# Inverse spectral problems and closed exponential systems 

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#### Abstract

Consider the inverse eigenvalue problem of the Schrödinger operator defined on a finite interval. We give optimal and almost optimal conditions for a set of eigenvalues to determine the Schrödinger operator. These conditions are simple closedness properties of the exponential system corresponding to the known eigenvalues. The statements contain nearly all former results of this topic. We give also conditions for recovering the Weyl-Titchmarsh $m$-function from its values $m\left(\lambda_{n}\right)$.


## 1. Introduction

Consider the Schrödinger operator

$$
\begin{equation*}
L y=-y^{\prime \prime}+q(x) y \tag{1.1}
\end{equation*}
$$

over the segment $[0, \pi]$ with a potential

$$
\begin{equation*}
q \in L_{1}(0, \pi) \quad \text { real-valued. } \tag{1.2}
\end{equation*}
$$

The eigenvalue problem

$$
\begin{align*}
& L y=\lambda y \quad \text { on } \quad[0, \pi],  \tag{1.3}\\
& y(0) \cos \alpha+y^{\prime}(0) \sin \alpha=0,  \tag{1.4}\\
& y(\pi) \cos \beta+y^{\prime}(\pi) \sin \beta=0 \tag{1.5}
\end{align*}
$$

defines a sequence of eigenvalues

$$
\begin{equation*}
\lambda_{0}<\lambda_{1}<\cdots<\lambda_{n}<\ldots, \quad \lambda_{n} \in \mathbf{R}, \quad \lambda_{n} \rightarrow+\infty \tag{1.6}
\end{equation*}
$$

they form together the spectrum $\sigma(q, \alpha, \beta)$.
In the inverse eigenvalue problems we aim to recover the potential $q$ from a given set of eigenvalues (not necessarily taken from the same spectrum). The first result of this type is given in

[^0]Theorem A (Ambarzumian [1]). Let $q \in C[0, \pi]$ and consider the Neumann eigenvalue problem

$$
\left.y^{\prime}(0)=y^{\prime}(\pi)=0 \quad \text { (i.e. } \alpha=\beta=\pi / 2\right) .
$$

If the eigenvalues are $\lambda_{n}=n^{2}, n \geq 0$ then $q \equiv 0$.
Later it was observed by G. Borg that the knowledge of the first eigenvalue $\lambda_{0}=0$ plays a crucial role here; he also found the general rule that in most cases two spectra are needed to recover the potential:

Theorem B (Borg [5]). Let $q \in L_{1}(0, \pi), \sigma_{1}=\sigma(q, 0, \beta), \sigma_{2}=$ $\sigma\left(q, \alpha_{2}, \beta\right), \sin \alpha_{2} \neq 0$ and

$$
\tilde{\sigma}_{2}= \begin{cases}\sigma_{2} & \text { if } \sin \beta=0 \\ \sigma_{2} \backslash\left\{\lambda_{0}\right\} & \text { if } \sin \beta \neq 0 .\end{cases}
$$

Then $\sigma_{1} \cup \tilde{\sigma}_{2}$ determines the potential a.e. and no proper subset has the same property.

Here determination means that there is no other potential $q^{*} \in L_{1}(0, \pi)$ with $\sigma_{1}=\sigma_{1}^{*}, \tilde{\sigma}_{2}=\tilde{\sigma}_{2}^{*}$. There is a related extension:

Theorem C (Levinson [16]). Let $q \in L_{1}(0, \pi)$. If $\sin \left(\alpha_{1}-\alpha_{2}\right) \neq 0$ then the two spectra $\sigma\left(q, \alpha_{1}, \beta\right)$ and $\sigma\left(q, \alpha_{2}, \beta\right)$ determine the potential a.e.

By an interesting observation of Hochstadt and Lieberman, if half of the potential is known then one spectrum is enough to recover the other half of $q$ :

Theorem D (Hochstadt and Lieberman [11]). If $q \in L_{1}(0, \pi)$, then $q$ on $(0, \pi / 2)$ and the spectrum $\sigma(q, \alpha, \beta)$ determine $q$ a.e. on $(0, \pi)$.

This idea has been further developed by Gesztesy and Simon:
Theorem E (Gesztesy, Simon [9]). Let $q \in L_{1}(0, \pi)$ and $\pi / 2<a<\pi$. Then $q$ on $(0, a)$ and a subset $S \subset \sigma=\sigma(q, \alpha, \beta)$ of eigenvalues satisfying

$$
\#\{\lambda \in S: \lambda \leq t\} \geq 2(1-a / \pi) \#\{\lambda \in \sigma: \lambda \leq t\}+a / \pi-1 / 2
$$

for sufficiently large $t>0$, uniquely determine $q$ a.e. on $(0, \pi)$.
Another statement of this type is given in
Theorem F (del Rio, Gesztesy, Simon [7]). Let $q \in L_{1}(0, \pi)$, let $\sigma_{i}=$ $\sigma\left(q, \alpha_{i}, \beta\right)$ be three different spectra and $S \subset \sigma_{1} \cup \sigma_{2} \cup \sigma_{3}$. If

$$
\#\{\lambda \in S: \lambda \leq t\} \geq 2 / 3 \#\left\{\lambda \in \sigma_{1} \cup \sigma_{2} \cup \sigma_{3}: \lambda \leq t\right\}
$$

for large $t$ then the eigenvalues in $S$ determine $q$.

In Horváth [12] a similar but more general sufficient condition is given for the case when the known eigenvalues are taken from N different spectra.

The following statement provides a necessary and sufficient condition for a set of eigenvalues to determine the potential; it is one of the major new results of this paper. Before its formulation it is useful to fix some terminology. Let $1 \leq p \leq \infty$ and $1 / p+1 / p^{\prime}=1$. A system $\left\{\varphi_{n}: n \geq 1\right\}, \varphi_{n} \in L_{p^{\prime}}(0, \pi)$ is called closed in $L_{p}(a, b)$ if $h \in L_{p}(a, b), \int_{0}^{\pi} h \varphi_{n}=0$ for all $n$ implies $h=0$. This is equivalent to the completeness of the $\varphi_{n}$ in $L_{p^{\prime}}(0, \pi)$ if $p>1$. Let $\beta \in \mathbf{R}$ be given and let $q^{*}, q \in L_{p}(0, \pi)$. We say that the (different) values $\lambda_{n} \in \mathbf{R}$ are common eigenvalues of $q^{*}$ and $q$ if there exist $\alpha_{n} \in \mathbf{R}$ with

$$
\lambda_{n} \in \sigma\left(q, \alpha_{n}, \beta\right) \cap \sigma\left(q^{*}, \alpha_{n}, \beta\right)
$$

So every eigenvalue $\lambda_{n}$ is allowed to belong to different spectra. The values $\cot \alpha_{n}$ are defined by $q, \lambda_{n}$ and $\beta$; see (1.12) below. In the above cited theorems the eigenvalues are taken from at most three spectra; in [12] the $\lambda_{n}$ belong to finitely many spectra.

Let $0 \leq a<\pi$ and $\lambda_{n} \in \mathbf{R}$ be different values. By the statement " $\beta, q$ on $(0, a)$ and the eigenvalues $\lambda_{n}$ determine $q$ in $L_{p}$ "
we mean that there are no two different potentials $q^{*}, q \in L_{p}(0, \pi)$ with $q^{*}=q$ a.e. on $(0, a)$ such that the $\lambda_{n}$ are common eigenvalues of $q^{*}$ and $q$. By the statement
" $\beta, q$ on $(0, a)$ and the eigenvalues $\lambda_{n}$ do not determine $q$ in $L_{p}$ "
we mean that for every $q \in L_{p}(0, \pi)$ there exists a different potential $q^{*} \in$ $L_{p}(0, \pi)$ with $q^{*}=q$ a.e. on $(0, a)$ such that the $\lambda_{n}$ are common eigenvalues of $q^{*}$ and $q$.

Theorem 1.1. Let $1 \leq p \leq \infty, q \in L_{p}(0, \pi), 0 \leq a<\pi$ and let $\lambda_{n} \in$ $\sigma\left(q, \alpha_{n}, 0\right)$ be real numbers with $\lambda_{n} \nrightarrow-\infty$. Then $\beta=0, q$ on $(0, a)$ and the eigenvalues $\lambda_{n}$ determine $q$ in $L_{p}$ if and only if the system

$$
\begin{equation*}
e(\Lambda)=\left\{e^{ \pm 2 i \mu x}, e^{ \pm 2 i \sqrt{\lambda_{n}} x}: n \geq 1\right\} \tag{1.7}
\end{equation*}
$$

is closed in $L_{p}(a-\pi, \pi-a)$ for some (for any) $\mu \neq \pm \sqrt{\lambda_{n}}$.
In case $\sin \beta \neq 0$ we find a different situation. First we state a sufficient condition:

Theorem 1.2. Let $1 \leq p \leq \infty, q \in L_{p}(0, \pi), \sin \beta \neq 0, \lambda_{n} \in \sigma\left(q, \alpha_{n}, \beta\right)$, $\lambda_{n} \nrightarrow-\infty$ and $0 \leq a<\pi$. If the set

$$
\begin{equation*}
e_{0}(\Lambda)=\left\{e^{ \pm 2 i \sqrt{\lambda_{n}} x}: n \geq 1\right\} \tag{1.8}
\end{equation*}
$$

is closed in $L_{p}(a-\pi, \pi-a)$ then $q$ on $(0, a)$ and the eigenvalues $\lambda_{n}$ determine $q$ in $L_{p}$.

The following example shows that the above closedness condition (1.8) is sharp in some cases:

Proposition 1.3. Let $\beta=\pi / 2$,

$$
\begin{aligned}
q(x) & = \begin{cases}0 & \text { on }(0, \pi / 2) \\
1 & \text { on }(\pi / 2, \pi),\end{cases} \\
q^{*}(x) & = \begin{cases}1 & \text { on }(0, \pi / 2) \\
0 & \text { on }(\pi / 2, \pi) .\end{cases}
\end{aligned}
$$

Then for the set of all common eigenvalues of $q^{*}$ and $q$, the system $e_{0}(\Lambda)$ has deficiency 1 in $L_{p}(-\pi, \pi), 1 \leq p<\infty$. In other words, the system $e_{1}(\Lambda)=$ $\left\{e^{2 i \mu x}, e^{ \pm 2 i \sqrt{\lambda_{n}} x}: n \geq 1\right\}$ with $\mu \neq \pm \sqrt{\lambda_{n}}$ is closed in $L_{p}(-\pi, \pi)$.

Remark. In the important special cases considered by Borg in Theorem B, however, the closedness of $e_{0}(\Lambda)$ is not an optimal condition in Theorem 1.2; in those situations the codimension of $e_{0}(\Lambda)$ is 1 for the set of eigenvalues defining the potential (see $\S 4$ ).

Remark. Denote by $v(x, \lambda)$ the solution of

$$
\begin{align*}
& -v^{\prime \prime}+q(x) v=\lambda v \quad \text { on }(0, \pi),  \tag{1.9}\\
& \quad v(\pi, \lambda)=\sin \beta, \quad v^{\prime}(\pi, \lambda)=-\cos \beta \tag{1.10}
\end{align*}
$$

and let $v^{*}(x, \lambda)$ be the same function defined by $q^{*}$ instead of $q$. Then the common eigenvalues of $q^{*}$ and $q$ under the boundary condition (1.5) are precisely the solutions $\lambda_{n} \in \mathbf{R}$ of the equation

$$
\begin{equation*}
v(0, \lambda) v^{* \prime}(0, \lambda)=v^{\prime}(0, \lambda) v^{*}(0, \lambda) . \tag{1.11}
\end{equation*}
$$

In this case $\lambda_{n} \in \sigma\left(q^{*}, \alpha_{n}, \beta\right) \cap \sigma\left(q, \alpha_{n}, \beta\right)$ with

$$
\begin{equation*}
\cot \alpha_{n}=-\frac{v^{\prime}\left(0, \lambda_{n}\right)}{v\left(0, \lambda_{n}\right)}=-\frac{v^{* \prime}\left(0, \lambda_{n}\right)}{v^{*}\left(0, \lambda_{n}\right)} . \tag{1.12}
\end{equation*}
$$

In looking for a necessary condition for $\sin \beta \neq 0$ we have to avoid the Ambarzumian-type exceptional cases where less than two spectra are enough to determine the potential. To this end, introduce the following minimality condition
(M) There exists $h \in L_{p}(a, \pi)$ such that

$$
\int_{a}^{\pi} h \neq 0 \quad \text { but } \quad \int_{a}^{\pi} h(x)\left[v^{2}\left(x, \lambda_{n}\right)-1 / 2 \sin ^{2} \beta\right] d x=0 \quad \forall n .
$$

For $1<p$ this condition can also be formulated in the following form: the closed subspace generated in $L_{p^{\prime}}(a, \pi)$ by the functions $v^{2}\left(x, \lambda_{n}\right)-1 / 2 \sin ^{2} \beta$ does not contain the constant function 1 ; here $1 / p+1 / p^{\prime}=1$.

ThEOREM 1.4. Let $\sin \beta \neq 0,0 \leq a<\pi, 1 \leq p \leq \infty$ and $\lambda_{n}, n \geq 1$ be different real numbers with $\lambda_{n} \nrightarrow-\infty$. Suppose (M) and that

$$
e(\Lambda)=\left\{e^{ \pm 2 i \mu x}, e^{ \pm 2 i \sqrt{\lambda_{n}} x}\right\}
$$

is not closed in $L_{p}(a-\pi, \pi-a)$, where $\mu \neq \pm \sqrt{\lambda_{n}}$. Then $q$ on $(0, a)$ and the eigenvalues $\lambda_{n}$ do not determine $q$ in $L_{p}$.

Define the Weyl-Titchmarsh m-function corresponding to the problem (1.3), (1.5) by

$$
\begin{equation*}
m_{\beta}(\lambda)=\frac{v^{\prime}(0, \lambda)}{v(0, \lambda)} \tag{1.13}
\end{equation*}
$$

where $v(x, \lambda)$ is given in (1.9), (1.10). It is a meromorphic function having poles at the zeros of $v(0, \lambda)$.

Theorem G (Borg [6], Marchenko [18]). The potential and the value $\tan \beta$ can be recovered from the $m$-function $m_{\beta}(\lambda)$.

In the context of the $m$-function Theorem 1.1 and Theorem 1.2 can be generalized in the following way:

THEOREM 1.5. Let $1 \leq p \leq \infty$ and $\lambda_{n}, n \geq 1$, be arbitrary different real numbers with $\lambda_{n} \nrightarrow-\infty$. Let $\beta_{1}, \beta_{2} \in \mathbf{R}, q^{*}, q \in L_{p}(0, \pi)$ and consider the m-functions $m_{\beta_{1}}$ and $m_{\beta_{2}}^{*}$, defined by $q$ and $q^{*}$ respectively.

- If the system $e_{0}(\Lambda)$ is closed in $L_{p}(-\pi, \pi)$ then

$$
\begin{equation*}
m_{\beta_{1}}\left(\lambda_{n}\right)=m_{\beta_{2}}^{*}\left(\lambda_{n}\right), \quad n \geq 1 \tag{1.14}
\end{equation*}
$$

implies $m_{\beta_{1}} \equiv m_{\beta_{2}}^{*}\left(\right.$ so $\tan \beta_{1}=\tan \beta_{2}$ and $\left.q^{*}=q\right)$.

- Let $\sin \beta_{1} \cdot \sin \beta_{2}=0$. Then (1.14) implies $\sin \beta_{1}=\sin \beta_{2}=0$. In this case (1.14) implies $m_{0}^{*} \equiv m_{0}$ if and only if the system $e(\Lambda)$ is closed in $L_{p}(-\pi, \pi)$.

Remark. We allow in (1.14) that both sides be infinite.
A former result of this type is given in
Theorem H (del Rio, Gesztesy, Simon [7]). Denote $c_{+}=\max (c, 0)$ and let $q \in L_{1}(0, \pi)$. If $\lambda_{n}>0$ are distinct numbers satisfying

$$
\begin{equation*}
\sum_{n=0}^{\infty} \frac{\left(\lambda_{n}-n^{2} / 4\right)_{+}}{1+n^{2}}<\infty \tag{1.15}
\end{equation*}
$$

then the values $m_{\beta}\left(\lambda_{n}\right)$ determine $m_{\beta}($ and $\tan \beta)$.

Since (1.15) implies the closedness of $e_{0}(\Lambda)$, this statement is a special case of Theorem 1.5; see Section 4.

Finally we mention the following localized version of Theorem G. It was first given in Simon [20]; see also Gesztesy and Simon [8], [10] and Bennewitz [4].

Theorem I ([20], [8], [10], [4]). Let $\beta_{1}, \beta_{2} \in \mathbf{R}, q^{*}, q \in L_{1}(0, \pi)$, $0 \leq a<\pi$. Then $q^{*}=q$ a.e. on $(0, a)$ if and only if for every $\varepsilon>0$

$$
\begin{equation*}
m_{\beta_{1}}(\lambda)-m_{\beta_{2}}^{*}(\lambda)=\mathbf{O}\left(e^{-2(a-\varepsilon)|\Im \sqrt{\lambda}|}\right) \tag{1.16}
\end{equation*}
$$

holds along a nonreal ray $\arg \lambda=\gamma, \sin \gamma \neq 0$.
From this statement the following generalization of Theorem 1.5 can be given:

Theorem 1.6. Let $1 \leq p \leq \infty$ and $\lambda_{n}, n \geq 1$ be arbitrary different real numbers with $\lambda_{n} \nrightarrow-\infty$. Let $\beta_{1}, \beta_{2} \in \mathbf{R}, q^{*}, q \in L_{p}(0, \pi)$ and suppose that (1.16) holds for every $\varepsilon>0$ along a nonreal ray.

- If the system $e_{0}(\Lambda)$ is closed in $L_{p}(a-\pi, \pi-a)$ then (1.14) implies $m_{\beta_{1}} \equiv m_{\beta_{2}}^{*}$.
- Let $\sin \beta_{1} \cdot \sin \beta_{2}=0$. Then (1.14) yields $\sin \beta_{1}=\sin \beta_{2}=0$. In this case (1.14) implies $m_{0}^{*} \equiv m_{0}$ if and only if the system $e(\Lambda)$ is closed in $L_{p}(a-\pi, \pi-a)$.

Remark. The statements of Theorems 1.1 and 1.5 for the Schrödinger operators on the half-line are investigated in the forthcoming paper [13]. It turns out that the inverse eigenvalue problem is closely related to the inverse scattering problem with fixed energy.

The organization of this paper is as follows. In Section 2 we provide the proof of Theorem 1.1; the main ingredient is Lemma 2.1. Some technical background needed in the proof is given only in Section 5 . Section 3 is devoted to prove Theorems $1.2,1.4,1.5$ and 1.6 by modifying the procedure presented in Section 2. The applications of the new results are collected in Section 4; we show how the above-mentioned former results can be presented as special cases of Theorems 1.1 to 1.6. This requires the use of some standard tools from the theory of nonharmonic Fourier series, more precisely, some closedness and basis tests for exponential systems. Finally at the end of Section 4 we check the properties of the counterexample formulated in Proposition 1.3.

## 2. Proof of Theorem 1.1

In this section we provide the proof of Theorem 1.1. We start with some lemmas.

Lemma 2.1. Let $B_{1}$ and $B_{2}$ be Banach spaces. For every $q \in B_{1}$ a continuous linear operator

$$
A_{q}: B_{1} \rightarrow B_{2}
$$

is defined so that for some $q_{0} \in B_{1}$

$$
\begin{equation*}
A_{q_{0}}: B_{1} \rightarrow B_{2} \text { is an (onto) isomorphism, } \tag{2.1}
\end{equation*}
$$

and the mapping $q \mapsto A_{q}$ is Lipschitzian in the sense that

$$
\begin{equation*}
\left\|\left(A_{q^{*}}-A_{q}\right) h\right\| \leq c\left(q_{0}\right)\left\|q^{*}-q\right\|\|h\| \quad \forall h, q, q^{*} \in B_{1},\|q\|,\left\|q^{*}\right\| \leq 2\left\|q_{0}\right\| \tag{2.2}
\end{equation*}
$$

the constant $c\left(q_{0}\right)$ being independent of $q, q^{*}$ and $h$. Then the set $\left\{A_{q}\left(q-q_{0}\right)\right.$ : $\left.q \in B_{1}\right\}$ contains a ball in $B_{2}$ with center at the origin.

Proof. Let $G_{0} \in B_{2}$ be an arbitrary element, the norm of which is small in a sense to be specified later. Our task is to find an element $q^{*} \in B_{1}$ such that

$$
\begin{equation*}
A_{q^{*}}\left(q^{*}-q_{0}\right)=G_{0} \tag{2.3}
\end{equation*}
$$

This will be done by the following iteration. The vector $q_{0}^{*}$ is defined by

$$
\begin{equation*}
A_{q_{0}}\left(q_{0}^{*}-q_{0}\right)=G_{0} \tag{2.4}
\end{equation*}
$$

and $q_{k+1}^{*}$ by

$$
\begin{equation*}
A_{q_{0}}\left(q_{k+1}^{*}-q_{0}\right)=G_{0}-\left(A_{q_{k}^{*}}-A_{q_{0}}\right)\left(q_{k}^{*}-q_{0}\right), \quad k \geq 0 \tag{2.5}
\end{equation*}
$$

This is justified by (2.1). We state that $q_{k}^{*} \rightarrow q^{*}$, a solution of (2.3). Indeed, consider the following corollary of (2.5):

$$
\begin{equation*}
A_{q_{0}}\left(q_{k+1}^{*}-q_{k}^{*}\right)=-\left(A_{q_{k}^{*}}-A_{q_{0}}\right)\left(q_{k}^{*}-q_{k-1}^{*}\right)-\left(A_{q_{k}^{*}}-A_{q_{k-1}^{*}}\right)\left(q_{k-1}^{*}-q_{0}\right) ; \tag{2.6}
\end{equation*}
$$

if $k=0$, we use instead

$$
A_{q_{0}}\left(q_{1}^{*}-q_{0}^{*}\right)=-\left(A_{q_{0}^{*}}-A_{q_{0}}\right)\left(q_{0}^{*}-q_{0}\right)
$$

Using the conditions (2.1), (2.2) we get from the formulae (2.4), (2.6') and (2.6) that

$$
\begin{gather*}
\left\|q_{0}^{*}-q_{0}\right\| \leq c_{1}\left\|G_{0}\right\|  \tag{2.7}\\
\left\|q_{1}^{*}-q_{0}^{*}\right\| \leq c_{1}\left\|q_{0}^{*}-q_{0}\right\|^{2} \text { if }\left\|q_{0}^{*}\right\| \leq 2\left\|q_{0}\right\| \tag{2.8}
\end{gather*}
$$

$$
\begin{array}{r}
\left\|q_{k+1}^{*}-q_{k}^{*}\right\| \leq c_{1}\left\|q_{k}^{*}-q_{k-1}^{*}\right\|\left(\left\|q_{k}^{*}-q_{0}\right\|+\left\|q_{k-1}^{*}-q_{0}\right\|\right),  \tag{2.9}\\
\text { if }\left\|q_{k}^{*}\right\| \leq 2\left\|q_{0}\right\|,\left\|q_{k-1}^{*}\right\| \leq 2\left\|q_{0}\right\|, k \geq 1
\end{array}
$$

with a constant $c_{1}$ independent of the $q_{k}^{*}, k \geq 0$, and of $G_{0}$. We suppose that $G_{0}$ is small enough to ensure

$$
\begin{equation*}
8 c_{1}^{2}\left\|G_{0}\right\| \leq 1, \quad c_{1}\left\|G_{0}\right\| \leq 1 / 2\left\|q_{0}\right\| \tag{2.10}
\end{equation*}
$$

and we prove that

$$
\begin{equation*}
\left\|q_{k+1}^{*}-q_{k}^{*}\right\| \leq 1 / 2\left\|q_{k}^{*}-q_{k-1}^{*}\right\|, \quad\left\|q_{k}^{*}\right\| \leq 2\left\|q_{0}\right\| \text { if } k \geq 1 \tag{2.11}
\end{equation*}
$$

Indeed, (2.7) and (2.10) imply $\left\|q_{0}^{*}\right\| \leq 3 / 2\left\|q_{0}\right\|$ and then by (2.8)

$$
\left\|q_{1}^{*}-q_{0}^{*}\right\| \leq c_{1}\left\|q_{0}^{*}-q_{0}\right\|^{2} \leq c_{1}^{2}\left\|G_{0}\right\| \cdot\left\|q_{0}^{*}-q_{0}\right\| \leq 1 / 2\left\|q_{0}^{*}-q_{0}\right\| \leq 1 / 4\left\|q_{0}\right\|
$$

and then

$$
\left\|q_{1}^{*}\right\| \leq\left\|q_{1}^{*}-q_{0}^{*}\right\|+\left\|q_{0}^{*}-q_{0}\right\|+\left\|q_{0}\right\| \leq(1 / 4+1 / 2+1)\left\|q_{0}\right\| .
$$

Consequently by (2.9)

$$
\begin{aligned}
\left\|q_{2}^{*}-q_{1}^{*}\right\| & \leq c_{1}\left\|q_{1}^{*}-q_{0}^{*}\right\|\left(\left\|q_{1}^{*}-q_{0}\right\|+\left\|q_{0}^{*}-q_{0}\right\|\right) \\
& \leq c_{1}\left\|q_{1}^{*}-q_{0}^{*}\right\|\left(\left\|q_{1}^{*}-q_{0}^{*}\right\|+2\left\|q_{0}^{*}-q_{0}\right\|\right) \\
& \leq\left\|q_{1}^{*}-q_{0}^{*}\right\|\left(c_{1}^{2}\left\|q_{0}^{*}-q_{0}\right\|^{2}+2 c_{1}\left\|q_{0}^{*}-q_{0}\right\|\right) \\
& \leq\left\|q_{1}^{*}-q_{0}^{*}\right\|\left(c_{1}^{4}\left\|G_{0}\right\|^{2}+2 c_{1}^{2}\left\|G_{0}\right\|\right) \leq 1 / 2\left\|q_{1}^{*}-q_{0}^{*}\right\|
\end{aligned}
$$

which is (2.11) for $k=1$. Now suppose (2.11) below a fixed value of $k$ and prove it for that $k$. We have

$$
\begin{aligned}
\left\|q_{i}^{*}-q_{0}\right\| & \leq\left\|q_{i}^{*}-q_{i-1}^{*}\right\|+\cdots+\left\|q_{1}^{*}-q_{0}^{*}\right\|+\left\|q_{0}^{*}-q_{0}\right\| \\
& \leq 2\left\|q_{1}^{*}-q_{0}^{*}\right\|+\left\|q_{0}^{*}-q_{0}\right\| \leq 2 c_{1}\left\|q_{0}^{*}-q_{0}\right\|^{2}+\left\|q_{0}^{*}-q_{0}\right\| \\
& \leq 2 c_{1}^{3}\left\|G_{0}\right\|^{2}+c_{1}\left\|G_{0}\right\| \leq 2 c_{1}\left\|G_{0}\right\| \leq\left\|q_{0}\right\|
\end{aligned}
$$

for $i \leq k$ and then

$$
\left\|q_{k}^{*}\right\| \leq\left\|q_{k}^{*}-q_{0}\right\|+\left\|q_{0}\right\| \leq 2\left\|q_{0}\right\| .
$$

Consequently

$$
\begin{aligned}
\left\|q_{k+1}^{*}-q_{k}^{*}\right\| & \leq c_{1}\left\|q_{k}^{*}-q_{k-1}^{*}\right\|\left(\left\|q_{k}^{*}-q_{0}\right\|+\left\|q_{k-1}^{*}-q_{0}\right\|\right) \\
& \leq\left\|q_{k}^{*}-q_{k-1}^{*}\right\|\left(4 c_{1}^{4}\left\|G_{0}\right\|^{2}+2 c_{1}^{2}\left\|G_{0}\right\|\right) \leq 1 / 2\left\|q_{k}^{*}-q_{k-1}^{*}\right\|
\end{aligned}
$$

and so (2.11) is proved and then $q_{k}^{*} \rightarrow q^{*}$ in $B_{1}$. Now

$$
\begin{equation*}
A_{q_{0}}\left(q_{k+1}^{*}-q_{0}\right)=G_{0}+\left(A_{q^{*}}-A_{q_{k}^{*}}\right)\left(q_{k}^{*}-q_{0}\right)-\left(A_{q^{*}}-A_{q_{0}}\right)\left(q_{k}^{*}-q_{0}\right) . \tag{2.12}
\end{equation*}
$$

Since

$$
\left\|\left(A_{q^{*}}-A_{q_{k}^{*}}\right)\left(q_{k}^{*}-q_{0}\right)\right\| \leq c\left\|q^{*}-q_{k}^{*}\right\| \cdot\left\|q_{k}^{*}-q_{0}\right\| \rightarrow 0 \quad k \rightarrow \infty,
$$

we can take the limit in (2.12) to obtain

$$
A_{q_{0}}\left(q^{*}-q_{0}\right)=G_{0}-\left(A_{q^{*}}-A_{q_{0}}\right)\left(q^{*}-q_{0}\right) .
$$

This is (2.3) so the proof is complete.

In the following statement the point a) (in a less general situation) and the formula (2.16) are due to Gesztesy and Simon [9], [10]. We give the whole proof for the sake of completeness.

Lemma 2.2. Let $0 \leq a<\pi, q, q^{*} \in L_{1}(0, \pi), q^{*}=q$ a.e. on $(0, a)$. Consider the function

$$
\begin{equation*}
F(z)=v^{*}(a, z) v^{\prime}(a, z)-v(a, z) v^{* \prime}(a, z) \tag{2.13}
\end{equation*}
$$

where $v$ and $v^{*}$ are defined by $q$ and $q^{*}$ respectively in (1.9), (1.10) with $\beta=0$. The derivatives in (2.13) refer to $x$. Then
a) The real zeros of $F(z)$ are precisely the common eigenvalues of $q$ and $q^{*}$; in other words, all values $z=\lambda \in \mathbf{R}$ for which there exists $\alpha \in \mathbf{R}$ with $\lambda \in \sigma\left(q^{*}, \alpha, 0\right) \cap \sigma(q, \alpha, 0)$.
b) If $\lambda_{n} \nrightarrow-\infty$ holds for the (infinitely many) common eigenvalues of $q^{*}$ and $q$ then

$$
\begin{equation*}
\int_{a}^{\pi}\left(q^{*}-q\right)=0 . \tag{2.14}
\end{equation*}
$$

Proof. $F(\lambda)=0$ if and only if the initial condition vectors $\left(v(a, \lambda), v^{\prime}(a, \lambda)\right)$ and $\left(v^{*}(a, \lambda), v^{* \prime}(a, \lambda)\right)$ are parallel. Since $q^{*}=q$ a.e. on $(0, a)$, this means that $v^{*}$ and $v$ are identical on $[0, a]$ up to a constant factor. In other words we have $\lambda \in \sigma\left(q^{*}, \alpha, 0\right) \cap \sigma(q, \alpha, 0)$ with $\tan \alpha=-\frac{v(0, \lambda)}{v^{\prime}(0, \lambda)}=-\frac{v^{*}(0, \lambda)}{v^{*}(0, \lambda)}$. This proves a). To show b) take the function

$$
\begin{equation*}
F(x, z)=v^{*}(x, z) v^{\prime}(x, z)-v(x, z) v^{* \prime}(x, z) . \tag{2.15}
\end{equation*}
$$

Now

$$
\begin{aligned}
\frac{\partial F}{\partial x}(x, z) & =v^{*}(x, z) v^{\prime \prime}(x, z)-v(x, z) v^{* \prime \prime}(x, z) \\
& =\left(q(x)-q^{*}(x)\right) v(x, z) v^{*}(x, z)
\end{aligned}
$$

which implies

$$
\begin{equation*}
F(z)=-\int_{a}^{\pi} \frac{\partial F}{\partial x}(x, z) d x=\int_{a}^{\pi}\left(q^{*}(x)-q(x)\right) v(x, z) v^{*}(x, z) d x \tag{2.16}
\end{equation*}
$$

If the zeros $\lambda_{n}$ have a finite accumulation point then the entire function $F(z)$ is identically zero, which implies $m^{*}=m$ and $q^{*}=q$; in this case (2.14) is
obvious. Otherwise the $\lambda_{n}$ have a subsequence tending to $+\infty$. By Lemma 5.2

$$
\begin{align*}
2\left(z^{2}-\mu^{2}\right) F\left(z^{2}\right)= & 2\left(z^{2}-\mu^{2}\right) \int_{a}^{\pi}\left(q^{*}(x)-q(x)\right) v\left(x, z^{2}\right) v^{*}\left(x, z^{2}\right) d x  \tag{2.17}\\
= & \int_{a}^{\pi}\left(q^{*}-q\right)-\int_{a}^{\pi}\left(q^{*}(x)-q(x)\right) \cos 2 z(\pi-x) d x \\
& -\int_{a}^{\pi}\left(q^{*}(x)-q(x)\right) \int_{0}^{2(\pi-x)} \cos z \tau M\left(\pi-x, \tau, \mu^{2}\right) d \tau d x \\
= & I_{1}-I_{2}-I_{3} .
\end{align*}
$$

Here $I_{3}$ has the form

$$
\begin{equation*}
I_{3}=\int_{0}^{2(\pi-a)} \cos z \tau \int_{a}^{\pi-\tau / 2}\left(q^{*}(x)-q(x)\right) M\left(\pi-x, \tau, \mu^{2}\right) d x d \tau \tag{2.18}
\end{equation*}
$$

This means that for the subsequence of values $z=\sqrt{\lambda_{n}}$ tending to $+\infty$ we have $I_{3} \rightarrow 0$. Since $I_{2} \rightarrow 0$ is obvious, from $F\left(\lambda_{n}\right)=0$ we infer (2.14) as asserted.

Proof of Theorem 1.1. We consider the closedness of the system

$$
\begin{equation*}
C(\Lambda)=\left\{\cos 2 \mu x, \cos 2 \sqrt{\lambda_{n}} x: n \geq 1\right\} \tag{2.19}
\end{equation*}
$$

in $L_{p}(0, \pi-a)$ instead of that of $e(\Lambda)$ in $L_{p}(a-\pi, \pi-a)$; this is justified in Lemma 5.4.

The if part. If the system $C(\Lambda)$ is closed in $L_{p}(0, \pi-a)$ then the eigenvalues $\lambda_{n}$ and $\left.q\right|_{(0, a)}$ determine $q$ on the whole $(0, \pi)$. Suppose indirectly that there exists another potential $q^{*} \in L_{p}$ with $q^{*}=q$ a.e. on $(0, a)$ and $\lambda_{n} \in \sigma\left(q^{*}, \alpha_{n}, 0\right) \cap \sigma\left(q, \alpha_{n}, 0\right)$ for some $\alpha_{n} \in \mathbf{R}$. Define $F(z)$ by (2.13); then $F\left(\lambda_{n}\right)=0(n \geq 1)$ and $F \not \equiv 0$. The function

$$
G(z)=-2\left(z^{2}-\mu^{2}\right) F\left(z^{2}\right)
$$

has zeros at $\pm \mu, \pm \sqrt{\lambda_{n}}$. From (2.14) we get

$$
\begin{equation*}
G(z)=\int_{a}^{\pi}\left(q^{*}(x)-q(x)\right)\left[1-2\left(z^{2}-\mu^{2}\right) v\left(x, z^{2}\right) v^{*}\left(x, z^{2}\right)\right] d x . \tag{2.20}
\end{equation*}
$$

Define the linear operators

$$
\begin{gathered}
A_{q^{*}}: L_{p}(a, \pi) \rightarrow L_{p}(a, \pi) \\
\left(A_{q^{*}} h\right)(x)=h(x)+2 \int_{a}^{x} h(\tau) M\left(\pi-\tau, 2(\pi-x), \mu^{2}, q, q^{*}\right) d \tau
\end{gathered}
$$

Then Lemma 5.2 gives, after an interchange of integrations,

$$
\begin{align*}
& \int_{a}^{\pi}\left(q^{*}(x)-q(x)\right)\left[1-2\left(z^{2}-\mu^{2}\right) v\left(x, z^{2}\right) v^{*}\left(x, z^{2}\right)\right] d x  \tag{2.21}\\
&=\int_{a}^{\pi} \cos 2 z(\pi-x)\left[A_{q^{*}}\left(q^{*}-q\right)\right](x) d x
\end{align*}
$$

Observe that

$$
\begin{equation*}
A_{q^{*}}: L_{p}(a, \pi) \rightarrow L_{p}(a, \pi) \text { is an isomorphism. } \tag{2.22}
\end{equation*}
$$

Indeed, the Volterra operator

$$
h \mapsto 2 \int_{a}^{x} h(\tau) M\left(\pi-\tau, 2(\pi-x), \mu^{2}, q, q^{*}\right) d \tau
$$

with continuous kernel is known to have the spectrum $\sigma=\{0\}$. In particular, $-1 \notin \sigma$ i.e. $A_{q^{*}}$ is an isomorphism. Now if $q^{*} \neq q$ then $A_{q^{*}}\left(q^{*}-q\right) \neq 0$; hence by $(2.20)$ and (2.21) the system $C(\Lambda)$ is not closed in $L_{p}(0, \pi-a)$. This contradiction proves the if part of Theorem 1.1.

The only if part. If $C(\Lambda)$ is not closed in $L_{p}(0, \pi-a)$ and if $\lambda_{n} \nrightarrow-\infty$ then for every $q \in L_{p}(0, \pi)$ there exists $q^{*} \in L_{p}(0, \pi), q^{*} \neq q$ but $q^{*}=q$ a.e. on $(0, a)$ and there exist values $\alpha_{n} \in \mathbf{R}$ with $\lambda_{n}=\sigma\left(q^{*}, \alpha_{n}, 0\right) \cap \sigma\left(q, \alpha_{n}, 0\right)$ for all $n \geq 1$. Indeed, since $C(\Lambda)$ is not closed, there exists a function $0 \neq h \in L_{p}(0, \pi-a)$ such that

$$
\begin{equation*}
G_{0}(z) \stackrel{\text { def }}{=} \int_{0}^{\pi-a} h(x) \cos 2 z x d x \tag{2.23}
\end{equation*}
$$

has zeros at $\pm \mu$ and $\pm \sqrt{\lambda_{n}}$. Our task is to show that for every $q \in L_{p}(0, \pi)$ there exists $q^{*} \in L_{p}(0, \pi), q^{*} \neq q, q^{*}=q$ a.e. on $(0, a)$ such that

$$
\begin{equation*}
\gamma G_{0}(z)=\int_{a}^{\pi}\left(q^{*}(x)-q(x)\right)\left[1-2\left(z^{2}-\mu^{2}\right) v\left(x, z^{2}\right) v^{*}\left(x, z^{2}\right)\right] d x \tag{2.24}
\end{equation*}
$$

holds for some constant $\gamma \neq 0$. Indeed, $G_{0}(\mu)=0$ and (2.24) gives (2.14) and then the function $F(z)$ defined in (2.13) has zeros $F\left(\lambda_{n}\right)=0$; i.e. the $\lambda_{n}$ are common eigenvalues of $q^{*}$ and $q$. Taking into account (2.21), (2.23) and (2.24), our task is to find $q^{*}$ with

$$
\begin{equation*}
\gamma h(\pi-x)=A_{q^{*}}\left(q^{*}-q\right)(x) \text { a.e. for some } \gamma \neq 0 . \tag{2.25}
\end{equation*}
$$

We check this representation by Lemma 2.1 applied with $B_{1}=B_{2}=L_{p}(a, \pi)$. The condition (2.1) is verified in (2.22) and (2.2) follows from Lemma 5.2, since
if $q, q^{*}, q^{* *} \in L_{p}$ with norms $\leq D$ then

$$
\begin{aligned}
&\left\|\left(A_{q^{* *}}-A_{q^{*}}\right) h\right\|=2\left\{\int_{a}^{\pi} \mid \int_{a}^{x} h(\tau)\left[M\left(\pi-\tau, 2(\pi-x), \mu^{2}, q, q^{* *}\right)\right.\right. \\
&\left.\left.-M\left(\pi-\tau, 2(\pi-x), \mu^{2}, q, q^{*}\right)\right]\left.d \tau\right|^{p} d x\right\}^{1 / p} \\
& \leq c(D)\left\|q^{* *}-q^{*}\right\|\left\{\int_{a}^{\pi}\left(\int_{a}^{x}|h|\right)^{p} d x\right\}^{1 / p} \leq c_{1}(D)\left\|q^{* *}-q^{*}\right\| \cdot\|h\|
\end{aligned}
$$

with straightforward modifications for $p=\infty$. So Lemma 2.1 applies and this shows the possibility of the representation (2.25) with sufficiently small $\gamma \neq 0$. The proof is complete.

## 3. Proofs of Theorems 1.2 to 1.6

In this part of the paper we give the proofs of the remaining new results. They are modifications of the proof of Theorem 1.1 or consequences of already proved results. The proof of Proposition 1.3 is deferred to Section 4.

Lemma 3.1. Let $1 \leq p \leq \infty, q, q^{*} \in L_{p}(0, \pi), 0 \leq a<\pi, q^{*}=q$ a.e. on $(0, a)$. Let $F(z)$ be defined by (2.13), where the functions $v$ and $v^{*}$ are as given in (1.9), (1.10) with $q$ and $q^{*}$. Let $\sin \beta \neq 0$. Then
a) The real zeros of $F(z)$ are precisely the common eigenvalues

$$
\lambda_{n} \in \sigma\left(q^{*}, \alpha_{n}, \beta\right) \cap \sigma(q, \alpha, \beta)
$$

of $q^{*}$ and $q$.
b) If $\lambda_{n} \nrightarrow-\infty$ holds for the (infinitely many) common eigenvalues of $q$ and $q^{*}$ then (2.14) holds.

Proof. The verification of Lemma 2.2 can be repeated, only (2.17) is replaced by

$$
\begin{align*}
F\left(z^{2}\right)= & \frac{\sin ^{2} \beta}{2} \int_{a}^{\pi}\left(q^{*}-q\right)+\frac{\sin ^{2} \beta}{2} \int_{a}^{\pi}\left(q^{*}(x)-q(x)\right) \cos 2 z(\pi-x) d x  \tag{3.1}\\
& +\int_{a}^{\pi}\left(q^{*}(x)-q(x)\right) \int_{0}^{2(\pi-x)} \cos z \tau L(\pi-x, \tau) d \tau d x
\end{align*}
$$

see Lemma 5.3. Consequently

$$
F\left(z^{2}\right) \rightarrow \frac{\sin ^{2} \beta}{2} \int_{a}^{\pi}\left(q^{*}-q\right) \text { if } z \rightarrow+\infty, z \in \mathbf{R}
$$

and the proof of (2.14) is finished as in Lemma 2.2.

Proof of Theorem 1.2. We must show that if the system

$$
\begin{equation*}
C_{0}(\Lambda)=\left\{\cos 2 \sqrt{\lambda_{n}} x: n \geq 1\right\} \tag{3.2}
\end{equation*}
$$

is closed in $L_{p}(0, \pi-a)$ then $\left.q\right|_{(0, a)}$ and the eigenvalues $\lambda_{n}$ determine $q$. Indeed, let $q^{*} \in L_{p}(0, \pi)$ be another potential with $q^{*}=q$ a.e. on $(0, a)$ such that $\lambda_{n} \in \sigma\left(q^{*}, \alpha_{n}, \beta\right) \cap \sigma(q, \alpha, \beta), n \geq 1$ for some $\alpha_{n} \in \mathbf{R}$. From Lemma 5.3 we infer for $h \in L_{p}(a, \pi)$

$$
\begin{equation*}
\int_{a}^{\pi} h(x)\left[v\left(x, z^{2}\right) v^{*}\left(x, z^{2}\right)-1 / 2 \sin ^{2} \beta\right] d x=\int_{a}^{\pi} \cos 2 z(\pi-x) A_{q^{*}} h(x) d x \tag{3.3}
\end{equation*}
$$

where

$$
\begin{equation*}
A_{q^{*}} h(x)=\frac{\sin ^{2} \beta}{2} h(x)+\int_{a}^{x} h(\tau) 2 L\left(\pi-\tau, 2(\pi-x), q, q^{*}\right) d \tau \tag{3.4}
\end{equation*}
$$

We observe that

$$
\begin{equation*}
A_{q^{*}}: L_{p}(a, \pi) \rightarrow L_{p}(a, \pi) \text { is an isomorphism } \tag{3.5}
\end{equation*}
$$

just as in the proof of Theorem 1.1. Let $F(z)$ be defined by (2.13), (2.16). It follows from (2.14) that

$$
\begin{equation*}
F\left(z^{2}\right)=\int_{a}^{\pi} \cos 2 z(\pi-x)\left[A_{q^{*}}\left(q^{*}-q\right)\right](x) d x \tag{3.6}
\end{equation*}
$$

Now if $q^{*} \neq q$ then $0 \neq h=A_{q^{*}}\left(q^{*}-q\right) \in L_{p}$ satisfies

$$
\int_{a}^{\pi} h(x) \cos 2 \sqrt{\lambda_{n}}(\pi-x) d x=0 \quad \forall n
$$

in contradiction to the closedness of $C_{0}(\Lambda)$ in $L_{p}(0, \pi-a)$.
The following statement is the counterpart of Lemma 2.1:
Lemma 3.2. Let $B_{1}$ and $B_{2}$ be Banach spaces, let $\varphi: B_{2} \rightarrow \mathbf{C}$ be a bounded linear functional and let $B_{21}$ be a closed subspace of $B_{2}$. For every $q \in B_{1}$ define a continuous linear operator

$$
A_{q}: B_{1} \rightarrow B_{2}
$$

Suppose (2.1), (2.2) and

$$
\begin{equation*}
\operatorname{dim} B_{21} \geq 2, \quad B_{21} \not \subset \operatorname{Ker} \varphi . \tag{3.7}
\end{equation*}
$$

Then the set $\left\{A_{q}\left(q-q_{0}\right): q \in B_{1}, q-q_{0} \in A_{q_{0}}^{-1}(\operatorname{Ker} \varphi)\right\}$ contains a nonzero element of $B_{21}$.

Proof. Take an element $0 \neq G_{0} \in B_{21} \cap \operatorname{Ker} \varphi$ and let $G_{00} \in B_{21} \backslash \operatorname{Ker} \varphi$ with $\varphi\left(G_{00}\right)=1$. Define the operator $P: B_{2} \rightarrow B_{2}$ by

$$
\begin{equation*}
P G=G-\varphi(G) G_{00} \tag{3.8}
\end{equation*}
$$

By definition we have

$$
\begin{equation*}
\operatorname{ImP} \subset \operatorname{Ker} \varphi \tag{3.9}
\end{equation*}
$$

The vector $q_{0}^{*}$ is defined by

$$
\begin{equation*}
A_{q_{0}}\left(q_{0}^{*}-q_{0}\right)=G_{0} \tag{3.10}
\end{equation*}
$$

and $q_{k+1}^{*}$ by

$$
\begin{equation*}
A_{q_{0}}\left(q_{k+1}^{*}-q_{0}\right)=G_{0}-P\left(\left(A_{q_{k}^{*}}-A_{q_{0}}\right)\left(q_{k}^{*}-q_{0}\right)\right), \quad k \geq 0 . \tag{3.11}
\end{equation*}
$$

Then we have for $k \geq 1$

$$
A_{q_{0}}\left(q_{k+1}^{*}-q_{k}^{*}\right)=-P\left[\left(A_{q_{k}^{*}}-A_{q_{0}}\right)\left(q_{k}^{*}-q_{k-1}^{*}\right)-\left(A_{q_{k}^{*}}-A_{q_{k-1}^{*}}\right)\left(q_{k-1}^{*}-q_{0}\right)\right] ;
$$

if $k=0$, we use instead

$$
A_{q_{0}}\left(q_{1}^{*}-q_{0}^{*}\right)=-P\left(\left(A_{q_{0}^{*}}-A_{q_{0}}\right)\left(q_{0}^{*}-q_{0}\right)\right) .
$$

These correspond to the formulae (2.6), (2.6'). Since the operator $P$ is bounded, the same estimation procedure can be executed (as in Lemma 2.1). So (2.11) holds and then $q_{k}^{*} \rightarrow q^{*} \in B_{1}$. Taking the limit in (3.11) we can verify again as in Lemma 2.2 that

$$
\begin{equation*}
A_{q_{0}}\left(q^{*}-q_{0}\right)=G_{0}-P\left(\left(A_{q^{*}}-A_{q_{0}}\right)\left(q^{*}-q_{0}\right)\right) ; \tag{3.12}
\end{equation*}
$$

i.e.,

$$
\begin{equation*}
A_{q^{*}}\left(q^{*}-q_{0}\right)=G_{0}+(I-P)\left(\left(A_{q^{*}}-A_{q_{0}}\right)\left(q^{*}-q_{0}\right)\right)=G_{0}+c G_{00} \tag{3.13}
\end{equation*}
$$

with some constant $c$. This shows that $0 \neq A_{q^{*}}\left(q^{*}-q_{0}\right) \in B_{21}$. From (3.12) and (3.9) we finally get $q^{*}-q_{0} \in A_{q_{0}}^{-1}(\operatorname{Ker} \varphi)$. Lemma 3.2 is proved.

Proof of Theorem 1.4. Let $q \in L_{p}(0, \pi)$; our task is to find a different $q^{*} \in L_{p}(0, \pi), q^{*}=q$ on $(0, a)$ such that the $\lambda_{n}$ are common eigenvalues of $q^{*}$ and $q$. This will be done by applying Lemma 3.2 with $B_{1}=B_{2}=L_{p}(a, \pi)$,

$$
\varphi: L_{p}(a, \pi) \rightarrow \mathbf{C}, \quad \varphi(h)=\int_{a}^{\pi} A_{q}^{-1} h
$$

and

$$
B_{21}=\left\{h \in L_{p}(a, \pi): \int_{a}^{\pi} h(x) \cos 2 \sqrt{\lambda_{n}}(\pi-x) d x=0 \quad \forall n\right\} .
$$

Now condition (2.1) is given in (3.5), (2.2) follows from Lemma 5.3. In order to check $\operatorname{dim} B_{21} \geq 2$ recall the following identity (see Young [21, Ch. III]):

Let $\alpha(t)$ belong to $L_{p}(-d, d)$ and suppose that

$$
f(z)=\int_{-d}^{d} \alpha(t) e^{i z t} d t \quad \text { satisfies } f(\mu)=0
$$

Then for every $\lambda \neq \mu, \lambda \in \mathbf{C}$ there exists $\beta(t) \in L_{p}(-d, d)$ with

$$
\frac{z-\lambda}{z-\mu} f(z)=\int_{-d}^{d} \beta(t) e^{i z t} d t
$$

namely,

$$
\beta(t)=\alpha(t)+i(\lambda-\mu) e^{-i \mu t} \int_{-d}^{t} \alpha(s) e^{i \mu s} d s
$$

This can be verified by direct substitution. A repeated application of this idea gives that if $f( \pm \mu)=0\left(\right.$ or $f(0)=f^{\prime}(0)=0$ for $\left.\mu=0\right)$, then for every $\lambda \neq \pm \mu$ there exists $\gamma(t) \in L_{p}(-d, d)$ with

$$
\frac{z^{2}-\lambda^{2}}{z^{2}-\mu^{2}} f(z)=\int_{-d}^{d} \gamma(t) e^{i z t} d t .
$$

Supposing that $\alpha(t)$ is even, $\alpha(-t)=\alpha(t)$, we see that $f(z)$ and thus $\gamma(t)$ is even. In other words, $f(z)=\int_{0}^{d} 2 \alpha(t) \cos z t d t, f\left(\mu^{2}\right)=0$ implies

$$
\frac{z^{2}-\lambda^{2}}{z^{2}-\mu^{2}} f(z)=\int_{0}^{d} 2 \gamma(t) \cos z t d t
$$

Since $C(\Lambda)$ is not closed in $L_{p}(0, \pi-a)$, there exists $0 \neq h \in L_{p}(0, \pi-a)$ with

$$
f(z)=\int_{0}^{\pi-a} h(t) \cos z t d t, \quad f\left(2 \sqrt{\lambda_{n}}\right)=0 \forall n, \quad f(2 \mu)=0 .
$$

Take any number $\mu_{1} \neq \pm \mu, \mu_{1} \neq \pm \sqrt{\lambda_{n}}$, then

$$
\frac{z^{2}-4 \mu_{1}^{2}}{z^{2}-4 \mu^{2}} f(z)=\int_{0}^{\pi-a} h_{1}(t) \cos z t d t \text { for some } h_{1} \in L_{p}(0, \pi-a)
$$

Consequently $h(\pi-t)$ and $h_{1}(\pi-t)$ are linearly independent elements of $B_{21}$; thus $\operatorname{dim} B_{21} \geq 2$ as asserted. Finally the minimality condition (M) implies by (3.3) that there exists a function $h \in L_{p}(a, \pi)$ satisfying

$$
\int_{a}^{\pi} A_{q} h(x) \cos 2 \sqrt{\lambda_{n}}(\pi-x) d x=0 \quad \forall n \text { but } \int_{a}^{\pi} h \neq 0 .
$$

Let $h_{1}=A_{q} h$, then $h_{1} \in B_{21} \backslash \operatorname{Ker} \varphi$ showing that $B_{21} \not \subset \operatorname{Ker} \varphi$. Thus all conditions formulated in Lemma 3.2 are fulfilled, so there exists $q^{*} \neq q, q^{*} \in$ $L_{p}(a, \pi)$ such that

$$
\begin{equation*}
A_{q^{*}}\left(q^{*}-q\right) \in B_{21} \text { and } A_{q}\left(q^{*}-q\right) \in \operatorname{Ker} \varphi \text { i.e. } \int_{a}^{\pi}\left(q^{*}-q\right)=0 \tag{3.14}
\end{equation*}
$$

Define $F(z)$ corresponding to $q^{*}$ and $q$. Putting together the formulae (2.16), (3.3) and (3.14) gives

$$
\begin{aligned}
F\left(z^{2}\right) & =\int_{a}^{\pi}\left(q^{*}(x)-q(x)\right)\left[v\left(x, z^{2}\right) v^{*}\left(x, z^{2}\right)-1 / 2 \sin ^{2} \beta\right] d x \\
& =\int_{a}^{\pi} \cos 2 z(\pi-x)\left[A_{q^{*}}\left(q^{*}-q\right)\right](x) d x
\end{aligned}
$$

and then $F\left(\lambda_{n}\right)=0$; i.e., the $\lambda_{n}$ are common eigenvalues of $q^{*}$ and $q$. The proof of Theorem 1.4 is complete.

Proofs of Theorems 1.5 and 1.6. To make explicit the dependence on the parameter $\beta$ we denote by $v(x, \lambda, \beta)$ the solution of (1.9), (1.10). Let

$$
F(x, z)=v^{\prime}\left(x, z, \beta_{1}\right) v^{*}\left(x, z, \beta_{2}\right)-v\left(x, z, \beta_{1}\right) v^{* \prime}\left(x, z, \beta_{2}\right) .
$$

We have $F\left(0, \lambda_{n}\right)=0$ by (1.14). The condition (1.16) means that $q^{*}=q$ a.e. on $(0, a)$ and then

$$
F\left(\lambda_{n}\right)=0 \text { if } F(z)=F(a, z) .
$$

If the values $\lambda_{n}$ have a finite accumulation point then $F(0, z) \equiv 0$ and $m^{*}=m$ follows. In this case $e_{0}(\Lambda)$ is also closed in $L_{p}(a-\pi, \pi-a)$. Indeed, if $G\left(\sqrt{\lambda_{n}}\right)$ $=0$ with $G(z)=\int_{0}^{\pi-a} h(x) \cos 2 z x d x$ where $h \in L_{p}(0, \pi-a)$ then $G \equiv 0$ and $h=0$. So in what follows we can suppose that $\lambda_{n_{k}} \rightarrow \infty$ for a subsequence. As in Lemma 2.2 we can verify that

$$
\begin{equation*}
F(z)=\int_{a}^{\pi}\left(q^{*}(x)-q(x)\right) v\left(x, z, \beta_{1}\right) v^{*}\left(x, z, \beta_{2}\right) d x+F(\pi, z) \tag{3.15}
\end{equation*}
$$

where

$$
F(\pi, z)=-\cos \beta_{1} \sin \beta_{2}+\cos \beta_{2} \sin \beta_{1}=\sin \left(\beta_{1}-\beta_{2}\right)
$$

Suppose first that

$$
\begin{equation*}
\sin \beta_{1} \cdot \sin \beta_{2} \neq 0 \tag{3.16}
\end{equation*}
$$

Analogously as in (3.3) we can check by Lemma 5.3' (last section) that

$$
\begin{align*}
\int_{a}^{\pi} h(x)\left[v\left(x, z^{2}, \beta_{1}\right) v^{*}\left(x, z^{2}, \beta_{2}\right)-1 / 2\right. & \left.\sin \beta_{1} \sin \beta_{2}\right] d x  \tag{3.17}\\
& =\int_{a}^{\pi} \cos 2 z(\pi-x) B_{q^{*}} h(x) d x
\end{align*}
$$

where

$$
\begin{align*}
B_{q^{*}} h(x)= & \frac{\sin \beta_{1} \sin \beta_{2}}{2} h(x)  \tag{3.18}\\
& +\int_{a}^{x} h(\tau) 2 L\left(\pi-\tau, 2(\pi-x), q, q^{*}, \beta_{1}, \beta_{2}\right) d \tau
\end{align*}
$$

Consequently

$$
\begin{aligned}
F\left(z^{2}\right)= & \sin \left(\beta_{1}-\beta_{2}\right)+1 / 2 \sin \beta_{1} \sin \beta_{2} \int_{a}^{\pi}\left(q^{*}-q\right)+ \\
& +\int_{a}^{\pi} \cos 2 z(\pi-x) B_{q^{*}}\left(q^{*}-q\right)(x) d x
\end{aligned}
$$

From $\lambda_{n_{k}} \rightarrow+\infty, F\left(\lambda_{n_{k}}\right)=0$, it follows that

$$
\sin \left(\beta_{1}-\beta_{2}\right)+1 / 2 \sin \beta_{1} \sin \beta_{2} \int_{a}^{\pi}\left(q^{*}-q\right)=0
$$

and then

$$
0=F\left(\lambda_{n}\right)=\int_{a}^{\pi} \cos 2 \sqrt{\lambda_{n}}(\pi-x) B_{q^{*}}\left(q^{*}-q\right)(x) d x \quad \forall n .
$$

Since $C_{0}(\Lambda)$ is closed in $L_{p}(0, \pi-a)$, we infer $B_{q^{*}}\left(q^{*}-q\right)=0$; i.e., $F \equiv 0$, and hence $m_{\beta_{1}} \equiv m_{\beta_{2}}^{*}$.

Now consider the case

$$
\begin{equation*}
\sin \beta_{1} \cdot \sin \beta_{2}=0 \tag{3.19}
\end{equation*}
$$

We see from (5.14) that for $\sin \beta=0$

$$
v\left(\pi-x, z^{2}, \beta\right)=\mathbf{O}\left(\frac{1}{|z|} e^{|\Im z| x}\right)
$$

uniformly in $z \in \mathbf{C}, z \neq 0$ and $0 \leq x \leq \pi-a$. Analogously from (5.25) we get for $\sin \beta \neq 0$

$$
v\left(\pi-x, z^{2}, \beta\right)=\mathbf{O}\left(e^{|\Im z| x}\right)
$$

This implies by (3.15) that

$$
F\left(z^{2}\right)=\sin \left(\beta_{1}-\beta_{2}\right)+\mathbf{O}\left(\frac{1}{|z|} e^{|\Im z| x}\right) .
$$

Now from $F\left(\lambda_{n_{k}}\right)=0, \lambda_{n_{k}} \rightarrow+\infty$, it follows that $\sin \left(\beta_{1}-\beta_{2}\right)=0$ and then $\sin \beta_{1}=\sin \beta_{2}=0$. So (1.14) has the form

$$
m_{0}\left(\lambda_{n}\right)=m_{0}^{*}\left(\lambda_{n}\right) ;
$$

in other words the $\lambda_{n}$ are common eigenvalues of $q^{*}$ and $q$. In this case $m_{0}^{*} \equiv m_{0}$ (i.e., $q^{*}=q$ ) follows if and only if $e(\Lambda)$ is closed in $L_{p}(a-\pi, \pi-a)$; this is stated in Theorem 1.1. The proofs of Theorems 1.5 and 1.6 are complete.

## 4. Applications

This section is devoted to demonstrate how the formerly known theorems listed in Section 1 can be deduced from the new results. At the end of this section we provide the proof of Proposition 1.3.

Consider an arbitrary sequence $\left\{\mu_{n}: n \geq 1\right\}$ of different complex numbers satisfying

$$
\begin{equation*}
\left|\mu_{n}\right| \rightarrow \infty \tag{4.1}
\end{equation*}
$$

Define the counting function

$$
\begin{equation*}
n(r)=\sum_{\left|\mu_{n}\right| \leq r} 1 \quad \text { for } \quad r>0 \tag{4.2}
\end{equation*}
$$

and the function

$$
\begin{equation*}
N(r)=\int_{1}^{r} \frac{n(t)}{t} d t \tag{4.3}
\end{equation*}
$$

Recall the following classical closedness test of Levinson:
Theorem 4.1 ([15], see also Young [21]). Let $0 \leq a<\pi, 1 \leq p<\infty$, $1 / p+1 / p^{\prime}=1$. If

$$
\begin{equation*}
\limsup _{r \rightarrow \infty}\left(N(r)-2(1-a / \pi) r+1 / p^{\prime} \ln r\right)>-\infty \tag{4.4}
\end{equation*}
$$

then the system $\left\{e^{i \mu_{n} x}: n \geq 1\right\}$ is closed in $L_{p}(a-\pi, \pi-a)$.
Remark. The original form of Theorem 4.1 in [21] refers to $1<p$ and to the case $a=0$ because it is formulated as a completeness test in $L_{p^{\prime}}$ and this is equivalent to closedness in $L_{p}$ only if $p>1$. However the proof given in [21] works also for $p=1$ and it can be transformed into the form (4.4).

Remark. We can easily extend Theorem 4.1 to those cases where there are values $\mu_{n}$ taken with multiplicities $1<m_{n}<\infty$; in this case (4.2) is accordingly modified and the exponential system contains the members $e^{i \mu_{n} x}, x e^{i \mu_{n} x}, \ldots, x^{m_{n}-1} e^{i \mu_{n} x}$.

In applying Theorem 4.1 we need the following estimates for the $N$ function corresponding to a complete spectrum.

Lemma 4.2. Denote by $N_{\sigma}$ the $N$-function for the set $\left\{ \pm 2 \sqrt{\lambda_{n}}: \lambda_{n} \in\right.$ $\sigma(q, \alpha, \beta)\}$; if $\lambda_{n}=0$ then the value 0 has multiplicity 2 in this set. Then, as $r \rightarrow+\infty$,

$$
\begin{equation*}
\sin \alpha=\sin \beta=0 \Rightarrow N_{\sigma}(r)=r-\ln r+\mathbf{O}(1) \tag{4.5}
\end{equation*}
$$

$$
\begin{align*}
\sin \alpha=0 & \neq \sin \beta \text { or } \sin \alpha \neq 0=\sin \beta  \tag{4.6}\\
& \Rightarrow N_{\sigma}(r)=r+\mathbf{O}(1),
\end{align*}
$$

$$
\begin{equation*}
\sin \alpha \neq 0, \sin \beta \neq 0 \Rightarrow N_{\sigma}(r)=r+\ln r+\mathbf{O}(1) \tag{4.7}
\end{equation*}
$$

Proof. Consider first the Dirichlet case (4.5). Recall the well-known eigenvalue asymptotics

$$
\begin{equation*}
\sqrt{\lambda_{n}}=n+\mathbf{O}\left(\frac{1}{n}\right), \quad n \geq 1 \tag{4.8}
\end{equation*}
$$

(see, e.g., [17]). Let $n^{(1)}$ and $N^{(1)}$ be the corresponding functions if we substitute the values $\pm 2 \sqrt{\lambda_{n}}$ by $\pm 2 n$. From (4.8) it follows that

$$
\begin{equation*}
N_{\sigma}(r)-N^{(1)}(r)=\mathbf{O}(1) . \tag{4.9}
\end{equation*}
$$

We can count $N^{(1)}(r)$ for $2 k \leq r \leq 2 k+2$ as follows

$$
\begin{align*}
N^{(1)}(r)= & \int_{1}^{r} \frac{n^{(1)}(t)}{t} d t=\sum_{i=2}^{k}(2 i-2)[\ln (2 i)-\ln (2 i-2)]  \tag{4.10}\\
& +2 k[\ln r-\ln (2 k)]=2 k \ln r-2 \sum_{i=1}^{k} \ln (2 i) \\
= & 2 k \ln r-2 k \ln 2-2 \ln (k!)=2 k(\ln r-\ln 2)-2 k(\ln k-1) \\
& -\ln k+\mathbf{O}(1)=2 k \ln \left(\frac{r}{2 k}\right)+2 k-\ln k+\mathbf{O}(1) \\
= & 2 k-\ln k+\mathbf{O}(1)=r-\ln r+\mathbf{O}(1) .
\end{align*}
$$

From (4.9) we get $N_{\sigma}(r)=r-\ln r+\mathbf{O}(1)$ as asserted.
In the second case (4.6) we start with the asymptotics

$$
\begin{equation*}
\sqrt{\lambda_{n}}=n+1 / 2+\mathbf{O}\left(\frac{1}{n+1}\right), \quad n \geq 0 . \tag{4.11}
\end{equation*}
$$

Define the functions $n^{(2)}$ and $N^{(2)}$ by the main term of (4.11). Taking into account $n^{(2)}(t)=n^{(1)}(t+1), n^{(1)}(t)=t+\mathbf{O}(1)$ we obtain

$$
\begin{aligned}
N^{(2)}(r) & =\int_{1}^{r} \frac{n^{(1)}(t+1)}{t} d t=\int_{1}^{r} \frac{n^{(1)}(t+1)}{t+1} d t+\int_{1}^{r} \frac{n^{(1)}(t+1)}{t(t+1)} d t \\
& =N^{(1)}(r+1)+\int_{1}^{r} \frac{t+\mathbf{O}(1)}{t(t+1)} d t+\mathbf{O}(1)=r+\mathbf{O}(1)
\end{aligned}
$$

With $N_{\sigma}(r)=N^{(2)}(r)+\mathbf{O}(1)$ this implies (4.6). Finally in case (4.7) we argue
similarly starting from

$$
\begin{equation*}
\sqrt{\lambda_{n}}=n+\mathbf{O}\left(\frac{1}{n+1}\right), \quad n \geq 0 \tag{4.12}
\end{equation*}
$$

and using the fact that $n^{(3)}(t)=n^{(1)}(t+2)$. The proof is complete.
Remark. Let $d>0$ and denote by $N_{d}$ the $N$-function corresponding to the set $\{ \pm 2 n d: n \geq 1\}$. Then we have

$$
\begin{equation*}
N_{d}(r)=\frac{r}{d}-\ln r+\mathbf{O}(1) \tag{4.13}
\end{equation*}
$$

Indeed, $n_{d}(t)=n^{(1)}(t / d)$ implies

$$
\begin{aligned}
N_{d}(r) & =\int_{1}^{r} \frac{n^{(1)}\left(\frac{t}{d}\right)}{t} d t=\int_{\frac{1}{d}}^{\frac{r}{d}} \frac{n^{(1)}(t)}{t} d t \\
& =N\left(\frac{r}{d}\right)+\mathbf{O}(1)=\frac{r}{d}-\ln r+\mathbf{O}(1)
\end{aligned}
$$

Proof of Theorem C. Note first that $\sigma\left(q, \alpha_{1}, \beta\right) \cap \sigma\left(q, \alpha_{2}, \beta\right)=\emptyset$. If $\sin \beta$ $=0$, we have two spectra of type (4.6) or one of type (4.6) and the other of type (4.5). Thus

$$
N_{\Lambda_{0}}(r) \geq 2 r-\ln r+\mathbf{O}(1)
$$

if $N_{\Lambda_{0}}$ is the $N$-function corresponding to the values $\pm 2 \sqrt{\lambda}$ where $\lambda$ runs over the eigenvalues from the two spectra. Adjoining the pair $\pm 2 \mu$ we get

$$
N_{\Lambda}(r) \geq 2 r+\ln r+\mathbf{O}(1)
$$

By Theorem 4.1 this implies that $e(\Lambda)$ is closed in $L_{1}(-\pi, \pi)$ and then the potential is determined by the two spectra. Now, if $\sin \beta \neq 0$, we have two spectra of type (4.7) or one of type (4.7) and one of type (4.6). Hence

$$
N_{\Lambda_{0}}(r) \geq 2 r+\ln r+\mathbf{O}(1)
$$

so that $e_{0}(\Lambda)$ is closed in $L_{1}(-\pi, \pi)$; thus the potential is again determined.
Proof of Theorem D. This follows from the estimates:

$$
\sin \beta \neq 0 \Rightarrow N_{\Lambda_{0}}=N_{\sigma} \geq r+\mathbf{O}(1),
$$

$$
\sin \beta=0 \Rightarrow N_{\Lambda}=N_{\sigma}+2 \ln r+\mathbf{O}(1) \geq r+\ln r+\mathbf{O}(1)
$$

Proof of Theorem E. Denote by $n_{S}(t)$ the $n$-function corresponding to the set $\left\{ \pm 2 \sqrt{\lambda_{n}}: \lambda_{n} \in S\right\}$. Then for large $t$

$$
\begin{aligned}
n_{S}(t)= & 2 \#\left\{\lambda_{n} \in S: \lambda_{n} \leq t^{2} / 4\right\} \geq 4(1-a / \pi) \#\left\{\lambda_{n} \in \sigma: \lambda_{n} \leq t^{2} / 4\right\} \\
& +2 a / \pi-1=2(1-a / \pi) n_{\sigma}(t)+2 a / \pi-1 .
\end{aligned}
$$

Hence

$$
\begin{aligned}
\sin \beta \neq 0 \Rightarrow & N_{\Lambda_{0}} \geq 2(1-a / \pi) r+(2 a / \pi-1) \ln r+\mathbf{O}(1), \\
\sin \beta=0 \Rightarrow & N_{\Lambda} \geq 2(1-a / \pi)(r-\ln r)+2 \ln r \\
& +(2 a / \pi-1) \ln r+\mathbf{O}(1)= \\
= & 2(1-a / \pi) r+2 a / \pi \ln r+(2 a / \pi-1) \ln r+\mathbf{O}(1)
\end{aligned}
$$

which verifies Theorem E even after deleting $a / \pi-1 / 2$ from the density condition on S .

## Proof of Theorem F.

$$
\begin{aligned}
\sin \beta \neq 0 \Rightarrow N_{\Lambda_{0}} & \geq 2 / 3\left(N_{\sigma_{1}}+N_{\sigma_{2}}+N_{\sigma_{3}}\right)+\mathbf{O}(1) \\
& \geq 2 / 3(3 r+2 \ln r)+\mathbf{O}(1), \\
\sin \beta=0 \Rightarrow N_{\Lambda} & \geq 2 / 3\left(N_{\sigma_{1}}+N_{\sigma_{2}}+N_{\sigma_{3}}\right)+2 \ln r+\mathbf{O}(1) \\
& \geq 2 / 3(3 r-\ln r)+2 \ln r+\mathbf{O}(1) .
\end{aligned}
$$

Proof of Theorem H. By Theorem 1.5 it is enough to verify that (1.15) implies the closedness of $e_{0}(\Lambda)$ in $L_{1}(-\pi, \pi)$. Consider the $N$-function for the set $\left\{ \pm 2 \sqrt{\lambda_{n}}: n \geq 0\right\}$; we have to check that

$$
\begin{equation*}
N(r)-2 r \nrightarrow-\infty \quad(r \rightarrow+\infty) \tag{4.14}
\end{equation*}
$$

We shift the values $\lambda_{n}<\frac{n^{2}}{4}$ into $\frac{n^{2}}{4}$. This will diminish $N(r)$ and it is enough to prove (4.14) for the diminished $N$. But we can also shift the values $\lambda_{n}>\frac{n^{2}}{4}$ into $\frac{n^{2}}{4}$ since this will grow $N$ by a bounded quantity. Indeed, the growth is at most

$$
\sum_{n<r} 2 \int_{n}^{2 \sqrt{\lambda_{n}}} \frac{d t}{t}=\sum_{n<r} \ln \frac{4 \lambda_{n}}{n^{2}}=\sum_{n<r} \ln \left(1+\frac{4\left(\lambda_{n}-n^{2} / 4\right)}{n^{2}}\right)=\mathbf{O}(1)
$$

by (1.15). For the shifted system $\left\{\lambda_{n}=n^{2} / 4: n \geq 0\right\}$ we have $N_{0}(r)=$ $2 r+\ln r+\mathbf{O}(1)$; hence $N(r) \geq 2 r+\ln r+\mathbf{O}(1)$ and (4.14) follows.

In order to check Theorem B we need a stability result of Riesz bases. By definition, a Riesz basis is an isomorphic image of an orthonormal basis of a Hilbert space. A famous result of Kadec [14] says that if $\lambda_{n}$ are arbitrary real numbers with $\left|\lambda_{n}-n\right| \leq L<1 / 4$ for all $n \in \mathbf{Z}$ then the system $\left\{e^{i \lambda_{n} x}\right.$ : $n \in \mathbf{Z}\}$ forms a Riesz basis in $L_{2}(-\pi, \pi)$. It has been previously known that the constant $1 / 4$ is best possible here; see e.g., Young [21]. Later on, S. A. Avdonin [2] realized that it is not necessary to impose the bound $L<1 / 4$ for every individual shift $\left|\lambda_{n}-n\right|$; instead, it is enough to take this bound only for the average shifts in the following sense:

Theorem 4.3 ([2]). Suppose that the shifts $\delta_{n} \in \mathbf{C}$ are bounded and the shifted exponents $\lambda_{n}=n+\delta_{n}$ are separated; i.e., $\inf _{n \neq m}\left|\lambda_{n}-\lambda_{m}\right|>0$. If the average Kadec condition

$$
\begin{equation*}
\left.\left.\limsup _{R \rightarrow \infty} \sup _{x \in \mathbf{R}} \frac{1}{R}\right|_{x<n+\Re \delta_{n}<x+R} \delta_{n} \right\rvert\,<\frac{1}{4} \tag{4.15}
\end{equation*}
$$

holds then the shifted system $\left\{e^{i \lambda_{n} x}: n \in \mathbf{Z}\right\}$ forms a Riesz basis in $L_{2}(-\pi, \pi)$.
Proof of Theorem B (in case $\sin \beta=0$ ). The sufficiency of two spectra is proved in Theorem C; we investigate the necessity. For the eigenvalues $\lambda_{n}^{(1)}$ of $\sigma(q, 0,0)$ and $\lambda_{n}^{(2)}$ of $\sigma\left(q, \alpha_{2}, 0\right)$ we have

$$
\begin{equation*}
\sqrt{\lambda_{n}^{(1)}}=n+\mathbf{o}(1)(n \geq 1), \quad \sqrt{\lambda_{n}^{(2)}}=n+1 / 2+\mathbf{o}(1)(n \geq 0) \tag{4.16}
\end{equation*}
$$

So the set of all values $\pm 2 \sqrt{\lambda}$ is an $\mathbf{o}(1)$-perturbation of $\mathbf{Z} \backslash\{0\}$. Since the eigenvalues are different, this means that the shifted exponents are separated and (4.15) holds with limsup $=0$. Consequently $e_{0}(\Lambda)$ is a Riesz basis of codimension 1. Hence $e(\Lambda)$ is complete in $L_{2}$ and after deleting an arbitrary eigenvalue it becomes incomplete (of codimension 1). In other words, after the deletion it is not closed in $L_{2}$, thus it is not closed in $L_{1}$. By Theorem 1.1 this proves Theorem B if $\sin \beta=0$.

Remark. The case $\sin \beta \neq 0$ cannot be dealt with in this general framework. Roughly speaking, we have "half an eigenvalue" deficiency and excess in $e_{0}(\Lambda)$ and $e(\Lambda)$, respectively. This prevents us from applying Theorems 1.2 and 1.4. It would be possible to give ad hoc modifications, based on the special structure of the set of eigenvalues in order to cover this special case; we do not give the details.

Our last topic in this section is the proof of Proposition 1.3. We need the following elementary

Lemma 4.4. In the domain $|w|>1$ the function

$$
f(w)=w \sin \frac{\pi}{2} w+\frac{1}{w} \sin \frac{\pi}{2 w}
$$

has only real zeros.
Proof. Since $f(-w)=f(w)$, we can suppose $\Re w \geq 0$. From the wellknown formula

$$
\begin{equation*}
|\sin (a+i b)|=\sqrt{\sin ^{2} a+\sinh ^{2} b} \tag{4.17}
\end{equation*}
$$

we can easily check that

$$
\begin{equation*}
\left|\sin \frac{\pi}{2} w\right|>\left|\sin \frac{\pi}{2 w}\right| \quad \text { if }|w|>1,0 \leq \Re w \leq 1, \tag{4.18}
\end{equation*}
$$

hence $f(w)$ has no zeros in this domain. Indeed, if $w=x+i y, 0 \leq x \leq 1$, $x^{2}+y^{2}>1$ then $\sin ^{2} \frac{\pi}{2} x \geq \sin ^{2} \frac{\pi}{2} \frac{x}{x^{2}+y^{2}}, \sinh ^{2} \frac{\pi}{2} y \geq \sinh ^{2} \frac{\pi}{2} \frac{y}{x^{2}+y^{2}}$ and equality cannot occur in both cases. Now consider the case $x=1+\varepsilon, \varepsilon>0$ being appropriately small. From

$$
\sin \frac{\pi}{2}(1+\varepsilon)=1-\mathbf{O}\left(\varepsilon^{2}\right)
$$

and (4.17) we get

$$
\left|\sin \frac{\pi}{2} w\right|>\left|\sin \frac{\pi}{2 w}\right|-\mathbf{O}\left(\varepsilon^{2}\right) \quad \text { if } x=1+\varepsilon
$$

Consequently

$$
\begin{equation*}
\left|w \sin \frac{\pi}{2} w\right|>\left|\frac{1}{w} \sin \frac{\pi}{2 w}\right| \quad \text { if } x=1+\varepsilon \tag{4.19}
\end{equation*}
$$

Indeed, this is trivial if $|y|$ is large enough, and for other values $y\left|\sin \frac{\pi}{2 w}\right|$ is not very small, so that

$$
\begin{aligned}
\left|w \sin \frac{\pi}{2} w\right| & >(1+\varepsilon)\left(\left|\sin \frac{\pi}{2 w}\right|-\mathbf{O}\left(\varepsilon^{2}\right)\right) \\
& >\left|\frac{1}{w} \sin \frac{\pi}{2 w}\right|+2 \varepsilon\left|\sin \frac{\pi}{2 w}\right|-\mathbf{O}\left(\varepsilon^{2}\right)>\left|\frac{1}{w} \sin \frac{\pi}{2 w}\right|
\end{aligned}
$$

Hence $f$ has no zeros on the line $x=1+\varepsilon$. We can simply check by (4.17) that

$$
\begin{equation*}
\left|w \sin \frac{\pi}{2} w\right|>\left|\frac{1}{w} \sin \frac{\pi}{2 w}\right| \quad \text { if } x=2 k+1(k=1,2, \ldots) \text { and } y \in \mathbf{R} \tag{4.20}
\end{equation*}
$$

and that for large $R>0$

$$
\begin{equation*}
\left|w \sin \frac{\pi}{2} w\right|>\left|\frac{1}{w} \sin \frac{\pi}{2 w}\right| \quad \text { if }|y| \geq R \tag{4.21}
\end{equation*}
$$

This means by Rouché's theorem that $f(w)$ has exactly one zero in each of the rectangles $[1+\varepsilon, 3] \times[-R, R]$ and $[2 k+1,2 k+3] \times[-R, R]$ for $k \geq 1$ and no other zeros exist. These zeros must be real since $f(\bar{w})=\overline{f(w)}$ and this completes the proof.

Proof of Proposition 1.3. On the interval $[\pi / 2, \pi]$ we have $v(x, z)=$ $\cos (\pi-x) \sqrt{z-1}$. Thus $v\left(\frac{\pi}{2}, z\right)=\cos \frac{\pi}{2} \sqrt{z-1}, v^{\prime}\left(\frac{\pi}{2}, z\right)=\sqrt{z-1} \sin \frac{\pi}{2} \sqrt{z-1}$ and then in $[0, \pi / 2]$

$$
\begin{aligned}
v(x, z)= & \cos \left(\frac{\pi}{2}-x\right) \sqrt{z} \cdot \cos \frac{\pi}{2} \sqrt{z-1} \\
& -\frac{\sin \left(\frac{\pi}{2}-x\right) \sqrt{z}}{\sqrt{z}} \sqrt{z-1} \sin \frac{\pi}{2} \sqrt{z-1} .
\end{aligned}
$$

Finally we get that

$$
\begin{align*}
v(0, z) & =\cos \frac{\pi}{2} \sqrt{z} \cdot \cos \frac{\pi}{2} \sqrt{z-1}-\frac{\sqrt{z-1}}{\sqrt{z}} \sin \frac{\pi}{2} \sqrt{z} \cdot \sin \frac{\pi}{2} \sqrt{z-1}  \tag{4.22}\\
v^{\prime}(0, z) & =\sqrt{z} \sin \frac{\pi}{2} \sqrt{z} \cdot \cos \frac{\pi}{2} \sqrt{z-1}+\sqrt{z-1} \cos \frac{\pi}{2} \sqrt{z} \cdot \sin \frac{\pi}{2} \sqrt{z-1} \tag{4.23}
\end{align*}
$$

Similarly

$$
\begin{gather*}
v^{*}(0, z)=\cos \frac{\pi}{2} \sqrt{z} \cdot \cos \frac{\pi}{2} \sqrt{z-1}-\frac{\sqrt{z}}{\sqrt{z-1}} \sin \frac{\pi}{2} \sqrt{z} \cdot \sin \frac{\pi}{2} \sqrt{z-1},  \tag{4.24}\\
v^{* \prime}(0, z)=v^{\prime}(0, z) \tag{4.25}
\end{gather*}
$$

Consider the function

$$
\begin{align*}
F(z) & =v^{\prime}(0, z) v^{*}(0, z)-v(0, z) v^{* \prime}(0, z)  \tag{4.26}\\
& =v^{\prime}(0, z)\left(v^{*}(0, z)-v(0, z)\right)=-v^{\prime}(0, z) \frac{\sin \frac{\pi}{2} \sqrt{z}}{\sqrt{z}} \cdot \frac{\sin \frac{\pi}{2} \sqrt{z-1}}{\sqrt{z-1}} .
\end{align*}
$$

Its real zeros are precisely the common eigenvalues of $q$ and $q^{*}$. In order to find the zeros of $v^{\prime}(0, z)$, consider the decomposition

$$
\begin{align*}
v^{\prime}(0, z)= & \sqrt{z} \sin \frac{\pi}{2}(\sqrt{z}+\sqrt{z-1})  \tag{4.27}\\
& -(\sqrt{z}-\sqrt{z-1}) \cos \frac{\pi}{2} \sqrt{z} \sin \frac{\pi}{2} \sqrt{z-1} \\
= & \sqrt{z} \sin \pi \sqrt{z}+\left[\sqrt{z}\left(\sin \pi \sqrt{z}-\sin \frac{\pi}{2}(\sqrt{z}+\sqrt{z-1})\right)\right. \\
& \left.-(\sqrt{z}-\sqrt{z-1}) \cos \frac{\pi}{2} \sqrt{z} \sin \frac{\pi}{2} \sqrt{z-1}\right] \\
\stackrel{\text { def }}{=} & g(z)+[h(z)] .
\end{align*}
$$

From

$$
\sin \pi \sqrt{z}-\sin \frac{\pi}{2}(\sqrt{z}+\sqrt{z-1})=2 \sin \frac{\pi}{2} \frac{\sqrt{z}-\sqrt{z-1}}{2} \cdot \cos \frac{\pi}{2} \frac{3 \sqrt{z}+\sqrt{z-1}}{2}
$$

we infer

$$
\begin{aligned}
\sqrt{z}\left(\sin \pi \sqrt{z}-\sin \frac{\pi}{2}(\sqrt{z}+\right. & \sqrt{z-1})) \\
& =\mathbf{O}\left(\sqrt{z} \frac{1}{\sqrt{z}} e^{\frac{\pi}{4}(3|\Im \sqrt{z}|+|\Im \sqrt{z-1}|)}\right)=\mathbf{O}\left(e^{\pi|\Im \sqrt{z}|}\right)
\end{aligned}
$$

Then

$$
\begin{equation*}
h(z)=\mathbf{O}\left(e^{\pi|\Im \sqrt{z}|}\right), \quad|z| \rightarrow \infty . \tag{4.28}
\end{equation*}
$$

The (simple) zeros of the function $g(z)$ are $z=k^{2}, k=0,1, \ldots$. We know that

$$
|g(z)| \geq c \sqrt{|z|} e^{\pi|\Im \sqrt{z}|} \quad \text { if } \quad|z|=(N+1 / 2)^{2}, n \in \mathbf{N}
$$

with $c>0$ independent of $z$ and $N$. Comparing this estimate with (4.28) we get from Rouché's theorem that $v^{\prime}(0, z)$ has precisely $N+1$ zeros in the disk $|z|<(N+1 / 2)^{2}$ and (again by Rouché's theorem) that the zeros satisfy

$$
\begin{equation*}
\sqrt{\lambda_{n}^{(1)}}=n+\mathbf{O}\left(\frac{1}{n+1}\right) \quad(n \geq 0, n \rightarrow \infty) . \tag{4.29}
\end{equation*}
$$

We have to check that these zeros are real. Apply trigonometric identities to obtain

$$
\begin{aligned}
v^{\prime}(0, z)= & \sqrt{z} \frac{\sin \frac{\pi}{2}(\sqrt{z}+\sqrt{z-1})+\sin \frac{\pi}{2}(\sqrt{z}-\sqrt{z-1})}{2} \\
& +\sqrt{z-1} \frac{\sin \frac{\pi}{2}(\sqrt{z}+\sqrt{z-1})-\sin \frac{\pi}{2}(\sqrt{z}-\sqrt{z-1})}{2} \\
= & \frac{\sqrt{z}+\sqrt{z-1}}{2} \sin \frac{\pi}{2}(\sqrt{z}+\sqrt{z-1}) \\
& +\frac{\sqrt{z}-\sqrt{z-1}}{2} \sin \frac{\pi}{2}(\sqrt{z}-\sqrt{z-1}) \\
= & 1 / 2 f(w), \quad w=\sqrt{z}+\sqrt{z-1} .
\end{aligned}
$$

Here we used $(\sqrt{z}+\sqrt{z-1})(\sqrt{z}-\sqrt{z-1})=1$. By appropriately defining the square roots we can suppose $|w|=|\sqrt{z}+\sqrt{z-1}| \geq 1$. If $|w|=1$, then $\overline{\sqrt{z}+\sqrt{z-1}}=\bar{w}=\frac{1}{w}=\sqrt{z}-\sqrt{z-1}$, hence $w+\bar{w}=2 \sqrt{z}$ is real, $w-\bar{w}=$ $2 \sqrt{z-1}$ is purely imaginary. This means that $0 \leq z \leq 1$. If $|w|>1$ and $f(w)=0$ then by Lemma 4.4 the root $w=\sqrt{z}+\sqrt{z-1}$ is real. Then $\frac{1}{w}=\sqrt{z}-\sqrt{z-1}$ and hence $\sqrt{z}$ and $\sqrt{z-1}$ are also real; i.e., $z \geq 1$. This shows indeed that $v^{\prime}(0, z)$ has only real zeros. The other two factors in (4.26) have the zeros $\lambda_{n}^{(2)}, \lambda_{n}^{(3)}$ satisfying

$$
\begin{align*}
& \sqrt{\lambda_{n}^{(2)}}=2 n \quad(n \geq 1)  \tag{4.30}\\
& \sqrt{\lambda_{n}^{(3)}}=\sqrt{4 n^{2}+1}=2 n+\mathbf{O}\left(\frac{1}{n}\right) \quad(n \geq 1) \tag{4.31}
\end{align*}
$$

So for the $N$-function of the sets of values $\pm 2 \sqrt{\lambda_{n}}$ we get by (4.13) that

$$
N^{(1)}(r)=r+\ln r+\mathbf{O}(1), \quad N^{(2)}(r), N^{(3)}(r)=\frac{r}{2}-\ln r+\mathbf{O}(1) .
$$

If $\mu \neq \pm \sqrt{\lambda_{n}^{(j)}}$, then the $N$-function of all values $\pm 2 \sqrt{\lambda_{n}^{(j)}}$ and $2 \mu$ satisfies

$$
N(r)=2 r+\mathbf{O}(1)
$$

which means by the Levinson test that the system

$$
e_{1}(\Lambda)=\left\{e^{2 i \mu x}, e^{ \pm 2 i \sqrt{\lambda} x}: F(\lambda)=0\right\}
$$

is closed in $L_{p}(-\pi, \pi)$. On the other hand $e_{0}(\Lambda)$ cannot be closed by Theorem 1.2, so it has deficiency 1 as asserted.

## 5. Technical background

In this last part of the paper we give the auxiliary results used in the above proofs. More precisely we provide integral representations for the products of eigenfunctions and a connection between the closedness of cosine and exponential systems. The first result is a refinement of the known representation (5.1) below; see, e.g., Marchenko [19].

Lemma 5.1. Let $1 \leq p \leq \infty, 0<d<\infty, q \in L_{p}(-d, d)$ and consider the solution $e(x, \lambda)$ of the initial value problem

$$
-y^{\prime \prime}+q y=\lambda^{2} y \text { on }(-d, d), \quad y(0)=1, \quad y^{\prime}(0)=i \lambda .
$$

It has a representation of the form

$$
\begin{equation*}
e(x, \lambda)=e^{i \lambda x}+\int_{-x}^{x} K_{1}(x, t) e^{i \lambda t} d t \tag{5.1}
\end{equation*}
$$

with a continuous kernel $K_{1}(x, t)$. If there exist two potentials $q^{*}, q \in L_{p}(-d, d)$ with norm $\leq D$ then

$$
\begin{gather*}
\left|K_{1}(x, t, q)\right| \leq c(D),  \tag{5.2}\\
\left|K_{1}\left(x, t, q^{*}\right)-K_{1}(x, t, q)\right| \leq c(D)\left\|q^{*}-q\right\|_{p} \tag{5.3}
\end{gather*}
$$

with a constant $c(D)=c(D, p, d)$ independent of $q, q^{*}, x$ and $t$.
Proof. Define $H(\alpha, \beta)=K_{1}(\alpha+\beta, \alpha-\beta)$ for $\alpha, \beta \geq 0$. Introduce the notation

$$
\sigma(u)=\int_{0}^{u}|q|, \quad \varrho(u, v)=\int_{0}^{u} \int_{0}^{v}|q(\alpha+\beta)| d \beta d \alpha .
$$

It is shown in Marchenko [19] that

$$
\begin{equation*}
H(u, v)=1 / 2 \int_{0}^{u} q+\int_{0}^{u} \int_{0}^{v} q(\alpha+\beta) H(\alpha, \beta) d \beta d \alpha \tag{5.4}
\end{equation*}
$$

and

$$
\begin{equation*}
|H(u, v)| \leq 1 / 2 \sigma(u) e^{\varrho(u, v)} \tag{5.5}
\end{equation*}
$$

From $\sigma(u) \leq c(D), \varrho(u, v) \leq c(D)$ we get $|H(u, v)| \leq c(D)$ which is (5.2). To show (5.3) consider the decomposition

$$
\begin{align*}
H^{*}(u, v)-H(u, v)= & 1 / 2 \int_{0}^{u}\left(q^{*}-q\right)  \tag{5.6}\\
& +\int_{0}^{u} \int_{0}^{v}\left(q^{*}(\alpha+\beta)-q(\alpha+\beta)\right) H^{*}(\alpha, \beta) d \beta d \alpha \\
& +\int_{0}^{u} \int_{0}^{v} q(\alpha+\beta)\left(H^{*}(\alpha, \beta)-H(\alpha, \beta)\right) d \beta d \alpha .
\end{align*}
$$

This implies

$$
\begin{align*}
& \left|H^{*}(u, v)-H(u, v)\right| \leq c\left\|q^{*}-q\right\|_{p}  \tag{5.7}\\
& \quad+c(D) \int_{0}^{u} \int_{0}^{v}\left|q^{*}(\alpha+\beta)-q(\alpha+\beta)\right| d \beta d \alpha \\
& \quad+\int_{0}^{u} \int_{0}^{v}\left|q(\alpha+\beta) \| H^{*}(\alpha, \beta)-H(\alpha, \beta)\right| d \beta d \alpha \\
& \leq \\
& \quad c(D)\left\|q^{*}-q\right\|_{p}+\int_{0}^{u} \int_{0}^{v}\left|q(\alpha+\beta) \| H^{*}(\alpha, \beta)-H(\alpha, \beta)\right| d \beta d \alpha .
\end{align*}
$$

Recall the following inequality of Wendroff (see, e.g., in [3]): Let $c \geq 0$, $u(s, r) \geq 0, v(s, r) \geq 0, u$ continuous, $v$ locally integrable in the domain $r, s \geq 0$. Now if

$$
\begin{equation*}
u(x, y) \leq c+\int_{0}^{x} \int_{0}^{y} v(r, s) u(r, s) d s d r, \quad x, y \geq 0 \tag{5.8}
\end{equation*}
$$

then

$$
\begin{equation*}
u(x, y) \leq c e^{\int_{0}^{x} \int_{0}^{y} v(r, s) d s d r}, \quad x, y \geq 0 \tag{5.9}
\end{equation*}
$$

Applying this to (5.7) gives

$$
\begin{align*}
\left|H^{*}(u, v)-H(u, v)\right| & \leq c(D)\left\|q^{*}-q\right\|_{p} e^{\int_{0}^{u} \int_{0}^{v}|q(\alpha+\beta)| d \beta d \alpha}  \tag{5.10}\\
& \leq c(D)\left\|q^{*}-q\right\|_{p}
\end{align*}
$$

which is equivalent to (5.3).
Our next topic is an integral representation for $v(x, \lambda) v^{*}(x, \lambda)$ :
Lemma 5.2. Let $\beta=0$ in (1.10) and $\mu \in \mathbf{C}, 1 \leq p, q^{*}, q \in L_{p}(0, \pi)$. Then for $z \in \mathbf{C}$

$$
\begin{align*}
1-2\left(z^{2}-\mu^{2}\right) v\left(\pi-x, z^{2}\right) v^{*}(\pi & \left.-x, z^{2}\right)  \tag{5.11}\\
& =\cos 2 z x+\int_{0}^{2 x} \cos z \tau M\left(x, \tau, \mu^{2}\right) d \tau
\end{align*}
$$

where the kernel function $M\left(x, \tau, \mu^{2}\right)$ is linear in $\mu^{2}$, continuous in $(x, \tau)$ and independent of $z$. Further if $q^{* *} \in L_{p}$ and $\|q\|_{p},\left\|q^{*}\right\|_{p},\left\|q^{* *}\right\|_{p} \leq D$ then

$$
\begin{equation*}
\left|M\left(x, \tau, \mu^{2}, q, q^{* *}\right)-M\left(x, \tau, \mu^{2}, q, q^{*}\right)\right| \leq c(D, \mu, p)\left\|q^{* *}-q^{*}\right\|_{p} \tag{5.12}
\end{equation*}
$$

Proof. It can be checked from (5.1) that there exists a continuous kernel $K(x, t), 0 \leq t \leq x$, satisfying $K(x, 0)=0$, the analogues of (5.2), (5.3) and

$$
\begin{equation*}
v\left(\pi-x, z^{2}\right)=\frac{\sin z x}{z}+\int_{0}^{x} K(x, t) \frac{\sin z t}{z} d t . \tag{5.14}
\end{equation*}
$$

Indeed, define $K_{1}$ for the potential $q(\pi-x)$; then $K(x, t)=K_{1}(x, t)-K_{1}(x,-t)$ satisfies (5.14). Consequently

$$
\begin{align*}
1- & 2 z^{2} v\left(\pi-x, z^{2}\right) v^{*}\left(\pi-x, z^{2}\right)  \tag{5.15}\\
= & 1-2 \sin ^{2} z x-2 \int_{0}^{x} K(x, t) \sin z x \sin z t d t \\
& -2 \int_{0}^{x} K^{*}(x, t) \sin z x \sin z t d t \\
& -2 \int_{0}^{x} \int_{0}^{x} K(x, t) K^{*}(x, u) \sin z t \sin z u d u d t \\
= & \cos 2 z x-\int_{0}^{x} K(x, t)[\cos z(x-t)-\cos z(x+t)] d t \\
& -\int_{0}^{x} K^{*}(x, t)[\cos z(x-t)-\cos z(x+t)] d t \\
& -\int_{0}^{x} \int_{0}^{x} K(x, t) K^{*}(x, u)[\cos z(t-u)-\cos z(t+u)] d u d t \\
= & \cos 2 z x-I_{1}-I_{1}^{*}-I_{2} .
\end{align*}
$$

We have to check that $I_{1}, I_{1}^{*}$ and $I_{2}$ have integral representations as in the right side of (5.11) with continuous kernels satisfying (5.12) and (5.13). In $I_{1}$

$$
\begin{aligned}
\int_{0}^{x} K(x, t) \cos z(x-t) d t & =\int_{0}^{x} K(x, x-\tau) \cos z \tau d \tau \\
\int_{0}^{x} K(x, t) \cos z(x+t) d t & =\int_{x}^{2 x} K(x, \tau-x) \cos z \tau d \tau
\end{aligned}
$$

i.e.,

$$
\begin{align*}
I_{1} & =\int_{0}^{2 x} \cos z \tau M_{1}(x, \tau) d \tau  \tag{5.16}\\
M_{1}(x, \tau) & = \begin{cases}K(x, x-\tau) & \text { if } 0 \leq \tau \leq x \\
-K(x, \tau-x) & \text { if } x \leq \tau \leq 2 x\end{cases}
\end{align*}
$$

The kernel $M_{1}$ is continuous since $K(x, 0)=0$ and the analogues of (5.12), (5.13) are also satisfied ((5.13) is trivial). In $I_{1}^{*}$ we argue similarly. In $I_{2}$ we change the order of integrations:

$$
\begin{align*}
\int_{0}^{x} & \int_{0}^{x} K(x, t) K^{*}(x, u) \cos z(t-u) d u d t  \tag{5.17}\\
& =\int_{0}^{x} \int_{t-x}^{t} K(x, t) K^{*}(x, t-\tau) \cos z \tau d \tau d t \\
= & \int_{0}^{x}\left(\int_{\tau}^{x} K(x, t) K^{*}(x, t-\tau) d t\right) \cos z \tau d \tau \\
& +\int_{0}^{x}\left(\int_{0}^{x-\tau} K(x, t) K^{*}(x, t+\tau) d t\right) \cos z \tau d \tau
\end{align*}
$$

and

$$
\begin{align*}
\int_{0}^{x} & \int_{0}^{x} K(x, t) K^{*}(x, u) \cos z(t+u) d u d t  \tag{5.18}\\
& =\int_{0}^{x} \int_{t}^{x+t} K(x, t) K^{*}(x, \tau-t) \cos z \tau d \tau d t \\
= & \int_{0}^{x}\left(\int_{0}^{\tau} K(x, t) K^{*}(x, \tau-t) d t\right) \cos z \tau d \tau \\
& +\int_{x}^{2 x}\left(\int_{\tau-x}^{x} K(x, t) K^{*}(x, \tau-t) d t\right) \cos z \tau d \tau
\end{align*}
$$

Consequently

$$
\begin{equation*}
I_{2}=\int_{0}^{2 x} \cos z \tau M_{2}(x, \tau) d \tau \tag{5.19}
\end{equation*}
$$

with the kernel

$$
M_{2}(x, \tau)= \begin{cases}\int_{\tau}^{x} K(x, t) K^{*}(x, t-\tau) d t &  \tag{5.20}\\ +\int_{0}^{x-\tau} K(x, t) K^{*}(x, t+\tau) d t & \\ -\int_{0}^{\tau} K(x, t) K^{*}(x, \tau-t) d t & \text { if } 0 \leq \tau \leq x \\ -\int_{\tau-x}^{x} K(x, t) K^{*}(x, \tau-t) d t & \text { if } x \leq \tau \leq 2 x\end{cases}
$$

continuous also at $\tau=x$ by definition. Here (5.12) and (5.13) also follow from Lemma 5.1. In order to complete the proof of (5.19) we have to find an integral representation of $2 \mu^{2} v\left(\pi-x, z^{2}\right) v^{*}\left(\pi-x, z^{2}\right)$. We apply the identities

$$
\begin{equation*}
\int_{0}^{t}(t-\tau) \cos z \tau d \tau=\frac{1-\cos z t}{z^{2}} \tag{5.21}
\end{equation*}
$$

and

$$
\begin{aligned}
v(\pi- & \left.x, z^{2}\right) v^{*}\left(\pi-x, z^{2}\right)=\frac{1-\cos 2 z x}{2 z^{2}} \\
& +\int_{0}^{x} K(x, t) \frac{\cos z(x-t)-\cos z(x+t)}{2 z^{2}} d t \\
& +\int_{0}^{x} K^{*}(x, t) \frac{\cos z(x-t)-\cos z(x+t)}{2 z^{2}} d t \\
& +\int_{0}^{x} \int_{0}^{x} K(x, t) K^{*}(x, u) \frac{\cos z(t-u)-\cos z(t+u)}{2 z^{2}} d u d t \\
= & \frac{1}{2} I_{3}+I_{4}+I_{4}^{*}+I_{5} .
\end{aligned}
$$

Now

$$
I_{3}=\int_{0}^{2 x}(2 x-\tau) \cos z \tau d \tau
$$

and in $I_{4}$

$$
\begin{array}{rl}
\int_{0}^{x} K & K(x, t) \frac{1-\cos z(x+t)}{z^{2}} d t \\
& =\int_{0}^{x} K(x, t) \int_{0}^{x+t}(x+t-\tau) \cos z \tau d \tau d t \\
& =\int_{0}^{2 x} \cos z \tau\left(\int_{\max (0, \tau-x)}^{x} K(x, t)(x+t-\tau) d t\right) d \tau, \\
\int_{0}^{x} K & K(x, t) \frac{1-\cos z(x-t)}{z^{2}} d t \\
& =\int_{0}^{x} \cos z \tau\left(\int_{0}^{x-\tau} K(x, t)(x-t-\tau) d t\right) d \tau .
\end{array}
$$

The kernel arising here is zero at $\tau=x$, so it can be continuously extended to $x \leq \tau \leq 2 x$. This proves an appropriate integral representation for $I_{4}$. The case of $I_{4}^{*}$ is similar. Finally in $I_{5}$ we get by twofold interchange of integrations

$$
\begin{aligned}
& \int_{0}^{x} \int_{0}^{x} K(x, t) K^{*}(x, u) \frac{1-\cos z(t+u)}{z^{2}} d u d t \\
&=\int_{0}^{2 x}\left(\int_{\max (\tau-x, 0)}^{x} K(x, t) \int_{\max (\tau-t, 0)}^{x} K^{*}(x, u)(t+u-\tau) d u d t\right) \\
& \quad \cdot \cos z \tau d \tau \\
& \int_{0}^{x} \int_{0}^{x} K(x, t) K^{*}(x, u) \frac{1-\cos z(t-u)}{z^{2}} d u d t \\
&=\int_{0}^{x} \cos z \tau\left(\int_{\tau}^{x} K(x, t) \int_{0}^{t-\tau} K^{*}(x, u)(t-u-\tau) d u d t\right) d \tau \\
& \quad+\int_{0}^{x} \cos z \tau\left(\int_{0}^{x-\tau} K(x, t) \int_{t+\tau}^{x} K^{*}(x, u)(u-t-\tau) d u d t\right) d \tau
\end{aligned}
$$

Since the last two kernels can be continuously extended by zero to the domain $x \leq \tau \leq 2 x$ and the analogue of (5.12), (5.13) is again a trivial corollary of (5.2) and (5.3), the proof of Lemma 5.3 is complete.

A similar statement holds for $\sin \beta \neq 0$ :
Lemma 5.3. Let $\sin \beta \neq 0,1 \leq p$ and $q, q^{*} \in L_{p}(0, \pi)$; then for $z \in \mathbf{C}$,

$$
\begin{align*}
v\left(\pi-x, z^{2}\right) v^{*}\left(\pi-x, z^{2}\right)-1 / 2 \sin ^{2} \beta= & 1 / 2 \sin ^{2} \beta \cos 2 z x  \tag{5.22}\\
& +\int_{0}^{2 x} L(x, t) \cos z t d t
\end{align*}
$$

with a kernel $L(x, t)$ continuous in $(x, t)$. Further if $q^{* *} \in L_{p}$ and $\|q\|_{p},\left\|q^{*}\right\|_{p}$, $\left\|q^{* *}\right\|_{p} \leq D$ then

$$
\begin{gather*}
\left|L\left(x, t, q, q^{*}\right)\right| \leq c(D, p),  \tag{5.23}\\
\left|L\left(x, t, q, q^{* *}\right)-L\left(x, t, q, q^{*}\right)\right| \leq c(D, p)\left\|q^{* *}-q^{*}\right\|_{p} . \tag{5.24}
\end{gather*}
$$

Proof. From Lemma 5.1 we know that

$$
\begin{equation*}
v\left(\pi-x, z^{2}\right)=\sin \beta \cos z x+\int_{0}^{x} N(x, t) \cos z t d t \tag{5.25}
\end{equation*}
$$

with a continuous kernel $N$ satisfying the analogue of (5.2), (5.3). Indeed, if we define the kernel $K_{1}$ for the potential $q(\pi-x)$ then

$$
N(x, t)=\sin \beta \frac{K_{1}(x, t)+K_{1}(x,-t)}{2}+\cos \beta \frac{K_{1}(x, t)-K_{1}(x,-t)}{2} .
$$

Now

$$
\begin{aligned}
v(\pi-x, & \left.z^{2}\right) v^{*}\left(\pi-x, z^{2}\right)-1 / 2 \sin ^{2} \beta \\
& =1 / 2 \sin ^{2} \beta \cos 2 z x+I_{1} \sin \beta+I_{1}^{*} \sin \beta+2 I_{2} \\
I_{1} & =\int_{0}^{x} N(x, t)(\cos z(x-t)+\cos z(x+t)) d t \\
I_{1}^{*} & =\int_{0}^{x} N^{*}(x, t)(\cos z(x-t)+\cos z(x+t)) d t \\
I_{2} & =\int_{0}^{x} \int_{0}^{x} N(x, t) N^{*}(x, u)(\cos z(t-u)+\cos z(t+u)) d u d t
\end{aligned}
$$

As above we can check that

$$
\begin{gathered}
I_{1}=\int_{0}^{2 x} \cos z \tau L_{1}(x, \tau) d \tau, L_{1}(x, \tau)= \begin{cases}N(x, x-\tau) & \text { if } 0 \leq \tau \leq x \\
N(x, \tau-x) & \text { if } x \leq \tau \leq 2 x\end{cases} \\
I_{2}=\int_{0}^{2 x} \cos z \tau L_{2}(x, \tau) d \tau+\int_{0}^{x} \cos z \tau L_{3}(x, \tau) d \tau
\end{gathered}
$$

with

$$
\begin{aligned}
& L_{2}(x, \tau)=\int_{\max (\tau-x, 0)}^{\min (\tau, x)} N(x, t) N^{*}(x, \tau-t) d t \\
& L_{3}(x, \tau)=\int_{0}^{x-\tau} N(x, t) N^{*}(x, t+\tau) d t+\int_{\tau}^{x} N(x, t) N^{*}(x, t-\tau) d t
\end{aligned}
$$

Since $L_{3}$ can be continuously extended by zero to the domain $x \leq \tau \leq 2 x$ and (5.12), (5.13) follow from Lemma 5.1, the proof is complete.

After obvious modifications we can also prove
Lemma 5.3'. Let $\sin \beta_{1} \neq 0, \sin \beta_{2} \neq 0$ and $q, q^{*} \in L_{1}(0, \pi)$; then for $z \in \mathbf{C}$

$$
\begin{align*}
& v\left(\pi-x, z^{2}, \beta_{1}\right) v^{*}\left(\pi-x, z^{2}, \beta_{2}\right)-1 / 2 \sin \beta_{1} \sin \beta_{2}  \tag{5.26}\\
& \quad=1 / 2 \sin ^{2} \beta_{1} \sin \beta_{2} \cos 2 z x+\int_{0}^{2 x} \cos z t L\left(x, t, q, q^{*}, \beta_{1}, \beta_{2}\right) d t
\end{align*}
$$

with a kernel $L(x, t)$ continuous in $(x, t)$.
Our final auxiliary result is a connection between the closedness of exponential systems and that of cosine systems.

Lemma 5.4. Let $z_{n}, n \geq 1$, be arbitrary different complex numbers and let $d>0,1 \leq p \leq \infty$. The system $\left\{\cos z_{n} x: n \geq 1\right\}$ is closed in $L_{p}(0, d)$ if and only if the system $\left\{e^{ \pm i z_{n} x}: n \geq 1\right\}$ is closed in $L_{p}(-d, d)$. If in case $z_{n}=0$, then 1 and $x$ are chosen instead of $e^{ \pm i z_{n} x}$.

Proof. The only if part. If the cosine system is not closed in $L_{p}(0, d)$, then there exists $0 \neq \overline{h \in L_{p}}(0, d)$ with

$$
\begin{equation*}
\int_{0}^{d} h(x) \cos z_{n} x d x=0, \quad n \geq 1 \tag{5.27}
\end{equation*}
$$

Define $h(-x)=h(x)$; then (5.27) implies

$$
0=\int_{-d}^{d} h(x) \cos z_{n} x d x=\int_{-d}^{d} h(x) e^{ \pm i z_{n} x} d x
$$

and in case $z_{n}=0$ we also have $\int_{-d}^{d} h(x) x d x=0$. Consequently $\left\{e^{ \pm i z_{n} x}: n \geq 1\right\}$ is not closed in $L_{p}(-d, d)$.

The if part. If the exponential system is not closed then there exists a function $0 \neq h \in L_{p}(-d, d)$ with

$$
\begin{equation*}
0=\int_{-d}^{d} h(x) e^{ \pm i z_{n} x} d x, \quad n \geq 1 \tag{5.28}
\end{equation*}
$$

Then

$$
\begin{aligned}
& 0=\int_{-d}^{d} h(-x) e^{ \pm i z_{n} x} d x \\
& 0=\int_{-d}^{d}(h(x)+h(-x)) e^{ \pm i z_{n} x} d x=2 \int_{0}^{d}(h(x)+h(-x)) \cos z_{n} x d x
\end{aligned}
$$

and this proves that the cosine system is not closed unless $h$ is odd. Now if $h$ is odd, we get from (5.28) that

$$
0=\int_{0}^{d} h(x) \sin z_{n} x d x, \quad n \geq 1 .
$$

Integrating by parts gives

$$
0=z_{n} \int_{0}^{d} \cos z_{n} x\left(\int_{x}^{d} h\right) d x
$$

In other words, $0 \neq \int_{x}^{d} h \in L_{p}(0, d)$ is orthogonal to all functions $\cos z_{n} x$, $z_{n} \neq 0$. If $z_{n}=0$, then

$$
\begin{aligned}
0 & =\int_{-d}^{d} x h(x) d x=2 \int_{0}^{d} x h(x) d x=2 \int_{0}^{d}\left(\int_{x}^{d} h\right) d x \\
& =2 \int_{0}^{d} \cos z_{n} x\left(\int_{x}^{d} h\right) d x
\end{aligned}
$$

Thus the cosine system is not closed in $L_{p}(0, d)$, which was to be proved.

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## References

[1] V. Ambarzumian, Über eine Frage der Eigenwerttheorie, Zeitschrift für Physik 53 (1929), 690-695.
[2] S. A. Avdonin, On the question of Riesz bases of exponential functions in $L^{2}$ (in Russian), Vestnik Leningrad Univ. Ser. Mat. 13 (1974), 5-12.
[3] E. F. Beckenbach and R. Bellmann, Inequalities, Springer-Verlag, New York, 1965.
[4] Ch. Bennewitz, A proof of the local Borg-Marchenko theorem, Comm. Math. Phys. 218 (2001), 131-132.
[5] G. Borg, Eine Umkehrung der Sturm-Liouvilleschen Eigenwertaufgabe, Acta Math. 78 (1946), 1-96.
[6] $\quad$, Uniqueness theorems in the spectral theory of $y^{\prime \prime}+(\lambda-q(x)) y=0$, Proc. $11^{\text {th }}$ Scandinavian Congress of Mathematicians (Trondheim, 1949) (Oslo), Johan Grundt Tanums Forlag, 1952, pp. 276-287.
[7] F. Gesztesy, R. del Rio, and B. Simon, Inverse spectral analysis with partial information on the potential, III. Updating boundary conditions, Internat. Math. Research Notices 15 (1997), 751-758.
[8] F. Gesztesy and B. Simon, A new approach to inverse spectral theory II. General real potentials and the connection to the spectral measure, Ann. of Math. 152 (2000), 593643.
[9] , Inverse spectral analysis with partial information on the potential, II. The case of discrete spectrum, Trans. Amer. Math. Soc. 352 (2000), 2765-2787.
[10] , On local Borg-Marchenko uniqueness results, Comm. Math. Phys. 211 (2000), 273-287.
[11] H. Hochstadt and B. Lieberman, An inverse Sturm-Liouville problem with mixed given data, SIAM J. Appl. Math. 34 (1978), 676-680.
[12] M. Horváth, On the inverse spectral theory of Schrödinger and Dirac operators, Trans. Amer. Math. Soc. 353 (2001), 4155-4171.
[13] M. Horváth, Inverse scattering with fixed energy and an inverse eigenvalue problem on the half-line, Trans. Amer. Math. Soc. (to appear).
[14] M. I. Kadec, The exact value of the Paley-Wiener constant, Sov. Math. Dokl. 5 (1964), 559-561.
[15] N. Levinson, Gap and Density Theorems, AMS Colloq. Publ., A. M. S., New York, 1940.
[16] ——, The inverse Sturm-Liouville problem, Mat. Tidsskr. B (1949), 25-30.
[17] B. M. Levitan and I. S. Sargsjan, Introduction to Spectral Theory (in Russian), Nauka, Moscow, 1970.
[18] V. A. Marchenko, Some questions in the theory of one-dimensional linear differential operators of the second order I (in Russian), Trudy Moskov Mat. Obsc. 1 (1952), 327420.
[19] , Sturm-Liouville Operators and Their Applications (in Russian), Naoukova Dumka, Kiev, 1977.
[20] B. Simon, A new approach to inverse spectral theory I. Fundamental formalism, Ann. of Math. 150 (1999), 1029-1057.
[21] R. M. Young, An Introduction to Nonharmonic Fourier Series, Academic Press, New York, 1980.
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