Calderón’s inverse conductivity problem in the plane

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Abstract

We show that the Dirichlet to Neumann map for the equation \( \nabla \cdot \sigma \nabla u = 0 \) in a two-dimensional domain uniquely determines the bounded measurable conductivity \( \sigma \). This gives a positive answer to a question of A. P. Calderón from 1980. Earlier the result has been shown only for conductivities that are sufficiently smooth. In higher dimensions the problem remains open.

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1. Introduction and outline of the method

Suppose that \( \Omega \subset \mathbb{R}^n \) is a bounded domain with connected complement and \( \sigma : \Omega \to (0, \infty) \) is measurable and bounded away from zero and infinity. Given the boundary values \( \phi \in H^{1/2}(\partial \Omega) \) let \( u \in H^1(\Omega) \) be the unique solution to

\[
\begin{align*}
\nabla \cdot \sigma \nabla u &= 0 \text{ in } \Omega, \\
u|_{\partial \Omega} &= \phi \in H^{1/2}(\partial \Omega).
\end{align*}
\]

This so-called conductivity equation describes the behavior of the electric potential in a conductive body.

*The research of both authors is supported by the Academy of Finland.
In 1980 A. P. Calderón [11] posed the problem whether one can recover the conductivity $\sigma$ from the boundary measurements, i.e. from the Dirichlet to Neumann map

$$\Lambda_{\sigma} : \phi \mapsto \sigma \frac{\partial u}{\partial \nu} |_{\partial \Omega}.$$ 

Here $\nu$ is the unit outer normal to the boundary and the derivative $\sigma \partial u / \partial \nu$ exists as an element of $H^{-1/2}(\partial \Omega)$, defined by

$$\langle \sigma \frac{\partial u}{\partial \nu}, \psi \rangle = \int_{\Omega} \sigma \nabla u \cdot \nabla \psi \, dm,$$

where $\psi \in H^1(\Omega)$ and $dm$ denotes the Lebesgue measure.

The aim of this paper is to give a positive answer to Calderón’s question in dimension two. More precisely, we prove

**Theorem 1.** Let $\Omega \subset \mathbb{R}^2$ be a bounded, simply connected domain and $\sigma_i \in L^\infty(\Omega)$, $i = 1, 2$. Suppose that there is a constant $c > 0$ such that $c^{-1} \leq \sigma_i \leq c$. If

$$\Lambda_{\sigma_1} = \Lambda_{\sigma_2}$$

then $\sigma_1 = \sigma_2$.

Note, in particular, that no regularity is required for the boundary. Our approach to Theorem 1 yields, in principle, also a method to construct $\sigma$ from the Dirichlet to Neumann operator $\Lambda_{\sigma}$. For this see Section 8. The case of an anisotropic conductivity has been fully analyzed in the follow-up paper with Lassas [6].

Calderón faced the above problem while working as an engineer in Argentina in the 1950’s. He was able to show that the linearized problem at constant conductivities has a unique solution. Decades later Alberto Grünbaum convinced Calderón to publish his result [11]. The problem rises naturally in geophysical prospecting. Indeed, the Slumberger–Doll company was founded to find oil by using electromagnetic methods.

In medical imaging Calderón’s problem is known as *Electrical Impedance Tomography*. It has been proposed as a valuable diagnostic tool especially for detecting pulmonary emboli [12]. One may find a review for medical applications in [13]; for statistical methods in electrical impedance tomography see [17].

That $\Lambda_{\sigma}$ uniquely determines $\sigma$ was established in dimension three and higher for smooth conductivities by J. Sylvester and G. Uhlmann [30] in 1987. In dimension two, A. Nachman [22] produced in 1995 a uniqueness result for conductivities with two derivatives. Earlier, the problem was solved for piecewise analytic conductivities by Kohn and Vogelius [19], [20] and the generic uniqueness was established by Sun and Uhlmann [29].
The regularity assumptions have since been relaxed by several authors (cf. [23], [24]) but the original problem of Calderón has still remained unsolved. In dimensions three and higher the uniqueness is known for conductivities in \( W^{3/2, \infty}(\Omega) \), see [26], and in two dimensions the best result so far was \( \sigma \in W^{1, p}(\Omega), \ p > 2 \), [10].

The original approach in [30] and [22] was to reduce the conductivity equation (1.1) to the Schrödinger equation by substituting \( v = \sigma^{1/2}u \). Indeed, after such a substitution \( v \) satisfies
\[
\Delta v - qv = 0
\]
where \( q = \sigma^{-1/2} \Delta \sigma^{1/2} \). This explains why in this method one needs two derivatives. For the numerical implementation of [22] see [27].

Following the ideas of Beals and Coifman [8], Brown and Uhlmann [10] found a first order elliptic system equivalent to (1.1). Indeed, by denoting
\[
\left( \begin{array}{c}
v \\
w
\end{array} \right) = \sigma^{1/2} \left( \begin{array}{c}
\partial u \\
\partial \bar{u}
\end{array} \right)
\]
one obtains the system
\[
D \left( \begin{array}{c}
v \\
w
\end{array} \right) = Q \left( \begin{array}{c}
v \\
w
\end{array} \right),
\]
where
\[
D = \left( \begin{array}{cc}
\bar{\partial} & 0 \\
0 & \partial
\end{array} \right), \quad Q = \left( \begin{array}{cc}
0 & q \\
q & 0
\end{array} \right)
\]
and \( q = -\frac{1}{2} \partial \log \sigma \). This allowed Brown and Uhlmann to work with conductivities with only one derivative. Note however, that the assumption \( \sigma \in W^{1, p}(\Omega), \ p > 2 \), necessary in [10], implies that \( \sigma \) is Hölder continuous. From the viewpoint of applications this is still not satisfactory. Our starting point is to replace (1.1) with an elliptic equation that does not require any differentiability of \( \sigma \).

We will base our argument on the fact that if \( u \in H^1(\Omega) \) is a real solution of (1.1) then there exists a real function \( v \in H^1(\Omega) \), called the \( \sigma \)-harmonic conjugate of \( u \), such that \( f = u + iv \) satisfies the \( \mathbb{R} \)-linear Beltrami equation
\[
\partial f = \mu \overline{\partial f},
\]
where \( \mu = (1 - \sigma)/(1 + \sigma) \). In particular, note that \( \mu \) is real-valued. The assumptions for \( \sigma \) imply that \( \| \mu \|_{L^\infty} \leq \kappa < 1 \), and the symbol \( \kappa \) will retain this role throughout the paper.

The structure of the paper is the following: Since the \( \sigma \)-harmonic conjugate is unique up to a constant we can define the \( \mu \)-Hilbert transform \( \mathcal{H}_\mu : H^{1/2}(\partial \Omega) \to H^{1/2}(\partial \Omega) \) by
\[
\mathcal{H}_\mu : u|_{\partial \Omega} \mapsto v|_{\partial \Omega}.
\]
We show in Section 2 that the Dirichlet to Neumann map $\Lambda_\sigma$ uniquely determines $\mathcal{H}_\mu$ and vice versa. Theorem 1 now implies the surprising fact that $\mathcal{H}_\mu$ uniquely determines $\mu$ in equation (1.4) in the whole domain $\Omega$.

Recall that a function $f \in H^1_{\text{loc}}(\Omega)$ satisfying (1.4) is called a \textit{quasiregular mapping}; if it is also a homeomorphism then it is called \textit{quasiconformal}. These have a well established theory, cf. [2], [5], [14], [21], that we will employ at several points in the paper. The $H^1_{\text{loc}}$-solutions $f$ to (1.4) are automatically continuous and admit a factorization $f = \psi \circ H$, where $\psi$ is $\mathbb{C}$-analytic and $H$ is a quasiconformal homeomorphism. Solutions with less regularity may not share these properties [14]. The basic tools to deal with the Beltrami equation are two linear operators, the Cauchy transform $P = \frac{\partial}{\partial z}$ and the Beurling transform $S = \frac{\partial}{\partial \bar{z}}$. In Section 3 we recall the basic properties of these operators with some useful preliminary results.

It is not difficult to see, cf. Section 2, that we can assume $\Omega = \mathbb{D}$, the unit disk of $\mathbb{C}$, and that outside $\Omega$ we can set $\sigma \equiv 1$, i.e., $\mu \equiv 0$.

In Section 4 we establish the existence of the geometric optics solutions $f = f_\mu$ of (1.4) that have the form

$$f_\mu(z, k) = e^{ikz} M_\mu(z, k),$$

where

$$M_\mu(z, k) = 1 + O\left(\frac{1}{z}\right) \text{ as } |z| \to \infty.$$  

(1.6)

As in the smooth case these solutions obey a $\bar{\partial}$-equation also in the $k$ variable. However, their asymptotics as $|k| \to \infty$ are now more subtle and considerably more difficult to handle.

It turns out that it is instructive to consider the conductivities $\sigma$ and $\sigma^{-1}$, or equivalently the Beltrami coefficients $\mu$ and $-\mu$, simultaneously. By defining

$$h_+ = \frac{1}{2}(f_\mu + f_{-\mu}), \quad h_- = \frac{i}{2}(f_\mu - f_{-\mu})$$

(1.7)

we show in Section 5 that with respect to the variable $k$, $h_+$ and $h_-$ satisfy the equations

$$\partial_k h_+ = \tau_\mu h_-, \quad \partial_k h_- = \tau_\mu h_+$$

(1.8)

where the scattering coefficient $\tau_\mu = \tau_\mu(k)$ is defined by

$$\tau_\mu(k) = \frac{i}{4\pi} \int \partial_{\bar{z}}(M_\mu - M_{-\mu}) \, dz \wedge d\bar{z}.$$  

(1.9)

The remarkable fact in the equations (1.8) is that the coefficient $\tau_\mu(k)$ does not depend on the space variable $z$; the idea of using such a phenomenon is due to Beals and Coifman [8] and in connection with the Dirichlet to Neumann operator to Nachman [22]. In Section 6 we show that $\Lambda_\sigma$ uniquely determines
the scattering coefficient $\tau_\mu(k)$ as well as the geometric optics solutions $f_\mu$ and $f_{-\mu}$ outside $\mathbb{D}$.

The crucial problem in the proof of Theorem 1 is the behavior of the function $M_\mu(z,k) - 1 = e^{-ikz}f_\mu(z,k) - 1$ with respect to the $k$-variable. In the case of [22] and [10] the behaviour is roughly like $|k|^{-1}$. In the $L^\infty$-case we cannot expect such good behavior. Instead, we can show that $M_\mu(z,k)$ grows at most subexponentially in $k$. This is the key tool to our argument and it takes a considerable effort to prove it. Precisely, we show in Section 7 that

$$f_\mu(z,k) = \exp(ik\varphi(z,k))$$

where $\varphi$ is a quasiconformal homeomorphism in the $z$-variable and satisfies the nonlinear Beltrami equation

$$(1.10) \quad \partial_z \varphi = -\frac{k}{k} \mu(z) e^{-k(\varphi(z))} \partial_{\bar{z}} \varphi$$

with the boundary condition

$$(1.11) \quad \varphi(z) = z + O\left(\frac{1}{z}\right)$$

at infinity. Here the unimodular function $e_k$ is given by

$$(1.12) \quad e_k(z) = e^{i(kz + k\tau)}.$$

The main result in Section 7 is that the unique solution of (1.10) and (1.11) satisfies

$$(1.13) \quad \varphi(z,k) - z \to 0 \quad \text{as} \quad |k| \to \infty,$$

uniformly in $z$.

Section 8 is devoted to the proof of Theorem 1. Since

$$\mu = \partial \bar{f}_\mu / \partial f_{\mu}$$

and $\partial f$ for a nonconstant quasiregular map $f$ can vanish only in a set of Lebesgue measure zero, we are reduced to determining the function $f_\mu$ in the interior of $\mathbb{D}$. As said before, we already know these functions outside of $\mathbb{D}$. To solve this problem we introduce the so-called transport matrix that transforms the solutions outside $\mathbb{D}$ to solutions inside. We show that this matrix is uniquely determined by $\Lambda_\sigma$. At this point one may work either with equation (1.1) or equation (1.4). We chose to go back to the conductivity equation since it slightly simplifies the formulas. More precisely, we set

$$(1.14) \quad u_1 = h_+ - ih_- \quad \text{and} \quad u_2 = i(h_+ + ih_-).$$

Then $u_1$ and $u_2$ are complex solutions of the conductivity equations

$$(1.15) \quad \nabla \cdot \sigma \nabla u_1 = 0 \quad \text{and} \quad \nabla \cdot \frac{1}{\sigma} \nabla u_2 = 0,$$
respectively, and of the $\partial_k$-equation

\begin{equation}
\frac{\partial}{\partial k} u_j = -i \tau_\mu(k) \overline{u_j}, \quad j = 1, 2,
\end{equation}

with the asymptotics $u_1 = e^{ikz}(1 + \mathcal{O}(1/z))$ and $u_2 = e^{ikz}(i + \mathcal{O}(1/z))$ in the $z$-variable. Uniqueness of (1.15) with these asymptotics gives that in the smooth case $u_1$ is exactly the exponentially growing solution of [22].

We then choose a point $z_0 \in \mathbb{C}, |z_0| > 1$. It is possible to write for each $z, k \in \mathbb{C}$

\begin{equation}
\begin{aligned}
  u_1(z, k) &= a_1 u_1(z_0, k) + a_2 u_2(z_0, k), \\
  u_2(z, k) &= b_1 u_1(z_0, k) + b_2 u_2(z_0, k)
\end{aligned}
\end{equation}

where $a_j = a_j(z, z_0; k)$ and $b_j = b_j(z, z_0; k)$ are real-valued. The transport matrix $T_{z, z_0}^\sigma(k)$ is now defined by

\begin{equation}
T_{z, z_0}^\sigma(k) = \begin{pmatrix} a_1 & a_2 \\ b_1 & b_2 \end{pmatrix}.
\end{equation}

It is an invertible $2 \times 2$ real matrix depending on $z, z_0$ and $k$. The proof of Theorem 1 is thus reduced to

**Theorem 2.** Assume that $\Lambda_\sigma = \Lambda_{\tilde{\sigma}}$ for two $L^\infty$-conductivities $\sigma$ and $\tilde{\sigma}$. Then for all $z, k \in \mathbb{C}$ and $|z_0| > 1$ the corresponding transport matrices $T_{z, z_0}^\sigma(k)$ and $T_{z, z_0}^{\tilde{\sigma}}(k)$ are equal.

The idea behind the proof is to use the Beals-Coifman method in an efficient manner and to show that the functions

\begin{equation}
\alpha(k) = a_1(k) + ia_2(k) \quad \text{and} \quad \beta(k) = b_1(k) + ib_2(k)
\end{equation}

both satisfy, with respect to the parameter $k$, the Beltrami equation

\begin{equation}
\partial_k \alpha = \nu_{z_0}(k) \partial_k \overline{\alpha}.
\end{equation}

Here the coefficient

\begin{equation}
\nu_{z_0}(k) = i \frac{h_-(z_0, k)}{h_+(z_0, k)}
\end{equation}

is determined by the data as proved in Section 6. Moreover it satisfies

\[|\nu_{z_0}(k)| \leq q < 1,
\]

where the number $q$ is independent of $k$ (or $z$). These facts and the subexponential growth of the solutions serve as the key elements for the proof of Theorem 2.
2. The Beltrami equation and the Hilbert transform

In a general domain $\Omega$ we identify $H^{1/2}(\partial \Omega) = H^1(\Omega)/H^1_0(\Omega)$. When $\partial \Omega$ has enough regularity, trace theorems and extension theorems [31] readily yield the standard interpretation of $H^{1/2}(\partial \Omega)$. The Dirichlet condition (1.2) is consequently defined in the Sobolev sense, requiring that $u - \phi \in H^1_0(\Omega)$ for the element $\phi \in H^{1/2}(\partial \Omega)$. Furthermore, $H^{-1/2}(\partial \Omega) = H^{1/2}(\partial \Omega)^*$ and via (1.3) it is then clear that $\Lambda_\sigma$ becomes a well-defined and bounded operator from $H^{1/2}(\partial \Omega)$ to $H^{-1/2}(\partial \Omega)$.

In this setup Theorem 1 quickly reduces to the case where the domain $\Omega$ is the unit disk. In fact, let $\Omega$ be a simply connected domain with $\overline{\Omega} \subset \mathbb{D}$ and let $\sigma$ and $\overline{\sigma}$ be two $L^\infty$-conductivities on $\Omega$ with $\Lambda_\sigma = \Lambda_{\overline{\sigma}}$. Continue both conductivities as the constant 1 outside $\Omega$ to obtain new $L^\infty$-conductivities $\sigma_0$ and $\overline{\sigma}_0$. Given $\phi \in H^{1/2}(\partial \mathbb{D})$, let $u_0 \in H^1(\mathbb{D})$ be the solution to the Dirichlet problem $\nabla \cdot \sigma_0 \nabla u_0 = 0$ in $\mathbb{D}$, $u_0|_{\partial \mathbb{D}} = \phi$. Assume also that $\tilde{u} \in H^1(\Omega)$ is the solution to

$$\nabla \cdot \overline{\sigma_0} \nabla \tilde{u} = 0 \quad \text{in} \quad \Omega, \quad \tilde{u} - u_0 \in H^1_0(\Omega).$$

Then $\tilde{u}_0 = \tilde{u} \chi_\Omega + u_0 \chi_{\mathbb{D}\setminus\Omega} \in H^1(\mathbb{D})$ since zero extensions of $H^1_0(\Omega)$ functions remain in $H^1$. Moreover, an application of the definition (1.3) to the condition $\Lambda_\sigma = \Lambda_{\overline{\sigma}}$ yields that $\tilde{u}_0$ satisfies

$$\nabla \cdot \overline{\sigma_0} \nabla \tilde{u}_0 = 0$$

in the disk $\mathbb{D}$. Since in $\mathbb{D}\setminus\Omega$ we have $u_0 \equiv \tilde{u}_0$ and $\sigma_0 \equiv \overline{\sigma}_0$, we obtain $\Lambda_{\overline{\sigma}_0} \phi = \Lambda_{\sigma_0} \phi$, and this holds for all $\phi \in H^{1/2}(\partial \mathbb{D})$. Thus if Theorem 1 holds for $\mathbb{D}$ we get $\overline{\sigma}_0 = \sigma_0$ and especially that $\overline{\sigma} = \sigma$.

From now on we assume that $\Omega = \mathbb{D}$, the unit disc in $\mathbb{C}$.

Let us then consider the complex analytic interpretation of (1.1). We will use the notation $\partial = \frac{1}{2}(\partial_x - i \partial_y)$ and $\overline{\partial} = \frac{1}{2}(\partial_x + i \partial_y)$; when clarity requires we may write $\overline{\partial} = \partial_{\overline{z}}$ or $\partial = \partial_z$. For derivatives with respect to the parameter $k$ we always use the notation $\partial_k$ and $\overline{\partial}_k$.

We start with a simple lemma:

**Lemma 2.1.** Assume $u \in H^1(\mathbb{D})$ is real-valued and satisfies the conductivity equation (1.1). Then there exists a function $v \in H^1(\mathbb{D})$, unique up to a constant, such that $f = u + iv$ satisfies the $\mathbb{R}$-linear Beltrami equation

$$\overline{\partial} f = \mu \overline{\partial} f,$$

where $\mu = (1 - \sigma)/(1 + \sigma)$.

Conversely, if $f \in H^1(\mathbb{D})$ satisfies (2.1) with an $\mathbb{R}$-valued $\mu$, then $u = \text{Re} \ f$ and $v = \text{Im} \ f$ satisfy

$$\nabla \cdot \sigma \nabla u = 0 \quad \text{and} \quad \nabla \cdot \frac{1}{\sigma} \nabla v = 0,$$

respectively, where $\sigma = (1 - \mu)/(1 + \mu)$. 

Proof. Denote by $w$ the vectorfield
\[ w = (-\sigma \partial_2 u, \sigma \partial_1 u) \]
where $\partial_1 = \partial/\partial x$ and $\partial_2 = \partial/\partial y$ for $z = x + iy \in \mathbb{C}$. Then by (1.1) the integrability condition $\partial_2 w_1 = \partial_1 w_2$ holds for the distributional derivatives. Therefore there exists $v \in H^1(\mathbb{D})$, unique up to a constant, such that
\begin{align*}
\partial_1 v &= -\sigma \partial_2 u, \\
\partial_2 v &= \sigma \partial_1 u.
\end{align*}
(2.3) (2.4)
A simple calculation shows that this is equivalent to (2.1).
\[ \square \]

We want to stress that every solution of (2.1) is also a solution of the standard $\mathbb{C}$-linear Beltrami equation
\[ \bar{\partial} f = \tilde{\mu} \partial f \]
(2.5)
but with a different $\mathbb{C}$-valued $\tilde{\mu}$ having, however, the same modulus as the old one. We note that the uniqueness properties of (2.1) and (2.5) are quite different (cf. [32], [5]) and that the conditions for $\sigma$ given in Theorem 1 imply the existence of a constant $0 \leq \kappa < 1$ such that
\[ |\mu(z)| \leq \kappa \]
holds for almost every $z \in \mathbb{C}$.

Since the function $v$ in Lemma 2.1 is defined only up to a constant we will normalize it by assuming
\[ \int_{\partial \mathbb{D}} v \, ds = 0. \]
(2.6)
This way we obtain a unique map $\mathcal{H}_\mu : H^{1/2}(\partial \mathbb{D}) \to H^{1/2}(\partial \mathbb{D})$ by setting
\[ \mathcal{H}_\mu : u |_{\partial \mathbb{D}} \mapsto v |_{\partial \mathbb{D}}. \]
(2.7)
The function $v$ satisfying (2.3), (2.4) and (2.6) is called the $\sigma$-harmonic conjugate of $u$ and $\mathcal{H}_\mu$ the Hilbert transform corresponding to equation (2.1).

Since $v$ is the real part of the function $g = -if$ satisfying $\bar{\partial} g = -\mu \partial g$, we have
\[ \mathcal{H}_\mu \circ \mathcal{H}_{-\mu} u = \mathcal{H}_{-\mu} \circ \mathcal{H}_\mu u = -u + \int_{\partial \mathbb{D}} u \, ds = -u + L(u) \]
(2.8)
where
\[ L(u) = \int_{\partial \mathbb{D}} u \, ds = \frac{1}{2\pi} \int_{\partial \mathbb{D}} u \, ds \]
is the average operator. In particular, $\mathcal{H}_{-\mu} = L - (\mathcal{H}_\mu + L)^{-1}$. 

So far we have only defined $\mathcal{H}_\mu(u)$ for real-valued $u$. By setting

$$H^\mu(u) = i H^{-\mu}(u)$$

we have extended the definition of $H^\mu(g)$ $\mathbb{R}$-linearly to all $\mathbb{C}$-valued $g \in H^{1/2}(\partial \mathbb{D})$. We also define $Q^\mu : H^{1/2}(\partial \mathbb{D}) \rightarrow H^{1/2}(\partial \mathbb{D})$ by

$$(2.9) \quad Q^\mu = \frac{1}{2} (I - i H^\mu).$$

Then $g \mapsto Q^\mu(g) - \frac{1}{2} \int_{\partial \mathbb{D}} g \, ds$ is a projection in $H^{1/2}(\partial \mathbb{D})$. In fact

$$(2.10) \quad Q^2\mu(g) = Q^\mu(g) - \frac{1}{4} \int_{\partial \mathbb{D}} g \, ds.$$ 

The proof of the following lemma is straightforward.

**Lemma 2.2.** If $g \in H^{1/2}(\partial \mathbb{D})$, the following conditions are equivalent:

a) $g = f |_{\partial \mathbb{D}}$, where $f \in H^1(\mathbb{D})$ and satisfies (2.1).

b) $Q^\mu(g)$ is a constant.

We close this section with

**Proposition 2.3.** The Dirichlet to Neumann map $\Lambda_\sigma$ uniquely determines $\mathcal{H}_\mu, \mathcal{H}_{-\mu}$ and $\Lambda_{\sigma^{-1}}$.

**Proof.** Choose the counter clockwise orientation for $\partial \mathbb{D}$ and denote by $\partial T$ the tangential (distributional) derivative on $\partial \mathbb{D}$ corresponding to this orientation. We will show for real-valued $u$ that

$$(2.11) \quad \partial T H^\mu(u) = \Lambda_\sigma(u)$$

holds in the weak sense. This will be enough as $\mathcal{H}_\mu$ uniquely determines $\mathcal{H}_{-\mu}$ by (2.8). Since $-\mu = (1 - \sigma^{-1})/(1 + \sigma^{-1})$ we also have $\Lambda_{\sigma^{-1}}(u) = \partial T H_{-\mu}(u)$. Note that the right-hand side of (2.11) is complex linear but the left-hand side is not.

By the definition of $\Lambda_\sigma$,

$$\int_{\partial \mathbb{D}} \varphi \Lambda_\sigma u \, ds = \int_{\mathbb{D}} \nabla \varphi \cdot \sigma \nabla u \, dm, \quad \varphi \in C^\infty(\mathbb{D}).$$

Thus, by (2.3), (2.4) and integration by parts, we get

$$\int_{\partial \mathbb{D}} \varphi \Lambda_\sigma u = \int_{\mathbb{D}} (\partial_1 \varphi \partial_2 v - \partial_2 \varphi \partial_1 v) \, dm$$

$$= - \int_{\partial \mathbb{D}} v \partial T \varphi \, ds$$

and (2.11) follows. \qed
3. Beltrami operators

The Beltrami differential equation (1.4) and its solutions are effectively governed and controlled by two basic linear operators, the Cauchy transform and the Beurling transform. Any analysis of (1.4) requires basic facts of these operators. We briefly recall those in this section.

The Cauchy transform

\[ P g(z) = -\frac{1}{\pi} \int_{C} \frac{g(\omega)}{\omega - z} \, dm(\omega) \]  

acts as the inverse operator to \( \partial \); \( P \partial g = \partial P g = g \) for \( g \in C_0^\infty(\mathbb{C}) \). We recall some mapping properties of \( P \) in appropriate Lebesgue, Sobolev and Lipschitz spaces. Below we denote

\[ L^p(\Omega) = \left\{ g \in L^p(\mathbb{C}) \mid g|_{\mathbb{C}\setminus\Omega} \equiv 0 \right\}. \]

**Proposition 3.1.** Let \( \Omega \subset \mathbb{C} \) be a bounded domain and let \( 1 < q < 2 \) and \( 2 < p < \infty \). Then

(i) \( P : L^p(\mathbb{C}) \to \text{Lip}_\alpha(\mathbb{C}) \), where \( \alpha = 1 - 2/p \);

(ii) \( P : L^p(\Omega) \to W^{1,p}(\mathbb{C}) \) is bounded;

(iii) \( P : L^p(\Omega) \to L^p(\mathbb{C}) \) is compact;

(iv) \( P : L^p(\mathbb{C}) \cap L^q(\mathbb{C}) \to C_0(\mathbb{C}) \) is bounded, where \( C_0 \) is the closure of \( C_0^\infty \) in \( L^\infty \).

For proof of Proposition 3.1 we refer to [32], but see also [22].

The Beurling transform is formally determined by \( S g = \partial P g \) and more precisely as a principal-value integral

\[ S g(z) = -\frac{1}{\pi} \int_{C} \frac{g(\omega)}{\omega - z} \frac{1}{(\omega - z)^2} \, dm(\omega). \]

It is a Calderón-Zygmund operator with a holomorphic kernel. Since \( S \) is a Fourier multiplier operator with symbol

\[ m(\xi) = \frac{\xi}{\xi^2}, \quad \xi = \xi_1 + i\xi_2 \]

in terms of the transform (5.16) below, we see in particular that

\[ S(\overline{\partial} \varphi) = \partial \varphi \quad \text{for} \quad \varphi \in C_0^\infty(\mathbb{C}). \]

Moreover, we have

\[ S = R_1^2 + 2iR_1R_2 - R_2^2, \]
where $R_i$'s denote the Riesz-transforms. Also, it follows (cf. [2], [28]) that

\begin{equation}
S : L^p(\mathbb{C}) \to L^p(\mathbb{C}), \quad 1 < p < \infty,
\end{equation}

and $\lim_{p \to 2} \|S\|_{L^p \to L^p} = \|S\|_{L^2 \to L^2} = 1$.

Because of (3.4), the mapping properties of the Beurling transform control
the solutions to the Beltrami equation (1.4). For instance, if $\text{supp}(\mu)$ is compact
as it is in our case, finding a solution to (1.4) with asymptotics

\begin{equation}
f(z) = \lambda z + O\left(\frac{1}{z}\right), \quad |z| \to \infty
\end{equation}

is equivalent to solving

$$g = \mu S g + \lambda \mu$$

and setting

$$f(z) = \lambda z + P g(z),$$

where $P$ is the Cauchy transform. Therefore, if we denote by $\overline{S}$ the $\mathbb{R}$-linear
operator $\overline{S}(g) = \overline{S(g)}$, we need to understand the mapping properties of $P$ and
the invertibility of the operator $I - \mu \overline{S}$ in appropriate $L^p$-spaces in order to
determine to which $L^p$-class the gradient of the solution to (1.4) belongs.

Recently, Astala, Iwaniec and Saksman established through the funda-
mental theory of quasiconformal mappings the precise $L^p$-invertibility range of
these operators.

**Theorem 3.2.** Let $\mu_1$ and $\mu_2$ be two $\mathbb{C}$-valued measurable functions such
that

\begin{equation}
|\mu_1(z)| + |\mu_2(z)| \leq \kappa
\end{equation}

holds for almost every $z \in \mathbb{C}$ with a constant $0 \leq \kappa < 1$. Suppose that $1 + \kappa < p < 1 + 1/\kappa$. Then the Beltrami operator

\begin{equation}
B = I - \mu_1 S - \mu_2 \overline{S}
\end{equation}

is bounded and invertible in $L^p(\mathbb{C})$, with norms of $B$ and $B^{-1}$ bounded by
constants depending only on $\kappa$ and $p$.

Moreover, the bound in $p$ is sharp; for each $p \leq 1 + \kappa$ and for each $p \geq 1 + 1/\kappa$ there are $\mu_1$ and $\mu_2$ as above such that $B$ is not invertible in $L^p(\mathbb{C})$.

For the proof see [4]. Since $\|S\|_{L^2 \to L^2} = 1$, all operators in (3.8) are
invertible in $L^2(\mathbb{C})$ as long as $\kappa < 1$. Thus Theorem 3.2 determines the interval
around the exponent $p = 2$ where the invertibility remains true. Note that it
is a famous open problem [16] whether

$$\|S\|_{L^p \to L^p} = \max \left\{ p - 1, \frac{1}{p - 1} \right\}.$$
If this turns out to be the case, then
\[ \|\mu S\|_{L^p \to L^p} \leq \|\mu\|_{L^\infty} \|S\|_{L^p \to L^p} < 1 \]
whenever \( p < 1 + 1/\|\mu\|_{L^\infty} \). This would then give an alternative proof to
Theorem 3.2.

Theorem 3.2 also has nonlinear counterparts [4] yielding solutions to non-linear uniformly elliptic PDE’s; here see also [15], [5]. On the other hand, in two dimensions the uniqueness of solutions to general nonlinear elliptic systems is typically reduced to the study of the pseudoanalytic functions of Bers (cf. [9], [32]). In the sequel we will need the following version of this principle.

**Proposition 3.3.** Let \( F \in W^1_{loc}(\mathbb{C}) \) and \( \gamma \in L^p_{loc}(\mathbb{C}) \) for some \( p > 2 \).

Suppose that for some constant \( 0 \leq \kappa < 1 \),
\[ \left| \partial F(z) \right| \leq \kappa \left| \partial F(z) \right| + \gamma(z) |F(z)| \]
holds for almost every \( z \in \mathbb{C} \). Then,

a) If \( F(z) \to 0 \) as \( |z| \to \infty \) and \( \gamma \) has a compact support then
\[ F(z) \equiv 0. \]

b) If for large \( |z| \), \( F(z) = \lambda z + \varepsilon(z) z \) where the constant \( \lambda \neq 0 \) and \( \varepsilon(z) \to 0 \)
as \( |z| \to \infty \), then \( F(z) = 0 \) exactly in one point \( z = z_0 \in \mathbb{C} \).

**Proof.** The result a) is essentially from [32]. For the convenience of the reader we will outline a proof for it after first proving b):

The continuity of \( F(z) = \lambda z + \varepsilon(z) z \) and an application of the degree theory [33] or an appropriate homotopy argument show that \( F \) is surjective and consequently there exists at least one point \( z_0 \in \mathbb{C} \) such that \( F(z_0) = 0 \).

To show that \( F \) cannot have more zeros, let \( z_1 \in \mathbb{C} \) and choose a large disk \( B = B(0, R) \) containing both \( z_1 \) and \( z_0 \). If \( R \) is so large that \( \varepsilon(z) < \lambda/2 \) for \( |z| = R \), then \( F\big|_{\{z\}=R} \) is homotopic to the identity relative to \( \mathbb{C} \setminus \{0\} \).

Next we express (3.9) in the form
\[ \bar{\partial}F = \nu(z) \partial F + A(z)F \]
where \( |\nu(z)| \leq \kappa < 1 \) and \( |A(z)| \leq \gamma(z) \) for almost every \( z \in \mathbb{C} \). Now \( A\chi_B \in L^r(\mathbb{C}) \) for all \( 1 \leq r \leq p' = \min\{p, 1 + 1/\kappa\} \) and we obtain from Theorem 3.2 that \( (I - \nu S)^{-1}(A\chi_B) \in L^r \) for all \( 1 + \kappa < r < p' \).

Next we define, cf. [32], \( \eta = P((I - \nu S)^{-1}(A\chi_B)) \). By Proposition 3.1, \( \eta \in C_0(\mathbb{C}) \) and clearly we have
\[ \bar{\partial}\eta - \nu \partial \eta = A(z), \quad z \in B. \]
By a differentiation we see that the function
\[ g = e^{-\eta F} \]
satisfies
\begin{equation}
\bar{\partial} g - \nu \partial g = 0, \quad z \in B.
\end{equation}

Since \( \eta \in W^{1,r}(\mathbb{C}) \) by Proposition 3.1, also \( g \in W^{1,r}_{\text{loc}}(\mathbb{C}) \) and thus \( g \) is quasiregular in \( B \). As such, see e.g. [14, Th. 1.1.1], \( g = h \circ \psi \), where \( \psi : B \to B \) is a quasiconformal homeomorphism and \( h \) is holomorphic, both continuous up to a boundary.

Since \( \eta \) is continuous, (3.12) shows that \( g \rvert \{ |z| = R \} \) is homotopic to the identity relative to \( \mathbb{C} \setminus \{0\} \), and so is the holomorphic function \( h \rvert \{ |z| = R \} \). Therefore, \( h \) has by the principle of the argument ([25, Ths. V.7.1 and VIII.3.5]) precisely one zero in \( B = B(0,R) \). As \( h(\psi(z_0)) = e^{-\eta(z_0)}F(z_0) = 0 \), there can be no further zeros for \( F \) either. This finishes the proof of b).

For the claim a) the condition \( F(z) = \varepsilon(z)z \) is too weak to guarantee \( F \equiv 0 \) in general. But if \( \gamma \) has a compact support we may choose \( \text{supp} \gamma \subset B(0,R) \) and thus the function \( \eta \) solves (3.11) for all \( z \in \mathbb{C} \). Consequently (3.13) holds in the whole plane and \( g \) in (3.11) is quasiregular in \( \mathbb{C} \). But since \( F \) and \( \eta \) are bounded, \( g \) also is bounded and thus constant by Liouville’s theorem. Now (3.12) gives

\begin{equation}
F = C_1 e^\eta, \quad \eta \in C_0(\mathbb{C}).
\end{equation}

With the assumption \( F(z) \to 0 \) as \( |z| \to \infty \) we then obtain \( C_1 = 0 \). \qed

Proposition 3.3 generalizes two classical theorems from complex analysis, Liouville’s theorem and the principle of the argument. Indeed, part b) implies that \( F \) is a homeomorphism. With the condition \( F(z) = \lambda z + \varepsilon(z)z \) the winding number of \( F \) around the origin is one. It is not difficult to find generalizations for arbitrary winding numbers.

We also have the following useful

**Corollary 3.4.** Suppose \( F \in W^{1,p}_{\text{loc}}(\mathbb{C}) \cap L^\infty(\mathbb{C}) \), \( p > 2 \), \( 0 \leq \kappa < 1 \) and that \( \gamma \in L^p(\mathbb{C}) \) has compact support. If

\[ |\bar{\partial}F(z)| \leq \kappa |\partial F(z)| + \gamma(z) |F(z)|, \quad z \in \mathbb{C}, \]

then

\[ F(z) = C_1 e^\eta \]

where \( C_1 \) is constant and \( \eta \in C_0(\mathbb{C}) \).

**Proof.** This is a reformulation of (3.14) from above. \qed
4. Complex geometric optics solutions

In this section we establish the existence of the solution to (1.4) of the form
\[ f_\mu(z, k) = e^{ikz} M_\mu(z, k) \]  
where
\[ M_\mu(z, k) - 1 = O \left( \frac{1}{z} \right) \quad \text{as} \quad |z| \to \infty. \]  
Moreover, we show that
\[ \text{Re} \left( \frac{M_\mu(z, k)}{M_{-\mu}(z, k)} \right) > 0, \quad \text{for all} \quad z, k \in \mathbb{C}. \]  
The importance of (4.3) lies for example in the fact that
\[ \nu_z(k) = -e^{-k(z)} \frac{M_\mu(z, k) - M_{-\mu}(z, k)}{M_\mu(z, k) + M_{-\mu}(z, k)} \]  
appears in Section 8 as the coefficient in a Beltrami equation in the \( k \)-variable. The result (4.3) clearly implies
\[ |\nu_z(k)| < 1 \quad \text{for all} \quad z, k \in \mathbb{C}. \]  
We start with

**Proposition 4.1.** Assume that \( 2 < p < 1 + 1/\kappa \), that \( \alpha \in L^\infty(\mathbb{C}) \) with \( \text{supp}(\alpha) \subset \mathbb{D} \) and that \( |\nu(z)| \leq \kappa \chi_\mathbb{D}(z) \) for almost every \( z \in \mathbb{D} \). Define the operator \( K : L^p(\mathbb{C}) \to L^p(\mathbb{C}) \) by
\[ Kg = P \left( I - \nu S \right)^{-1} (\alpha g). \]  
Then \( K : L^p(\mathbb{C}) \to W^{1,p}(\mathbb{C}) \) and \( I - K \) is invertible in \( L^p(\mathbb{C}) \).

**Proof.** First we note that by Theorem 3.2, \( I - \nu S \) is invertible in \( L^p \) and by Proposition 3.1 (iii) the operator \( K : L^p(\mathbb{C}) \to L^p(\mathbb{C}) \) is well-defined and compact. We also have \( \text{supp} (I - \nu S)^{-1}(\alpha g) \subset \mathbb{D} \). Thus, by Fredholm’s alternative, we need to show that \( I - K \) is injective. So suppose that \( g \in L^p(\mathbb{C}) \) satisfies
\[ g = P \left( (I - \nu S)^{-1} (\alpha g) \right). \]  
By Proposition 3.1 (ii) \( g \in W^{1,p} \) and thus by (4.6)
\[ \overline{\partial} g = (I - \nu S)^{-1} (\alpha \overline{g}) \]  
or equivalently
\[ \overline{\partial} g - \nu \overline{\partial} g = \alpha \overline{g}. \]
Finally, from (4.7) it follows that $g$ is analytic outside the unit disk. This together with $g \in L^p(\mathbb{C})$ implies

$$g(z) = O\left(\frac{1}{z}\right) \text{ for } z \to \infty.$$ 

Thus the assumptions of Proposition 3.3 a) are fulfilled and we must have $g \equiv 0$. $\square$

It is not difficult to find examples showing that Proposition 4.1 fails for $p \leq 2$ and for $p \geq 1 + 1/\kappa$.

We are now ready to establish the existence of the complex geometric optics solutions to (1.4).

**Theorem 4.2.** For each $k \in \mathbb{C}$ and for each $2 < p < 1 + 1/\kappa$ the equation (1.4) admits a unique solution $f \in W^{1,p}_{\text{loc}}(\mathbb{C})$ of the form (1.5) such that the asymptotic formula (1.6) holds true.

In particular, $f(z,0) \equiv 1$.

**Proof.** If we write $f_\mu(z,k) = e^{ikz}M_\mu(z,k) = e^{ikz}(1 + \eta(z))$

and plug this into (1.4) we obtain

$$\bar{\partial}\eta - e_{-k} \mu \bar{\partial}\eta = \alpha \eta + \alpha \tag{4.8}$$

where $e_{-k}$ is as defined in (1.12) and

$$\alpha(z) = -ik e_{-k}(z) \mu(z). \tag{4.9}$$

Hence

$$\bar{\partial}\eta = (I - e_{-k} \mu \overline{S})^{-1}(\alpha \overline{\eta} + \alpha). \tag{4.10}$$

If now $K$ is defined as in Proposition 4.1 with $\nu = e_{-k} \mu$ we get

$$\eta - K \eta = K(\chi_D) \in L^p(\mathbb{C}). \tag{4.11}$$

Since by Proposition 4.1 the operator $I - K$ is invertible in $L^p(\mathbb{C})$, and by (4.8) $\eta$ is analytic in $\mathbb{C} \setminus \overline{D}$ the claims follow by (4.10) and (4.11). $\square$

Next, let $f_\mu(z,k) = e^{ikz}M_\mu(z,k)$ and $f_{-\mu}(z,k) = e^{ikz}M_{-\mu}(z,k)$ be the solutions given by Theorem 4.2 corresponding to conductivities $\sigma$ and $\sigma^{-1}$, respectively.

**Proposition 4.3.** For all $k,z \in \mathbb{C}$,

$$\text{Re}\left(\frac{M_\mu(z,k)}{M_{-\mu}(z,k)}\right) > 0. \tag{4.12}$$
Proof. Firstly, note that (1.4) implies for $M_{\pm\mu}$

\[(4.13) \quad \bar{\partial}M_{\pm\mu} \mp \mu e_{-k}\bar{\partial}M_{\pm\mu} = \mp i\mu e_{-k}M_{\pm\mu}.\]

Thus we may apply Corollary 3.4 to get

\[(4.14) \quad M_{\pm\mu}(z) = \exp(\eta_{\pm}(z)) \neq 0\]

and consequently $M_{\pm\mu}/M_{-\mu}$ is well defined. Secondly, if (4.12) is not true the continuity of $M_{\pm\mu}$ and the fact $\lim_{z \to \infty} M_{\pm\mu}(z,k) = 1$ imply the existence of $z_0 \in \mathbb{C}$ such that

\[
M_{\mu}(z_0,k) = it M_{-\mu}(z_0,k)
\]

for some $t \in \mathbb{R} \setminus \{0\}$. But then $g = M_{\mu} - it M_{-\mu}$ satisfies

\[
\bar{\partial}g = \mu \bar{\partial}(e_{-k}g),
\]

\[
g(z) = 1 - it + O\left(\frac{1}{z}\right), \quad \text{as } z \to \infty.
\]

According to Corollary 3.4 this implies

\[
g(z) = (1 - it) \exp(\eta(z)) \neq 0,
\]

contradicting the assumption $g(z_0) = 0$.

\[\square\]

5. $\partial_k$-equations

We will prove in this section the $\partial_k$-equation (1.8) for the complex geometric optics solutions. We begin by writing (1.4)–(1.6) in the form

\[(5.1) \quad \bar{\partial}M_{\mu} = \mu \bar{\partial}(e_{-k}M_{\mu}), \quad M_{\mu} - 1 \in W^{1,p}(\mathbb{C}).\]

By introducing an $\mathbb{R}$-linear operator $L_{\mu}$,

\[
L_{\mu}g = P\left(\mu \bar{\partial}(e_{-k}g)\right),
\]

we see that (5.1) is equivalent to

\[(5.2) \quad (I - L_{\mu})M_{\mu} = 1.
\]

The following refinement of Proposition 4.1 will serve as the main tool in proving (1.8). Below we will study functions of the form $f = \text{constant} + f_0$, where $f_0 \in W^{1,p}(\mathbb{C})$, and use the notation $W^{1,p}(\mathbb{C}) \oplus \mathbb{C}$ for the corresponding Banach space.

**Theorem 5.1.** Assume that $k \in \mathbb{C}$ and $\mu \in L^\infty_{\text{comp}}(\mathbb{C})$ with $\|\mu\|_{\infty} \leq \kappa < 1$. Then for $2 < p < 1 + \frac{1}{1/\kappa}$ the operator

\[
I - L_{\mu} : W^{1,p}(\mathbb{C}) \oplus \mathbb{C} \to W^{1,p}(\mathbb{C}) \oplus \mathbb{C}
\]

is bounded and invertible.
Proof. We write $L_\mu(g)$ as

$$L_\mu(g) = P \left( \mu e_{-k}\overline{\partial g} - ik\mu e_{-k}\overline{g} \right).$$

Proposition 3.1 (ii) now yields that

$$L_\mu : W^{1,p}(\mathbb{C}) \oplus \mathbb{C} \to W^{1,p}(\mathbb{C})$$

is bounded. Thus we need to show that $I - L_\mu$ is bijective on $W^{1,p}(\mathbb{C}) \oplus \mathbb{C}$. To this end, assume

$$(I - L_\mu)(g + C_0) = h + C_1$$

for $g, h \in W^{1,p}(\mathbb{C})$ and for constants $C_0, C_1$. This yields

$$C_1 - C_0 = g - h - L_\mu(g + C_0)$$

which by (5.4) gives $C_0 = C_1$. By differentiating, rearranging and by using the operator $K = K_\mu$ from Proposition 4.1 with $\alpha = -i\overline{\mu}e_{-k}$ and $\nu = \mu e_{-k}$ we see that (5.5) is equivalent to

$$g - K_\mu(g) = K_\mu(C_0\chi_D) + P \left[ (I - \mu e_{-k}\overline{S})^{-1}\partial h \right].$$

Since the right-hand side belongs to $L^p(\mathbb{C})$ for each $h \in W^{1,p}(\mathbb{C})$, this equation has a unique solution $g \in W^{1,p}(\mathbb{C})$ by Proposition 4.1.

As an immediate corollary we get the following important

**Corollary 5.2.** The operator $I - L^2_\mu$ is invertible on $W^{1,p}(\mathbb{C}) \oplus \mathbb{C}$.

**Proof.** Since $L_\mu = -L_{-\mu}$ we have $I - L^2_\mu = (I - L_\mu)(I - L_{-\mu})$.

Next, we make use of the differentiability properties of the operator $L_\mu$. For later purposes it will be better to work with $L^2_\mu$ which can be written in the following convenient form

$$L^2_\mu g = P \left( \mu \overline{\partial}(\partial + ik)^{-1}\mu(\partial + ik)g \right)$$

where the operator $(\partial + ik)^{-1}$ is defined by

$$(\partial + ik)^{-1}g = e_{-k}\partial^{-1}(e_k g), \quad g \in L^p(\mathbb{C}).$$

Note that many mapping properties of this operator follow from Proposition 3.1. Moreover, we have

**Lemma 5.3.** Let $p > 2$. Then the operator valued map $k \mapsto (\partial + ik)^{-1}$ is continuously differentiable in $\mathbb{C}$, in the uniform operator topology: $L^p(\mathbb{D}) \to W^{1,p}_{\text{loc}}(\mathbb{C})$.

**Proof.** The lemma is a straightforward reformulation of [22, Lemma 2.2], where slightly different function spaces were used. Note that $W^{1,p}_{\text{loc}}(\mathbb{C})$ has the topology given by the seminorms $\|f\|_n = \|f\|_{W^{1,p}(B(0,n))}$, $n \in \mathbb{N}$. 

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Combination of Lemma 5.3 with (5.7) shows that \( k \rightarrow L^2_\mu \) is a \( C^1 \)-family of operators \( L^2_\mu : W^{1,p}(\mathbb{C}) \oplus \mathbb{C} \rightarrow W^{1,p}(\mathbb{C}) \oplus \mathbb{C} \) in the uniform operator topology. If we iterate the equation (5.2) once, we get

\[
M_\mu = 1 + P(\mu \partial e_k - k) + L^2_\mu(M_\mu).
\]

Therefore the above lemma shows that \( k \mapsto M_\mu(z, k) \) is a continuously differentiable family of functions in \( W^{1,p}(\mathbb{C}) \oplus \mathbb{C}, \ p > 2 \). In particular, for each fixed \( z \in \mathbb{C} \), \( M_\mu(z, k) \) is continuously differentiable in \( k \). An alternative way to see this is to note that \( k \mapsto L_\mu \) is smooth in the operator norm topology of \( L(W^{1,p}(\mathbb{C}) \oplus \mathbb{C}) \) and then use Theorem 5.1. This gives by (5.2) that for fixed \( z \) the map \( k \mapsto M_\mu(z, k) \) is, indeed, \( C^\infty \)-smooth.

Furthermore, by (5.1), with respect to the first variable, \( M_\mu(z, k) \) is complex analytic in \( \mathbb{C} \setminus \overline{D} \) with development

\[
M_\mu(z, k) = 1 + \sum_{n=1}^{\infty} b_n(k)z^{-n}, \quad \text{for } |z| > 1.
\]

We define the scattering amplitude corresponding to \( M_\mu \) to be

\[
t_\mu(k) = \frac{1}{\pi} \int_{\mathbb{D}} \mu \partial(e_kM_\mu) \ dm.
\]

Beals and Coifman [8] introduced the idea of studying the \( k \)-dependence of operators associated to complex geometric optics solutions. We will use the Beals-Coifman principle in the following form:

**Lemma 5.4.** Suppose \( g \in W^{1,p}(\mathbb{C}) \oplus \mathbb{C} \) is fixed. Then

\[
\frac{\partial}{\partial k}(e_{-k} \partial^{-1} \mu \partial e_kg) = -it_\mu(g; k)e_{-k}
\]

where

\[
t_\mu(g; k) = \frac{1}{\pi} \int_{\mathbb{C}} \mu \partial(e_kg) \ dm.
\]

**Proof.** For \( f \in L^p_{\text{comp}}(\mathbb{C}) \) we have

\[
(e_{-k} \partial^{-1} e_kf)(z) = -\frac{1}{\pi} \int_{\mathbb{C}} \frac{e_k(\xi - z)}{\xi - \overline{z}} f(\xi) \ dm(\xi).
\]

Using this representation [22, Lemma 2.2] shows that

\[
\frac{\partial}{\partial k}(e_{-k} \partial^{-1} e_kf)(z) = \frac{\partial}{\partial k}((\partial + ik)^{-1}f)(z) = -i \hat{f}(k)e_{-k}(z)
\]

where

\[
\hat{f}(k) = \frac{1}{\pi} \int_{\mathbb{C}} e_k(\xi) f(\xi) \ dm(\xi).
\]

By rewriting the left-hand side of (5.13) in the form \( \frac{\partial}{\partial k}((\partial + ik)^{-1} \mu(\partial + ik)g) \) we see that the claim follows from (5.15). \( \square \)
To get rid of the second term on the right-hand side of (5.9) we introduce

$$F_+ = \frac{1}{2} (M_\mu + M_{-\mu}),$$

$$F_- = \frac{ie_{-k}}{2} (M_\mu - M_{-\mu}).$$

In particular, (5.9) gives

$$F_+ = 1 + L_\mu^2 F_+.$$  

From Lemma 5.4 and (5.7) one has $\partial_k L_\mu^2 (g) = -it_\mu (g; k) P(\mu \partial e_{-k})$ for every $g \in W^{1,p}(\mathbb{C}) \oplus \mathbb{C}$. Hence a differentiation of (5.19) yields

$$L_\mu^2 (\partial_k F_+) = -i \tau_\mu (k) P(\mu \partial e_{-k})$$

where the scattering coefficient $\tau_\mu (k)$ is

$$\tau_\mu (k) \equiv t_\mu (F_+; k) = \frac{1}{2} (t_\mu (k) - t_{-\mu} (k)).$$

Note that this is consistent with (1.9).

One way to identify $\partial_k F_+$ is by readily observing that the unique solution to (5.20) also has other realizations. Namely, if one subtracts the equation (5.9) applied to $M_\mu$ from the corresponding equation for $M_{-\mu}$, one obtains, after using $L_\mu^2 = L_{-\mu}^2$,

$$(I - L_\mu^2) (\partial_k F_+) = -i \tau_\mu (k) P(\mu \partial e_{-k}).$$

Thus by Corollary 5.2, (5.20) and (5.22) we have proved the first part of

**Theorem 5.5.** For each fixed $z \in \mathbb{C}$, the functions $k \mapsto F_\pm (z, k)$ are continuously differentiable with

a) $\partial_k F_+ (z, k) = \tau_\mu (k) e_{-k}(z) F_- (z, k),$ 

b) $\partial_k F_- (z, k) = \tau_\mu (k) e_{-k}(z) F_+ (z, k).$

**Proof.** The differentiability is clear since $M_{\pm \mu} (z, k)$ are continuously differentiable in $k$. Hence we are left proving b). We start by adding and subtracting (5.2) for $M_\mu$ and $M_{-\mu}$ to arrive at the equations

$$F_+ = 1 - i \overline{\partial}^{-1} \mu \overline{\partial} F_-,$$

$$F_- = ie_{-k} \partial^{-1} \mu \partial e_k F_+.$$  

By differentiating the second equation with respect to $\overline{k}$ and by applying Lemma 5.4 we get

$$\partial_{\overline{k}} F_- = \tau_\mu (k) e_{-k} + ie_{-k} \partial^{-1} \mu \partial (e_k \partial_{\overline{\overline{k}}} F_+).$$

Combining this with part a) we have

$$\partial_{\overline{k}} F_- = \tau_\mu (k) e_{-k} (1 + i \partial^{-1} \mu \partial F_-).$$

This together with (5.23) yields b).
We close this section by returning to the functions
\[ h_+ = \frac{1}{2}(f_\mu + f_{-\mu}) = e^{ikz}F_+ \]
and
\[ h_- = \frac{i}{2}(f_{\bar{\mu}} - f_{-\bar{\mu}}) = e^{ikz}F_- . \]
These expressions with Theorem 5.5 immediately give the identities (1.8). Note that by Theorem 5.5, \( k \mapsto h_\pm(z, k) \) is \( C^1 \) in \( \mathbb{C} \), for each fixed \( z \).

6. From \( \Lambda_\sigma \) to \( \tau \)

We next prove that the Dirichlet to Neumann operator \( \Lambda_\sigma \) uniquely determines \( f_\mu(z) \) and \( f_{-\mu}(z) \) at the points \( z \) that lie outside the unit disk \( \mathbb{D} \). This will also show that \( \Lambda_\sigma \) determines \( \tau_\mu(k) \) for all \( k \in \mathbb{C} \).

**Proposition 6.1.** If \( \sigma \) and \( \bar{\sigma} \) are two conductivities satisfying the assumptions of Theorem 1 with \( \Lambda_\sigma = \Lambda_{\bar{\sigma}} \), and if \( \mu \) and \( \bar{\mu} \) are the corresponding Beltrami coefficients, then
\[ f_\mu(z) = f_{\bar{\mu}}(z) \quad \text{and} \quad f_{-\mu}(z) = f_{-\bar{\mu}}(z) \]
for all \( z \in \mathbb{C} \setminus \mathbb{D} \).

**Proof.** We assume \( \Lambda_\sigma = \Lambda_{\bar{\sigma}} \) which by Proposition 2.3 implies that \( \mathcal{H}_\mu = \mathcal{H}_{\bar{\mu}} \). Since \( \Lambda_\sigma \) by the same proposition determines \( \Lambda_{\sigma-} \), it is enough to prove the first claim of (6.1).

From (2.9) we see firstly that the projections \( Q_\mu = Q_{\bar{\mu}} \) and thus by Lemma 2.2
\[ Q_\mu(f - \bar{f}) = \text{constant} \]
where we have written \( f = f_\mu|_{\partial \mathbb{D}} \) and \( \bar{f} = f_{\bar{\mu}}|_{\partial \mathbb{D}} \). By using Lemma 2.2 again we see that there exists a function \( G \in H^1(\mathbb{D}) \) with \( \overline{\partial G} = \mu \overline{\partial G} \) in \( \mathbb{D} \) and
\[ G|_{\partial \mathbb{D}} = f - \bar{f} . \]
Define then \( G \) outside \( \mathbb{D} \) by
\[ G(z) = f_\mu(z) - f_{\bar{\mu}}(z), \quad |z| \geq 1, \]
to get a global solution to
\[ \overline{\partial G}(z) - \mu(z)\overline{\partial G}(z) = 0, \quad z \in \mathbb{C} . \]
Since \( G \) is an \( H^1_{\text{loc}} \)-solution to (6.2), the general smoothness properties of quasiregular mappings [3] give \( G \in W^{1,p}_{\text{loc}}(\mathbb{C}) \) for all \( p < 1 + 1/\kappa \). This regularity can also readily be seen from Theorem 3.2, since the compactly supported function \( h = f_\mu - f_{\bar{\mu}} - G \) satisfies
\[ \overline{\partial h} = -(1 - \mu S)^{-1} (\chi_{\mathbb{D}} \overline{\partial f_{\bar{\mu}}} - \mu \overline{\partial f_\mu}) . \]
Finally, from the above we obtain that the function $G_0$, defined by
\[ G_0(z) = e^{-ikz}G(z), \]
belongs to $W^{1,p}(\mathbb{C})$ and satisfies $G_0(z) = O(1/z)$ with
\[ \partial G_0 - e_{-k\mu} \overline{\partial G_0} = -i k e_{-k\mu} \overline{G_0}. \]
By Proposition 3.3 a) the function $G_0$ must hence vanish identically, which proves (6.1).

**Corollary 6.2.** The operator $\Lambda_\sigma$ uniquely determines $t_\mu$, $t_{-\mu}$ and $\tau_\mu$.

**Proof.** The claim follows immediately from Proposition 6.1, (1.5) and from the definitions (5.11) of the scattering coefficients.

From the results of Section 5 it follows that the coefficient $\tau_\mu$ is continuously differentiable in $k$ and vanishes at the origin. For the global properties of $\tau_\mu$ we apply Schwarz’s lemma.

**Proposition 6.3.** The complex geometric optics solutions
\[ f_{\pm\mu}(z, k) = e^{ikz}M_{\pm\mu}(z, k) \]
satisfy for $|z| > 1$ and for all $k \in \mathbb{C}$
\[ (6.3) \quad \left| \frac{M_\mu(z, k) - M_{-\mu}(z, k)}{M_\mu(z, k) + M_{-\mu}(z, k)} \right| \leq \frac{1}{|z|}. \]
Moreover, for the scattering coefficient $\tau_\mu(k)$,
\[ (6.4) \quad |\tau_\mu(k)| \leq 1 \quad \text{for all } k \in \mathbb{C}. \]

**Proof.** Fix the parameter $k \in \mathbb{C}$ and denote
\[ m(z) = \frac{M_\mu(z, k) - M_{-\mu}(z, k)}{M_\mu(z, k) + M_{-\mu}(z, k)}. \]
Then by Proposition 4.3, $|m(z)| < 1$ for all $z \in \mathbb{C}$. Moreover, $m$ is $\mathbb{C}$-analytic in $z \in \mathbb{C} \setminus \overline{D}$, $m(\infty) = 0$, and thus by Schwarz’s lemma we have $|m(z)| \leq 1/|z|$ for all $z \in \mathbb{C} \setminus \overline{D}$. Since (5.11) and (5.21) give $\lim_{z \to -\infty} zm(z) = \overline{\tau_\mu}$, both claims of the proposition follow.

**7. Subexponential growth**

We know from Section 4 and (4.1), (4.14) that the complex geometric optics solution $f_\mu$ of (1.4) can be written in the exponential form. Here we begin with a more detailed analysis of this fact. For later purposes we also need to generalize the situation a bit by considering complex Beltrami coefficients $\mu_\lambda$ of the form $\mu_\lambda = \lambda \mu$, where the constant $\lambda \in \partial\overline{D}$ and $\mu = (1 - \sigma)/(1 + \sigma)$ is
as before. Precisely as in Section 4 one can show the existence and uniqueness of $f_{\lambda \mu} \in W^1_{\text{loc}}(\mathbb{C})$ satisfying
\begin{equation}
\overline{\partial} f_{\lambda \mu} = \lambda \mu \frac{\partial f_{\lambda \mu}}{\partial \mu}
\end{equation}
and
\begin{equation}
f_{\lambda \mu}(z, k) = e^{ikz} \left(1 + O \left(\frac{1}{z}\right)\right) \text{ as } |z| \to \infty.
\end{equation}

**Lemma 7.1.** The function $f_{\lambda \mu}$ admits a representation
\begin{equation}
f_{\lambda \mu}(z, k) = e^{ik\varphi_\lambda(z, k)},
\end{equation}
where for each fixed $k \in \mathbb{C} \setminus \{0\}$ and $\lambda \in \partial \mathbb{D}$, the function $\varphi_\lambda(\cdot, k) : \mathbb{C} \to \mathbb{C}$ is a quasiconformal homeomorphism that satisfies
\begin{equation}
\varphi_\lambda(z, k) = z + O \left(\frac{1}{z}\right) \text{ for } z \to \infty
\end{equation}
and
\begin{equation}
\overline{\partial} \varphi_\lambda(z, k) = -\frac{k}{\lambda} \mu_\lambda(z) (e^{-k \circ \varphi_\lambda})(z, k) \overline{\partial \varphi_\lambda(z, k)}, \quad z \in \mathbb{C}.
\end{equation}

**Proof.** Since the parameter $k$ is fixed we drop it from the notation and write simply $f_{\lambda \mu}(z, k) = f_{\lambda \mu}(z), \varphi_\lambda(z, k) = \varphi_\lambda(z)$, etc. Denote
\begin{equation}
\mu_1(z) = \mu_\lambda(z) \frac{\partial f_{\mu}(z)}{\partial f_{\mu}(z)}.
\end{equation}
Then
\begin{equation}
\overline{\partial} f_{\lambda \mu} = \mu_1 \partial f_{\lambda \mu}.
\end{equation}
On the other hand, by the general theory of quasiconformal maps ([2], [14], [21]) there exists a unique quasiconformal homeomorphism $\varphi_\lambda \in H^1_{\text{loc}}(\mathbb{C})$ satisfying
\begin{equation}
\overline{\partial} \varphi_\lambda = \mu_1 \partial \varphi_\lambda
\end{equation}
and having the asymptotics
\begin{equation}
\varphi_\lambda(z) = z + O \left(\frac{1}{z}\right) \text{ as } z \to \infty.
\end{equation}
Moreover, any $H^1_{\text{loc}}$-solution to (7.6) is obtained from $\varphi_\lambda$ by post-composing with an analytic function ([14, Th. 11.1.2]). In particular,
\begin{equation}
f_{\lambda \mu}(z) = h \circ \varphi_\lambda(z)
\end{equation}
where $h : \mathbb{C} \to \mathbb{C}$ is an entire analytic function. But
\begin{equation}
\frac{h \circ \varphi_\lambda(z)}{\exp(ik\varphi_\lambda(z))} = \frac{f_{\lambda \mu}(z)}{\exp(ik\varphi_\lambda(z))}
\end{equation}
has by (1.5), (1.6) and (7.8) the limit 1 as the variable $z \to \infty$. Thus
\begin{equation}
h(z) \equiv e^{ikz}.
\end{equation}
Finally, (7.5) follows immediately from (1.4) and (7.3). \qed
Note that the results in Section 4 show that (7.4) and (7.5) have a unique solution. The existence of such a solution can also be directly verified by using Schauder’s fixed point theorem [15], [5]. The result of Lemma 7.1 demonstrates that after a change of coordinates \( z \mapsto \varphi(z) \) the complex geometric optics solution \( f_\mu \) is simply an exponential function.

The main goal of this section is to show

**Theorem 7.2.** Let \( \varphi_\lambda \in H^1_{\text{loc}}(\mathbb{C}) \) satisfy (7.4) and (7.5). Then as \( k \to \infty \),

\[
\varphi_\lambda(z, k) \to z
\]

uniformly in \( z \in \mathbb{C} \) and \( \lambda \in \partial \mathbb{D} \).

We have split the proof of Theorem 7.2 into several lemmas.

**Lemma 7.3.** Suppose \( \varepsilon > 0 \) is given. Suppose also that for \( \mu_\lambda(z) = \lambda \mu(z) \),

\[
(7.9) \quad f_n = \mu_\lambda S_n \mu_\lambda S_{n-1} \mu_\lambda \cdots \mu_\lambda S_1 \mu_\lambda
\]

where \( S_j : L^2(\mathbb{C}) \to L^2(\mathbb{C}) \) are Fourier multiplier operators, each with a unimodular symbol. Then there is a number \( R_n = R_n(\kappa, \varepsilon) \) depending only on \( \mu \), \( n \) and \( \varepsilon \) such that

\[
(7.10) \quad |\hat{f}_n(\xi)| < \varepsilon \text{ for } |\xi| > R_n.
\]

**Proof.** Clearly it is enough to prove the claim for \( \lambda = 1 \).

Recall that for the Fourier transform \( \hat{\phi} \) we use the definition (5.16). By assumption

\[
\hat{S}_j g(\xi) = m_j(\xi) \hat{g}(\xi)
\]

where \( |m_j(\xi)| = 1 \) for \( \xi \in \mathbb{C} \). We have by (7.9)

\[
(7.11) \quad \|f_n\|_{L^2} \leq \|\mu\|_{L^\infty} \|\mu\|_{L^2} \leq \sqrt{\pi} \kappa^{n+1}
\]

since \( \text{supp}(\mu) \subset \mathbb{D} \). Choose first \( \rho_n \) so that

\[
(7.12) \quad \int_{|\xi| > \rho_n} |\hat{\mu}(\xi)|^2 d\mu(\xi) < \varepsilon^2.
\]

After this choose \( \rho_{n-1}, \rho_{n-2}, \ldots, \rho_1 \) inductively so that for \( l = n - 1, \ldots, 1 \)

\[
(7.13) \quad \pi \int_{|\xi| > \rho_l} |\hat{\mu}(\xi)|^2 d\mu(\xi) \leq \varepsilon^2 \left( \prod_{j=l+1}^n \pi \rho_j \right)^{-2}.
\]

Finally, choose \( \rho_0 \) so that

\[
(7.14) \quad |\hat{\mu}(\xi)| < \varepsilon \pi^{-n} \left( \prod_{j=1}^n \rho_j \right)^{-1}, \text{ when } |\xi| > \rho_0.
\]

All these choices are possible since \( \mu \in L^1 \cap L^2 \).
Now, we set \( R_n = \sum_{j=0}^{n} \rho_j \) and claim that (7.10) holds for this choice of \( R_n \). Hence assume that \( |\xi| > \sum_{j=0}^{n} \rho_j \). Now,

\[
(7.15) \quad |\hat{f}_n(\xi)| \leq \int_{|\xi-\eta| \leq \rho_n} |\hat{\mu}(\xi-\eta)| |\hat{f}_{n-1}(\eta)| \, dm(\eta) + \int_{|\xi-\eta| \geq \rho_n} |\hat{\mu}(\xi-\eta)| |\hat{f}_{n-1}(\eta)| \, dm(\eta).
\]

But if \( |\xi-\eta| \leq \rho_n \) then \( |\eta| > \sum_{j=0}^{n-1} \rho_j \). Thus, if we denote

\[
\Delta_n = \sup \left\{ |\hat{f}_n(\xi)| : |\xi| > \sum_{j=0}^{n} \rho_j \right\}
\]

it follows from (7.15) and (7.11) that

\[
\Delta_n \leq \Delta_{n-1}(\pi \rho_n^2)^{1/2} \|\hat{\mu}\|_{L^2} + \left( \int_{|\eta| \geq \rho_n} |\hat{\mu}(\eta)|^2 \, dm(\eta) \right)^{1/2} \|\hat{f}_{n-1}\|_{L^2}
\]

\[
\leq \pi \rho_n \kappa \Delta_{n-1} + \kappa^n \left( \pi \int_{|\eta| \geq \rho_n} |\hat{\mu}(\eta)|^2 \, dm(\eta) \right)^{1/2}
\]

for \( n \geq 2 \). Moreover, the same argument shows that

\[
\Delta_1 \leq \pi \rho_1 \kappa \sup \{ |\hat{\mu}(\xi)| : |\xi| > \rho_0 \} + \kappa \left( \pi \int_{|\eta| \geq \rho_1} |\hat{\mu}(\eta)|^2 \, dm(\eta) \right)^{1/2}.
\]

In conclusion, after an iteration we have

\[
\Delta_n \leq (\kappa \pi)^n \left( \prod_{j=1}^{n} \rho_j \right) \sup \{ |\hat{\mu}(\xi)| : |\xi| > \rho_0 \} + \kappa^n \sum_{l=1}^{n} \left( \prod_{j=l+1}^{n} \pi \rho_j \right) \left( \pi \int_{|\eta| \geq \rho_l} |\hat{\mu}(\eta)|^2 \, dm(\eta) \right)^{1/2},
\]

where we define \( \prod_{j=n+1}^{n} \pi \rho_j = 1 \). With the choices (7.12)–(7.14) this gives

\[
\Delta_n \leq (n + 1) \kappa^n \varepsilon \leq \frac{\varepsilon}{1 - \kappa},
\]

which proves the claim. \( \square \)

Our next goal is to use Lemma 7.3 to obtain a similar asymptotic result as in Theorem 7.2 for a solution to a linear equation somewhat similar to (7.5).
Proposition 7.4. Suppose $\psi \in H^1_{\text{loc}}(\mathbb{C})$ satisfies

$$
\overline{\partial}\psi = \frac{\mp}{k}e_{-k}(z)\partial\psi, \quad \text{and}
$$

$$
\psi(z) = z + O\left(\frac{1}{z}\right) \quad \text{as } z \to \infty.
$$

Then $\psi(z,k) \to z$, uniformly in $z \in \mathbb{C}$ and $\lambda \in \partial\Omega$, as $k \to \infty$.

For Proposition 7.4 we need some preparations. First, as $\|S\|_{L^p \to L^p} \to 1$ when $p \to 2$, we can choose a $\delta_k > 0$ so that $\kappa\|S\|_{L^p \to L^p} < 1$ for every $2 - \delta_k \leq p \leq 2 + \delta_k$. With this notation we then have

Lemma 7.5. Let $\psi \in H^1_{\text{loc}}(\mathbb{C})$ be the unique solution to (7.16) and let $\varepsilon > 0$. Then one can decompose $\overline{\partial}\psi$ in the following way: $\overline{\partial}\psi = g + h$ where

(i) $\|h(\cdot,k)\|_{L^p} < \varepsilon$ for $2 - \delta_k \leq p \leq 2 + \delta_k$, uniformly in $k$,

(ii) $\|g(\cdot,k)\|_{L^p} \leq C_0 = C_0(\kappa)$, uniformly in $k$ and

(iii) $\hat{g}(\xi,k) \to 0$ as $k \to \infty$,

where in (iii) the convergence is uniform on compact subsets of the $\xi$-plane and also uniform in $\lambda \in \partial\Omega$. The Fourier transform is with respect to the first variable only.

Proof. We may solve (7.16) by Born-series which converge in $L^p$,

$$
\overline{\partial}\psi = \sum_{n=0}^{\infty} \left(\frac{\mp}{k}\mu e_{-k} S\right)^n \left(\frac{\mp}{k}\mu e_{-k}\right).
$$

Let

$$
h = \sum_{n=n_0}^{\infty} \left(\frac{\mp}{k}\mu e_{-k} S\right)^n \left(\frac{\mp}{k}\mu e_{-k}\right).
$$

Then

$$
\|h\|_{L^p} \leq \pi^{1/p} \kappa^{n_0+1} \|S\|_{L^p \to L^p}^{n_0} \frac{1 - \kappa\|S\|_{L^p \to L^p}}{1 - \kappa\|S\|_{L^p \to L^p}}.
$$

We obtain (i) by choosing $n_0$ large enough.

The remaining part clearly satisfies (ii) with a constant $C_0$ that is independent of $k$ and $\lambda$. To prove (iii) we first note that

$$
S(e_{-k}\phi) = e_{-k}S_k\phi
$$

where $(S_k\phi)\hat{\phi}(\xi) = m(\xi + k)\hat{\phi}(\xi)$ and $m(\xi) = \xi/\xi$. Consequently,

$$
(\mu e_{-k} S)^n e_{-k} = e_{-(n+1)k}\mu S_{nk}\mu S_{(n-1)k} \cdots \mu S_k\mu
$$
and so
\[ g = \sum_{j=1}^{n_0} e^{-jk\mu S(j-1)k \mu} \cdots \mu S_{j\mu}. \]

Therefore
\[ g = \sum_{j=1}^{n_0} e^{-jk G_j} \]

where by Lemma 7.3, \(|\hat{G}_j(\xi)| < \tilde{\varepsilon}\) whenever \(|\xi| > R = \max_{j \leq n_0} R_j\). As \((e^{-jkG_j})(\xi) = \hat{G}_j(\xi - jk)\), for any fixed compact set \(K_0\) we can take \(k\) so large that \(K_0 - jk \subset \mathbb{C} \setminus B(0, R)\) for each \(1 \leq j \leq n_0\). Then
\[ \sup_{\xi \in K_0} |\hat{g}(\xi, k)| \leq n_0 \tilde{\varepsilon}. \]

This proves (iii).

Proof of Proposition 7.4. We show first that when \(k \to \infty\), \(\overline{\partial} \psi \to 0\) weakly in \(L^p\), \(2 - \delta_\kappa \leq p \leq 2 + \delta_\kappa\). For this, suppose \(f_0 \in L^q\), \(q = p/(p-1)\), is fixed and choose \(\varepsilon > 0\). Then there exists \(f \in C^\infty_0(\mathbb{C})\) such that \(\|f_0 - f\|_{L^q} < \varepsilon\) and so by Lemma 7.5
\[ |\langle f_0, \overline{\partial} \psi \rangle| \leq \varepsilon C_1(f_0) + \int |\hat{f}(\xi)\hat{g}(\xi, k)\, dm(\xi)|. \]

First, choose \(R\) so large that
\[ \int_{\mathbb{C} \setminus B(0, R)} |\hat{f}(\xi)|^2 \, dm(\xi) \leq \varepsilon^2 \]

and then \(|k|\) so large that \(|\hat{g}(\xi, k)| \leq \varepsilon/(\sqrt{\pi}R)\) for all \(\xi \in B(0, R)\). Now
\[ \left| \int \hat{f}(\xi)\hat{g}(\xi, k) \, dm(\xi) \right| \leq \int_{B(0, R)} |\hat{f}(\xi)| \hat{g}(\xi, k) \, dm(\xi) + \int_{\mathbb{C} \setminus B(0, R)} |\hat{f}(\xi)| \hat{g}(\xi, k) \, dm(\xi) \]
\[ \leq \varepsilon (\|f\|_{L^2} + \|g\|_{L^2}) \leq \varepsilon C_2(f_0). \]

The bound is the same for all \(\lambda\); hence \(\sup_{\lambda \in \partial \mathbb{D}} |\langle f_0, \overline{\partial} \psi \rangle| \to 0\) as \(|k| \to \infty\).

To prove the uniform convergence of \(\psi\) itself we write
\[ \psi(z, k) = z - \frac{1}{\pi} \int_{\mathbb{D}} \frac{1}{\omega - z} \overline{\partial} \psi(\omega, k) \, dm(\omega). \]

Here note that \(\text{supp}(\overline{\partial} \psi) \subset \mathbb{D}\) and \(\chi_\mathbb{D}(\omega)/(\omega - z) \in L^q\) for all \(q < 2\). Thus by the weak convergence we get
\[ \psi(z, k) \to z \text{ as } k \to \infty, \]

for each fixed \(z \in \mathbb{C}\), but uniformly in \(\lambda \in \partial \mathbb{D}\). On the other hand as \(\sup_k \|\overline{\partial} \psi\|_{L^p} \leq C_0(\kappa) < \infty\), for all \(z\) sufficiently large \(|\psi(z, k) - z| < \varepsilon\), uniformly in \(k \in \mathbb{C}\) and \(\lambda \in \partial \mathbb{D}\). Thus (7.17) with Proposition 3.1.(i) shows that
the family \( \{ \psi(\cdot, k) : k \in \mathbb{C}, \lambda \in \partial D \} \) is equicontinuous. Combination of all these observations shows that the convergence in (7.18) is uniform in \( z \in \mathbb{C} \) and \( \lambda \in \partial D \).

Finally we proceed to the nonlinear case: assume that \( \varphi_\lambda \) satisfies (7.4) and (7.5). Since \( \varphi \) is a (quasiconformal) homeomorphism we may consider its inverse \( \psi_\lambda : \mathbb{C} \to \mathbb{C} \),

\[
\psi_\lambda \circ \varphi_\lambda (z) = z,
\]
which also is quasiconformal. By differentiating (7.19) with respect to \( z \) and \( \lambda \) one obtains that \( \psi \) satisfies

\[
\partial \psi_\lambda = -\lambda \frac{k}{k} (\mu \circ \psi_\lambda) e^{\frac{-k}{k}} \partial \psi_\lambda \quad \text{and} \quad \psi_\lambda (z, k) = z + O \left( \frac{1}{z} \right) \quad \text{as} \quad z \to \infty.
\]

Proof of Theorem 7.2. It is enough to show that

\[
\psi_\lambda (z, k) \to z
\]
uniformly in \( z \) and \( \lambda \) as \( k \to \infty \). To prove this we need to recall some further facts from quasiconformal mappings. Let us use the notation

\[
\Sigma_\kappa = \{ g \in H^1_{\text{loc}}(\mathbb{C}) : \partial g = \nu g, \ |\nu| \leq \kappa |\chi|_D \ \text{and} \ \ g = z + O \left( \frac{1}{z} \right) \ \text{as} \ z \to \infty \}. \tag{7.23}
\]

Note that with the above normalization at \( \infty \) all elements \( g \in \Sigma_\kappa \) are homeomorphisms \([2]\). Also, the use of \( \chi_D \) will become clear soon.

Lemma 7.6. a) The family \( \Sigma_\kappa \) is compact in the topology of uniform convergence on \( \mathbb{C} \).

b) Suppose that \( f, g \in \Sigma_\kappa, \ 1 + \kappa < p < 1 + 1/\kappa \) and that \( \varepsilon > 0 \) is so small that \( (1 + \varepsilon)p < 1 + 1/\kappa \). Then

\[
\int_{\mathbb{C}} |\bar{\partial} f - \bar{\partial} g|^p \, dm \leq C(p, \varepsilon) \left( \int_{\mathbb{C}} |\nu_f - \nu_g|^{\frac{p+1}{\kappa}} \, dm \right)^{\frac{\kappa}{p+1}}
\]

where

\[
\nu_f = \frac{\partial f}{\partial f} \ \text{and} \ \nu_g = \frac{\partial g}{\partial g}.
\]

Proof. The claim a) follows from \([21, \text{Th. II.51}]\). Furthermore, since \( \partial g = 1 + S(\bar{\partial} g) \) for each \( g \in \Sigma_\kappa \), we have

\[
\bar{\partial} f - \bar{\partial} g = (I - \nu_f S)^{-1} \left( \nu_f - \nu_g + (\nu_f - \nu_g) S(\bar{\partial} g) \right), \quad f, g \in \Sigma_\kappa.
\]
Applying Theorem 3.2 and Hölder’s inequality then gives the claim b). \( \Box \)
The support of \( \mu \circ \psi \) need not anymore be contained in \( \mathbb{D} \). However, by Koebe’s 1/4-theorem, see e.g. [1, Cor. 5.3], \( \varphi(\mathbb{D}) \subset 4\mathbb{D} \) and thus \( \text{supp}(\mu \circ \psi) \subset 4\mathbb{D} \). Therefore by Lemma 7.6 a) we have sequences \( k_n \to \infty \) and sequences \( \lambda_n \to \lambda \in \partial \mathbb{D} \) such that \( \psi_{\lambda_n}(\cdot, k_n) \to \psi_{\infty} \) uniformly, with \( \psi_{\infty} \in \Sigma_\kappa \). To prove Theorem 7.2 it is enough to show that for any such sequence \( \psi_{\infty}(z) \equiv z \).

Hence assume that we have such a limit function \( \psi_{\infty} \). We then consider the \( H^{1}_{\text{loc}} \)-solution \( \Phi(z) = \Phi_{\lambda}(z, k) \) of

\[
\overline{\partial} \Phi = -\frac{k}{k} \lambda (\mu \circ \psi_{\infty}) e_{-k} \partial \Phi,
\]

\[
\Phi(z) = z + O \left( \frac{1}{z} \right) \quad \text{as} \quad z \to \infty.
\]

This is now a linear Beltrami equation which [2] has a unique solution \( \Phi \in \Sigma_\kappa \) for each \( k \in \mathbb{C} \). According to Proposition 7.4

\[
(7.24) \quad \Phi_{\lambda}(z, k) \to z \quad \text{as} \quad k \to \infty.
\]

Secondly, when \( 2 < p < 1 + 1/\kappa \), by Lemma 7.6 b)

\[
(7.25) \quad |\psi_{\lambda_n}(z, k_n) - \Phi_{\lambda}(z, k_n)| = \frac{1}{\pi} \left| \int_{\partial \mathbb{D}} \frac{1}{\omega - z} \overline{\partial} \left( \psi_{\lambda_n}(\omega, k_n) - \Phi_{\lambda}(\omega, k_n) \right) \, d\mu(\omega) \right|
\]

\[
\leq C_1 \| \overline{\partial} \left( \psi_{\lambda_n}(\cdot, k_n) - \Phi_{\lambda}(\cdot, k_n) \right) \|_{L^p}
\]

\[
\leq C_2 |\lambda_n - \lambda| + C_2 \left( \int_{\partial \mathbb{D}} \left| \mu(\psi_{\lambda_n}(\omega, k_n)) - \mu(\psi_{\infty}(\omega)) \right|^{\frac{p(1+\epsilon)}{1-\epsilon}} \, d\mu(\omega) \right)^{\frac{1}{p(1+\epsilon)}}.
\]

Lastly, we use the higher integrability of the Jacobian [3]; for all \( 2 < p < 1 + 1/\kappa \) and for all \( g = \psi^{-1} \), \( \psi \in \Sigma_\kappa \),

\[
(7.26) \quad \int_{\mathbb{D}} |J_g(z)|^{p/2} \, dm(z) \leq \int_{\mathbb{D}} |\partial g|^p \, dm \leq C(\kappa) < \infty,
\]

where \( C(\kappa) \) depends only on \( \kappa \). Again, the bound can also be deduced from Theorem 3.2 since \( \overline{\partial} g = (I - \nu g S)^{-1} \nu g \) and \( \partial g = 1 + S(\overline{\partial} g) \). We use this estimate in the cases \( \psi(z) = \psi_{\lambda_n}(z, k_n) \) and \( \psi = \psi_{\infty} \). Namely, we have for each \( \eta \in C_0^\infty(\mathbb{D}) \) that

\[
(7.27) \quad \int_{\partial \mathbb{D}} |\mu(\psi) - \eta(\psi)|^{\frac{p(1+\epsilon)}{1-\epsilon}} = \int_{\partial \mathbb{D}} |\mu - \eta|^{\frac{p(1+\epsilon)}{1-\epsilon}} J_g
\]

\[
\leq \left( \int_{\partial \mathbb{D}} |\mu - \eta|^{\frac{p^2(1+\epsilon)}{p(1-\epsilon)(p-2)}} (p-2)^{p/2} \left( \int_{\partial \mathbb{D}} J_g^{p/2} \right)^{2/p} \right).
\]

Since \( \mu \) can be approximated in the mean by smooth \( \eta \), the last term in (7.27) can be made arbitrarily small. Since by uniform convergence \( \eta(\psi_{\lambda_n}(z, k_n)) \to \eta(\psi_{\infty}(z)) \) we see that the last bound in (7.25) converges to zero as \( \lambda_n \to \lambda \) and \( k_n \to k \). In view of (7.24) and (7.25) we have established that

\[
\psi_{\lambda_n}(z, k_n) \to z
\]

and that \( \psi_{\infty}(z) \equiv z \). The theorem is proved.
The transport matrix

The gradient of a quasiregular map can vanish only on a set of Lebesgue measure zero ([2, p. 34]). By the equation (1.4) the derivative $\partial f_\mu$ can vanish only on points $z$ where the whole gradient of $f_\mu$ vanishes. This means that we can recover $\mu$ and hence $\sigma$ from $f_\mu$ by the formulae

$$\mu = (\partial f_\mu)^{-1} \partial f_\mu$$

and

$$\sigma = \frac{1 - \mu}{1 + \mu}.$$  \hfill (8.1)

Next, let $u_1 = h_+ - ih_-$ and $u_2 = i(h_+ + ih_-)$ be defined as in (1.14). From Lemma 2.1 and (1.7) we see that $u_1$, $u_2$ satisfy the conductivity equations (1.15). With respect to the parameter $k$ the $u_j$’s are $C^1$-mappings; cf. Theorem 5.5. A straightforward derivation using (1.8) shows also that at each point $z \in \mathbb{C}$, both $u_1(z,k)$ and $u_2(z,k)$ satisfy the $\partial_k$-equation

$$\frac{\partial}{\partial k} u(z,k) = -i \tau_\mu(k) u(z,k).$$  \hfill (8.2)

In addition, $u_1(z,0) \equiv 1$ and $u_2(z,0) \equiv i$.

It is clear that the pair $\{u_1(z,k), u_2(z,k)\}$ determines the pair $\{f_\mu(z,k), f_{-\mu}(z,k)\}$ and vice versa. Thus by Proposition 6.1 the Dirichlet to Neumann map $\Lambda_\sigma$ uniquely determines the functions $u_1$ and $u_2$ outside $\mathbb{D}$. To prove Theorem 1 it therefore suffices to transport these solutions from outside to inside $\mathbb{D}$ by using $\Lambda_\sigma$. For this purpose we will employ (8.2) and the fact from Corollary 6.2 that $\Lambda_\sigma$ uniquely determines the scattering coefficient $\tau_\mu$.

We have only shown that $\tau_\mu$ is bounded and that $u_j(z,k)e^{-ikz}$ behaves subexponentially as $k \to \infty$, and these facts alone are much too weak to guarantee the uniqueness of solutions of the pseudoanalytic equation (8.2). To remedy this we need to understand the transport matrix from (1.18).

We start by arguing that $u_1$ and $u_2$ are $\mathbb{R}$-linearly independent, i.e. that

$$u_2(z,k) \neq 0 \text{ and } u_1(z,k)/u_2(z,k) \notin \mathbb{R}$$  \hfill (8.3)

holds for all $z, k \in \mathbb{C}$. By Proposition 4.3 we have

$$\left| \frac{h_-(z,k)}{h_+(z,k)} \right| < 1 \text{ and } h_+(z,k) \neq 0 \text{ for all } z, k \in \mathbb{C}. $$

Since

$$\frac{u_1}{u_2} = -\frac{h_+ - ih_-}{h_+ + ih_-},$$

this proves (8.3) and enables us to define the transport matrix $T_{z,z_0}^\sigma(k)$ as in (1.18). It also follows that $T_{z,z_0}^\sigma(k)$ is invertible.

It was discovered by Bers [9] that the coefficients connecting different solutions of the same pseudoanalytic equation as in (1.17) give rise to a quasiregular mapping. In [9] these mappings are called pseudoanalytic functions of the
second kind. In our case this means that differentiating
\[ u_1(z, k) = a_1(z, z_0; k)u_1(z_0, k) + a_2(z, z_0; k)u_2(z_0, k) \]
with respect to \( k \) and using (8.2) with (1.14) give for \( \alpha = a_1 + ia_2 \)
\[ \partial_k \alpha(z, z_0; k) = \nu_{z_0}(k) \partial_k \alpha(z, z_0; k), \quad \nu_{z_0}(k) = \frac{h_-(z_0, k)}{h_+(z_0, k)}. \] (8.4)
Note that \( \alpha(z, z_0; k) \) inherits the continuous differentiability with respect to \( k \) from the \( u_j \)'s. For an explicit expression for \( \alpha \) in terms of \( u_1 \) and \( h \pm \) see (8.8) below.

Moreover, the second row of \( T_{z, z_0}(k) \) gives similarly a solution \( \beta = b_1 + ib_2 \)
for the same equation (8.4). In fact, if we write the \( \mu \) dependence explicitly
for \( u_i: u_i = u_i(\mu), i = 1, 2, \) then
\[ u_2(\mu) = iu_1(-\mu) \quad \text{and} \quad u_1(\mu) = -iu_2(-\mu). \] (8.5)
Thus \( \beta_\mu = i\alpha_{-\mu}. \)

We have now shown that the rows of the transport matrix produce quasi-
regular mappings, with respect to \( k \), satisfying (8.4). Our next task is to
determine their asymptotic behaviour at \( \infty \).

**Proposition 8.1.** Suppose \( z_0 \in \mathbb{C}, |z_0| > 1. \) Then:

a) For each fixed \( k \neq 0 \) and \( z_0, \) with respect to \( z, \)
\[ \alpha(z, z_0; k) = \exp(ikz + v(z)) \]
where \( v \in L^\infty. \)

b) For each fixed \( z \) and \( z_0, \) with respect to \( k, \)
\[ \alpha(z, z_0; k) = \exp(ik(z - z_0) + k\varepsilon(k)) \]
where \( \varepsilon(k) \to 0 \) as \( k \to \infty. \)

Up to the factor \( i, \beta_\mu = i\alpha_{-\mu} \) has the same asymptotics.

**Proof.** For a) we write
\[ u_1 = \frac{1}{2} \left( f_\mu + f_{-\mu} + f_\mu - f_{-\mu} \right) \]
\[ = f_\mu \left( 1 + \frac{f_\mu - f_{-\mu}}{f_\mu + f_{-\mu}} \right)^{-1} \left( 1 + \frac{f_{-\mu}}{f_\mu + f_{-\mu}} \right). \] (8.7)
All factors in the product are nonvanishing. Taking the logarithm and applying
(1.5), (1.6) lead to
\[ u_1(z, k) = \exp \left( ikz + O_k \left( \frac{1}{|z|} \right) \right). \]
On the other hand, dropping temporarily the fixed $k$ from the notation, we have

$$u_1(z) = \alpha h_+(z_0) - i\alpha h_-(z_0)$$

or

$$\frac{u_1(z)}{h_+(z_0)} = \alpha - i\frac{h_-(z_0)}{h_+(z_0)}.$$

This gives

$$\alpha = \left(1 - \left|\frac{h_-(z_0)}{h_+(z_0)}\right|^2\right)^{-1}\left(1 + i\frac{h_-(z_0) u_1(z)}{h_+(z_0) u_1(z)}\right)\frac{u_1(z)}{h_+(z_0)}.$$  (8.8)

According to Proposition 6.3, $|h_-(z_0)/h_+(z_0)| \leq 1/|z_0| < 1$ while $h_+(z_0)$ is constant for fixed $k$ and $z_0$. This proves a).

To prove b) note that

$$u_1 = h_+ \left(1 - i\frac{h_-}{h_+}\right) \quad \text{and} \quad h_+ = f_\mu \left(1 + i\frac{h_-}{h_+}\right)^{-1}.$$  

Again, the factors are continuous and pointwise nonvanishing. Therefore the identity (8.8) and Theorem 7.2 reduce the proof of b) to

**Lemma 8.2.** For each fixed $z \in \mathbb{C},$

$$\left|\frac{h_-(z, k)}{h_+(z, k)}\right| \leq 1 - e^{-|k|}|z(k).$$

**Proof.** By the definition of $h_+$ and $h_-$ it suffices to show

$$\inf_t \left|\frac{f_\mu - f_-}{f_\mu + f_-} + e^{it}\right| \geq e^{-|k|}|z(k).$$  (8.9)

For this, define

$$\Phi_t = e^{-it/2}(f_\mu \cos t/2 + if_- \sin t/2).$$

Then for each fixed $k,$

$$\Phi_t(z, k) = e^{ikz} \left(1 + O_k \left(\frac{1}{z}\right)\right) \quad \text{as} \quad z \to \infty$$

and

$$\partial_t \Phi_t = \mu e^{-it} \partial \Phi_t.$$  

Thus $\Phi_t = f_\lambda \mu$, $\lambda = e^{-it}$, the exponentially growing solution of (7.1) and (7.2) from Section 7. But

$$\frac{f_\mu - f_-}{f_\mu + f_-} + e^{it} = \frac{2e^{it} \Phi_t}{f_\mu + f_-} = \frac{f_\lambda \mu}{f_\mu} \frac{2e^{it}}{1 + M_\mu/M_-}.$$  (8.10)
By Theorem 7.2,
\begin{equation}
\label{eq:8.11}
e^{-|k|\varepsilon_1(k)} \leq |M_{\pm\mu}(z, k)| \leq e^{|k|\varepsilon_1(k)}
\end{equation}
and
\begin{equation}
\label{eq:8.12}
e^{-|k|\varepsilon_2(k)} \leq \inf_{\lambda \in \partial D} \left| \frac{f_{\lambda\mu}(z, k)}{f_{\mu}(z, k)} \right| \leq \sup_{\lambda \in \partial D} \left| \frac{f_{\lambda\mu}(z, k)}{f_{\mu}(z, k)} \right| \leq e^{|k|\varepsilon_2(k)}
\end{equation}
where \( \varepsilon_j(k) \to 0 \) as \( k \to \infty \). Since \( \text{Re}(M_{-\mu}/M_{\mu}) > 0 \) the inequality (8.9) follows from (8.10), (8.11) and (8.12). This finishes the proof of Lemma 8.2 and thus also the proof of Proposition 8.1.

Now the last remaining obstacle is Theorem 2; that is, we need to prove that the data determine the transport matrices \( T^\sigma_{z, z_0}(k) \). We know that in the \( k \)-variable the rows \( \alpha \) and \( \beta \) of \( T^\sigma_{z, z_0} \) are quasiregular mappings satisfying (8.4) and having the asymptotics given by Proposition 8.1 b).

It is not clear if the asymptotics (8.6) are strong enough to determine the individual solution. However, if we consider the entire family \( \{T^\sigma_{z, z_0} : z \in \mathbb{C} \} \), then the uniqueness does hold:

**Proof of Theorem 2.** Choose \( |z_0| > 1 \) and write \( \mu = (1-\sigma)/(1+\sigma) \), \( \mu = (1-\sigma)/(1+\sigma) \). Let \( \alpha_\mu(z, z_0, k) \) and \( \alpha_{\mu}(z, z_0, k) \) be defined by (1.17) and (1.19). Since neither \( \alpha_\mu \) or \( \alpha_{\mu} \) vanishes at any point we can define the corresponding logarithms \( \delta_\mu \) and \( \delta_{\mu} \) by
\begin{equation}
\delta_\mu(z, z_0; k) = \log \alpha_\mu(z, z_0, k) = ik(z - z_0) + k\varepsilon_1(k)
\end{equation}
\begin{equation}
\delta_{\mu}(z, z_0; k) = \log \alpha_{\mu}(z, z_0, k) = ik(z - z_0) + k\varepsilon_2(k)
\end{equation}
where for \( |k| \to \infty, \varepsilon_j(k) \to 0 \) by Proposition 8.1 b). Moreover, by Theorem 4.2,
\[ \delta_\mu(z, z_0; 0) \equiv \delta_{\mu}(z, z_0; 0) \equiv 0 \]
for all \( z \in \mathbb{C}, |z_0| > 1 \).

In addition, \( z \mapsto \delta_\mu(z, z_0; k) \) is continuous, \( \delta_\mu(z_0, z_0; k) = 0 \) and we have
\begin{equation}
\delta_\mu(z, z_0; k) = ikz \left( 1 + \frac{v_k(z)}{ikz} \right), \quad k \neq 0,
\end{equation}
where by Proposition 8.1 a), \( v_k \in L^\infty(\mathbb{C}) \) for each fixed \( k \in \mathbb{C} \). The use of degree theory or homotopy argument [33] gives that \( z \mapsto \delta_\mu(z, z_0; k) \) is surjective \( \mathbb{C} \to \mathbb{C} \).

To prove the theorem it suffices to show that if \( \Lambda_\sigma = \Lambda_{\sigma} \), then we have
\begin{equation}
\delta_{\mu}(z, z_0; k) \neq \delta_{\mu}(w, z_0; k), \quad \text{for} \ z \neq w \text{ and } k \neq 0.
\end{equation}
Namely then (8.15) and the surjectivity of \( z \mapsto \delta_\mu(z, z_0; k) \) show that necessarily we have \( \delta_{\mu}(z, z_0; k) = \delta_\mu(z, z_0, k) \) for all \( k, z \in \mathbb{C} \) and \( |z_0| > 1 \). Hence \( \alpha_\mu \equiv \alpha_{\mu} \). By (8.5) we have \( \beta_{\mu} = \beta_{\mu} \) as well and hence that \( T^\sigma_{z, z_0}(k) \equiv T^\sigma_{z, z_0}(k) \).
To show (8.15) fix \( z \neq w \) and note that by (8.4) \( \delta_\mu \) and \( \delta_\mu^\dagger \) satisfy
\[
\partial_\tau \delta = \nu_{z_0}(k)e^{\bar{\delta}-\delta k \partial_\delta}, \quad k \in \mathbb{C},
\]
where by Proposition 6.1 and the assumption \( \Lambda_\sigma = \Lambda_{\bar{\sigma}} \), the coefficient \( \nu_{z_0} \) is the same for both \( \delta_\mu \) and \( \delta_\mu^\dagger \). The difference
\[
g(k) = \delta_\mu(w, z_0; k) - \delta_\mu^\dagger(z, z_0; k)
\]
satisfies the equation
\[
\partial_\tau g - \nu_{z_0}e^{(\delta_\mu - \delta_\mu^\dagger)} \partial_\delta g = \nu_{z_0} \partial_\delta \delta_\mu \left( e^{(\delta_\mu - \delta_\mu^\dagger)} - e^{(\delta_\mu^\dagger - \delta_\mu)} \right).
\]
In other words, there exist functions \( \eta \) and \( \gamma \) such that
\[
\partial_\tau g - \eta \partial_\delta g = \gamma g
\]
with \( |\eta| \leq |\nu_{z_0}| \leq 1/|z_0| < 1 \) and \( |\gamma| \leq 2|\nu_{z_0}| ||\partial_\delta \delta_\mu|| \leq 2||\partial_\delta \delta_\mu|| \). From (8.13) we have \( g(k) = i(kw - z)k + k\varepsilon(k) \). Since \( \alpha_\mu, \alpha_\mu^\dagger \) are \( C^1 \) with respect to \( k \), we see that \( \gamma \) is locally bounded with respect to \( k \) and we may apply Proposition 3.3 b) (with respect to \( k \)) to obtain that \( g \) vanishes only at \( k = 0 \). This shows (8.15).

The proof of Theorem 1 is now immediate. If \( \Lambda_\sigma = \Lambda_{\bar{\sigma}} \), then by Proposition 6.1 we have \( u_j^\sigma(z) = u_j^\bar{\sigma}(z) \) for \( |z| > 1 \) and \( j = 1, 2 \). Theorem 2 with (1.17) then gives \( u_j^\sigma \equiv u_j^\bar{\sigma} \) and by (8.1) that \( \sigma \equiv \bar{\sigma} \).

Lastly, we describe how our uniqueness proof leads to a constructive procedure for recovering \( \sigma \). A somewhat similar algorithm is given in [18] for \( C^{1+\varepsilon} \)-conductivities. First we need to construct the geometric optics solutions \( M_{\pm \mu}(z, k) \) on \( \partial \mathbb{D} \). There is a Fredholm equation of the second kind on \( \partial \mathbb{D} \) described in terms of the operators \( Q_{\pm \mu} \) and having \( M_{\pm \mu} \big|_{\partial \mathbb{D}} \) as its unique solution; for details see [7]. Since by (2.9) and (2.11) these projection operators can be calculated from \( \Lambda_\sigma \) we obtain the boundary values of \( M_{\pm \mu} \). Another way to obtain these is to notice that by Lemma 2.2 and equation (5.1) the solution \( M_\mu(\cdot, k) \) belongs to the space \( \text{Range}(P_\mu^k) \cap \text{Range}(Q_0) \) where the operators \( P_\mu^k \) and \( Q_0 \) are defined by
\[
Q_0(g)(z) = \frac{1}{2}(I - i\mathcal{H}_0)(g)(z)
\]
\[
P_\mu^k(g)(z) = e^{-i k z}(I - Q_\mu)(e^{i k \cdot} g)(z).
\]
By Theorem 4.2 the space \( \text{Range}(P_\mu^k) \cap \text{Range}(Q_0) \) is one dimensional. Since it is defined by the data, we may use the asymptotics (1.6) to construct \( M_\mu \) on \( \partial \mathbb{D} \) from the equations
\[
(Q_0 - P_\mu^k)g = 0, \quad Q_0g = c M_\mu.
\]
By (2.8) this procedure also gives \( M_{- \mu} \) on \( \partial \mathbb{D} \).
The following step is to use the Fourier coefficients of $M_{\pm\mu}$ on $\partial D$ to construct $M_{\pm\mu}$ in the exterior of $D$. This gives by (1.7) and (1.21) the Beltrami coefficient $\nu_{z_0}(k)$ in (1.20). Finally by Theorem 2 we can uniquely solve equation (1.20) with the asymptotics (8.6) to obtain the transport matrix $T^\sigma_{z,z_0}(k)$ and hence $f_\mu(z,k)$ for $z \in D$. Formula (8.1) yields then the conductivity $\sigma$.

Acknowledgements. We are grateful for J. Kinnunen, M. Lassas, P. Ola and J. Väisälä for useful discussions. We also thank M. Orispää for helping in the preparation of the final form of the manuscript.

References


(Received August 8, 2003)