdot K_2$.

Suppose now that $D \neq 0$ and $AD - BC \neq 0$, and write[2]

$$\phi(z) = -\frac{A + Bz}{C + Dz} \qquad , \qquad \psi(z) = \frac{A}{Dz},$$

---

[2] This form of the argument was communicated to me by F. J. Dyson.

where ϕ, ψ are considered as mappings of the Riemann sphere (add a point at infinity to $\mathbf{C}$, K_1, K_2). If we write $\omega = \phi\psi^{-1}$, $z_2 = \omega z_1$ is equivalent to

$$AB + ADz_1 + ADz_2 + CDz_1z_2 = 0,$$

showing that ω is an involution. Since K_2 is a proper closed set, ωK_2 cannot be interior to K_2 (otherwise $\omega^2 K_2$ would be interior to K_2). Thus

$$(\omega K_2) \cap (\mathbf{C}\backslash K_2)^- \neq \varnothing$$

where $(\mathbf{C}\backslash K_2)^-$ denotes the closure of the complement of K_2. By assumption $\mathbf{C}\backslash K_2 \subset \phi K_1$ and, since ϕK_1 is closed, $(\mathbf{C}\backslash K_2)^- \subset \phi K_1$; hence $\omega K_2 \cap \phi K_1 \neq \varnothing$, or $\phi\psi^{-1}K_2 \cap \phi K_1 \neq \varnothing$, or $K_2 \cap \psi K_1 \neq \varnothing$, i.e.,

$$(\exists z)\left(\frac{A}{Dz} \in K_1 \text{ and } z \in K_2\right)$$

so that $A/D \in K_1 \cdot K_2$ and finally $-A/D \in -K_1 \cdot K_2$. $\qquad\square$

PROPOSITION A.2 ([7]). . *Let K_{ij} be a closed subset of $\mathbf{C}$ with $K_{ij} \not\ni 0$, for $i = 1, 2$ and $j = 1, \ldots, n$. If $\Phi_1, \Phi_2 \in \mathcal{A}_n$, and $\Phi_i(z_1, \ldots, z_n) \neq 0$ whenever $z_1 \notin K_{i1}, \ldots, z_n \notin K_{in}$, then $\Phi_1 * \Phi_2(z_1, \ldots, z_n) \neq 0$ whenever $z_1 \notin -K_{11} \cdot K_{21}, \ldots, z_n \notin -K_{1n} \cdot K_{2n}$.*

To show this, start from the element of $\mathcal{A}_{2n}$ defined by the product $\Phi_1(z_{11}, \ldots \ldots, z_{1n}) \times \Phi_2(z_{21}, \ldots, z_{2n})$, different from 0 whenever $z_{ij} \notin K_{ij}$ for $i = 1, 2; j = 1, \ldots, n$. Then perform successive Asano contractions $(z_{1k}, z_{2k}) \to z_k$ for $k = 1, \ldots, n$. At the k-th step we have an element of $\mathcal{A}_{2n-k}$. Induction on k and use of Lemma A.1 yield that this element of $\mathcal{A}_{2n-k}$ is $\neq 0$ when $z_1 \notin -K_{11} \cdot K_{21}, \ldots, z_k \notin -K_{1k} \cdot K_{2k}, z_{1,k+1} \notin K_{1,k+1}, z_{2,k+1} \notin K_{2,k+1}, \ldots, z_{1n} \notin K_{1n}, z_{2n} \notin K_{2n}$. Taking $k = n$ proves the proposition. $\qquad\square$

COROLLARY A.3. *If $\Phi_1, \Phi_2 \in \mathcal{A}_n$, then $r(\Phi_1 * \Phi_2) \geq r(\Phi_1)r(\Phi_2)$.*

This follows by taking $K_{ij} = \{z : |z| \geq r(\Phi_i)\}$. $\qquad\square$

A.4. *Remark.* Asano [1] originally considered the situation where $K_1, K_2 = \{z : |z| \geq 1\}$ in Lemma A.1. In that case, the fact that

$$A + Bz_1 + Cz_2 + Dz_1z_2 \neq 0$$

if $z_1 \notin K_1, z_2 \notin K_2$, implies that the roots of $A + (B + C)z + Dz^2$ have absolute value ≥ 1, and the same is true of their product A/D, proving the lemma.

THEOREM A.5 (Grace's theorem). *Let P be a complex polynomial of degree n in one variable and $\Phi \in \mathcal{A}_n$ be the only polynomial symmetric in its n arguments such that*

$$\Phi(z, \ldots, z) = P(z).$$

If the n roots of P are contained in a closed circular region K and $z_1 \notin K, \ldots, z_k \notin K$, then $\Phi(z_1, \ldots, z_n) \neq 0$.

A closed circular region is a closed subset K of **C** bounded by a circle or a straight line. We allow the coefficients of $z^n, z^{n-1}, \ldots$ in P to vanish: we then say that some of the roots of P are at ∞, and we take K noncompact.

For a proof see Polya and Szegö [6, V, Ex. 145].   □

References

[1] T. ASANO, Theorems on the partition functions of the Heisenberg ferromagnets., *J. Phys. Soc. Japan* **29** (1970), 350–359. MR 43 #5848

[2] B. BEAUZAMY, On complex Lee and Yang polynomials, *Comm. Math. Phys.* **182** (1996), 177–184. MR 99e:82004 Zbl 0877.12002

[3] J. BORCEA and P. BRÄNDÉN, The Lee-Yang and Polya-Schur programs. I. Linear operators preserving stability, *Invent. Math.* **177** (2009), 541–569.

[4] _____ , The Lee-Yang and Polya-Schur programs. II. Theory of stable polynomials and applications, *Comm. Pure Appl. Math.* **62** (2009), 1595–1631.

[5] T. D. LEE and C. N. YANG, Statistical theory of equations of state and phase transitions. II. Lattice gas and Ising model, *Phys. Rev.* **87** (1952), 410–419. MR 14,711c Zbl 0048.43401

[6] G. PÓLYA and G. SZEGŐ, *Problems and Theorems in Analysis.* II, *Classics in Mathematics*, Springer-Verlag, New York, 1998. MR 1492448

[7] D. RUELLE, Extension of the Lee-Yang circle theorem, *Phys. Rev. Lett.* **26** (1971), 303–304. MR 45 #1527

[8] _____ , Grace-like polynomials, in *Foundations of Computational Mathematics. Proc. of Smale-fest* 2000 (Hong Kong, 2000), World Sci. Publ., River Edge, NJ, 2002, Cucker, F. and Rojas, J. M. (eds.), pp. viii+469. MR 2004g:65003

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