# The sharp quantitative isoperimetric inequality 

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#### Abstract

A quantitative sharp form of the classical isoperimetric inequality is proved, thus giving a positive answer to a conjecture by Hall.


## 1. Introduction

The classical isoperimetric inequality states that if $E$ is a Borel set in $\mathbb{R}^{n}$, $n \geq 2$, with finite Lebesgue measure $|E|$, then the ball with the same volume has a lower perimeter, or, equivalently, that

$$
\begin{equation*}
n \omega_{n}^{1 / n}|E|^{(n-1) / n} \leq P(E) \tag{1.1}
\end{equation*}
$$

Here $P(E)$ denotes the distributional perimeter of $E$ (which coincides with the classical ( $n-1$ )-dimensional measure of $\partial E$ when $E$ has a smooth boundary) and $\omega_{n}$ is the measure of the unit ball $B$ in $\mathbb{R}^{n}$. It is also well known that equality holds in (1.1) if and only if $E$ is a ball.

The history of the various proofs and different formulations of the isoperimetric inequality is definitely a very long and complex one. Therefore we shall not even attempt to sketch it here, but we refer the reader to the many review books and papers (e.g. [3], [18], [5], [21], [7], [13]) available on the subject and to the original paper by De Giorgi [8] (see [9] for the English translation) where (1.1) was proved for the first time in the general framework of sets of finite perimeter.

In this paper we prove a quantitative version of the isoperimetric inequality. Inequalities of this kind have been named by Osserman [19] Bonnesen type inequalities, following the results proved in the plane by Bonnesen in 1924 (see [4] and also [2]). More precisely, Osserman calls in this way any inequality of the form

$$
\lambda(E) \leq P(E)^{2}-4 \pi|E|,
$$

valid for smooth sets $E$ in the plane $\mathbb{R}^{2}$, where the quantity $\lambda(E)$ has the following three properties: (i) $\lambda(E)$ is nonnegative; (ii) $\lambda(E)$ vanishes only when $E$ is a ball; (iii) $\lambda(E)$ is a suitable measure of the "asymmetry" of $E$.

In particular, any Bonnesen inequality implies the isoperimetric inequality as well as the characterization of the equality case.

The study of Bonnesen type inequalities in higher dimension has been carried on in recent times in [12], [16], [15]. In order to describe these results let us introduce, for any Borel set $E$ in $\mathbb{R}^{n}$ with $0<|E|<\infty$, the isoperimetric deficit of $E$

$$
D(E):=\frac{P(E)}{n \omega_{n}^{1 / n}|E|^{(n-1) / n}}-1=\frac{P(E)-P(r B)}{P(r B)}
$$

where $r$ is the radius of the ball having the same volume as $E$, that is $|E|=$ $r^{n}|B|$.

The paper [12] by Fuglede deals with convex sets. Namely, he proves that if $E$ is a convex set having the same volume of the unit ball $B$ then

$$
\min \left\{\delta_{H}(E, x+B): x \in \mathbb{R}^{n}\right\} \leq C(n) D(E)^{\alpha(n)}
$$

where $\delta_{H}(\cdot, \cdot)$ denotes the Hausdorff distance between two sets and $\alpha(n)$ is a suitable exponent depending on the dimension $n$. This result is sharp, in the sense that in [12] examples are given showing that the exponent $\alpha(n)$ found in the paper cannot be improved (at least if $n \neq 3$ ).

When dealing with general nonconvex sets, we cannot expect the isoperimetric deficit to control the Hausdorff distance from $E$ to a ball. To see this it is enough to take, in any dimension, the union of a large ball and a far away tiny one or, if $n \geq 3$, a connected set obtained by adding to a ball an arbitrarily long (and suitably thin) "tentacle". It is then clear that in this case a natural notion of asymmetry is the so-called Fraenkel asymmetry of $E$, defined by

$$
\lambda(E):=\min \left\{\frac{d(E, x+r B)}{r^{n}}: x \in \mathbb{R}^{n}\right\}
$$

where $r>0$ is again such that $|E|=r^{n}|B|$ and $d(E, F)=|E \Delta F|$ denotes the measure of the symmetric difference between any two Borel sets $E, F$.

This kind of asymmetry has been considered by Hall, Hayman and Weitsman in [16] where it is proved that if $E$ is a smooth open set with a sufficiently small deficit $D(E)$, then there exists a suitable straight line such that, denoting by $E^{*}$ the Steiner symmetral of $E$ with respect to the line (see definition in Section 3), one has

$$
\begin{equation*}
\lambda(E) \leq C(n) \sqrt{\lambda\left(E^{*}\right)} \tag{1.2}
\end{equation*}
$$

Later on Hall proved in [15] that for any axially symmetric set $F$

$$
\begin{equation*}
\lambda(F) \leq C(n) \sqrt{D(F)} \tag{1.3}
\end{equation*}
$$

and thus, combining (1.3) (applied with $F=E^{*}$ ) with (1.2), he was able to conclude that

$$
\begin{equation*}
\lambda(E) \leq C(n) D\left(E^{*}\right)^{1 / 4} \leq C(n) D(E)^{1 / 4} \tag{1.4}
\end{equation*}
$$

where the last inequality is immediate when one recalls that Steiner symmetrization lowers the perimeter, hence the deficit. Though both estimates (1.2) and (1.3) are sharp, in the sense that one cannot replace the square root on the right-hand side by any better power, the exponent $1 / 4$ appearing in (1.4) does not seem to be optimal. And in fact Hall himself conjectured that the term $D(E)^{1 / 4}$ should be replaced by the smaller term $D(E)^{1 / 2}$. If so, the resulting inequality would be optimal, as one can easily check by taking an ellipsoid $E$ with $n-1$ semiaxes of length 1 and the last one of length larger than 1.

We give a positive answer to Hall's conjecture by proving the following estimate.

ThEOREM 1.1. Let $n \geq 2$. There exists a constant $C(n)$ such that for every Borel set $E$ in $\mathbb{R}^{n}$ with $0<|E|<\infty$

$$
\begin{equation*}
\lambda(E) \leq C(n) \sqrt{D(E)} \tag{1.5}
\end{equation*}
$$

A few remarks are in order. As we have already observed, the exponent $1 / 2$ in the above inequality is optimal and cannot be replaced by any bigger power. Notice also that both $\lambda(E)$ and $D(E)$ are scale invariant; therefore it is enough to prove (1.5) for sets of given measure. Thus, throughout the paper we shall assume that

$$
|E|=|B|
$$

Moreover, since $\lambda(E) \leq 2|B|$, it is clear that one needs to prove Theorem 1.1 only for sets with a small isoperimetric deficit. In fact, if $D(E) \geq \delta>0$, (1.5) is trivially satisfied by taking a suitably large constant $C(n)$. Finally, a more or less standard truncation argument (see Lemma 5.1) shows that in order to prove Theorem 1.1 it is enough to assume that $E$ is contained in a suitably large cube.

Let us now give a short description of how the proof goes. Apart from the isoperimetric property of the sphere, we do not use any sophisticated technical tool. On the contrary, the underlying idea is to reduce the problem, by means of suitable geometric constructions, to the case of more and more symmetric sets.

To be more precise, let us introduce the following definition, which will play an important role in the sequel. We say that a Borel set $E \subseteq \mathbb{R}^{n}$ is $n$-symmetric if $E$ is symmetric with respect to $n$ orthogonal hyperplanes $H_{1}, \ldots, H_{n}$. A simple, but important property of $n$-symmetric sets is that the Fraenkel asymmetry $\lambda(E)$ is equivalent to the distance from $E$ to the ball centered at the intersection $x_{0}$ of the $n$ hyperplanes $H_{i}$. In fact, we have (see Lemma 2.2)

$$
\begin{equation*}
\lambda(E) \leq d\left(E, x_{0}+B\right) \leq 2^{n} \lambda(E) \tag{1.6}
\end{equation*}
$$

Coming back to the proof, the first step is to pass from a general set $E$ to a set $E^{\prime}$ symmetric with respect to a hyperplane, without losing too much in terms of isoperimetric deficit and asymmetry, namely

$$
\begin{equation*}
\lambda(E) \leq C(n) \lambda\left(E^{\prime}\right) \quad \text { and } \quad D\left(E^{\prime}\right) \leq C(n) D(E) \tag{1.7}
\end{equation*}
$$

A natural way to do this, could be to take any hyperplane dividing $E$ in two parts of equal measure and then to reflect one of them. In fact, calling $E^{+}$and $E^{-}$the two resulting sets (see Figure 1.a), it is easily checked that

$$
D\left(E^{+}\right)+D\left(E^{-}\right) \leq 2 D(E),
$$

but, unfortunately, it is not true in general that

$$
\lambda(E) \leq C(n) \max \left\{\lambda\left(E^{+}\right), \lambda\left(E^{-}\right)\right\}
$$



Figure 1: The sets $E, E^{+}$and $E^{-}$
This is clear if we take, for instance, $E$ equal to the union of two slightly shifted half-balls, as in Figure 1.b. However, if we take, instead, two orthogonal hyperplanes, each one dividing $E$ in two parts of equal volume, at least one of the four sets thus obtained by reflection will satisfy (1.7) for a suitable constant $C(n)$ (see Lemma 2.5). Thus, iterating this procedure, we obtain a set with ( $n-1$ ) symmetries and eventually, using a variant of this argument to get the last symmetry, an $n$-symmetric set $E^{\prime}$ satisfying (1.7).

Once we have reduced the proof of Theorem 1.1 to the case of an $n$-symmetric set $E$, equivalently to a set symmetric with respect to all coordinate hyperplanes, all we have to do, thanks to (1.6), is to estimate $d(E, B)$ by $\sqrt{D(E)}\left(\right.$ as $\left.x_{0}=0\right)$.

To this aim we compare $E$ with its Steiner symmetral $E^{*}$ with respect to one of the coordinate axes, say $x_{1}$. Simplifying a bit, the idea is to estimate each one of the two terms appearing on the right-hand side of the triangular inequality

$$
\begin{equation*}
d(E, B) \leq d\left(E, E^{*}\right)+d\left(E^{*}, B\right) \tag{1.8}
\end{equation*}
$$

by the square root of the isoperimetric deficit. Concerning the first term, by Fubini's theorem we can write

$$
\begin{equation*}
d\left(E, E^{*}\right)=\int_{\mathbb{R}} \mathscr{H}^{n-1}\left(E_{t} \Delta E_{t}^{*}\right) d t \tag{1.9}
\end{equation*}
$$

where, for any set $F, F_{t}$ stands for $\left\{x \in F: x_{1}=t\right\}$. Since $E_{t}^{*}$ is the ( $n-1$ )-dimensional ball with the same measure of $E_{t}$, centered on the axis $x_{1}$, and since $E_{t}$ is symmetric with respect to the remaining $(n-1)$ coordinate hyperplanes, by applying (1.6) -in one dimension less- to the set $E_{t}$ suitably rescaled we get

$$
\begin{equation*}
\mathscr{H}^{n-1}\left(E_{t} \Delta E_{t}^{*}\right) \leq \frac{2^{n-1}}{\omega_{n-1}} \mathscr{H}^{n-1}\left(E_{t}\right) \lambda_{n-1}\left(E_{t}\right) \leq C \lambda_{n-1}\left(E_{t}\right) \tag{1.10}
\end{equation*}
$$

where $\lambda_{n-1}\left(E_{t}\right)$ denotes the Fraenkel asymmetry of $E_{t}$ in $\mathbb{R}^{n-1}$. Then, assuming that Theorem 1.1 holds in dimension $n-1$, we can estimate $\lambda_{n-1}\left(E_{t}\right)$ by the deficit $D_{n-1}\left(E_{t}\right)$ of $E_{t}$ in $\mathbb{R}^{n-1}$, thus getting from (1.9) and (1.10)

$$
d\left(E, E^{*}\right) \leq C \int_{\mathbb{R}} \lambda_{n-1}\left(E_{t}\right) d t \leq C \int_{\mathbb{R}} \sqrt{D_{n-1}\left(E_{t}\right)} d t
$$

Finally, by a suitable choice of the symmetrization axis $x_{i}$, we are able to prove that

$$
\int_{\mathbb{R}} \sqrt{D_{n-1}\left(E_{t}\right)} d t \leq C \sqrt{D(E)}
$$

thus concluding from (1.6) and (1.8) that

$$
\lambda(E) \leq d(E, B) \leq C \sqrt{D(E)}+d\left(E^{*}, B\right) \leq C \sqrt{D(E)}+2^{n} \lambda\left(E^{*}\right)
$$

At this point, in order to control $\lambda\left(E^{*}\right)$ by $\sqrt{D\left(E^{*}\right)}$, which in turn is smaller than $\sqrt{D(E)}$, we could have relied on Hall's inequality (1.3). However, we have preferred to do otherwise. In fact in our case, since we may assume that $E$ is $n$-symmetric (and thus $E^{*}$ is $n$-symmetric too) we can give a simpler, selfcontained proof, ultimately reducing the required estimate to the case of two overlapping balls with the same radii (see the proof of Theorem 4.1). And this particular case can be handled by elementary one-dimensional calculations.

The methods developed in this paper, besides giving a positive answer to the question posed by Hall, can also be used to obtain an optimal quantitative version of the Sobolev inequality. This application is contained in the forthcoming paper [14] by the same authors.

## 2. Reduction to $n$-symmetric sets

In this section, we aim to reduce ourselves to the case of a set with wide symmetry, namely an $n$-symmetric one. Since, as will shall see in Section 5, we may always reduce the proof of Theorem 1.1 to the case where $E$ is contained
in a suitably large cube $Q_{l}=(-l, l)^{n}$ and $|E|=|B|$, we shall work here and in the next two sections with uniformly bounded sets contained in

$$
X:=\left\{E \subseteq \mathbb{R}^{n}: E \text { is Borel, }|E|=|B|\right\}
$$

And thus we shall use the convention that $C=C(n, l)$ denotes a sufficiently large constant, that may change from line to line, and that depends uniquely on the dimension $n$ and on $l$.

The whole section is devoted to show the following result.
Theorem 2.1. For every $E \in X, E \subseteq Q_{l}$, there exists a set $F \in X$, $F \subseteq Q_{3 l}$, symmetric with respect to $n$ orthogonal hyperplanes and such that,

$$
\lambda(E) \leq C(n, l) \lambda(F), \quad D(F) \leq 2^{n} D(E)
$$

This section is divided into two subsections: in the first one, we collect some technical properties needed later, and in the second we prove Theorem 2.1.
2.1. Some technical facts. In this subsection we collect some technical facts to be used throughout the paper. Even though the ball centered in the center of symmetry is in general not optimal for an $n$-symmetric set, the next lemma states that this is true apart from a constant factor. In the sequel, for any two sets $E, F \subseteq \mathbb{R}^{n}$ we shall denote by $\lambda(E \mid F)$ the Fraenkel asymmetry relative to $F$, that is

$$
\lambda(E \mid F):=\min \left\{\frac{d(E, x+r B)}{r^{n}}: x \in F\right\},
$$

again being $|E|=r^{n}|B|$.
Lemma 2.2. Let $E \in X$ be a set symmetric with respect to $k$ orthogonal hyperplanes $H_{j}=\left\{x \in \mathbb{R}^{n}: x \cdot \nu_{j}=0\right\}$ for $1 \leq j \leq k$. Then one has

$$
\lambda\left(E \mid \bigcap_{j=1}^{k} H_{j}\right) \leq 2^{k} \lambda(E) .
$$

Proof. Let us set

$$
\begin{aligned}
& Q^{-}:=\left\{x \in \mathbb{R}^{n}: x \cdot \nu_{j} \leq 0 \quad \forall 1 \leq j \leq k\right\}, \\
& Q^{+}:=\left\{x \in \mathbb{R}^{n}: x \cdot \nu_{j} \geq 0 \quad \forall 1 \leq j \leq k\right\}
\end{aligned}
$$

By definition and by symmetry, $\lambda(E)=d(E, p+B)$ for some point $p \in \mathbb{R}^{n}$ belonging to $Q^{-}$; as an immediate consequence, denoting by $p_{0}$ the orthogonal projection of $p$ on $\bigcap_{j=1}^{k} H_{j}$, one has that $\left(p_{0}+B\right) \cap Q^{+} \supseteq(p+B) \cap Q^{+}$. Hence,

$$
\left(E \backslash\left(p_{0}+B\right)\right) \cap Q^{+} \subseteq(E \backslash(p+B)) \cap Q^{+}
$$

The conclusion follows just by noticing that, since both $p_{0}+B$ and $E$ are symmetric with respect to the hyperplanes $H_{j}$,

$$
\begin{aligned}
\lambda\left(E \mid \bigcap_{j=1}^{k} H_{j}\right) & \leq d\left(E,\left(p_{0}+B\right)\right)=2\left|E \backslash\left(p_{0}+B\right)\right| \\
& =2 \cdot 2^{k}\left|\left(E \backslash\left(p_{0}+B\right)\right) \cap Q^{+}\right| \\
& \leq 2 \cdot 2^{k}\left|(E \backslash(p+B)) \cap Q^{+}\right| \\
& \leq 2 \cdot 2^{k}|E \backslash(p+B)|=2^{k} d(E, p+B) \\
& =2^{k} \lambda(E)
\end{aligned}
$$

The second result we present shows the stability of the isoperimetric inequality, but without any estimate about the rate of convergence; keep in mind that the goal of this paper is exactly to give a precise and sharp quantitative estimate about this convergence. This weak result is easy and very well known; we present a proof only for the reader's convenience, and to keep this paper self-contained. The proof is based on a simple compactness argument.

Lemma 2.3. Let $l>0$. For any $\varepsilon>0$ there exists $\delta=\delta(n, l, \varepsilon)>0$ such that if $E \in X, E \subseteq Q_{l}$, and $D(E) \leq \delta$ then $\lambda(E) \leq \varepsilon$.

Proof. We argue by contradiction. If the assertion were not true, there would be a sequence $\left\{E_{j}\right\} \subseteq X$ with $E_{j} \subseteq Q_{l}, D\left(E_{j}\right) \rightarrow 0$ and $\lambda\left(E_{j}\right) \geq \varepsilon>0$ for all $j \in \mathbb{N}$. Since each set $E_{j}$ is contained in the same cube $Q_{l}$, thanks to a well-known embedding theorem (see for instance Theorem 3.39 in [1]) we can assume, up to a subsequence, that $\chi_{E_{j}} \xrightarrow{L^{1}} \chi_{E_{\infty}}$ for some set $E_{\infty}$ of finite perimeter; we deduce that $E_{\infty}$ is a set with $\left|E_{\infty}\right|=|B|$, and by the lower semicontinuity of the perimeters $P\left(E_{\infty}\right) \leq P(B)$, then $E_{\infty}$ is a ball. The fact that $\chi_{E_{j}}$ strongly converges in $L^{1}$ to $\chi_{E_{\infty}}$ immediately implies that $\left|E_{j} \Delta E_{\infty}\right| \rightarrow 0$, against the assumption $\lambda\left(E_{j}\right) \geq \varepsilon$. The contradiction concludes the proof.

The last result is an estimate about the distance of two sets obtained via translations of half-balls; the proof that we present was suggested by Sergio Conti.

Lemma 2.4. Let $H_{1}$ and $H_{2}$ be two orthogonal hyperplanes and let $H_{i}^{ \pm}$be the corresponding two pairs of half-spaces. Consider two points $x_{1}, \sigma_{1} \in H_{1}$, two points $x_{2}, \sigma_{2} \in H_{2}$ and the sets

$$
B_{i}:=x_{i}+B, \quad B_{i}^{ \pm}:=B_{i} \cap H_{i}^{ \pm}, \quad D_{i}:=B_{i}^{+} \cup\left(B_{i}^{-}+\sigma_{i}\right)
$$

There are two constants $\varepsilon=\varepsilon(n)$ and $C=C(n)$ such that, provided $\left|x_{1}-x_{2}\right|$ $\leq \varepsilon$ and $\left|\sigma_{1}\right|,\left|\sigma_{2}\right| \leq \varepsilon$, then

$$
\max \left\{\left|\sigma_{1}\right|,\left|\sigma_{2}\right|\right\} \leq C d\left(D_{1}, D_{2}\right)
$$

Proof. For suitable constants $\delta(n)$ and $C(n)$, given two unitary balls $F_{1}$ and $F_{2}$ with the centers lying at a distance $\delta \leq \delta(n)$, we have

$$
\delta \leq C(n) d\left(F_{1}, F_{2}\right) .
$$

In particular, up to changing $C(n)$, if $Q$ denotes any intersection of two orthogonal half-spaces of $\mathbb{R}^{n}$, with the property that

$$
\begin{equation*}
\min \left\{\left|F_{1} \cap Q\right|,\left|F_{2} \cap Q\right|\right\} \geq \frac{|B|}{8} \tag{2.1}
\end{equation*}
$$

then

$$
\delta \leq C(n) d\left(F_{1} \cap Q, F_{2} \cap Q\right)
$$

We now apply this statement twice to prove the lemma. In the first instance we choose $F_{1}=B_{1}, F_{2}=B_{2}$ and $Q=H_{1}^{+} \cap H_{2}^{+}$. Note that, provided $\varepsilon(n)$ is small enough, by construction, condition (2.1) is satisfied. Thus, if $2 \varepsilon(n) \leq \delta(n)$,

$$
d\left(D_{1}, D_{2}\right) \geq d\left(D_{1} \cap Q, D_{2} \cap Q\right)=d\left(B_{1} \cap Q, B_{2} \cap Q\right) \geq C^{-1}\left|x_{1}-x_{2}\right| .
$$

In the second instance we choose $F_{1}=\sigma_{1}+B_{1}, F_{2}=B_{2}$ and $Q=H_{1}^{-} \cap H_{2}^{+}$, and find similarly that

$$
\begin{aligned}
d\left(D_{1}, D_{2}\right) & \geq d\left(D_{1} \cap Q, D_{2} \cap Q\right) \\
& =d\left(\left(\sigma_{1}+B_{1}\right) \cap Q, B_{2} \cap Q\right) \geq C^{-1}\left|x_{1}+\sigma_{1}-x_{2}\right|
\end{aligned}
$$

Thus $\left|\sigma_{1}\right| \leq 2 C d\left(D_{1}, D_{2}\right)$, and by symmetry we have the analogous estimate on $\sigma_{2}$.
2.2. The proof of Theorem 2.1. We show first a technique to perform a single symmetrization, and the claim will then be proved by successive applications of this main step. We need also a bit of notation: given a set $E \in X$ and a unit vector $\nu \in \mathbb{S}^{n-1}$, we denote by $H_{\nu}^{+}=\left\{x \in \mathbb{R}^{n}: x \cdot \nu>t\right\}$ an open half-space orthogonal to $\nu$ where $t \in \mathbb{R}$ is chosen in such a way that

$$
\left|E \cap H_{\nu}^{+}\right|=\frac{|E|}{2} ;
$$

we also denote by $r_{\nu}: \mathbb{R}^{n} \rightarrow \mathbb{R}^{n}$ the reflection with respect to $H_{\nu}=\partial H_{\nu}^{+}$, and by $H_{\nu}^{-}=r_{\nu}\left(H_{\nu}^{+}\right)$the open half-space complementary to $H_{\nu}^{+}$. Finally, we write $E_{\nu}^{ \pm}=E \cap H_{\nu}^{ \pm}$.

Lemma 2.5. There exist two constants $C$ and $\delta$, depending only on $n$ and $l$ such that, given $E \in X, E \subseteq Q_{l}$, and two orthogonal vectors $\nu_{1}$ and $\nu_{2}$, there are $i \in\{1,2\}$ and $s \in\{+,-\}$ such that, setting $E^{\prime}=E_{\nu_{i}}^{s} \cup r_{\nu_{i}}\left(E_{\nu_{i}}^{s}\right)$, one has

$$
\begin{equation*}
\lambda(E) \leq C \lambda\left(E^{\prime}\right), \quad D\left(E^{\prime}\right) \leq 2 D(E) \tag{2.2}
\end{equation*}
$$

provided that $D(E) \leq \delta$.

Proof. First of all, given any unit vector $\nu$, let us denote by $B_{\nu}^{+}$a half-ball with center on $H_{\nu}$ that best approximates $E_{\nu}^{+}$, i.e. we set $B_{\nu}^{+}=(p+B) \cap H_{\nu}^{+}$ for some $p$ realizing

$$
\min \left\{d\left(E_{\nu}^{+},(p+B) \cap H_{\nu}^{+}\right): p \in H_{\nu}\right\} ;
$$

analogously, we let $B_{\nu}^{-}$be a half-ball with center in $H_{\nu}$ which best approximates $E_{\nu}^{-}$.

We consider now the sets

$$
F_{\nu}^{1,2}:=E_{\nu}^{ \pm} \cup r_{\nu}\left(E_{\nu}^{ \pm}\right), \quad T_{\nu}:=B_{\nu}^{+} \cup B_{\nu}^{-},
$$

(the first equation must be understood as $F_{\nu}^{1}=E_{\nu}^{+} \cup r_{\nu}\left(E_{\nu}^{+}\right)$and $F_{\nu}^{2}=E_{\nu}^{-} \cup$ $\left.r_{\nu}\left(E_{\nu}^{-}\right)\right)$. Note that $T_{\nu}$ is the union of two half-balls with (possibly different) centers on $H_{\nu}$, and that $F_{\nu}^{1,2}$, with $\nu=\nu_{1}$ or $\nu=\nu_{2}$, are the four sets among which we need to select $E^{\prime}$. Notice also that by a compactness argument similar to the one used in proving Lemma 2.3 it is clear that, if $\varepsilon>0$ is chosen as in Lemma 2.4, the centers of the four half-balls $B_{\nu_{i}}^{ \pm}$are at distance less than $\varepsilon$, provided that $D(E)$ is smaller than a suitable $\delta$.

The following two remarks will be useful. First we note that clearly, by construction of $B_{\nu}^{ \pm}$and by symmetry of $F_{\nu}^{1,2}$, we have

$$
\begin{equation*}
\lambda\left(F_{\nu}^{1,2} \mid H_{\nu}\right)=d\left(F_{\nu}^{1,2}, B_{\nu}^{ \pm} \cup r_{\nu}\left(B_{\nu}^{ \pm}\right)\right)=2 d\left(E_{\nu}^{ \pm}, B_{\nu}^{ \pm}\right) \tag{2.3}
\end{equation*}
$$

It can be easily checked that $P\left(F_{\nu}^{i} \mid H_{\nu}\right)=0$, being $P\left(F_{\nu}^{i} \mid H_{\nu}\right)$ the perimeter of $F_{\nu}^{i}$ relative to $H_{\nu}$ (see the appendix); hence

$$
P(E) \geq P\left(E \mid H_{\nu}^{+}\right)+P\left(E \mid H_{\nu}^{-}\right)=\frac{P\left(F_{\nu}^{1}\right)+P\left(F_{\nu}^{2}\right)}{2}
$$

so that we have always

$$
\begin{equation*}
\max \left\{D\left(F_{\nu}^{1}\right), D\left(F_{\nu}^{2}\right)\right\} \leq 2 D(E) \tag{2.4}
\end{equation*}
$$

As shown by (2.4), all the four sets among which we have to choose $E^{\prime}$ satisfy the estimate on the right in (2.2), so that we need only to take care of the one on the left.

Assume now for the moment that, for some constant $K=K(n)$ to be determined later and for some unit vector $\nu \in S^{n-1}$,

$$
\begin{equation*}
d\left(B_{\nu}^{-}, r_{\nu}\left(B_{\nu}^{+}\right)\right) \leq K\left(d\left(E_{\nu}^{+}, B_{\nu}^{+}\right)+d\left(E_{\nu}^{-}, B_{\nu}^{-}\right)\right) \tag{2.5}
\end{equation*}
$$

Then we can easily estimate, also recalling (2.3),

$$
\begin{aligned}
\lambda(E) & \leq d\left(E, B_{\nu}^{+} \cup r_{\nu}\left(B_{\nu}^{+}\right)\right)=d\left(E_{\nu}^{+}, B_{\nu}^{+}\right)+d\left(E_{\nu}^{-}, r_{\nu}\left(B_{\nu}^{+}\right)\right) \\
& \leq d\left(E_{\nu}^{+}, B_{\nu}^{+}\right)+d\left(E_{\nu}^{-}, B_{\nu}^{-}\right)+d\left(B_{\nu}^{-}, r_{\nu}\left(B_{\nu}^{+}\right)\right) \\
& \leq(K+1)\left(d\left(E_{\nu}^{+}, B_{\nu}^{+}\right)+d\left(E_{\nu}^{-}, B_{\nu}^{-}\right)\right) \\
& =\frac{K+1}{2}\left(\lambda\left(F_{\nu}^{1} \mid H_{\nu}\right)+\lambda\left(F_{\nu}^{2} \mid H_{\nu}\right)\right) .
\end{aligned}
$$

Therefore, up to swapping $F_{\nu}^{1}$ and $F_{\nu}^{2}$, we have that

$$
\lambda\left(F_{\nu}^{1} \mid H_{\nu}\right) \geq \frac{1}{K+1} \lambda(E)
$$

Since $F_{\nu}^{1}$ is symmetric with respect to the hyperplane $H_{\nu}$, by Lemma 2.2 we conclude

$$
\lambda\left(F_{\nu}^{1}\right) \geq \frac{1}{2(K+1)} \lambda(E) .
$$

Keeping in mind (2.4), the proof of this lemma will be concluded once we show that (2.5) holds either with $\nu=\nu_{1}$ or with $\nu=\nu_{2}$.

Suppose this is not true; then, define $\sigma_{1}$ as the vector connecting the centers of $B_{\nu_{1}}^{+}$and $B_{\nu_{1}}^{-}$. Since $\left|\sigma_{1}\right|<\varepsilon$

$$
d\left(B_{\nu_{1}}^{-}, r_{\nu_{1}}\left(B_{\nu_{1}}^{+}\right)\right) \leq C(n)\left|\sigma_{1}\right| .
$$

Therefore, the assumption that (2.5) does not hold with $\nu=\nu_{1}$ implies that

$$
d\left(E, T_{\nu_{1}}\right)=d\left(E_{\nu_{1}}^{+}, B_{\nu_{1}}^{+}\right)+d\left(E_{\nu_{1}}^{-}, B_{\nu_{1}}^{-}\right) \leq \frac{1}{K} d\left(B_{\nu_{1}}^{-}, r_{\nu_{1}}\left(B_{\nu_{1}}^{+}\right)\right) \leq \frac{C(n)}{K}\left|\sigma_{1}\right| .
$$

Analogously, assuming that (2.5) does not hold with $\nu=\nu_{2}$ yields

$$
d\left(E, T_{\nu_{2}}\right) \leq \frac{C(n)}{K}\left|\sigma_{2}\right| .
$$

We deduce, by the triangular inequality, that

$$
d\left(T_{\nu_{1}}, T_{\nu_{2}}\right) \leq \frac{C(n)}{K}\left(\left|\sigma_{1}\right|+\left|\sigma_{2}\right|\right) .
$$

By Lemma 2.4, this leads to a contradiction provided the constant $K$ is chosen sufficiently large; and, as already noticed, this contradiction completes the proof.

We can now show the main result of this section.
Proof of Theorem 2.1. Let us assume for the moment that $D(E)<$ $\delta / 2^{n-2}$, where $\delta$ is the constant appearing in Lemma 2.5. Let us take the standard orthonormal basis $\left\{e_{i}\right\}_{i=1}^{n}$ : we will prove the existence of a set $F$ of volume $|F|=|E|$, symmetric with respect to $n$ hyperplanes $H_{1}, H_{2}, \ldots, H_{n}$ such that each hyperplane $H_{i}$ is orthogonal to $e_{i}$, and with the property that

$$
\lambda(E) \leq C \lambda(F), \quad D(F) \leq 2^{n} D(E)
$$

We start with the versors $e_{1}$ and $e_{2}$ : thanks to Lemma 2.5, up to a permutation of $e_{1}$ and $e_{2}$ we can find a hyperplane $H_{1}$ orthogonal to $e_{1}$ and a set $F_{1}$ with $\left|F_{1}\right|=|E|$, symmetric with respect to $H_{1}$ and with the property that

$$
\lambda(E) \leq C \lambda\left(F_{1}\right), \quad D\left(F_{1}\right) \leq 2 D(E)
$$

Consider now the versors $e_{2}$ and $e_{3}$, and apply Lemma 2.5 to the set $F_{1}$; up to a permutation, we find a hyperplane $H_{2}$ orthogonal to $e_{2}$ and a set $F_{2}$ symmetric with respect to $\mathrm{H}_{2}$ with the property that

$$
\lambda(E) \leq C \lambda\left(F_{1}\right) \leq C^{2} \lambda\left(F_{2}\right), \quad D\left(F_{2}\right) \leq 2 D\left(F_{1}\right) \leq 4 D(E)
$$

Moreover, by Lemma 2.5 and by the fact that the hyperplanes $H_{1}$ and $H_{2}$ are orthogonal, the set $F_{2}$ is symmetric also with respect to $H_{1}$. By an immediate iteration, we arrive at a set $F_{n-1}$, symmetric with respect to $n-1$ orthogonal hyperplanes $H_{1}, H_{2}, \ldots, H_{n-1}$ and such that

$$
\begin{equation*}
\lambda(E) \leq C^{n-1} \lambda\left(F_{n-1}\right), \quad D\left(F_{n-1}\right) \leq 2^{n-1} D(E) \tag{2.6}
\end{equation*}
$$

To find the last hyperplane of symmetry we need a different argument, since we no longer have two different hyperplanes among which to choose. We then let $H_{n}$ be a hyperplane orthogonal to $e_{n}$ such that $\left|F_{n-1} \cap H_{n}^{+}\right|=$ $\left|F_{n-1} \cap H_{n}^{-}\right|$, being $H_{n}^{ \pm}$the two half-spaces corresponding to $H_{n}$, and we define $q=\cap_{i=1}^{n} H_{i} ; q$ is clearly a point since the $H_{i}$ 's are $n$ orthogonal hyperplanes. Defining now
$F_{n}^{1}:=\left(F_{n-1} \cap H_{n}^{+}\right) \cup r_{\nu}\left(F_{n-1} \cap H_{n}^{+}\right), \quad F_{n}^{2}:=\left(F_{n-1} \cap H_{n}^{-}\right) \cup r_{\nu}\left(F_{n-1} \cap H_{n}^{-}\right)$,
with $\nu=e_{n}$, we first notice that, with the same argument used to obtain (2.4),

$$
\begin{equation*}
\max \left\{D\left(F_{n}^{1}\right), D\left(F_{n}^{2}\right)\right\} \leq 2 D\left(F_{n-1}\right) . \tag{2.7}
\end{equation*}
$$

Moreover, by definition we have

$$
d\left(F_{n-1}, q+B\right)=\frac{d\left(F_{n}^{1}, q+B\right)+d\left(F_{n}^{2}, q+B\right)}{2} ;
$$

then applying Lemma 2.2 to $F_{n}^{1}$ and $F_{n}^{2}$-which are symmetric with respect to the $n$ orthogonal hyperplanes $H_{i^{-}}$we deduce

$$
\begin{align*}
\lambda\left(F_{n-1}\right) & \leq d\left(F_{n-1}, q+B\right)=\frac{d\left(F_{n}^{1}, q+B\right)+d\left(F_{n}^{2}, q+B\right)}{2}  \tag{2.8}\\
& =\frac{\lambda\left(F_{n}^{1} \mid \cap_{i=1}^{n} H_{i}\right)+\lambda\left(F_{n}^{2} \mid \cap_{i=1}^{n} H_{i}\right)}{2} \leq 2^{n-1}\left(\lambda\left(F_{n}^{1}\right)+\lambda\left(F_{n}^{2}\right)\right) \\
& \leq 2^{n} \max \left\{\lambda\left(F_{n}^{1}\right), \lambda\left(F_{n}^{2}\right)\right\} .
\end{align*}
$$

Putting together (2.6), (2.7) and (2.8), we find $F \in\left\{F_{n}^{1}, F_{n}^{2}\right\}$ such that

$$
\lambda(E) \leq C^{n-1} \lambda\left(F_{n-1}\right) \leq 2^{n} C^{n-1} \lambda(F), \quad D(F) \leq 2 D\left(F_{n-1}\right) \leq 2^{n} D(E)
$$

and conclude the proof under the assumption $D(E) \leq \delta / 2^{n-2}$, since the inclusion $F \subseteq Q_{3 l}$ is obvious by construction.

Finally, to conclude also when $D(E)>\delta / 2^{n-2}$, it is enough to take as $F$, independently from $E$, any $n$-symmetric set contained in $Q_{3 l}$ such that $0<D(F) \leq 4 \delta$, and possibly to modify $C(n, l)$ so that $C(n, l) \lambda(F) \geq 2 \omega_{n}$.

Remark 2.6. A quick inspection to the proofs of Lemma 2.5 and Theorem 2.1 shows that if we assume that the set $E$ satisfies the following condition

$$
\begin{equation*}
\mathscr{H}^{n-1}\left(\left\{x \in \partial^{*} E: \nu^{E}(x)= \pm e_{i}\right\}\right)=0 \quad \forall i=1, \ldots, n, \tag{2.9}
\end{equation*}
$$

then the same property is inherited by the set $F$ constructed above (here $\partial^{*} E$ denotes the reduced boundary of $E$; see the appendix). To see this, it is enough to check that if we take a set $E$ satisfying condition (2.9), split it in two parts by a hyperplane $H$ orthogonal with respect to one of the $e_{i}$ 's and reflect it with respect to $H$ (these are the only operations we performed in the previous proofs), then we obtain another set $E^{\prime}$ still satisfying (2.9).

## 3. Reduction to axially symmetric sets

In this section we show how to reduce the $n$-symmetric case to the axially symmetric one. The goal is to show Theorem 3.1 below, which will be the starting point for the induction argument over the dimension $n$ used in Section 5 to prove Theorem 1.1.

In the following, given any set $E$ of finite perimeter, for any $t \in \mathbb{R}$, we denote by $E_{t}$ the $(n-1)$-dimensional section $\left\{x \in E: x_{1}=t\right\}$ and by $E^{*}$ the Steiner symmetrization of $E$ with respect to the axis $e_{1}$, that is, the set $E^{*} \subseteq \mathbb{R}^{n}$ such that for any $t \in \mathbb{R}$ the section $E_{t}^{*}$ is the $(n-1)$-dimensional ball centered at $(t, 0,0, \ldots, 0)$ with $\mathscr{H}^{n-1}\left(E_{t}^{*}\right)=\mathscr{H}^{n-1}\left(E_{t}\right)$. We also set

$$
v_{E}(t):=\mathscr{H}^{n-1}\left(E_{t}\right),
$$

the $(n-1)$-dimensional measure of the section $E_{t}$, and denote by $p_{E}(t)$ the perimeter of $E_{t}$ in $\mathbb{R}^{n-1}$, i.e., for $\mathscr{H}^{1}$-a.e. $t$,

$$
p_{E}(t):=\mathscr{H}^{n-2}\left(\partial^{*} E_{t}\right) .
$$

Theorem 3.1. Let $E \in X, E \subseteq Q_{l}$ be symmetric with respect to the $n$ coordinate hyperplanes and assume that (2.9) holds. If $n=2$, or $n \geq 3$ and Theorem 1.1 holds in dimension $n-1$, then, up to a rotation of the coordinate axes,

$$
d(E, B) \leq 4 d\left(E^{*} \cap Z, B \cap Z\right)+C(n, l) \sqrt{D(E)},
$$

where $Z=\left\{x \in \mathbb{R}^{n}:\left|x_{1}\right| \leq \sqrt{2} / 2\right\}$.
Our first step will be to select one of the hyperplanes which will have a particular role in the following construction; we need to ensure that the symmetric difference between $E$ and $B$ is not, roughly speaking, too concentrated close to the "poles" of $B$ (i.e., the regions of $B$ having greatest distance from the hyperplane).

Lemma 3.2. Up to a rotation one can assume

$$
d(E, B) \leq 4 d(E \cap Z, B \cap Z)
$$

Proof. Define the sets

$$
Z_{1}:=\left\{x \in \mathbb{R}^{n}:\left|x_{1}\right| \leq \sqrt{2} / 2\right\}, \quad Z_{2}:=\left\{x \in \mathbb{R}^{n}:\left|x_{2}\right| \leq \sqrt{2} / 2\right\}
$$

so that $B \subseteq Z_{1} \cup Z_{2}$; as a consequence,

$$
B \backslash E=\left((B \backslash E) \cap Z_{1}\right) \cup\left((B \backslash E) \cap Z_{2}\right)
$$

Therefore, up to interchanging the axis $e_{1}$ with the axis $e_{2}$, we may assume that

$$
\left|(B \backslash E) \cap Z_{1}\right| \geq \frac{|B \backslash E|}{2}
$$

and finally conclude

$$
d(E, B)=2|B \backslash E| \leq 4\left|(B \backslash E) \cap Z_{1}\right| \leq 4 d\left(E \cap Z_{1}, B \cap Z_{1}\right)
$$

It is well known that the Steiner symmetrization lowers the perimeter (see e.g. [17]); that is,

$$
P\left(E^{*}\right) \leq P(E)
$$

In turn, the above inequality can be deduced also from the following estimate for the perimeters. This is immediate if $E$ is bounded and condition (3.1) holds, and follows by a simple approximation argument in the general case.

Lemma 3.3. Let $E \in X$ be a set of finite perimeter such that

$$
\begin{equation*}
\mathscr{H}^{n-1}\left(\left\{x \in \partial^{*} E: \nu^{E}(x)= \pm e_{1}\right\}\right)=0 \tag{3.1}
\end{equation*}
$$

Then $v_{E} \in W^{1,1}(\mathbb{R})$ and

$$
\begin{equation*}
P(E) \geq \int_{-\infty}^{+\infty} \sqrt{p_{E}(t)^{2}+v_{E}^{\prime}(t)^{2}} d t \tag{3.2}
\end{equation*}
$$

Moreover, for an axially symmetric set $E$ the preceding inequality is in fact an equality, and can be written as

$$
\begin{equation*}
P(E)=\int_{-\infty}^{+\infty} \sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+v_{E}^{\prime}(t)^{2}} d t \tag{3.3}
\end{equation*}
$$

where $\tau=(n-1)^{2} \omega_{n-1}^{\frac{2}{n-1}}$.

Proof. Since $E$ satisfies (3.1), from the co-area formula (6.1) we have that

$$
\begin{aligned}
P(E)=\mathscr{H}^{n-1}\left(\partial^{*} E\right) & =\int_{\partial^{*} E} \frac{\sqrt{1-\left|\nu_{1}^{E}\right|^{2}}}{\sqrt{1-\left|\nu_{1}^{E}\right|^{2}}} d \mathscr{H}^{n-1} \\
& =\int_{-\infty}^{+\infty} \int_{\left(\partial^{*} E\right)_{t}} \frac{1}{\sqrt{1-\left|\nu_{1}^{E}\right|^{2}}} d \mathscr{H}^{n-2} d t .
\end{aligned}
$$

Let us now recall that, by Theorem 6.2, $E_{t}$ is a set of finite perimeter for $\mathscr{H}^{1}$-a.e. $t$ and $\partial^{*} E_{t}$ equals $\left(\partial^{*} E\right)_{t}$ up to a $\mathscr{H}^{n-2}$-negligible set. Thus, from the equality above we get, using also Jensen's inequality,

$$
\begin{aligned}
P(E) & =\int_{-\infty}^{+\infty} \int_{\partial^{*} E_{t}} \frac{1}{\sqrt{1-\left|\nu_{1}^{E}\right|^{2}}} d \mathscr{H}^{n-2} d t \\
& =\int_{-\infty}^{+\infty} p_{E}(t) f_{\partial^{*} E_{t}} \sqrt{1+\frac{\left|\nu_{1}^{E}\right|^{2}}{1-\left|\nu_{1}^{E}\right|^{2}}} d \mathscr{H}^{n-2} d t \\
& \geq \int_{-\infty}^{+\infty} p_{E}(t) \sqrt{1+\left(f_{\partial^{*} E_{t}} \frac{\left|\nu_{1}^{E}\right|}{\sqrt{1-\left|\nu_{1}^{E}\right|^{2}}} d \mathscr{H}^{n-2}\right)^{2}} d t .
\end{aligned}
$$

Then, (3.2) immediately follows from the expression of $v_{E}^{\prime}$ given in (6.4). Finally, if $E$ is axially symmetric, then $\nu_{1}^{E}$ is clearly constant on each boundary $\partial^{*} E_{t}$. Hence the inequality above is indeed an equality and (3.3) follows.

Notice that by Theorem 6.3, if $E$ satisfies (3.1), then the same condition holds for $E^{*}$.

The second main ingredient that we give now is an $L^{\infty}$ estimate for $v_{E}-v_{B}$.
Lemma 3.4. For any $\rho>0$ there exists $\delta>0$ such that, if $E$ is as in Theorem 3.1 and $D(E) \leq \delta$, then

$$
\left\|v_{E}-v_{B}\right\|_{L^{\infty}} \leq \rho .
$$

Proof. Let us fix $\rho>0$. By Lemma 2.2, since $v_{E}=v_{E^{*}}$,

$$
\begin{equation*}
\left\|v_{E}-v_{B}\right\|_{L^{1}}=d\left(E^{*}, B\right) \leq 2^{n} \lambda\left(E^{*}\right) \tag{3.4}
\end{equation*}
$$

Therefore, given $\varepsilon>0$ (to be chosen later), by Lemma 2.3, if $D\left(E^{*}\right) \leq D(E) \leq \delta$, for $\delta$ small, then $\lambda\left(E^{*}\right) \leq \rho \varepsilon / 2^{n+2}$. Thus, from (3.4) we get

$$
\begin{equation*}
\mathscr{H}^{1}\left(\left\{t:\left|v_{E}-v_{B}\right|>\rho / 4\right\}\right)<\varepsilon . \tag{3.5}
\end{equation*}
$$

Recall that by assumption (2.9) and by Theorem 6.1, $v_{E}$ is continuous. Thus, if $\left\|v_{E}-v_{B}\right\|_{L^{\infty}}>\rho$, there exists $\bar{t}$ such that

$$
\left|v_{E}(\bar{t})-v_{B}(\bar{t})\right|>\rho .
$$

By the uniform continuity of $v_{B}$, provided $\varepsilon$ is sufficiently small, for $|t-\bar{t}|<\varepsilon$ one has $\left|v_{B}(t)-v_{B}(\bar{t})\right|<\rho / 4$. Then, by (3.5), and the continuity of $v_{E}$, there exist

$$
\bar{t}-\varepsilon<t^{-}<\bar{t}<t^{+}<\bar{t}+\varepsilon
$$

such that

$$
v_{E}\left(t^{-}\right)=v_{E}\left(t^{+}\right), \quad\left|v_{E}\left(t^{ \pm}\right)-v_{B}(\bar{t})\right|=\frac{\rho}{2} .
$$

Let us now define an axially symmetric set $F$ by letting $v_{F}(t)=v_{E}(t)$ if $t \notin\left(t^{-}, t^{+}\right)$and $v_{F}(t)=v_{E}\left(t^{-}\right)=v_{E}\left(t^{+}\right)$otherwise. Clearly

$$
\begin{equation*}
|F| \geq\left|E^{*}\right|-2^{n} l^{n-1} \varepsilon \tag{3.6}
\end{equation*}
$$

and

$$
\begin{equation*}
P(F)=P\left(E^{*}\right)+P\left(F \mid\left\{x: t^{-}<x_{1}<t^{+}\right\}\right)-P\left(E^{*} \mid\left\{x: t^{-}<x_{1}<t^{+}\right\}\right) . \tag{3.7}
\end{equation*}
$$

Moreover, by (3.3),

$$
\begin{align*}
& P\left(F \mid\left\{x: t^{-}<x_{1}<t^{+}\right\}\right)=\left(t^{+}-t^{-}\right) \sqrt{\tau} v_{E}\left(t^{+}\right)^{\frac{n-2}{n-1}} \leq C(n, l) \varepsilon,  \tag{3.8}\\
& P\left(E^{*} \mid\left\{x: t^{-}<x_{1}<t^{+}\right\}\right) \geq \int_{t^{-}}^{t^{+}}\left|v_{E}^{\prime}(t)\right| d t \geq \rho
\end{align*}
$$

From (3.7) and (3.8) we get

$$
P(F) \leq P\left(E^{*}\right)+C \varepsilon-\rho \leq P(E)+C \varepsilon-\rho .
$$

On the other hand, by the isoperimetric inequality and (3.6), if $\varepsilon$ is small enough, recalling that $\left|E^{*}\right|=|B|$, we have

$$
P(F) \geq n \omega_{n}^{1 / n}\left(\left|E^{*}\right|-2^{n} l^{n-1} \varepsilon\right)^{\frac{n-1}{n}} \geq P(B)-C(n, l) \varepsilon
$$

Therefore,

$$
D(E) \geq \frac{\rho-C \varepsilon}{P(B)}
$$

Hence, the assertion follows by choosing $\varepsilon<\rho / 2 C$ and $\delta=\rho / 2 P(B)$.
We can finally give the proof of Theorem 3.1.

Proof of Theorem 3.1. Since $d(E, B) \leq 2|B|$, by choosing $C(n, l)$ sufficiently large, we may assume $D(E)$ as small as we wish.

Thanks to Lemma 3.2, we can estimate

$$
\begin{align*}
d(E, B) & \leq 4 d(E \cap Z, B \cap Z)=4 \int_{-\sqrt{2} / 2}^{\sqrt{2} / 2} \mathscr{H}^{n-1}\left(E_{t} \Delta B_{t}\right) d t  \tag{3.9}\\
& =4 \int_{I} \mathscr{H}^{n-1}\left(E_{t} \Delta B_{t}\right) d t+4 \int_{J} \mathscr{H}^{n-1}\left(E_{t} \Delta B_{t}\right) d t
\end{align*}
$$

where we divide the interval $[-\sqrt{2} / 2, \sqrt{2} / 2]$ in two subsets $I$ and $J$ defined as

$$
\begin{aligned}
& I:=\left\{t \in\left[-\frac{\sqrt{2}}{2}, \frac{\sqrt{2}}{2}\right]:\left|v_{E}^{\prime}(t)\right| \leq M\right\} \\
& J:=\left\{t \in\left[-\frac{\sqrt{2}}{2}, \frac{\sqrt{2}}{2}\right]:\left|v_{E}^{\prime}(t)\right|>M\right\}
\end{aligned}
$$

for a constant $M$ to be determined later.
We will consider separately the situation in the sets $I$ and $J$. Let us then start working on $I$. By the triangular inequality we have immediately

$$
\begin{equation*}
\int_{I} \mathscr{H}^{n-1}\left(E_{t} \Delta B_{t}\right) d t \leq \int_{I} \mathscr{H}^{n-1}\left(E_{t}^{*} \Delta B_{t}\right) d t+\int_{I} \mathscr{H}^{n-1}\left(E_{t} \Delta E_{t}^{*}\right) d t . \tag{3.10}
\end{equation*}
$$

Concerning $\mathscr{H}^{n-1}\left(E_{t}^{*} \Delta B_{t}\right)$ we have easily

$$
\begin{equation*}
\int_{I} \mathscr{H}^{n-1}\left(E_{t}^{*} \Delta B_{t}\right) d t \leq \int_{-\sqrt{2} / 2}^{\sqrt{2} / 2} \mathscr{H}^{n-1}\left(E_{t}^{*} \Delta B_{t}\right) d t=d\left(E^{*} \cap Z, B \cap Z\right) \tag{3.11}
\end{equation*}
$$

On the other hand, since $E_{t}^{*}$ is an ( $n-1$ )-dimensional ball of ( $n-1$ )-dimensional volume $v_{E}(t)$, its perimeter is

$$
p_{E^{*}}(t)=\mathscr{H}^{n-2}\left(\partial E_{t}^{*}\right)=\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}} .
$$

Of course, $p_{E}(t) \geq p_{E^{*}}(t)$ for every $t$. Thus we can give the following definition, which will be very useful in the sequel,

$$
d(t):=p_{E}(t)^{2}-p_{E^{*}}(t)^{2}=p_{E}(t)^{2}-\tau v_{E}(t)^{\frac{2 n-4}{n-1}} \geq 0
$$

We now claim the inequality

$$
\mathscr{H}^{n-1}\left(E_{t} \Delta E_{t}^{*}\right) \leq C(n, l) \sqrt{p_{E}(t)-\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}}}
$$

If $n \geq 3$, this can be proved recalling again Lemma 2.2 and the symmetries of $E_{t}$, and applying Theorem 1.1 in dimension $n-1$ to $E_{t}$ (whose $\mathscr{H}^{n-1}$ measure is bounded by $2^{n-1} l^{n-1}$ ). On the other hand, for $n=2$ the inequality is true since
when $\mathscr{H}^{1}\left(E_{t} \Delta E_{t}^{*}\right)>0$, we have $\mathscr{H}^{1}\left(E_{t} \Delta E_{t}^{*}\right) \leq 2 l$ and $\sqrt{p_{E}(t)-p_{E^{*}}(t)} \geq 1$. Therefore,

$$
\begin{align*}
\int_{I} \mathscr{H}^{n-1}\left(E_{t} \Delta E_{t}^{*}\right) & \leq C \int_{I} \sqrt{p_{E}(t)-\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}}}  \tag{3.12}\\
& =C \int_{I} \sqrt{\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+d(t)}-\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}}} \\
& =C \int_{I} \sqrt{\frac{d(t)}{\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+d(t)}+\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}}}} \\
& \leq C\left(\int_{I} \frac{d(t)}{\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+d(t)}+\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}}}\right)^{1 / 2}
\end{align*}
$$

where the last inequality comes from Hölder inequality when we recall that the length of $I$ is of course a priori bounded by $\sqrt{2}$.

Let us estimate now the difference $P(E)-P\left(E^{*}\right)$. By applying formula (3.3) to the axially symmetric set $E^{*}$, we have

$$
\begin{align*}
P(E)-P\left(E^{*}\right) & \geq \int_{-\infty}^{+\infty} \sqrt{p_{E}(t)^{2}+v_{E}^{\prime}(t)^{2}}-\sqrt{p_{E^{*}}(t)^{2}+v_{E}^{\prime}(t)^{2}}  \tag{3.13}\\
& \geq \int_{I} \sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+v_{E}^{\prime}(t)^{2}+d(t)}-\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+v_{E}^{\prime}(t)^{2}} \\
& =\int_{I} \frac{d(t)}{\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+v_{E}^{\prime}(t)^{2}+d(t)}+\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+v_{E}^{\prime}(t)^{2}}} .
\end{align*}
$$

We use now the following fact: there exists a constant $C$, depending only on $n, l, M$, such that for any $t \in I$ one has

$$
\begin{equation*}
\frac{\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+v_{E}^{\prime}(t)^{2}+d(t)}+\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+v_{E}^{\prime}(t)^{2}}}{\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+d(t)}+\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}}} \leq C(n, l, M) . \tag{3.14}
\end{equation*}
$$

To show this estimate, we first remark that $v_{E}(t)$ is bounded from below by a strictly positive constant inside $[-\sqrt{2} / 2, \sqrt{2} / 2] \supseteq I$ (indeed, $v_{B}(t) \geq$ $\omega_{n-1} / 2^{(n-1) / 2}$ inside $[-\sqrt{2} / 2, \sqrt{2} / 2]$, so we apply Lemma 3.4). Then, we recall that by definition $\left|v_{E}^{\prime}\right| \leq M$ inside $I$, and this immediately concludes (3.14). Finally, putting together the estimates (3.12) and (3.13) and making use of (3.14)
we obtain

$$
\begin{align*}
& \int_{I} \mathscr{H}^{n-1}\left(E_{t} \Delta E_{t}^{*}\right)  \tag{3.15}\\
& \leq C(n, l, M)\left(\int_{I} \frac{d(t)}{\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+d(t)}+\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}}}\right)^{1 / 2} \\
& \leq C\left(\int_{I} \frac{d(t)}{\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+v_{E}^{\prime}(t)^{2}+d(t)}+\sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+v_{E}^{\prime}(t)^{2}}}\right)^{1 / 2} \\
& \quad \leq C \sqrt{P(E)-P\left(E^{*}\right)} \leq C(n, l, M) \sqrt{D(E)}
\end{align*}
$$

where the last inequality (up to a redefinition of the constant $C$ ) is true since $P\left(E^{*}\right) \geq P(B)$. Now, putting together (3.9), (3.10), (3.11) and (3.15), we finally obtain

$$
\begin{equation*}
d(E, B) \leq 4 d\left(E^{*} \cap Z, B \cap Z\right)+C(n, l, M) \sqrt{D(E)}+4 \int_{J} \mathscr{H}^{n-1}\left(E_{t} \Delta B_{t}\right) d t \tag{3.16}
\end{equation*}
$$

Let us now consider the situation in the set $J$. Notice that the above argument requires an a priori upper bound on $\left|v_{E}^{\prime}\right|$. On the other hand, it is clear that a region where $\left|v_{E}^{\prime}\right|$ is very large is far from being optimal from the point of view of the perimeter. Therefore, it is not surprising that an even stronger estimate than (3.15) holds in $J$, namely

$$
\begin{equation*}
\int_{J} \mathscr{H}^{n-1}\left(E_{t} \Delta B_{t}\right) d t \leq C D\left(E^{*}\right) \tag{3.17}
\end{equation*}
$$

Notice that this inequality together with (3.16) concludes the proof of the theorem.

To prove (3.17) observe that, since

$$
\mathscr{H}^{n-1}\left(E_{t} \Delta B_{t}\right) \leq \mathscr{H}^{n-1}\left(E_{t}\right)+\mathscr{H}^{n-1}\left(B_{t}\right)=v_{E}(t)+v_{B}(t)
$$

and $E \subseteq Q_{l}$, we have

$$
\int_{J} \mathscr{H}^{n-1}\left(E_{t} \Delta B_{t}\right) d t \leq C \mathscr{H}^{1}(J)
$$

Thus, (3.17) will be established once we prove that there exists a constant $K$ depending only on $n$ such that

$$
\begin{equation*}
\mathscr{H}^{1}(J) \leq \frac{K}{M} D\left(E^{*}\right) . \tag{3.18}
\end{equation*}
$$

To this aim, let us first assume that $v_{E}^{\prime}$ is continuous and $J=(a, a+\varepsilon)$ is an interval. As in the proof of Lemma 3.4, we introduce a new set $F$, not


Figure 2: Construction of $F$
necessarily of the same volume as $E^{*}$ and $B$, and prove (3.18) by applying the isoperimetric inequality $P(F) \geq P(\widetilde{B})$, where $\widetilde{B}$ is the ball having the same volume as $F$. As in Figure 2, let $h$ be the jump of $v_{E}$ inside $J$; we introduce a set $F$, axially symmetric with respect to the axis $e_{1}$, by defining $v_{F}$. To this aim, we consider the interval $\widetilde{J}=(a, a+h / N)$, where $N$ is an integer to be determined later depending only on the dimension $n$ of the ambient space. We set

$$
v_{F}(t):= \begin{cases}v_{E}(t) & \text { if } t \leq a \\ v_{E}(a) \pm N(t-a) & \text { if } a \leq t \leq a+h / N \\ v_{E}(t+\varepsilon-h / N) & \text { if } t \geq a+h / N\end{cases}
$$

In the above definition, the sign in the second row coincides with the sign of $v_{E}^{\prime}$ inside $J:$ since $\left|v_{E}^{\prime}\right|>M$ in $J$, by continuity this sign is well-defined. For instance, in the situation of the figure this sign is negative. Notice that, by definition, $v_{F}$ is a continuous function. Let us start by recalling the information we have so-far:

$$
\begin{equation*}
\mathscr{H}^{1}(J)=\varepsilon, \quad h \geq M \varepsilon, \quad \mathscr{H}^{1}(\widetilde{J})=\frac{h}{N} \tag{3.19}
\end{equation*}
$$

In order to evaluate both $|F|$ and $P(F)$, we set $J_{*}=\left\{x \in \mathbb{R}^{n}: x_{1} \in J\right\}$ and $\widetilde{J}_{*}=\left\{x \in \mathbb{R}^{n}: x_{1} \in \widetilde{J}\right\}$. Notice that $\omega_{n-1} / 2^{(n-1) / 2} \leq v_{B}(t) \leq \omega_{n-1}$ for $-\sqrt{2} / 2 \leq t \leq \sqrt{2} / 2$; thus by Lemma 3.4 and by construction we have

$$
\frac{\omega_{n-1}}{2^{n / 2}} \leq v_{E}(t) \leq 2 \omega_{n-1}, \quad \frac{\omega_{n-1}}{2^{n / 2}} \leq v_{F}(t) \leq 2 \omega_{n-1}
$$

This immediately ensures

$$
\left|E^{*} \cap J_{*}\right| \leq 2 \omega_{n-1} \mathscr{H}^{1}(J)=2 \omega_{n-1} \varepsilon, \quad\left|F \cap \widetilde{J}_{*}\right| \geq \frac{\omega_{n-1}}{2^{n / 2}} \mathscr{H}^{1}(\widetilde{J})=\frac{h \omega_{n-1}}{2^{n / 2} N},
$$

from which we deduce
$|F|=\left|E^{*}\right|+\left|F \cap \widetilde{J}_{*}\right|-\left|E^{*} \cap J_{*}\right| \geq\left|E^{*}\right|+\frac{h \omega_{n-1}}{2^{n / 2} N}-2 \omega_{n-1} \varepsilon \geq\left|E^{*}\right|+\frac{h \omega_{n-1}}{2 \cdot 2^{n / 2} N}$.
The last inequality holds provided $h \geq 4 \cdot 2^{n / 2} N \varepsilon$, and this can be achieved with a suitable choice of $M$ in view of the second estimate in (3.19) - recall
that the constant $M$ can be chosen as big as we need, while the integer $N$ (still to be determined) depends only on the dimension of the space.

We turn now to the evaluation of the perimeter $P(F)$ of $F$. Clearly, we have

$$
\begin{equation*}
P(F)=P\left(E^{*}\right)-P\left(E^{*} \mid J_{*}\right)+P\left(F \mid \widetilde{J}_{*}\right) . \tag{3.21}
\end{equation*}
$$

Moreover, since both $E^{*}$ and $F$ are axially symmetric, (3.3) allows us to calculate their perimeters in terms of $v_{E}$ and $v_{F}$. Keep in mind that our aim is to show that $F$ has a low perimeter, because this will ensure us that $F$ has an isoperimetric deficit quite a bit lower than $E^{*}$, which means that $E^{*}$ has quite a high deficit. Therefore, first of all we estimate $P\left(F \mid \widetilde{J}_{*}\right)$ from above, recalling again that $\left\|v_{F}\right\|_{L^{\infty}} \leq 2 \omega_{n-1}$, that $v_{F}^{\prime} \equiv N$ inside $\widetilde{J}$ and (3.19),

$$
\begin{align*}
P\left(F \mid \widetilde{J}_{*}\right) & =\int_{\widetilde{J}} \sqrt{\tau v_{F}(t)^{\frac{2 n-4}{n-1}}+v_{F}^{\prime}(t)^{2}}=\int_{\widetilde{J}} \sqrt{\tau v_{F}(t)^{\frac{2 n-4}{n-1}}+N^{2}} \\
& \leq \int_{\widetilde{J}} \sqrt{\tau\left(2 \omega_{n-1}\right)^{\frac{2 n-4}{n-1}}+N^{2}}=\int_{\widetilde{J}} \sqrt{\tau^{\prime}+N^{2}} \leq \int_{\widetilde{J}} N+\frac{\tau^{\prime}}{2 N}  \tag{3.22}\\
& =\left(N+\frac{\tau^{\prime}}{2 N}\right) \mathscr{H}^{1}(\widetilde{J})=h\left(1+\frac{\tau^{\prime}}{2 N^{2}}\right),
\end{align*}
$$

where we have denoted by $\tau^{\prime}$ the new constant (also depending only on $n$ ),

$$
\tau^{\prime}:=\tau\left(2 \omega_{n-1}\right)^{\frac{2 n-4}{n-1}}=2^{\frac{2 n-4}{n-1}}(n-1)^{2} \omega_{n-1}^{2}
$$

Now, we easily estimate $P\left(E^{*} \mid J_{*}\right)$ from below:

$$
\begin{equation*}
P\left(E^{*} \mid J_{*}\right)=\int_{J} \sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+v_{E}^{\prime}(t)^{2}} \geq \int_{J}\left|v_{E}^{\prime}(t)\right|=h \tag{3.23}
\end{equation*}
$$

Finally, gathering (3.21), (3.22) and (3.23) we obtain

$$
\begin{equation*}
P(F) \leq P\left(E^{*}\right)+h \frac{\tau^{\prime}}{2 N^{2}} \tag{3.24}
\end{equation*}
$$

Now we compare the volume estimate (3.20) with (3.24) thus getting, by the isoperimetric inequality, a perimeter estimate

$$
\begin{align*}
P(F) & \geq n \omega_{n}^{1 / n}|F|^{\frac{n-1}{n}} \geq n \omega_{n}^{1 / n}\left(\left|E^{*}\right|+\frac{h \omega_{n-1}}{2 \cdot 2^{n / 2} N}\right)^{\frac{n-1}{n}} \\
& =n \omega_{n}^{1 / n}\left|E^{*}\right|^{\frac{n-1}{n}}\left(1+\frac{h \omega_{n-1}}{2 \cdot 2^{n / 2} N\left|E^{*}\right|}\right)^{\frac{n-1}{n}}  \tag{3.25}\\
& =P(B)\left(1+\frac{h \omega_{n-1}}{2 \cdot 2^{n / 2} N\left|E^{*}\right|}\right)^{\frac{n-1}{n}} \geq P(B)+C_{0}(n) \frac{h}{N}
\end{align*}
$$

notice that the last inequality holds since $h$ is bounded from above (as a consequence of Lemma 3.4).

Finally, comparing (3.24) with (3.25) we get

$$
\begin{aligned}
P(B) D\left(E^{*}\right) & =P\left(E^{*}\right)-P(B) \geq P(F)-h \frac{\tau^{\prime}}{2 N^{2}}-P(B) \\
& \geq C_{0}(n) \frac{h}{N}-h \frac{\tau^{\prime}}{2 N^{2}}=h\left(\frac{C_{0}(n)}{N}-\frac{\tau^{\prime}}{2 N^{2}}\right)
\end{aligned}
$$

If we choose $N$ so that the expression in parentheses is bounded from below by a positive constant, the constant $M$, which as we noticed before depends only on $N$, will ultimately depend only on $n$ and

$$
D\left(E^{*}\right) \geq \frac{1}{K} h
$$

In view of (3.19), we then get

$$
\mathscr{H}^{1}(J)=\varepsilon \leq \frac{h}{M} \leq \frac{K}{M} D\left(E^{*}\right)
$$

where the constant $K$ again depends only on $n$. This proves (3.18) when $v_{E}^{\prime}$ is continuous and $J$ is an interval. When $v_{E}^{\prime}$ is continuous, in general $J$ is a countable union of pairwise disjoint open intervals and it is easily checked that the above proof still works, due to the additivity of both perimeter and volume.

If $v_{E}^{\prime}$ is not continuous, by assumptions (3.1) and $E \subseteq Q_{l}$, we have from Theorem 6.1 that $v_{E}$ is a nonnegative function in $W^{1,1}(\mathbb{R})$ with support in $[-l, l]$. Thus we may construct a sequence of smooth nonnegative functions $v_{h}$ converging to $v_{E}$ in $W^{1,1}(\mathbb{R})$, whose supports are all contained in $[-l, l]$, and such that

$$
\begin{equation*}
\int_{-\infty}^{+\infty} v_{h} d t=\int_{-\infty}^{+\infty} v_{E} d t \text { for all } h, v_{h} \rightarrow v_{E} \text { uniformly, } v_{h}^{\prime} \rightarrow v_{E}^{\prime} \text { a.e. . } \tag{3.26}
\end{equation*}
$$

Let us denote by $E_{h}$ the axially symmetric set defined by $v_{E_{h}}=v_{h}$. Clearly, from (3.26) and Lemma 3.3 we get easily that $\left|E_{h}\right|=\left|E^{*}\right|$ for all $h$ and $P\left(E_{h}\right) \rightarrow$ $P\left(E^{*}\right)$; hence $D\left(E_{h}\right) \rightarrow D\left(E^{*}\right)$. The result then follows by applying the above estimate to each set $J_{h}=\left\{t \in[-\sqrt{2} / 2, \sqrt{2} / 2]:\left|v_{h}^{\prime}(t)\right| \geq M\right\}$ and observing that $\mathscr{H}^{1}\left(J_{h}\right) \rightarrow \mathscr{H}^{1}(J)$.

This proves (3.18); hence the result follows.

## 4. The axially symmetric case

This section is entirely devoted to the proof of the following estimate, which refers to the particular case of an axially symmetric set.

Theorem 4.1. Let $E \in X$ be axially symmetric with respect to the first axis and symmetric with respect to the hyperplane $\left\{x_{1}=0\right\}$. Assume also that
$E \subseteq Q_{l}$ and that (3.1) holds. Then,

$$
\begin{equation*}
d(E \cap Z, B \cap Z) \leq C(n, l) \sqrt{D(E)}, \tag{4.1}
\end{equation*}
$$

where $Z=\left\{x \in \mathbb{R}^{n}:\left|x_{1}\right| \leq \sqrt{2} / 2\right\}$.
Notice that the claim above states, for an axially symmetric set, an inequality similar to the one we aim to prove, but only the volume of the internal part of the symmetric difference (that is, inside $Z$ ) is estimated.

Let us explain now the strategy we will adopt to prove Theorem 4.1. The first step will be to show the existence of a section of $E$ such that the ( $n-1$ )-dimensional symmetric difference between this section of $E$ and the corresponding section of $B$ (in the sense of Lemma 4.2 below) gives a bound to the left term in (4.1). The second step will be to prove that it is not restrictive to assume that this section is the central one. Finally, the last step will be a careful but not difficult comparison between the symmetric difference of the central sections of $E$ and $B$ and the isoperimetric deficit of $E$.

Lemma 4.2. There exist $\delta(n, l), C(n)>0$, with the property that if $E$ is as in Theorem 4.1 and $D(E)<\delta$, there exists a number $0 \leq \bar{x} \leq \sqrt{2} / 2$ such that, if $x^{\prime}$ is defined through the equality

$$
\begin{equation*}
\int_{0}^{x^{\prime}} v_{B}(t) d t=\int_{0}^{\bar{x}} v_{E}(t) d t \tag{4.2}
\end{equation*}
$$

then

$$
\begin{equation*}
\left|v_{E}(\bar{x})-v_{B}\left(x^{\prime}\right)\right| \geq \frac{1}{C} d(E \cap Z, B \cap Z) \tag{4.3}
\end{equation*}
$$

Notice that the number $x^{\prime}$ is chosen in such a way that the volume of the set $E$ in the half-space $\left\{x_{1} \geq \bar{x}\right\}$ coincides with the volume of the ball $B$ in the half-space $\left\{x_{1} \geq x^{\prime}\right\}$.

Proof of Lemma 4.2. Recalling that both $E$ and $B$ are axially symmetric, we can write

$$
d(E \cap Z, B \cap Z)=\int_{-\sqrt{2} / 2}^{\sqrt{2} / 2} \mathscr{H}^{n-1}\left(E_{t} \Delta B_{t}\right) d t=\int_{-\sqrt{2} / 2}^{\sqrt{2} / 2}\left|v_{E}(t)-v_{B}(t)\right| d t
$$

so that by symmetry there exists $0 \leq x_{0} \leq \sqrt{2} / 2$ with the property that

$$
\begin{equation*}
\left|v_{E}\left(x_{0}\right)-v_{B}\left(x_{0}\right)\right| \geq \frac{\sqrt{2}}{2} d(E \cap Z, B \cap Z) \tag{4.4}
\end{equation*}
$$

Let us assume for the moment that $v_{B}\left(x_{0}\right)>v_{E}\left(x_{0}\right)$. In this case, let us first assume that the inequality

$$
\begin{equation*}
\int_{0}^{x_{0}} v_{B}(t) d t \geq \int_{0}^{x_{0}} v_{E}(t) d t \tag{4.5}
\end{equation*}
$$

holds. Then, defining $\bar{x}=x_{0}$, by the symmetry of $E$ and $B$ with respect to $\left\{x_{1}=0\right\}$, we have that $0 \leq x^{\prime} \leq \bar{x} \leq 1-$ with $x^{\prime}$ defined according to (4.2) so that

$$
v_{B}\left(x^{\prime}\right) \geq v_{B}(\bar{x}) \geq v_{E}(\bar{x})+\frac{\sqrt{2}}{2} d(E \cap Z, B \cap Z),
$$

thus (4.3) follows, and so the thesis is obtained by choosing $C(n) \geq \sqrt{2}$.
When (4.5) does not hold, we consider two cases, depending on whether or not the following inequality holds

$$
\begin{equation*}
\int_{0}^{x_{0}} v_{B}(t) d t \geq \int_{0}^{x_{0}} v_{E}(t) d t-\frac{1}{2 \cdot 2^{n / 2}(n-1)} d(E \cap Z, B \cap Z) . \tag{4.6}
\end{equation*}
$$

Case I. If the inequality (4.6) holds. In this case, let again $\bar{x}=x_{0}$ and let $x^{\prime}$ be defined according to (4.2). Assuming $\delta(n, l)$ sufficiently small, by Lemmas 2.2 and 2.3 we have that $d(E, B)$ is small, and since $v_{B} \geq \omega_{n-1} / 2^{(n-1) / 2}$ in $[-\sqrt{2} / 2, \sqrt{2} / 2]$, we obtain that $\left|x^{\prime}-\bar{x}\right|$ is small; therefore, provided $\delta$ is sufficiently small, we have $v_{B}(t) \geq \omega_{n-1} / 2^{n / 2}$ for $t$ in $\left[\bar{x}, x^{\prime}\right]$. This implies, by (4.6)

$$
\begin{aligned}
\left|x^{\prime}-\bar{x}\right| & =\int_{\bar{x}}^{x^{\prime}} 1 d t \leq \frac{2^{n / 2}}{\omega_{n-1}} \int_{\bar{x}}^{x^{\prime}} v_{B}(t) d t=\frac{2^{n / 2}}{\omega_{n-1}}\left(\int_{0}^{x^{\prime}} v_{B}(t) d t-\int_{0}^{\bar{x}} v_{B}(t) d t\right) \\
& =\frac{2^{n / 2}}{\omega_{n-1}} \int_{0}^{\bar{x}} v_{E}(t)-v_{B}(t) d x \leq \frac{1}{2(n-1) \omega_{n-1}} d(E \cap Z, B \cap Z) .
\end{aligned}
$$

Notice now that, in $[0, \sqrt{2} / 2]$,

$$
0 \geq v_{B}^{\prime}(t)=-(n-1) \omega_{n-1} t\left(1-t^{2}\right)^{\frac{n-3}{2}} \geq-(n-1) \omega_{n-1}
$$

therefore, recalling also (4.4), the assumption $v_{B}(\bar{x})>v_{E}(\bar{x})$ and the fact that $\left|\bar{x}-x^{\prime}\right|$ is small enough if $\delta \ll 1$, we have

$$
\begin{aligned}
v_{B}\left(x^{\prime}\right) & =v_{B}(\bar{x})+\int_{\bar{x}}^{x^{\prime}} v_{B}^{\prime}(t) d t \geq v_{B}(\bar{x})-\frac{6}{5}\left|x^{\prime}-\bar{x}\right|(n-1) \omega_{n-1} \\
& \geq v_{B}(\bar{x})-\frac{3 d(E \cap Z, B \cap Z)}{5} \geq v_{E}(\bar{x})+\left(\frac{\sqrt{2}}{2}-\frac{3}{5}\right) d(E \cap Z, B \cap Z) ;
\end{aligned}
$$

hence also in this case (4.3), and so the thesis, is obtained, by choosing $C(n)$ large enough.

Case II. If the inequality (4.6) does not hold. In this case we have

$$
\int_{0}^{x_{0}} v_{E}(t)-v_{B}(t) d t>\frac{1}{2 \cdot 2^{n / 2}(n-1)} d(E \cap Z, B \cap Z)
$$

Since $0 \leq x_{0} \leq \sqrt{2} / 2$, this implies

$$
f_{0}^{x_{0}} v_{E}(t)-v_{B}(t) d t>\frac{\sqrt{2}}{2 \cdot 2^{n / 2}(n-1)} d(E \cap Z, B \cap Z)
$$

Now let $\bar{x}$ be the largest number in $\left[0, x_{0}\right]$ such that

$$
v_{E}(\bar{x})-v_{B}(\bar{x}) \geq \frac{\sqrt{2}}{2 \cdot 2^{n / 2}(n-1)} d(E \cap Z, B \cap Z) .
$$

Therefore, we deduce

$$
\begin{aligned}
\int_{\bar{x}}^{x_{0}} v_{E}(t)-v_{B}(t) & \leq\left(x_{0}-\bar{x}\right) \frac{\sqrt{2}}{2 \cdot 2^{n / 2}(n-1)} d(E \cap Z, B \cap Z) \\
& \leq \frac{1}{2 \cdot 2^{n / 2}(n-1)} d(E \cap Z, B \cap Z)
\end{aligned}
$$

the fact that (4.6) does not hold, together with this estimate, ensures

$$
\int_{0}^{\bar{x}} v_{E}(t)-v_{B}(t) d t>0
$$

and then $x^{\prime} \geq \bar{x}$, defining as usual $x^{\prime}$ according to (4.2). Then

$$
v_{B}\left(x^{\prime}\right) \leq v_{B}(\bar{x}) \leq v_{E}(\bar{x})-\frac{\sqrt{2}}{2 \cdot 2^{n / 2}(n-1)} d(E \cap Z, B \cap Z) ;
$$

so (4.3) is established also in this case, after the choice of a suitable constant $C(n)$.

Finally, it remains to consider what happens if $v_{B}\left(x_{0}\right) \leq v_{E}\left(x_{0}\right)$ : the proof in this case is very similar to the preceding one; it suffices to replace (4.5) with

$$
\int_{0}^{x_{0}} v_{B}(t) d t \leq \int_{0}^{x_{0}} v_{E}(t) d t
$$

and (4.6) with

$$
\int_{0}^{x_{0}} v_{B}(t) d t \leq \int_{0}^{x_{0}} v_{E}(t) d t+\frac{1}{2 \cdot 2^{n / 2}(n-1)} d(E \cap Z, B \cap Z)
$$

and to argue in the very same way as before.
Lemma 4.3. Let $E, \delta(n, l)$ and $\bar{x}$ be as in Lemma 4.2 and in Theorem 4.1; then there exists a set $E^{\prime} \in X, E^{\prime} \subseteq Q_{2 l}$, such that (3.1) holds, axially symmetric with respect to the first axis and symmetric with respect to the hyperplane $\left\{x_{1}=0\right\}$, with the property that

$$
P\left(E^{\prime}\right) \leq P(E), \quad\left|v_{E^{\prime}}(0)-v_{B}(0)\right| \geq \frac{1}{C}\left|v_{E}(\bar{x})-v_{B}\left(x^{\prime}\right)\right|
$$

for some constant $C=C(n)$.
Proof. Let us write $E=E^{l} \cup E^{c} \cup E^{r}$, where

$$
\begin{gathered}
E^{l}:=\left\{x \in E: x_{1} \leq-\bar{x}\right\}, \quad E^{c}:=\left\{x \in E:\left|x_{1}\right|<\bar{x}\right\}, \\
E^{r}:=\left\{x \in E: x_{1} \geq \bar{x}\right\} .
\end{gathered}
$$

Notice that there exists a unique pair of numbers $r \in \mathbb{R}$ and $x_{c}>0$ such that the intersection of the ball of radius $1+r$ with the strip $\left\{-x_{c} \leq x_{1} \leq x_{c}\right\}$,

$$
B^{c}:=\left\{x \in \mathbb{R}^{n}:|x| \leq 1+r,\left|x_{1}\right| \leq x_{c}\right\},
$$

satisfies

$$
\begin{equation*}
\left|B^{c}\right|=\left|E^{c}\right|, \quad \mathscr{H}^{n-1}\left(B_{x_{c}}^{c}\right)=\mathscr{H}^{n-1}\left(E_{\bar{x}}\right)=v_{E}(\bar{x}) . \tag{4.7}
\end{equation*}
$$

Let us now define $\widetilde{E}^{l}$ and $\widetilde{E}^{r}$, the following two horizontal translations of $E^{l}$ and $E^{r}$ :

$$
\widetilde{E}^{l}:=\left(\bar{x}-x_{c}\right) e_{1}+E^{l}, \quad \widetilde{E}^{r}:=\left(x_{c}-\bar{x}\right) e_{1}+E^{r},
$$

and finally set $E^{\prime}=\widetilde{E}^{l} \cup B^{c} \cup \widetilde{E}^{r}$. Notice that the translations of $E^{l}$ and $E^{r}$ are made in such a way that the three parts $\widetilde{E}^{l}, B^{c}$ and $\widetilde{E}^{r}$ match along their boundaries. We will show that $E^{\prime}$ has the desired properties. It is immediate from the construction that $\left|E^{\prime}\right|=|E|$-so that $E^{\prime} \in X$ - and that $E^{\prime}$ is axially symmetric and symmetric with respect to the hyperplane $\left\{x_{1}=0\right\}$. We show now that $P\left(E^{\prime}\right) \leq P(E)$. Define

$$
\begin{aligned}
B^{l} & :=\left\{x \in \mathbb{R}^{n}:|x| \leq 1+r, x_{1}<-x_{c}\right\}, \\
B^{r} & :=\left\{x \in \mathbb{R}^{n}:|x| \leq 1+r, x_{1}>x_{c}\right\},
\end{aligned}
$$

so that $B^{l} \cup B^{c} \cup B^{r}$ is the ball of radius $1+r$; moreover, let $\widetilde{B}^{l}$ and $\widetilde{B}^{r}$ be the horizontal translations of $B^{l}$ and $B^{r}$,

$$
\widetilde{B}^{l}:=\left(x_{c}-\bar{x}\right) e_{1}+B^{l}, \quad \widetilde{B}^{r}:=\left(\bar{x}-x_{c}\right) e_{1}+B^{r}
$$

and notice that $\widetilde{B}^{l} \cup E^{c} \cup \widetilde{B}^{r}$ is a set of the same volume as the ball $B^{l} \cup B^{c} \cup B^{r}$; this implies $P\left(\widetilde{B}^{l} \cup E^{c} \cup \widetilde{B}^{r}\right) \geq P\left(B^{l} \cup B^{c} \cup B^{r}\right)$, hence $P\left(E^{c}\right) \geq P\left(B^{c}\right)$ and so $P(E) \geq P\left(E^{\prime}\right)$. Then, to conclude the thesis, we need only to show that

$$
\left|v_{E^{\prime}}(0)-v_{B}(0)\right| \geq \frac{1}{C}\left|v_{E}(\bar{x})-v_{B}\left(x^{\prime}\right)\right|
$$

Let us suppose, just to fix the ideas, that $v_{E}(\bar{x})>v_{B}\left(x^{\prime}\right)$-in the opposite case the proof works in a very similar way. We claim that $r>0$ and $0<x_{c}<x^{\prime}$ : indeed, if $r \leq 0$ then by (4.2) and the fact that $\left|E^{c}\right|=\left|B^{c}\right|$ we would infer $x_{c} \geq x^{\prime}$, so that $\mathscr{H}^{n-1}\left(B_{x_{c}}^{c}\right) \leq \mathscr{H}^{n-1}\left(B_{x^{\prime}}^{c}\right)=v_{B}\left(x^{\prime}\right)<v_{E}(\bar{x})$ against (4.7). This implies $r>0$, and again by the fact that $\left|E^{c}\right|=\left|B^{c}\right|$ and by (4.2) we deduce $x_{c}<x^{\prime}$.

Since $r>0$,

$$
\begin{aligned}
\left|v_{E^{\prime}}(0)-v_{B}(0)\right| & =\left|\mathscr{H}^{n-1}\left(B_{0}^{c}\right)-\mathscr{H}^{n-1}\left(B_{0}\right)\right| \\
& =\omega_{n-1}\left((1+r)^{n-1}-1\right) \geq(n-1) \omega_{n-1} r .
\end{aligned}
$$

Hence, to obtain the thesis it is sufficient to show that

$$
\begin{equation*}
\left|v_{E}(\bar{x})-v_{B}\left(x^{\prime}\right)\right|=\left|v_{E^{\prime}}\left(x_{c}\right)-v_{B}\left(x^{\prime}\right)\right| \leq C r . \tag{4.8}
\end{equation*}
$$



Figure 3: Situation in the proof of Lemma 4.3

To this aim, consider Figure 3, which shows the two balls of radii 1 and $1+r$, the points $X_{C} \equiv\left(x_{c}, 0, \ldots, 0\right)$ and $X^{\prime} \equiv\left(x^{\prime}, 0, \ldots, 0\right)$ on the first axis, and the points $A, B, H$ and $K$ as drawn. By construction, the two shaded regions have the same volume; moreover, since $0 \leq \bar{x} \leq \sqrt{2} / 2,\left|x^{\prime}-\bar{x}\right|=O(D(E))$ and $0 \leq x_{c} \leq x^{\prime}$, there exists a constant $\alpha=\alpha(n)$ such that $\overline{A B} \leq \alpha r$. As an easy consequence, the volume of the lightly shaded region is less than $C(n) r$. Recalling again that $x^{\prime} \leq \sqrt{2} / 2+O(D(E))$ we deduce that the $(n-1)-$ dimensional measures $\mathscr{H}^{n-1}\left(B_{t}\right)$ of the sections $B_{t}$ of the ball $B$ with $x_{c} \leq t \leq$ $x^{\prime}$ are all greater than $\left((\sqrt{2} / 2)^{n-1} / 2\right) \omega_{n-1}=2^{-(n+1) / 2} \omega_{n-1}$. Therefore, the volume of the darker region (which equals that of the lighter one) is greater than $2^{-(n+1) / 2} \omega_{n-1}\left|x^{\prime}-x_{c}\right|$. Thus,

$$
\left|x^{\prime}-x_{c}\right|<C r .
$$

Since the angle $\widehat{A K H}$ is less than $\pi / 4+O(D(E))$,

$$
\overline{B H}=\overline{A B}+\overline{A H} \leq \alpha r+2 \overline{H K}=\alpha r+2\left|x^{\prime}-x_{c}\right| \leq C r .
$$

Finally, keeping in mind that

$$
v_{E^{\prime}}\left(x_{c}\right)=\omega_{n-1}{\overline{B X_{c}}}^{n-1}, \quad v_{B}\left(x^{\prime}\right)=\omega_{n-1}{\overline{K X^{\prime}}}^{n-1}=\omega_{n-1} \overline{H X}_{c}{ }^{n-1}
$$

and that $\overline{H X_{c}} \geq \sqrt{2} / 2+O(D(E))$, we finally get (4.8), thus the proof of Lemma 4.3 is concluded.

Thanks to the above lemmas, our last step is to perform a precise estimate of the isoperimetric deficit of an axially symmetric set as in the claim of Lemma 4.3; having reduced the problem to a very specific case, we see that the calculation will be involved but elementary, while an analogous calculation for a general set as in Theorem 4.1 would be extremely complicated.

Proof of Theorem 4.1. Thanks to Lemmas 4.2 and 4.3, the proof of Theorem 4.1 will be achieved once we show that, for any set $E$ as in Theorem 4.1,

$$
\begin{equation*}
\left|v_{E}(0)-v_{B}(0)\right| \leq C \sqrt{D(E)} \tag{4.9}
\end{equation*}
$$

First of all, let $\widetilde{B}$ be the ball centered on the first axis such that $\widetilde{B}_{0}=E_{0}$ and

$$
\left|\widetilde{B} \cap\left\{x \in \mathbb{R}^{N}: x_{1}<0\right\}\right|=\left|E \cap\left\{x \in \mathbb{R}^{N}: x_{1}<0\right\}\right|=\frac{|B|}{2} ;
$$

in words, in the open half-space $R_{-}^{n}=\left\{x_{1}<0\right\}$ the sets $\widetilde{B}$ and $E$ have the same volume, and their sections at $x_{1}=0$ coincide. As an immediate consequence, with the same argument used in Lemma 4.3, we obtain $P\left(\widetilde{B} \mid \mathbb{R}_{-}^{n}\right) \leq P\left(E \mid \mathbb{R}_{-}^{n}\right)$.



Figure 4: The sets $E$ and $E^{\prime}$

Define $E^{\prime}$ as the set which coincides with the ball $\widetilde{B}$ in $\mathbb{R}_{-}^{n}$ and which is symmetric with respect to the hyperplane $\left\{x_{1}=0\right\}$, as shown in Figure 4. By construction,

$$
v_{E^{\prime}}(0)=v_{E}(0), \quad P\left(E^{\prime}\right) \leq P(E),
$$

so that to show the claim of this lemma we can consider the set $E^{\prime}$ instead of $E$; more precisely, the proof will be concluded once we establish that

$$
\left|v_{E^{\prime}}(0)-v_{B}(0)\right| \leq C \sqrt{P\left(E^{\prime}\right)-P(B)} .
$$

Let us set

$$
\begin{equation*}
\varepsilon:=\frac{v_{E^{\prime}}(0)-v_{B}(0)}{\omega_{n-1}}, \tag{4.10}
\end{equation*}
$$

and let us denote the radius of the ball $\widetilde{B}$ by $1+r$ and the center of the ball $\widetilde{B}$ by $(-\delta, 0, \ldots, 0)$. Notice that the set $E^{\prime}$ is fully determined by $\varepsilon$, and then $r=r(\varepsilon)$ and $\delta=\delta(\varepsilon)$; moreover, when $\varepsilon=0$ one of course has $r=\delta=0$ since $E^{\prime} \equiv \widetilde{B} \equiv B$. We will now evaluate $r$ and $\delta$ in terms of $\varepsilon$.

To do that we first write the following two conditions:

$$
\left\{\begin{array}{l}
\left((1+r)^{2}-\delta^{2}\right)^{\frac{n-1}{2}}=1+\varepsilon  \tag{4.11}\\
\int_{-(1+r)}^{\delta}\left((1+r)^{2}-t^{2}\right)^{\frac{n-1}{2}} d t=\int_{-1}^{0}\left(1-t^{2}\right)^{\frac{n-1}{2}} d t
\end{array}\right.
$$

Let us prove these equations: concerning the first one, if we denote, as in Figure 4 , by $P$ the center of $\widetilde{B}$ and by $Q$ some point of $\partial E^{\prime} \cap\left\{x_{1}=0\right\}$, we have $\overline{P Q}^{2}=\overline{P O}^{2}+\overline{O Q}^{2}$ by Pythagoras Theorem; thus $\overline{O Q}=\sqrt{(1+r)^{2}-\delta^{2}}$. Since $E_{0}^{\prime}$ is a ball in $\mathbb{R}^{n-1}$ of radius $\overline{O Q}$, its volume $v_{E^{\prime}}(0)$ is $\omega_{n-1} \overline{O Q}^{n-1}$; recalling (4.10) and the fact that $v_{B}(0)=\omega_{n-1}$ by definition, we recover the first equation. The second one comes from the fact that $\left|E^{\prime}\right|=|B|$ and $E^{\prime}$ is symmetric with respect to the hyperplane $\left\{x_{1}=0\right\}$; indeed the first and the second integrals, up to the factor $\omega_{n-1}$, are the volumes of $E^{\prime}$ and of $B$ inside $\mathbb{R}_{-}^{n-1}$ respectively.

Thus, (4.9) follows once we show that

$$
\begin{equation*}
P\left(E^{\prime}\right)-P(B) \geq \frac{1}{C} \varepsilon^{2} \tag{4.12}
\end{equation*}
$$

To this aim, let us set

$$
a:=\int_{-1}^{0}\left(1-t^{2}\right)^{\frac{n-3}{2}} d t=\int_{0}^{\pi / 2} \sin ^{n-2} \theta d \theta
$$

Evaluating the first equation in (4.11) in $r$ and $\delta$ at the second order, we get

$$
\begin{equation*}
\varepsilon=(n-1) r+\frac{(n-1)(n-2)}{2} r^{2}-\frac{n-1}{2} \delta^{2}+o\left(r^{2}+\delta^{2}\right) \tag{4.13}
\end{equation*}
$$

The second order Taylor expansion in $r$ and $\delta$ of the second equation in (4.11) yields

$$
\begin{equation*}
0=(n-1) r a+\delta+\frac{(n-1)^{2} a}{2} r^{2}+(n-1) r \delta+o\left(r^{2}+\delta^{2}\right) \tag{4.14}
\end{equation*}
$$

Finally, (4.13) and (4.14) allow us to obtain a first order estimate of $r$ and $\delta$ :

$$
\left\{\begin{array}{l}
r=\frac{1}{n-1} \varepsilon+o(r+\delta) \\
\delta=-a \varepsilon+o(r+\delta)
\end{array}\right.
$$

This ensures that both $r$ and $\delta$ are of order $\varepsilon$, so that in particular $o(r+\delta)=$ $o(\varepsilon)$ and $o\left(r^{2}+\delta^{2}\right)=o\left(\varepsilon^{2}\right)$. However, we need a more precise estimate, namely an estimate of $r$ and $\delta$ up to the second order in $\varepsilon$. Setting

$$
r=\frac{1}{n-1} \varepsilon+\tilde{r}, \quad \delta=-a \varepsilon+\tilde{\delta}
$$

and plugging these expressions into (4.13) and (4.14), we find the second order expansion of $\tilde{r}$ and $\tilde{\delta}$, thus concluding that

$$
\left\{\begin{aligned}
r & =\frac{1}{n-1} \varepsilon+\frac{\varepsilon^{2}}{2(n-1)}\left((n-1) a^{2}-\frac{n-2}{n-1}\right)+o\left(\varepsilon^{2}\right) \\
\delta & =-a \varepsilon-\frac{a \varepsilon^{2}}{2}\left((n-1) a^{2}-\frac{n-2}{n-1}\right)+\frac{a \varepsilon^{2}}{2}+o\left(\varepsilon^{2}\right)
\end{aligned}\right.
$$

Once we know $r$ and $\delta$ with a second order precision, we pass to the calculation of the perimeter of $E^{\prime}$ keeping in mind that our aim is to show (4.12). Since $E^{\prime}$ is a part of a ball, it is immediate to write down its perimeter in terms of $r$ and $\delta$; that is,

$$
\frac{P\left(E^{\prime}\right)}{2(n-1) \omega_{n-1}}=\int_{-(1+r)}^{\delta}\left((1+r)^{2}-t^{2}\right)^{\frac{n-3}{2}}(1+r) d t=: \xi(r, \delta) .
$$

The second order Taylor expansion of $\xi$ gives

$$
\begin{aligned}
\xi(r, \delta) & =\xi(0,0)+r(n-1) a+\delta+\frac{(n-2)(n-1) a}{2} r^{2}+(n-2) r \delta+o\left(r^{2}+\delta^{2}\right) \\
& =\xi(0,0)+\frac{a}{2(n-1)} \varepsilon^{2}+o\left(\varepsilon^{2}\right) .
\end{aligned}
$$

Since

$$
\xi(0,0)=\frac{P(B)}{2(n-1) \omega_{n-1}},
$$

we conclude that

$$
P\left(E^{\prime}\right)-P(B)=a \omega_{n-1} \varepsilon^{2}+o\left(\varepsilon^{2}\right)
$$

and (4.12) follows.

## 5. The proof of the main theorem

We are now ready to give the proof of Theorem 1.1.
Proof of Theorem 1.1. As already observed, it is enough to prove the theorem under the assumption that $|E|=|B|$. Moreover we may also assume that $D(E) \leq 1$ (otherwise the assertion is trivial).

From Lemma 5.1 below it is clear that (with no loss of generality) we may assume that $E$ is contained in a fixed cube $Q_{l}$, where $l$ is a constant depending only on $n$. Let $n$ be either 2 or greater than 2 and such that Theorem 1.1 is true in dimension $n-1$.

Let us fix a set of finite perimeter $E \subseteq Q_{l}$. Since there are only countably many directions $\nu$ such that $\mathscr{H}^{n-1}\left(\left\{x \in \partial^{*} E: \nu^{E}(x)= \pm \nu\right\}\right)>0$, by rotating $E$ and slightly increasing $l$ if necessary, we may assume without loss of generality that for all $i=1, \ldots, n$

$$
\begin{equation*}
\mathscr{H}^{n-1}\left(\left\{x \in \partial^{*} E: \nu^{E}(x)= \pm e_{i}\right\}\right)=0 . \tag{5.1}
\end{equation*}
$$

By Theorem 2.1 and Remark 2.6, there exists a set $F$, symmetric with respect to the $n$ orthogonal coordinate hyperplanes, satisfying (5.1) and contained in the cube $Q_{3 l}$, such that

$$
\begin{equation*}
\lambda(E) \leq C \lambda(F), \quad D(F) \leq 2^{n} D(E) \tag{5.2}
\end{equation*}
$$

Thus, combining Theorems 3.1 and 4.1 we get that, up to a rotation of the coordinate axes,

$$
\lambda(F) \leq d(F, B) \leq 4 d\left(F^{*} \cap Z, B \cap Z\right)+C \sqrt{D(F)} \leq C \sqrt{D(F)}
$$

From this inequality and (5.2), the assertion immediately follows by a trivial induction argument, for any dimension $n \geq 2$.

The next lemma is a variant of a similar result proved in [10] (see also [11]).
Lemma 5.1. There exist two constants, $l=l(n)$ and $C=C(n)$ such that, if $E \in X$, then there is a set $E^{\prime} \subseteq Q_{l}$, such that $\left|E^{\prime}\right|=|B|$ and

$$
\begin{equation*}
\lambda(E) \leq \lambda\left(E^{\prime}\right)+C D(E), \quad D\left(E^{\prime}\right) \leq C D(E) \tag{5.3}
\end{equation*}
$$

Proof. First of all we observe that, by rotating $E$ if necessary, we may assume that

$$
\begin{equation*}
\mathscr{H}^{n-1}\left(\left\{x \in \partial^{*} E: \nu^{E}(x)= \pm e_{i}\right\}\right)=0 \tag{5.4}
\end{equation*}
$$

for all $i=1, \ldots, n$. By Lemma 6.1 we deduce that $v_{E} \in W^{1,1}(\mathbb{R})$, so that in particular $v_{E}$ is continuous.

Now, setting $E_{t}^{-}=\left\{x \in E: x_{1}<t\right\}$ for all $t \in \mathbb{R}$, we claim that

$$
\begin{equation*}
P\left(E_{t}^{-}\right) \leq P\left(E \mid\left\{x_{1}<t\right\}\right)+v_{E}(t), \quad P\left(E \backslash E_{t}^{-}\right) \leq P\left(E \mid\left\{x_{1}>t\right\}\right)+v_{E}(t) \tag{5.5}
\end{equation*}
$$

To prove these inequalities we first notice that, if $E$ is a smooth open set, then by (5.4) the inequalities (5.5) hold as equalities. Assuming now that $E$ is any set of finite perimeter satisfying (5.4), we may take a sequence of smooth open sets $E_{j}$ such that $\left|E_{j} \Delta E\right| \rightarrow 0, P\left(E_{j}\right) \rightarrow P(E)$ and (5.4) holds for each $E_{j}$. Since the sets $E_{j}$ are smooth, we already know that

$$
\begin{aligned}
P\left(\left(E_{j}\right)_{t}^{-}\right) & =P\left(E_{j} \mid\left\{x_{1}<t\right\}\right)+v_{E_{j}}(t) \\
P\left(E_{j} \backslash\left(E_{j}\right)_{t}^{-}\right) & =P\left(E_{j} \mid\left\{x_{1}>t\right\}\right)+v_{E_{j}}(t)
\end{aligned}
$$

Since $\left\|v_{E_{j}}-v_{E}\right\|_{L^{1}} \rightarrow 0$ and $v_{E}$ is continuous, one easily obtains by construction that $v_{E_{j}}(t) \rightarrow v_{E}(t)$ for every $t \in \mathbb{R}$; moreover, $P\left(E_{j} \mid \Omega\right) \rightarrow P(E \mid \Omega)$ for every open set $\Omega \subseteq \mathbb{R}^{n}$ such that $P(E \mid \partial \Omega)=0$. Clearly, for any $t \in \mathbb{R}$, $\left|\left(E_{j}\right)_{t}^{-} \Delta E_{t}^{-}\right| \rightarrow 0$; hence by the lower semicontinuity of the distributional perimeter we get

$$
\begin{aligned}
P\left(E_{t}^{-}\right) & \leq \liminf _{j \rightarrow \infty} P\left(\left(E_{j}\right)_{t}^{-}\right) \\
& =\lim _{j \rightarrow \infty} P\left(E_{j} \mid\left\{x_{1}<t\right\}\right)+v_{E_{j}}(t)=P\left(E \mid\left\{x_{1}<t\right\}\right)+v_{E}(t)
\end{aligned}
$$

this proves the first inequality in (5.5), and the second one is fully analogous.

Let us now define the function $g: \mathbb{R} \rightarrow \mathbb{R}^{+}$as

$$
g(t):=\frac{\left|E_{t}^{-}\right|}{|B|}
$$

$g$ is a nondecreasing, $\mathrm{C}^{1}$ function such that $g^{\prime}(t)=v_{E}(t) /|B|$. Let $-\infty \leq a<$ $b \leq+\infty$ be the numbers such that $\{t: 0<g(t)<1\}=(a, b)$; pick now any $a<t<b$ and notice that by definition

$$
\left|g(t)^{-\frac{1}{n}} E_{t}^{-}\right|=|B|
$$

therefore, $P\left(g(t)^{-1 / n} E_{t}^{-}\right) \geq P(B)$, which gives

$$
P\left(E_{t}^{-}\right)=g(t)^{\frac{n-1}{n}} P\left(g(t)^{-\frac{1}{n}} E_{t}^{-}\right) \geq g(t)^{\frac{n-1}{n}} P(B)
$$

Similarly,

$$
P\left(E \backslash E_{t}^{-}\right) \geq(1-g(t))^{\frac{n-1}{n}} P(B)
$$

Therefore, from (5.5) and assumption (5.4) we get that

$$
P(E)+2 v_{E}(t) \geq P(B)\left(g(t)^{\frac{n-1}{n}}+(1-g(t))^{\frac{n-1}{n}}\right)
$$

for all $t \in(a, b)$. On the other hand, by definition of isoperimetric deficit we have $P(E)=P(B)(1+D(E))$, and then

$$
\begin{equation*}
v_{E}(t) \geq \frac{1}{2} P(B)\left(g(t)^{\frac{n-1}{n}}+(1-g(t))^{\frac{n-1}{n}}-1-D(E)\right) \tag{5.6}
\end{equation*}
$$

Let us now define the function $\psi:[0,1] \rightarrow \mathbb{R}^{+}$, whose graph is drawn in Figure 5, as

$$
\psi(t):=t^{\frac{n-1}{n}}+(1-t)^{\frac{n-1}{n}}-1
$$

notice that $\psi(0)=\psi(1)=0$, that $\psi(1 / 2)=2^{1 / n}-1$ is the maximum, and that $\psi$ is concave, so that

$$
\begin{equation*}
\psi(t) \geq 2\left(2^{1 / n}-1\right) t \quad \forall 0 \leq t \leq \frac{1}{2} \tag{5.7}
\end{equation*}
$$



Figure 5: The graph of $\psi$

We may assume, without loss of generality, that $2 D(E)<\psi(1 / 2)$. Let $a<t_{1}<t_{2}<b$ be two numbers such that $g\left(t_{1}\right)=1-g\left(t_{2}\right)$ and $\psi\left(g\left(t_{1}\right)\right)=$ $\psi\left(g\left(t_{2}\right)\right)=2 D(E)$. Then

$$
\begin{equation*}
\psi(g(t)) \geq 2 D(E) \quad \forall t \in\left(t_{1}, t_{2}\right) \tag{5.8}
\end{equation*}
$$

and, by (5.7),

$$
\begin{equation*}
g\left(t_{1}\right)=1-g\left(t_{2}\right) \leq \frac{D(E)}{2^{1 / n}-1} . \tag{5.9}
\end{equation*}
$$

Thanks to (5.6), for any $t_{1} \leq t \leq t_{2}$ we have

$$
\begin{align*}
v_{E}(t) & \geq \frac{1}{2} P(B)(\psi(g(t))-D(E))  \tag{5.10}\\
& \geq \frac{1}{4} P(B) \psi(g(t))+\frac{1}{4} P(B)(\psi(g(t))-2 D(E)) \\
& \geq \frac{n \omega_{n}}{4} \psi(g(t)) .
\end{align*}
$$

Therefore, recalling that $v_{E}(t)=|B| g^{\prime}(t)$, we have

$$
\begin{equation*}
t_{2}-t_{1} \leq \frac{4}{n} \int_{t_{1}}^{t_{2}} \frac{g^{\prime}(t)}{\psi(g(t))} d t=\frac{4}{n} \int_{g\left(t_{1}\right)}^{g\left(t_{2}\right)} \frac{1}{\psi(s)} d s \leq \frac{4}{n} \int_{0}^{1} \frac{1}{\psi(s)} d s=\alpha \tag{5.11}
\end{equation*}
$$

for some dimensional constant $\alpha=\alpha(n)$. Let us now set

$$
\begin{aligned}
& \tau_{1}=\max \left\{t \in\left(a, t_{1}\right]: v_{E}(t) \leq \frac{n \omega_{n} D(E)}{2}\right\}, \\
& \tau_{2}=\min \left\{t \in\left[t_{2}, b\right): v_{E}(t) \leq \frac{n \omega_{n} D(E)}{2}\right\} .
\end{aligned}
$$

Notice that $\tau_{1}$ and $\tau_{2}$ are well defined since $v_{E}$ is continuous and $v_{E}(a)=$ $v_{E}(b)=0$; moreover, by (5.8) and (5.10), $v_{E}\left(\tau_{1}\right)=v_{E}\left(\tau_{2}\right)=\left(n \omega_{n} D(E)\right) / 2$. Moreover, from (5.9), we have

$$
t_{1}-\tau_{1} \leq \frac{2}{n \omega_{n} D(E)} \int_{\tau_{1}}^{t_{1}} v_{E}(t) d t=\frac{2}{n D(E)} \int_{\tau_{1}}^{t_{1}} g^{\prime}(t) d t \leq \frac{2 g\left(t_{1}\right)}{n D(E)} \leq \frac{2}{n\left(2^{1 / n}-1\right)}
$$

and a similar estimate for $\tau_{2}-t_{2}$.
Let us now set $\widetilde{E}=E \cap\left\{x: \tau_{1}<x_{1}<\tau_{2}\right\}$. From the above estimate and (5.11), we have that $\tau_{2}-\tau_{1}<\beta(n)$. Moreover, (5.9), the definition of $\tau_{1}, \tau_{2}$ and (5.5) immediately yield

$$
\begin{equation*}
|\widetilde{E}| \geq|B|\left(1-2 \frac{D(E)}{2^{1 / n}-1}\right), \quad P(\widetilde{E}) \leq P(E)+n \omega_{n} D(E) . \tag{5.12}
\end{equation*}
$$

Let us now assume that $D(E)<\left(2^{1 / n}-1\right) / 4$ and set

$$
\sigma=\left(\frac{|B|}{|\widetilde{E}|}\right)^{1 / n}, \quad E^{\prime}=\sigma \widetilde{E}
$$

Clearly, $\left|E^{\prime}\right|=|B|$ and $E^{\prime}$ is contained in a strip $\left\{\tau_{1}^{\prime}<x_{1}<\tau_{2}^{\prime}\right\}$, with $\tau_{2}^{\prime}-\tau_{1}^{\prime} \leq$ $\beta^{\prime}$, where $\beta^{\prime}$ is a constant depending only on $n$. Let us now show that $E^{\prime}$ satisfies (5.3) for a suitable constant $C$ depending only on $n$.

To this aim, notice that since we are assuming $D(E)$ small, from (5.12) we get that $1 \leq \sigma \leq 1+C_{0} D(E)$, with $C_{0}=C_{0}(n)$. Thus, from (5.12), we get

$$
\begin{aligned}
P\left(E^{\prime}\right) & =\sigma^{n-1} P(\widetilde{E}) \leq \sigma^{n-1}(P(E)+P(B) D(E))=\sigma^{n-1} P(B)(1+2 D(E)) \\
& \leq P(B)(1+C(n) D(E))
\end{aligned}
$$

Hence, the second inequality in (5.3) follows. To prove the first inequality, let us denote by $p+B$ a ball such that $\lambda\left(E^{\prime}\right)=\left|E^{\prime} \Delta(p+B)\right|$. From the first inequality in (5.12), we then get

$$
\begin{aligned}
\lambda(E) & \leq|E \Delta((p / \sigma)+B)| \\
& \leq|E \Delta \widetilde{E}|+\left|\widetilde{E} \Delta\left((p / \sigma)+B_{1 / \sigma}\right)\right|+\left|\left((p / \sigma)+B_{1 / \sigma}\right) \Delta((p / \sigma)+B)\right| \\
& =|E \backslash \widetilde{E}|+\frac{\lambda\left(E^{\prime}\right)}{\sigma^{n}}+\left|B \backslash B_{1 / \sigma}\right| \\
& \leq C(n) D(E)+\lambda\left(E^{\prime}\right)+C(n)(1-\sigma) \leq \lambda\left(E^{\prime}\right)+C D(E)
\end{aligned}
$$

Thus, the set $E^{\prime}$ satisfies (5.5). Starting from this set, we may repeat the same construction with respect to the $x_{2}$ axis, thus getting a new set, still denoted by $E^{\prime}$, uniformly bounded with respect to the first two coordinate directions and satisfying (5.5) with a new constant, still depending only on $n$. Thus, the assertion follows by repeating this argument with respect to all the remaining coordinate directions.

## 6. Appendix: Some definitions and technical facts

In this section we first briefly recall the notion of distributional perimeter, then we collect a couple of technical facts concerning sets of finite perimeter. These properties are certainly known to the experts in the field, but we have not been able to find a precise reference in the literature. Therefore, for the benefit of the nonexpert reader and for the sake of completeness, we have decided to present here a complete proof of these facts.
6.1. Distributional perimeter. A Borel set $E$ is said to be a set of finite perimeter if the distributional derivative $D \chi_{E}$ of its characteristic function $\chi_{E}$ is an $\mathbb{R}^{n}$-valued Radon measure with finite total variation. If $E$ is a set of finite perimeter then the total variation $\left|D \chi_{E}\right|\left(\mathbb{R}^{n}\right)$ is denoted by $P(E)$, and it is called the perimeter of $E$. Moreover, if $A$ is any Borel set in $\mathbb{R}^{n}$, the perimeter of $E$ in $A$ is defined by setting $P(E \mid A)=\left|D \chi_{E}\right|(A)$.

This distributional notion of perimeter, due to Caccioppoli and De Giorgi, extends the usual one in the sense that $P(E)=\mathscr{H}^{n-1}(\partial E)$ whenever $\partial E$ is a

Lipschitz manifold. Notice that, by definition of the distributional derivative, for any smooth function $f$ with compact support,

$$
\int_{E} \nabla f(x) d x=-\int_{\mathbb{R}^{n}} f(x) d\left[D \chi_{E}\right](x) .
$$

Moreover, De Giorgi's Rectifiability Theorem (see Theorem 3.59 in [1]) states that $D \chi_{E}=\nu^{E} d \mathscr{H}^{n-1}\left\llcorner\partial^{*} E\right.$, i.e. that

$$
\int_{E} \nabla f(x) d x=-\int_{\partial^{*} E} f(x) \nu^{E}(x) d \mathscr{H}^{n-1}, \quad P(E \mid A)=\mathscr{H}^{n-1}\left(\partial^{*} E \cap A\right) .
$$

Here $\partial^{*} E$ is the reduced boundary of $E$, i.e. the set of those points $x \in \mathbb{R}^{n}$ such that

$$
\nu^{E}(x)=\lim _{r \rightarrow 0} \frac{D \chi_{E}(x+r B)}{\left|D \chi_{E}\right|(x+r B)}
$$

and $\left|\nu^{E}(x)\right|=1$. The density $\nu^{E}$ is called the generalized inner normal to $E$, and for all $x \in \partial^{*} E$
$\lim _{r \rightarrow 0} \frac{|E \cap(x+r B)|}{r^{n}|B|}=\frac{1}{2}, \quad \lim _{r \rightarrow 0} \frac{\left|\left\{y \in E \cap(x+r B):(y-x) \cdot \nu^{E}(x)<0\right\}\right|}{r^{n}}=0$.
The validity of the isoperimetric inequality (1.1) in the class of sets of finite perimeter, together with the characterization of the equality case mentioned in the introduction, are due to De Giorgi ([8], see also [9] and [13]).
6.2. Some technical facts about sets of finite perimeter. In the following, we shall denote the generic point $x$ of $\mathbb{R}^{n}$ also by $(t, z)$, where $t \in \mathbb{R}, z \in \mathbb{R}^{n-1}$. Given a set of finite perimeter $E$, we denote by $v_{E}(t)$ the $\mathscr{H}^{n-1}$ measure of its section $E_{t}=\{(t, z) \in E\}$. The same symbol $F_{t}$ will be used to denote the sections of any set $F$.

Recalling that the reduced boundary $\partial^{*} E$ of a set of finite perimeter is a $\mathscr{H}^{n-1}$-rectifiable set (see [1, Th. 3.59]), from the coarea formula for rectifiable sets ( $\left[1\right.$, Th. 2.39 and Rem. 2.40]) we get that if $g: \mathbb{R}^{n} \rightarrow[0,+\infty]$ is a Borel function, then

$$
\begin{equation*}
\int_{\partial^{*} E} g(x) \sqrt{1-\left|\nu_{1}^{E}(x)\right|^{2}} d \mathscr{H}^{n-1}(x)=\int_{-\infty}^{+\infty} \int_{\left(\partial^{*} E\right)_{t}} g(t, z) d \mathscr{H}^{n-2}(z) d t \tag{6.1}
\end{equation*}
$$

The first result we present is the version in codimension $n-1$ of a similar result proved in codimension 1 in [6] (see Lemmas 3.1, 3.2 and Proposition 1.2).

To this aim, we recall that by definition of generalized inner normal to a set of finite perimeter we have that for all $\psi \in C_{0}^{1}\left(\mathbb{R}^{n}\right)$ and $i=1, \ldots, n$,

$$
\begin{equation*}
\int_{E} \frac{\partial \psi}{\partial x_{i}} d x=-\int_{\partial^{*} E} \psi \nu_{i}^{E} d \mathscr{H}^{n-1} \tag{6.2}
\end{equation*}
$$

Theorem 6.1. Let $E$ be a set of finite perimeter, with finite measure. Then $v_{E} \in B V(\mathbb{R})$. Moreover, if

$$
\begin{equation*}
\mathscr{H}^{n-1}\left(\left\{x \in \partial^{*} E: \nu_{1}^{E}= \pm 1\right\}=0\right. \tag{6.3}
\end{equation*}
$$

then $v_{E} \in W^{1,1}(\mathbb{R})$ and for $\mathscr{H}^{1}$-a.e. $t \in \mathbb{R}$

$$
\begin{equation*}
v_{E}^{\prime}(t)=\int_{\left(\partial^{*} E\right)_{t}} \frac{\nu_{1}^{E}(t, z)}{\sqrt{1-\left|\nu_{1}^{E}(t, z)\right|^{2}}} d \mathscr{H}^{n-2}(z) \tag{6.4}
\end{equation*}
$$

Proof. The fact that $v_{E} \in L^{1}(\mathbb{R})$ is a simple consequence of Fubini's theorem and of the assumption $|E|<\infty$. To prove that the distributional derivative of $v_{E}$ is a Radon measure, let us fix $\varphi \in C_{0}^{1}(\mathbb{R})$ and take a sequence $\psi_{h}$ of nonnegative functions from $C_{0}^{1}\left(\mathbb{R}^{n-1}\right)$, such that $\psi_{h}(z) \nearrow 1$. Then, by Fubini's theorem and (6.2), we get

$$
\begin{align*}
\int_{-\infty}^{+\infty} v_{E}(t) \varphi^{\prime}(t) d t & =\int_{-\infty}^{+\infty} \varphi^{\prime}(t) \int_{E_{t}} d z d t=\lim _{h \rightarrow \infty} \int_{E} \varphi^{\prime}(t) \psi_{h}(z) d x  \tag{6.5}\\
& =-\lim _{h \rightarrow \infty} \int_{\partial^{*} E} \varphi(t) \psi_{h}(z) \nu_{1}^{E}(t, z) d \mathscr{H}^{n-1} \\
& =-\int_{\partial^{*} E} \varphi(t) \nu_{1}^{E}(t, z) d \mathscr{H}^{n-1}
\end{align*}
$$

Taking the supremum of the left-hand side integral among all $\varphi$ such that $\|\varphi\|_{\infty} \leq 1$, we get immediately that $D v_{E}$ is a Radon measure. Notice that once we know this, by using (6.5) and the definition of distributional derivative, we obtain that for all $\varphi \in C_{0}^{1}(\mathbb{R})$

$$
\begin{equation*}
\int_{-\infty}^{+\infty} \varphi(t) d\left[D v_{E}\right](t)=\int_{\partial^{*} E} \varphi(t) \nu_{1}^{E}(t, z) d \mathscr{H}^{n-1} \tag{6.6}
\end{equation*}
$$

From this equation we get immediately that if $G$ is an open subset in $\mathbb{R}$ then

$$
\begin{equation*}
\left|D v_{E}\right|(G) \leq \int_{\partial^{*} E \cap\left(G \times \mathbb{R}^{n-1}\right)}\left|\nu_{1}^{E}(x)\right| d \mathscr{H}^{n-1}(x) . \tag{6.7}
\end{equation*}
$$

Again, we can deduce by approximation that the same inequality holds for any Borel set $G$ in $\mathbb{R}$.

Let us now assume that (6.3) holds and let us apply (6.7) to a null Borel set $G$. Since now $\sqrt{1-\left|\nu_{1}^{E}(x)\right|^{2}}>0$ for $\mathscr{H}^{n-1}$-a.e. $x \in \partial^{*} E$, from the co-area
formula (6.1) we get

$$
\begin{aligned}
\left|D v_{E}\right|(G) & \leq \int_{\partial^{*} E \cap\left(G \times \mathbb{R}^{n-1}\right)}\left|\nu_{1}^{E}\right| d \mathscr{H}^{n-1} \\
& =\int_{\partial^{*} E \cap\left(G \times \mathbb{R}^{n-1}\right)} \frac{\left|\nu_{1}^{E}\right|}{\sqrt{1-\left|\nu_{1}^{E}\right|^{2}}} \sqrt{1-\left|\nu_{1}^{E}\right|^{2}} d \mathscr{H}^{n-1} \\
& =\int_{G} d t \int_{\left(\partial^{*} E\right)_{t}} \frac{\left|\nu_{1}^{E}\right|}{\sqrt{1-\left|\nu_{1}^{E}\right|^{2}}} d \mathscr{H}^{n-2}=0 .
\end{aligned}
$$

thus proving that $v_{E} \in W^{1,1}(\mathbb{R})$. To conclude the proof, let us fix $\varphi \in C_{0}^{1}(\mathbb{R})$ and use (6.6) and the co-area formula again to get

$$
\begin{aligned}
\int_{-\infty}^{+\infty} \varphi(t) v_{E}^{\prime}(t) d t & =\int_{\partial^{*} E} \varphi(t) \nu_{1}^{E}(t, z) d \mathscr{H}^{n-1} \\
& =\int_{\partial^{*} E} \frac{\varphi(t)\left|\nu_{1}^{E}(t, z)\right|}{\sqrt{1-\left|\nu_{1}^{E}\right|^{2}}} \sqrt{1-\left|\nu_{1}^{E}\right|^{2}} d \mathscr{H}^{n-1} \\
& =\int_{-\infty}^{+\infty} \varphi(t) \int_{\left(\partial^{*} E\right)_{t}} \frac{\nu_{1}^{E}(t, z)}{\sqrt{1-\left|\nu_{1}^{E}(t, z)\right|^{2}}} d \mathscr{H}^{n-2}(z) d t .
\end{aligned}
$$

Hence (6.4) follows, by the arbitrariness of $\varphi$.
The next result is an extension to our situation of a well known result by Vol'pert, concerning one-dimensional sections of sets of finite perimeter (see [22] and also [1, Th. 3.108]).

Theorem 6.2. Let $E$ be a set of finite perimeter in $\mathbb{R}^{n}$. Then, for $\mathscr{H}^{1}$ a.e. $t \in \mathbb{R}$ the section $E_{t}$ is a set of finite perimeter in $\mathbb{R}^{n-1}$. Moreover, the $\partial^{*} E_{t}$ coincide with $\left(\partial^{*} E\right)_{t}$ up to a set of zero $\mathscr{H}^{n-2}$ measure.

Proof. Let us denote by $u_{h}$ a sequence of mollifiers of $\chi_{E}$. Then, $u_{h} \rightarrow \chi_{E}$ in $L_{\text {loc }}^{1}\left(\mathbb{R}^{n}\right)$ and it is well known that

$$
\begin{equation*}
\int_{\mathbb{R}^{n}}\left|\nabla u_{h}\right| d x \rightarrow\left|D \chi_{E}\right|\left(\mathbb{R}^{n}\right)=P(E), \quad \int_{\mathbb{R}^{n}}\left|\nabla_{z} u_{h}\right| d x \rightarrow\left|D_{z} \chi_{E}\right|\left(\mathbb{R}^{n}\right) \tag{6.8}
\end{equation*}
$$

where $\nabla_{z}=\left(\nabla_{x_{2}}, \ldots, \nabla_{x_{n}}\right)$ and $D_{z}$ is defined similarly. Moreover, by Fubini's theorem we get that for any compact subset $K \subseteq \mathbb{R}^{n-1}$ and any $R>0$

$$
\lim _{h \rightarrow \infty} \int_{K} \int_{-R}^{R}\left|u_{h}(t, z)-\chi_{E}(t, z)\right| d t d z=0 .
$$

From this equality, by a diagonalization argument, it follows that up to a (not relabeled) subsequence, $u_{h}(t, \cdot) \rightarrow \chi_{E}(t, \cdot)=\chi_{E_{t}}$ in $L_{\text {loc }}^{1}\left(\mathbb{R}^{n-1}\right)$ for $\mathscr{H}^{1}$-a.e. $t \in \mathbb{R}$. Hence, by the lower semicontinuity of the total variation, we have

$$
P_{n-1}\left(E_{t}\right) \leq \liminf _{h \rightarrow \infty} \int_{\mathbb{R}^{n-1}}\left|\nabla_{z} u_{h}(t, z)\right| d z,
$$

where by $P_{n-1}$ we denote the perimeter in $\mathbb{R}^{n-1}$. Integrating this inequality and using Fatou's lemma we have, from (6.8),

$$
\begin{align*}
\int_{\mathbb{R}} P_{n-1}\left(E_{t}\right) d t & \leq \int_{\mathbb{R}} \liminf _{h \rightarrow \infty} \int_{\mathbb{R}^{n-1}}\left|\nabla_{z} u_{h}(t, z)\right| d z d t  \tag{6.9}\\
& \leq \liminf _{h \rightarrow \infty} \int_{\mathbb{R}^{n}}\left|\nabla_{z} u_{h}(x)\right| d x=\left|D_{z} \chi_{E}\right|\left(\mathbb{R}^{n}\right)<\infty
\end{align*}
$$

Hence, we get in particular that, for a.e. $t \in \mathbb{R}, P_{n-1}\left(E_{t}\right)<\infty$; i.e., $E_{t}$ has finite perimeter.

Let us fix $\varphi \in C_{0}^{1}\left(\mathbb{R}^{n} ; \mathbb{R}^{n-1}\right)$ with $\|\varphi\| \leq 1$. We have

$$
\begin{aligned}
\int_{\mathbb{R}^{n}} \varphi(x) d\left[D_{z} \chi_{E}\right](x) & =-\int_{\mathbb{R}^{n}} \chi_{E} \operatorname{div}_{z} \varphi(x) d x \\
& =-\int_{\mathbb{R}} \int_{\mathbb{R}^{n-1}} \chi_{E_{t}}(z) \operatorname{div}_{z} \varphi(t, z) d z d t \leq \int_{\mathbb{R}} P_{n-1}\left(E_{t}\right) d t
\end{aligned}
$$

and, passing to the supremum over all $\varphi$, we get that

$$
\begin{equation*}
\left|D_{z} \chi_{E}\right|\left(\mathbb{R}^{n}\right) \leq \int_{\mathbb{R}} P_{n-1}\left(E_{t}\right) d t \tag{6.10}
\end{equation*}
$$

Therefore from this inequality and from (6.9) we conclude that in (6.10) the equality holds. With exactly the same argument it can be proved that if $\Omega$ is any open set in $\mathbb{R}^{n}$, then

$$
\left|D_{z} \chi_{E}\right|(\Omega)=\int_{\mathbb{R}} P_{n-1}\left(E_{t} \mid \Omega_{t}\right) d t
$$

and, by a simple approximation argument, that this equality still holds if we replace $\Omega$ by any Borel set $G \subseteq \mathbb{R}^{n}$,

$$
\left|D_{z} \chi_{E}\right|(G)=\int_{\mathbb{R}} P_{n-1}\left(E_{t} \mid G_{t}\right) d t
$$

Since $\sqrt{1-\left|\nu_{1}^{E}\right|^{2}}=\left|\nu_{z}^{E}\right|=\frac{d\left|D_{z} \chi_{E}\right|}{d\left|D \chi_{E}\right|}=\frac{d\left|D_{z} \chi_{E}\right|}{d \mathscr{H}^{n-1}\left\llcorner\partial^{*} E\right.}$, this equality can be written in the form

$$
\int_{\partial^{*} E \cap G} \sqrt{1-\left|\nu_{1}^{E}(x)\right|^{2}} d \mathscr{H}^{n-1}(x)=\int_{\mathbb{R}} \mathscr{H}^{n-2}\left(\partial^{*} E_{t} \cap G\right) d t
$$

Since this equality holds for any Borel set we may conclude that if $f: \mathbb{R}^{n} \rightarrow \mathbb{R}$ is any bounded Borel function, then

$$
\begin{equation*}
\int_{\partial^{*} E} f(x) \sqrt{1-\left|\nu_{1}^{E}(x)\right|^{2}} d \mathscr{H}^{n-1}(x)=\int_{\mathbb{R}} \int_{\partial^{*} E_{t}} f(t, z) d \mathscr{H}^{n-2}(z) d t \tag{6.11}
\end{equation*}
$$

On the other hand, from the co-area formula (6.1) we have also that

$$
\begin{equation*}
\int_{\partial^{*} E} f(x) \sqrt{1-\left|\nu_{1}^{E}(x)\right|^{2}} d \mathscr{H}^{n-1}(x)=\int_{\mathbb{R}} \int_{\left(\partial^{*} E\right)_{t}} f(t, z) d \mathscr{H}^{n-2}(z) d t \tag{6.12}
\end{equation*}
$$

Choosing $f(t, z)=\varphi(t) \psi(z)$, and comparing the two equations (6.11), (6.12) we then easily get that for $\mathscr{H}^{1}$-a.e. $t \in \mathbb{R}$

$$
\mathscr{H}^{n-2}\left\llcorner\partial^{*} E_{t}=\mathscr{H}^{n-2}\left\llcorner\left(\partial^{*} E\right)_{t} .\right.\right.
$$

Hence, the result follows.
We conclude this section with the proof of a result which has been used in Section 3.

Theorem 6.3. Let $E$ be a set of finite perimeter with finite measure. If $E$ satisfies (6.3), then the same is true for $E^{*}$.

Proof. First, notice that if $F$ is any set of finite perimeter, by applying (6.11) to the characteristic function of the set $C_{F}=\left\{x \in \partial^{*} F:\left|\nu_{1}^{F}(x)\right|=1\right\}$, we have that $\mathscr{H}^{1}\left(\left\{t \in \mathbb{R}: \mathscr{H}^{n-2}\left(C_{F} \cap \partial^{*} F_{t}\right)>0\right\}\right)=0$. From this equality, using Theorem 6.1 and arguing exactly as in the proof of [6, Lemma 3.2], we get that for $\mathscr{H}^{1}$-a.e. $t \in \mathbb{R}$

$$
v_{F}^{\prime}(t)=\int_{\left(\partial^{*} F\right)_{t}} \frac{\nu_{1}^{F}(t, z)}{\sqrt{1-\left|\nu_{1}^{F}(t, z)\right|^{2}}} d \mathscr{H}^{n-2}(z)
$$

where $v_{F}^{\prime}$ denotes the absolutely continuous part of the distributional derivative $D v_{F}$. From this formula, the same argument used to prove (3.2) immediately yields

$$
\begin{equation*}
P(F) \geq \mathscr{H}^{n-1}\left(\partial^{*} F \backslash C_{F}\right) \geq \int_{-\infty}^{+\infty} \sqrt{p_{F}(t)^{2}+v_{F}^{\prime}(t)^{2}} d t \tag{6.13}
\end{equation*}
$$

In particular, from this inequality we get that, if $E$ is any set of finite perimeter,

$$
\begin{equation*}
P\left(E^{*}\right) \geq \mathscr{H}^{n-1}\left(\partial^{*} E^{*} \backslash C_{E^{*}}\right) \geq \int_{-\infty}^{+\infty} \sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+v_{E}^{\prime}(t)^{2}} d t \tag{6.14}
\end{equation*}
$$

where $\tau$ is as in (3.3). On the other hand, if $E$ satisfies (6.3), then $v_{E} \in$ $W^{1,1}(\mathbb{R})$; hence there exists a sequence of smooth nonnegative functions $v_{h}$, with compact supports, converging to $v_{E}$ in $W^{1,1}(\mathbb{R})$. Denoting by $E_{h}$ the axially symmetric set defined by $v_{E_{h}}=v_{h}$, by the lower semicontinuity of perimeters and the fact that (3.3) clearly holds for each set $E_{h}$, we have

$$
\begin{aligned}
P\left(E^{*}\right) \leq \liminf _{h \rightarrow \infty} P\left(E_{h}\right) & =\lim _{h \rightarrow \infty} \int_{-\infty}^{+\infty} \sqrt{\tau v_{h}(t)^{\frac{2 n-4}{n-1}}+v_{h}^{\prime}(t)^{2}} d t \\
& =\int_{-\infty}^{+\infty} \sqrt{\tau v_{E}(t)^{\frac{2 n-4}{n-1}}+v_{E}^{\prime}(t)^{2}} d t
\end{aligned}
$$

From this inequality and (6.14) we immediately get that $\mathscr{H}^{n-1}\left(C_{E^{*}}\right)=0$, i.e. $E^{*}$ satisfies (6.3).

Notice that (6.13) generalizes (3.2) to any set of finite perimeter.
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